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Processes controlling the distribution of aerosol particles in the lower marine boundary layer during the First Aerosol Characterization Experiment (ACE 1)

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Abstract. The goals of the International Global Atmospheric Chemistry (IGAC) Program's First Aerosol Characterization Experiment (ACE 1) are to determine and understand the properties and controlling factors of the aerosol in the remote marine atmosphere that are relevant to radiative forcing and climate. A key question in terms of this goal and the overall biogeochemical sulfur cycle is what factors control the formation, growth, and evolution of particles in the marine boundary layer (MBL). To address this question, simultaneous measurements of dimethylsulfide (DMS), sulfur dioxide (SO₂), the aerosol chemical mass size distribution, and the aerosol number size distribution from 5 to 10,000 nm diameter were made on the National Oceanic and Atmospheric Administration (NOAA) ship *Discoverer*. From these data we conclude that the background MBL aerosol during ACE 1 often was composed of four distinct modes: an ultrafine (UF) mode ($D_p = 5\text{--}20$ nm), an Aitken mode ($D_p = 20\text{--}80$ nm), an accumulation mode ($D_p = 80\text{--}300$ nm), and a coarse mode ($D_p > 300$ nm). The presence of UF mode particles in the MBL could be explained by convective mixing between the free troposphere (FT) and the MBL associated with cloud pumping and subsidence following cold frontal passages. There was no evidence of major new particle production in the MBL. Oceanic emissions of DMS appeared to contribute to the growth of Aitken and accumulation mode particles. Coarse mode particles were comprised primarily of sea salt. Although these particles result from turbulence at the air-sea interface, the instantaneous wind speed accounted for only one third of the variance in the coarse mode number concentration in this region.

1. Introduction

Global aerosol size distributions of mass, number, and chemical composition are needed to calculate the direct and indirect radiative forcing due to aerosol particles.

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Given the short lifetimes of particles in the troposphere and their resulting spatial (vertical and horizontal) and temporal variability, it is not economically feasible to develop an observational network to map global aerosol distributions at the resolution needed for climate models. Accurate assessments of current radiative forcing due to aerosols must therefore rely on chemical transport models to generate the needed global aerosol distributions. Chemical transport models, developed and verified under present-day conditions, can then be used to generate global aerosol distributions using different emission scenarios and thus provide the needed input for prognostic analyses of future radiative forcing and climate response.

To generate accurate global aerosol distributions, chemical transport models must incorporate the physical and chemical processes that transform natural and anthropogenically derived emissions of aerosols and their precursors into the heterogeneously dispersed aerosol distributions that exist in the atmosphere. Although global chemical transport models [*Taylor and Penner,*

1994; Pham et al., 1995; Chin et al., 1996; Feichter et al., 1996; Kasibhatla et al., 1997] and coupled climate chemistry models [Chuang et al., 1997] continue to refine estimates of tropospheric aerosol distributions and the resulting radiative effects, reducing the uncertainties in these model estimates requires “more detailed and comprehensive data sets” [Kasibhatla et al., 1997, p. 3757] of aerosol properties and processes. These data sets can be acquired with a long-term global observation network and short-term intensive campaigns.

The spatial variability in aerosol properties and the interactions between the various chemical components require aerosol characterization and process studies in a globally representative range of natural and anthropogenically perturbed environments. Natural emissions and the resulting background aerosol are essential components of both chemical transport models and radiative forcing calculations. The reasons for this are several fold. First, the emissions of natural aerosols and their gaseous precursors can contribute significantly to the total aerosol burden over large areas of the Earth [Chin and Jacob, 1996]. Second, gaseous emissions from human sources can interact with the existing background aerosol to modify the aerosol mass and number distribution [i.e., Sievering et al., 1995; Chuang et al., 1997]. Third, since the different aerosol components interact with each other, the total aerosol distribution and resulting radiative effect are not necessarily the sum of the individual background and anthropogenic components. Fourth, radiative forcing is defined as the total

radiative effect due to aerosols minus the radiative effect due to background aerosols. Without a knowledge of the background aerosol properties, it is not possible to accurately quantify the direct and indirect forcing resulting from anthropogenically derived emissions.

The gas and aerosol measurements made aboard *Discoverer* during the First Aerosol Characterization Experiment (ACE 1) provide a means to characterize aerosol properties and assess the processes controlling the aerosol distribution in the lower marine boundary layer (MBL) south of Tasmania (Figure 1). In this minimally polluted marine environment the background atmospheric aerosol is composed largely of only two aerosol types, a sea-salt component and a non-sea-salt (nss) sulfate component. The primary source of both components is thought to be the sea surface with sea-salt particles being injected directly into the atmosphere and the nss sulfate component coming from the air-sea exchange of one primary gaseous precursor, dimethylsulfide (DMS). Less certain are the controlling processes that determine the aerosol size or spatial distribution at any given time. In this paper we focus on the processes controlling the physical and chemical aerosol distribution in the lower MBL. In a companion paper [Quinn et al., this issue] we discuss the effects of these distributions on the aerosol optical properties. Complementary descriptive aerosol data analyses from Macquarie Island [Brechtel et al., this issue; Kreidenweis et al., this issue] and the free troposphere [Clarke et al., this issue] are included in this special section.

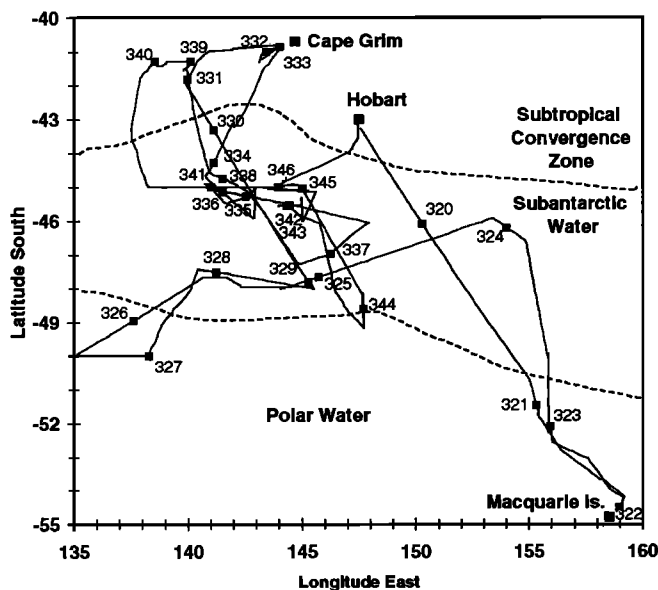


Figure 1. Discoverer cruise track during the intensive operations of ACE 1. The numbers (day of year in 1995) represent the ship's position at the beginning of each day (UTC). The three surface water masses are defined by salinity with subtropical convergence zone salinities >34.8 practical salinity unit (psu), Subantarctic water salinities >34.2 psu but <34.8 psu, and polar water salinities <34.2 psu.

2. Methods

Aerosol particles were sampled at 18 m above sea level through a heated mast [Bates et al., this issue]. The mast extended 6 m above the aerosol measurement container and was capped with a rotating cone-shaped inlet nozzle that was positioned into the relative wind. Air was pulled through this 5-cm-diameter inlet nozzle at $1 \text{ m}^3 \text{ min}^{-1}$ and down the 20-cm-diameter mast. The lower 1.5 m of the mast were heated to dry the aerosol to a relative humidity (RH) of 40 to 50%. Fifteen 1.9-cm-diameter conductive tubes extending into this heated zone were used to isokinetically subsample the air stream for the various aerosol sizing instruments and impactors at flows of 30 L min^{-1} . Comparisons of the total particle count ($D_p > 5 \text{ nm}$) during intercomparisons with the National Center for Atmospheric Research (NCAR) C-130 agreed to within 20%, suggesting minimal loss of particles in the inlet system.

The number size distribution between 5 and 600 nm was measured every 10 min with two differential mobility particle sizers (DMPS) at an RH of approximately 10%. The mobility distribution from the DMPSs was inverted to a number distribution by assuming that a Fuchs-Boltzman charge distribution resulted from the Kr^{85} charge neutralizer. The data were corrected for diffusional losses [Covert et al., 1997] and size depen-

dent counting efficiencies [Wiedensohler *et al.*, 1997] based on the pre-ACE 1 intercalibration exercises. The number distribution between 0.6 and 9.6 μm was measured with an aerodynamic particle sizer (APS) at an RH of approximately 40%. Data at diameters larger than 5 μm were discarded due to interferences from phantom counts and uncertainties in large particle collection efficiencies. The APS diameters were converted to geometric diameters by dividing by the square root of the particle density and dried to 10% RH assuming a sea-salt growth factor of 1.5 between 40 and 10% RH [Berg *et al.*, this issue]. An interactive routine was used to fit lognormal curves to the different modes of the number size distribution.

Aerosol chemical sampling was conducted with Berner-type multijet cascade impactors. Air was sampled only when the relative wind direction was forward of the beam, the relative wind speed was greater than 3 m s^{-1} , and the total particle count indicated the air was free of ship contamination. Further details of the sampling and ion chromatographic and gravimetric analysis are reported elsewhere [Quinn and Coffman, this issue]. The 50% aerodynamic cutoff diameters of the impactors [Quinn *et al.*, this issue] were converted to geometric diameters by dividing by the square root of the particle density and dried to 10% RH assuming an ammonium bisulfate growth factor of 1.2 [Berg *et al.*, this issue].

Air samples for DMS analysis were collected through a Teflon line which ran approximately 60 m from the top of the aerosol sampling mast to the analytical system. One hundred mL min^{-1} of the 4 L min^{-1} flow were pulled through a KI solution at the analytical system to eliminate oxidant interferences [Cooper and Saltzman, 1993]. The air sample volume ranged from 0.5 to 1.5 L depending on the DMS concentration. Seawater samples were collected from the ship's seawater pumping system at a depth of approximately 5 m. The seawater line ran to the analytical system where 5.1 mL of sample were valved into a Teflon gas stripper. The samples were purged with hydrogen at 80 mL min^{-1} for 5 min. Water vapor in either the air or purged seawater sample stream was removed by passing the flow through a -25°C Teflon tube filled with silanized glass wool. DMS was then trapped in a -25°C Teflon tube filled with Tenax. During the sample trapping period, 6.2 pmol of methylethyl sulfide (MES) were valved into the hydrogen stream as an internal standard. At the end of the sampling/purge period the coolant was pushed away from the trap, and the trap was electrically heated. DMS was desorbed onto a DB-1 megabore fused silica column where the sulfur compounds were separated isothermally at 50°C and quantified with a sulfur chemiluminescence detector. The system was calibrated using gravimetrically calibrated DMS and MES permeation tubes. The precision of the analysis, based on both replicate analyses of a single water sample and replicate analyses of a standard introduced at the inlet of the air sample line, was typically $\pm 8\%$. The detection limit during ACE 1

was approximately 0.8 pmol. The performance of the system was monitored regularly by running blanks and standards through the entire system. Values reported here have been corrected for recovery losses. Losses within the system were $<10\%$. System blanks were below detection limit. Air samples are reported in units of nmol m^{-3} at standard temperature (25°C) and pressure (1013 mbar) such that 1 nmol m^{-3} equals 24.5 ppt.

Air was sampled for SO_2 analysis from the ship's flying bridge, 18 m above sea level [Bates *et al.*, this issue]. The inlet was heated and connected to a nafion dryer to minimize losses. The gaseous SO_2 in dry air was absorbed into an aqueous scrubbing solution in a gas/liquid exchange coil, converted to sulfite, and then derivatized to isoindole. The derivative was separated from interferences by isocratic reverse phase HPLC and detected by fluorescence. Further details of the instrument are reported elsewhere [Saltzman *et al.*, 1993; De Bruyn *et al.*, this issue]. The system was calibrated using a two-stage gas-dilution system with a low-level permeation device. The precision of replicate measurements was typically better than 5%, and the detection limit during ACE 1 was approximately 3–5 parts per trillion by volume (pptv). The performance of the system was monitored regularly by running blanks and standards through the entire system. Losses within the system were $<15\%$. Values reported here have been corrected for recovery losses.

Additional measurements made aboard the ship included atmospheric temperature, pressure, and humidity, surface seawater temperature, salinity and nitrate concentrations, wind speed and direction, rainfall rates, solar radiation, atmospheric radon [Whittlestone and Zahorowski, this issue], and vertical profiles of temperature, dew point temperature, and wind speed and direction from radiosondes. All references to time are reported here in UTC which is 10 hours behind local standard time.

3. Results and Discussion

3.1. Characteristics of the MBL Aerosol

During the 29-day field experiment the ship encountered four episodes of continental air with high radon ($>100 \text{ mBq m}^{-3}$) and total particle concentrations (2000–4000 cm^{-3}). These were generally on the back side of high-pressure systems which brought flow off the Australian continent. During the remainder of the month the air masses sampled at the ship were typical of background conditions in the midlatitudes [Covert *et al.*, 1996a] with very low radon concentrations ($<100 \text{ mBq m}^{-3}$) and a mean total ($D_p > 5 \text{ nm}$) particle number concentration of 500 cm^{-3} (Figure 2).

The particle number size distributions, $dN/d\log D_p$, fell into three distinct categories and could be modeled with generally three and sometimes four modes (Figure 3). All the distributions had a similar coarse mode ($D_p = 300\text{--}5000 \text{ nm}$) comprised primarily of sea salt

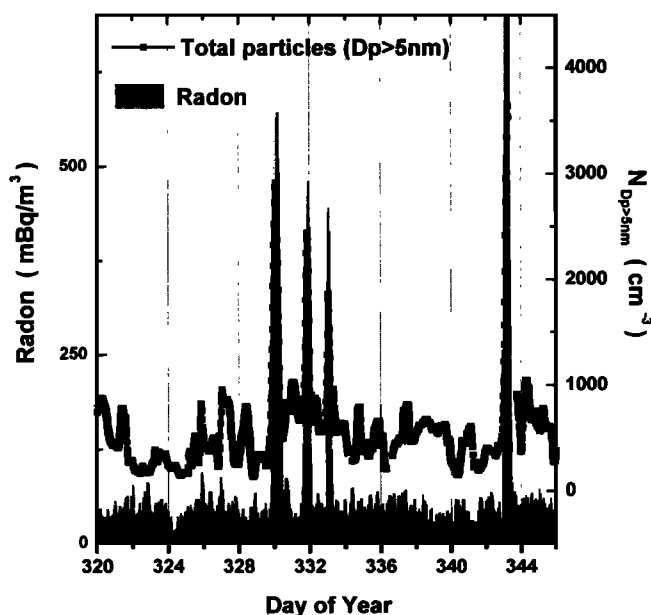


Figure 2. Total particle number concentration ($D_p > 5$ nm) and atmospheric radon concentrations measured aboard *Discoverer* as a function of day of year (DOY).

(Table 1). The upper size on this mode was limited by instrumental uncertainties and sample inlet losses. The background marine distributions generally had distinct Aitken ($D_p = 20\text{--}80$ nm) and accumulation modes ($D_p = 80\text{--}300$ nm) with a minimum between the modes (60–100 nm). The aerosol mass in this size range ($D_p < 450$ nm) was comprised primarily of nss sulfate, ammonium, methanesulfonate (MSA), and sea salt (Table 1). The sum of these ions ($D_p < 450$ nm) were, on average, $100 \pm 10\%$ of the total gravimetric mass of particles with $D_p < 450$ nm [Quinn and Coffman, this issue]. The average NH_4^+ to nss SO_4^- molar ratio in this size range was 0.72 ± 0.23 , suggesting only partly neutralized sulfuric acid particles. The presence of sea salt in this size range does not necessarily imply an additional sea-salt mode, but more likely represents the tail of the coarse mode extending into this size range [Clarke and Porter, 1993]. The presence of appreciable sea-salt mass in this size range has important implications for sea-salt aerosol lifetimes and the climatic effects of these particles [Quinn et al., this issue].

The second most frequent size distribution (Figure 3) also occurred in background marine conditions, but had a fourth mode ($D_p = 5\text{--}20$ nm) which we refer to here as the ultrafine (UF) mode. This smallest mode is a short-lived transition from newly nucleated particles of a few nanometers diameter to the more stable Aitken mode centered around 30–40 nm [Covert et al., 1996a, b; Wiedensohler et al., 1996]. The UF mode was encountered at the ship 13 times, for periods ranging from 0.5 to 36 hours, during the 29-day experiment (Figure 4). The third type of size distribution occurred in continental air masses (Figure 3). There were once again

three modes, although the number and mass concentrations in the Aitken and accumulation modes were much higher than in the background marine distributions and there was no clear concentration minimum between the modes.

3.2. Sources of Aerosol and Aerosol Precursors to the MBL

The ocean is a source of both sea-salt aerosol and nss sulfur to the MBL. The nss sulfur is emitted to the MBL primarily in the form of dimethylsulfide (DMS), a biologically produced trace gas [Bates et al., 1994]. During ACE 1, seawater DMS concentrations (Table 2), which drive the ocean to atmosphere DMS flux, were highest in the subtropical convergence zone west of Tasmania ($2.8 \text{ nM} \pm 1.2 \text{ nM}$). DMS concentrations in the subantarctic water (Figure 1) increased during the experiment. Initial concentrations ($1.1 \pm 0.3 \text{ nM}$) were very similar to the low concentrations observed in the polar water ($0.9 \pm 0.4 \text{ nM}$). During the latter part of the experiment, as the seasonal biological cycle progressed, Subantarctic water DMS concentrations ($2.4 \pm 1.0 \text{ nM}$) increased to levels similar to the subtropical convergence zone.

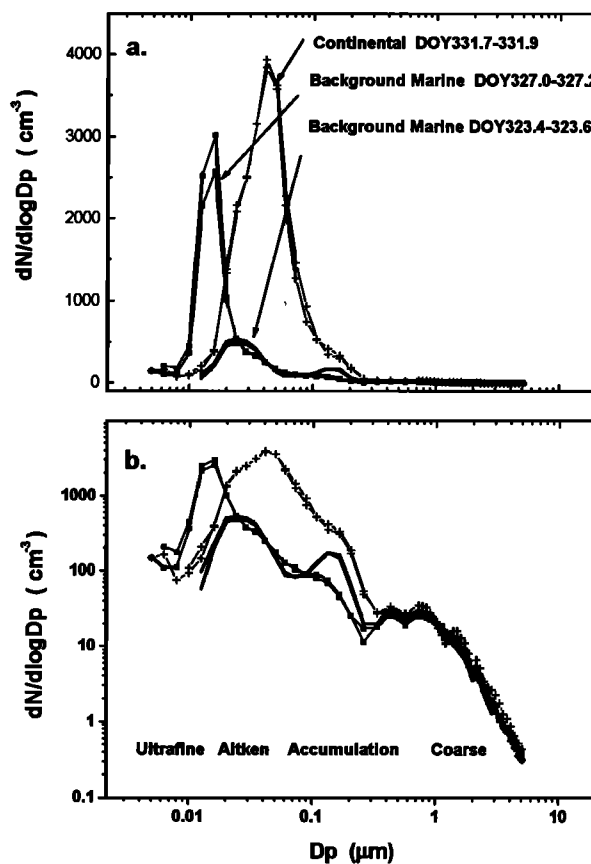


Figure 3. Typical number size distributions measured during ACE 1 showing the four distinct modes. (a) Linear and (b) logarithmic concentration plots are shown to present the entire range of data in perspective.

Table 1. Aerosol Number and Chemical Mass Characteristics Measured Aboard the NOAA Ship *Discoverer* During the ACE 1 Intensive (November 15 to December 14, 1995)

	Modes			
	Ultrafine	Aitken	Accumulation	Coarse
<i>Background Marine</i>				
N , cm^{-3}	190	210	74	15
D_{gn} , nm	16	33	110	540
σ_g	1.45	1.40	1.41	2.02
Sea salt, ng m^{-3}			280	9,000
Range, ng m^{-3}			130–610	3,000–20,000
nss $\text{SO}_4^- + \text{NH}_4^+$, ng m^{-3}			170	44
Range, ng m^{-3}			93–320	0–130
MSA, ng m^{-3}			31	23
Range, ng m^{-3}			10–180	5–67
<i>Continental</i>				
N , cm^{-3}	500	720	75	23
D_{gn} , nm	18	37	130	530
σ_g	1.42	1.50	1.32	2.08

The mean number concentration, number mean geometric diameter at 10% RH, and standard deviation of the averaged data were determined from a log normal fit to the modes. The chemical concentrations are given as <450 nm geometric diameter at 10% RH (first three impactor stages [Quinn *et al.*, this issue] and >450 nm mass concentrations.

The flux of DMS to the atmosphere is calculated from the seawater DMS concentration and wind speed. Two wind speed transfer velocity relationships [Liss and Merlivat, 1986; Wanninkhof, 1992] currently used to calculate air-sea gas exchange result in fluxes that vary by roughly a factor of 2 (Table 2). Using these two re-

lationships to estimate the range of DMS fluxes in the ACE 1 region results in calculated median DMS fluxes of 2.3 to 4.1 $\mu\text{mol m}^{-2} \text{d}^{-1}$ in the polar waters, 3.0 to 7.3 $\mu\text{mol m}^{-2} \text{d}^{-1}$ in the Subantarctic water, and 5.4 to 8.9 $\mu\text{mol m}^{-2} \text{d}^{-1}$ in the subtropical convergence zone (Table 2). The ocean-atmosphere flux of DMS in the

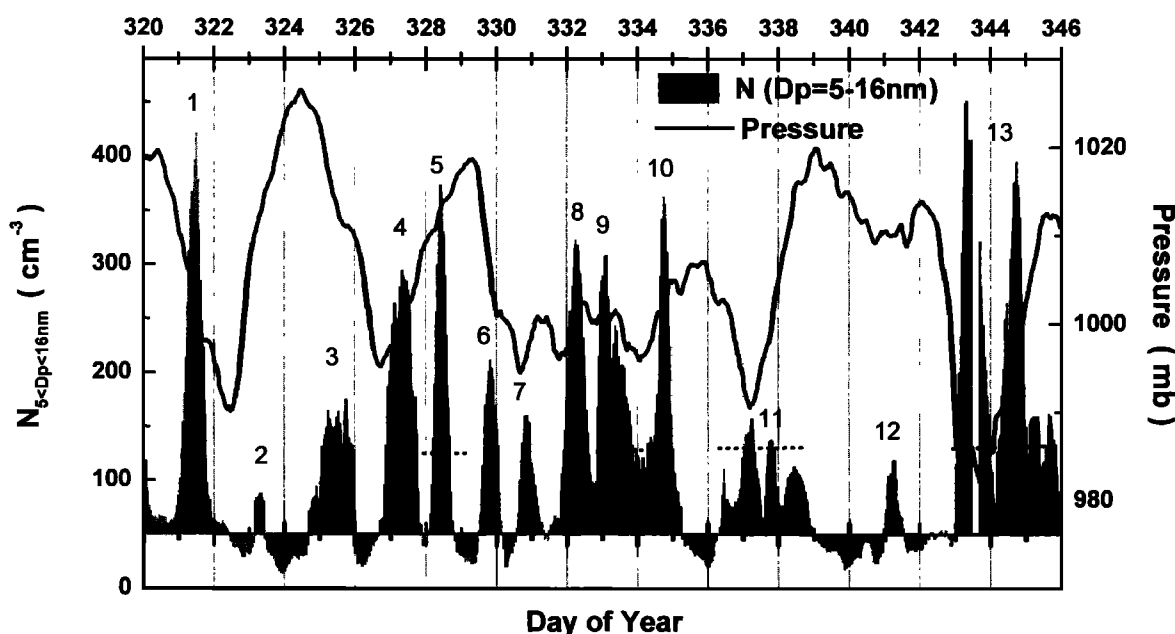


Figure 4. Total number concentration of particles with diameters between 5 and 16 nm plotted with atmospheric pressure as a function of DOY. A particle threshold of 50 cm^{-3} is used to more clearly delineate the 13 UF particle events.

Table 2. Measured Seawater and Atmospheric DMS Concentrations Associated With the Three Surface Water Masses in the ACE 1 Study Area (Figure 1)

	Seawater DMS, nM	Atmospheric DMS, nmol m ⁻³	DMS Flux [Liss and Merlivat, 1986], μmol m ⁻² d ⁻¹	DMS Flux [Wanninkhof, 1992], μmol m ⁻² d ⁻¹
Polar water				
Median	0.79	3.3	2.3	4.1
Average	0.90	4.1	3.4	6.1
Standard deviation (1σ)	0.39	2.3	2.7	5.2
Range	0.5–2.9	1.7–9.7	0–11	0.3–21
Subantarctic water (DOY 320–338)				
Median	1.09	3.0	3.0	5.2
Average	1.10	3.3	3.0	5.6
Standard Deviation (1σ)	0.31	1.4	2.0	3.9
Range	0.4–2.4	1.0–9.2	0–9.7	0–19
Subantarctic water (DOY 338–346)				
Median	2.24	5.6	4.5	7.3
Average	2.43	6.0	6.6	11.6
Standard Deviation (1σ)	1.00	2.9	7.3	13.7
Range	0.9–5.6	1.0–14.2	0–46	0–85
Subtropical Convergence Zone				
Median	2.60	4.9	5.4	8.9
Average	2.84	5.5	6.9	12.6
Standard Deviation (1σ)	1.21	2.6	5.7	10.8
Range	0.8–6.8	1.7–14.3	0–33	0–61

DMS fluxes were calculated for each seawater measurement, using the in-situ wind speed and the wind speed transfer velocity relationships of both Liss and Merlivat [1986] and Wanninkhof [1992].

polar waters is typical of winter fluxes at this latitude, while the fluxes calculated for the Subantarctic waters late in the experiment and the subtropical convergence zone waters during the entire ACE 1 were more typical of summer conditions [Bates *et al.*, 1987].

The ocean-atmosphere flux of DMS during ACE 1 maintained an atmospheric DMS concentration of 1 to 14 nmol m⁻³ (Table 2) with an overall mean concentration of 4.5 ± 2.5 nmol m⁻³. This concentration is nearly identical to the mean December DMS concentration (4.9 nmol m⁻³) measured at Cape Grim from 1988 to 1992 [Ayers *et al.*, 1995]. DMS in the MBL has two major sinks, transport to the free troposphere (FT) and oxidation in the MBL by OH [Ayers *et al.*, 1995; Bandy *et al.*, 1996; Yvon *et al.*, 1996]. The oxidation of DMS by OH proceeds by two pathways and can yield a variety of gas phase products including dimethylsulfoxide, dimethylsulfone, SO₂, MSA, and sulfuric acid [Berresheim *et al.*, 1995]. SO₂, the only gas phase DMS oxidation product measured aboard *Discoverer*, was present at an average concentration of 0.5 ± 0.3 nmol m⁻³. Some fraction of the gas phase DMS oxidation products will nucleate new particles or condense/react onto existing particles in the MBL depending primarily on the existing particle surface area [Covert *et al.*, 1996a].

The ocean is a source of sea-salt particles to the atmosphere through the bursting of air bubbles associated with breaking waves (film and jet droplets) and the mechanical tearing (spume droplets) and spilling

over (splash droplets) of wave crests (reviewed by Andreas *et al.* [1995]). While the total number concentration of these mechanically produced sea-salt particles is relatively small even in the remote MBL (Figure 2 and Table 1), they can dominate the mass size distribution (Table 1 and Quinn *et al.* [1996]) and thus have a significant effect on chemical reactions occurring in the MBL [Sievering *et al.*, 1992], the formation of new particles [Covert *et al.*, 1996a, b], the production of cloud condensation nuclei [O'Dowd *et al.*, 1997], and the light scattering by aerosols [Quinn *et al.*, 1996; this issue]. The magnitude of the ocean-atmosphere flux of sea-salt particles is a function of the wind-dependent sea state [Monahan *et al.*, 1986], while the instantaneous atmospheric sea-salt distribution is a function of the ocean-atmosphere flux, mixing between the MBL and FT, size-dependent removal processes, and advection.

The FT also can be a source of aerosol and aerosol precursors to the MBL. Long-range transport of SO₂ from the continents can bring anthropogenically and volcanically derived SO₂ to remote regions where it can be entrained into the MBL [Andreae *et al.*, 1988; Thornton *et al.*, 1996]. This did not appear to be the case during ACE 1 as SO₂ concentrations generally decreased with altitude above the MBL (A. Bandy, personal communication, 1997). The low concentrations of SO₂ in the FT and the presence of DMS in the FT (A. Bandy, personal communication, 1997) suggest a MBL source for the FT sulfur in this region.

Aerosol concentrations in the remote FT can also be

high due to long-range transport of continental aerosols [Clarke *et al.*, 1996] and the in situ production of UF particles [Hegg *et al.*, 1990; Clarke, 1993; Perry and Hobbs, 1994]. Aerosol distributions, soot concentrations, and volatility measurements in the FT during ACE 1 [Clarke *et al.*, this issue] indicate a low concentration of continental aerosol. The FT therefore appears to be a minimal source of aerosol mass to the MBL. However, high concentrations of newly formed

UF particles were often observed in the FT in the vicinity of both low and high level clouds [Clarke *et al.*, this issue]. The large number concentrations of these UF particles could have a major impact on the MBL particle concentrations during periods of strong subsidence and entrainment from the FT. Modeling studies [Raes and Van Dingenen, 1992; Raes, 1995] and previous observations [Clarke, 1993; Covert *et al.*, 1996a, b; Bigg *et al.*, 1996] have shown that this process can be a ma-

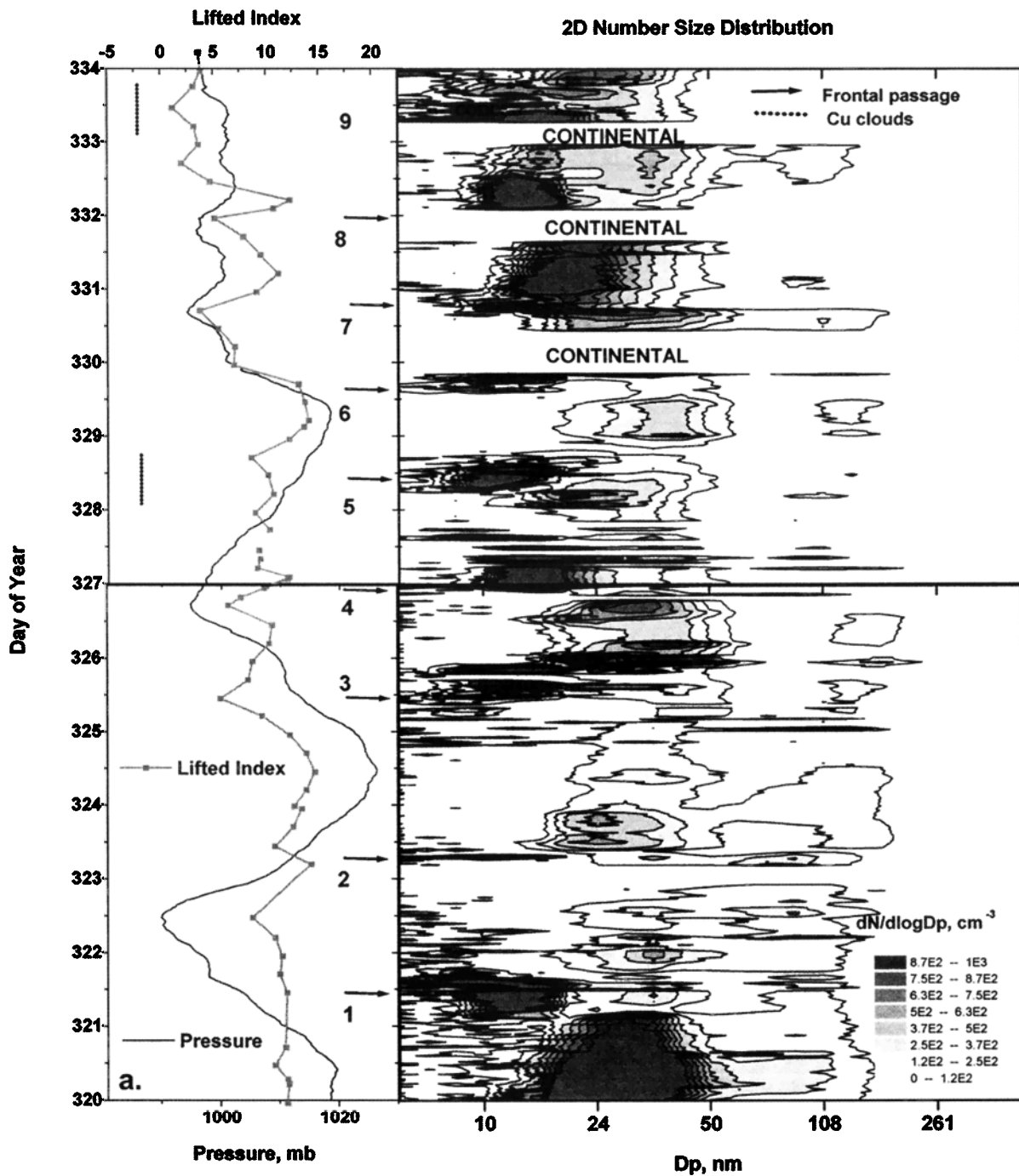


Figure 5. Contour plots of the number size distributions measured during ACE 1 during background marine conditions. The 13 UF particle events shown in Figure 4 are labeled. Each event can be attributed to either the passage of a cold front (arrow) or a cumulus cloud field (dots). Also plotted are atmospheric pressure and the lifted index (LI) which was generated from the shipboard radiosonde data. (a) DOY 320–334, (b) DOY 334–347.

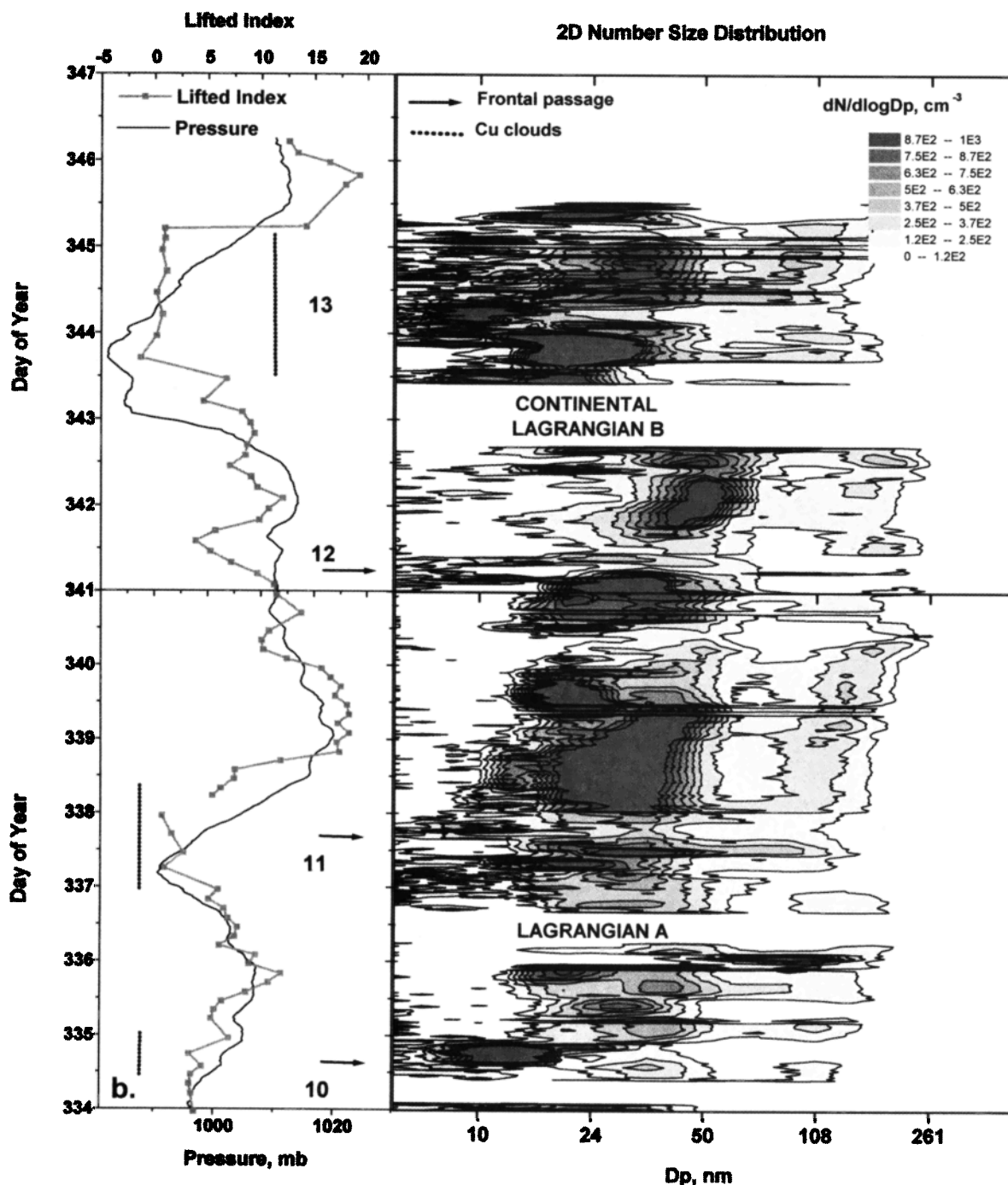


Figure 5. (continued)

major source of new submicrometer aerosol to the remote MBL.

3.3. Processes Controlling the MBL Aerosol

High concentrations (>50 particles cm^{-3}) of UF mode particles ($D_p = 5\text{--}20$ nm) were observed aboard *Discoverer* approximately 50% of the time during ACE 1, occurring as 13 separate events (Figures 4 and 5). Measurements during ACE 1 [Clarke *et al.*, this issue; Weber *et al.*, this issue], previous studies [Shaw, 1989; Covert *et al.*, 1992; Weber *et al.*, 1995, 1996, 1997], and model calculations [Hegg *et al.*, 1992; Russell *et al.*, 1994; Raes,

1995] have shown that UF mode particle production is critically dependent upon H_2SO_4 concentrations and preexisting aerosol surface area concentrations. Model calculations using typical MBL conditions suggest that preexisting surface areas must be less than $5 \mu\text{m}^2 \text{cm}^{-3}$ [Hegg *et al.*, 1992] or $25 \mu\text{m}^2 \text{cm}^{-3}$ [Raes, 1995] for in situ particle formation to occur. During ACE 1 there was no correlation between the concentration of the H_2SO_4 precursors (DMS and SO_2) measured aboard *Discoverer* and the observed UF mode particle concentrations. Furthermore, the average preexisting aerosol surface area measured 18 m above the sea surface dur-

ing the 13 UF events ranged from 40 to $130 \mu\text{m}^2 \text{cm}^{-3}$ at ambient RH. This surface area concentration would likely suppress new particle production and high UF mode particle concentrations by scavenging condensable vapors, clusters, and newly formed particles [Shaw, 1989; Covert *et al.*, 1992; Raes, 1995; Weber *et al.*, 1995].

Conditions in the FT and the cloud-topped boundary between the MBL and FT are often more favorable for new particle production than the conditions which generally exist in the MBL [Hegg *et al.*, 1990; Raes and Van Dingenen, 1992; Clarke, 1993; Hoppel *et al.*, 1994a; Raes, 1995; Covert *et al.*, 1996a, b]. The UF particles produced aloft can be entrained into the MBL. During ACE 1, most of the observed UF events (1–8 and 10–12, Figures 4 and 5) coincided with the passage of cold fronts and the postfrontal subsidence associated with them [Bond and Fleagle, 1988]. For example (Figure 6), the prefrontal air mass sampled at the ship at the end of day 331 had passed over the Australian continent picking up radon. The passage of the front, which began with a brief rain event, was marked by a change in wind

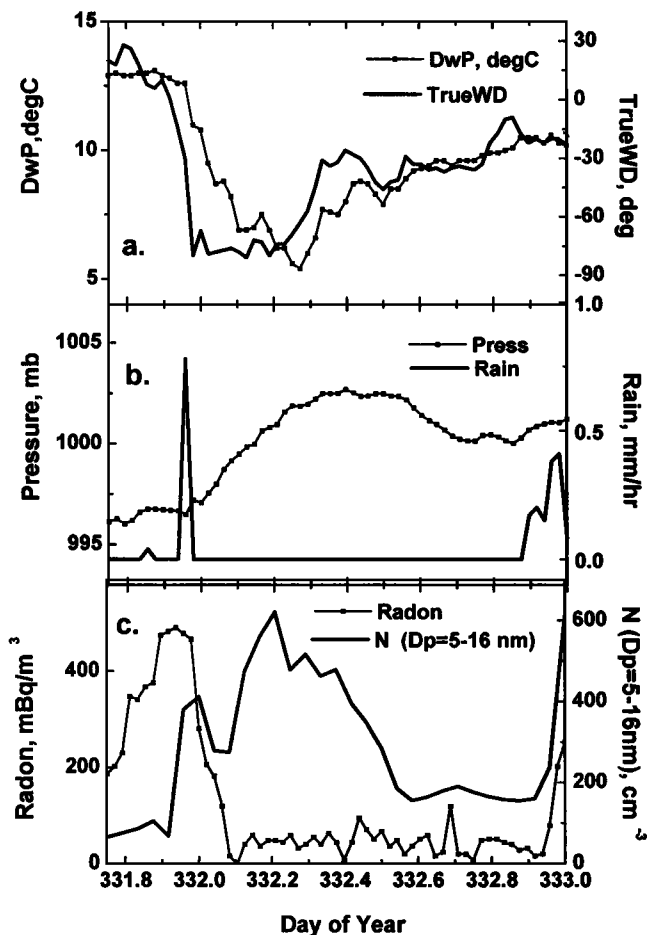


Figure 6. Aerosol and meteorological data for a frontal passage on DOY 332. (a) True wind direction and dew point temperature, (b) atmospheric pressure and rainfall rate, (c) radon concentrations and number of particles with diameters between 5 and 16 nm.

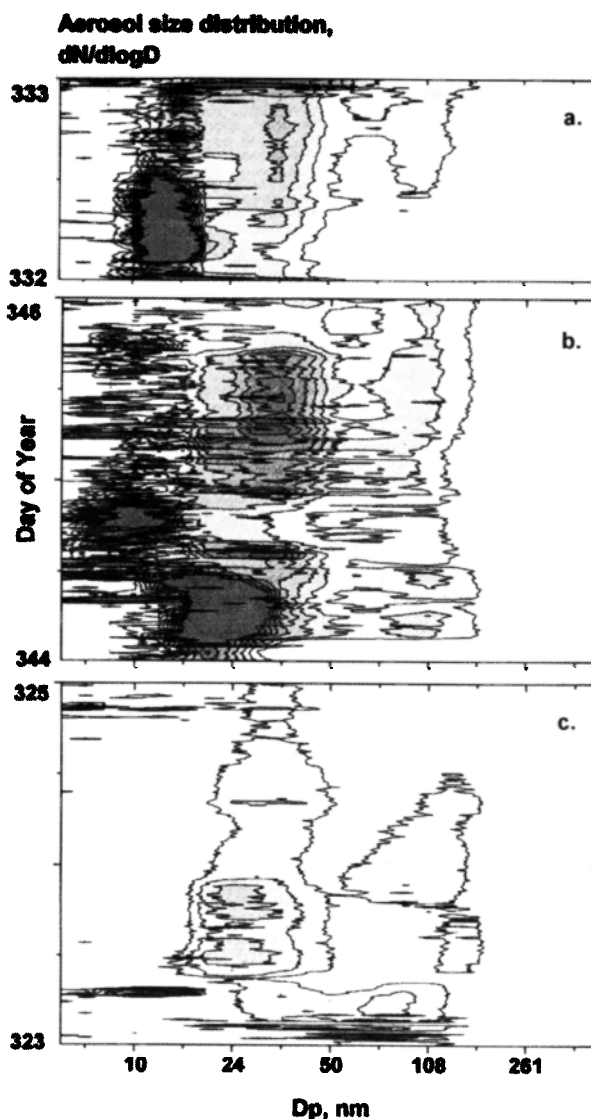


Figure 7. Contour plots of the number size distribution for (a) the frontal passage on DOY 332, (b) the cumulus cloud mixing period during DOY 344–346, and (c) the stratus cloud period during DOY 323–325. The scale is the same as in Figure 5.

direction and depression in dew point temperature. The air mass behind the front had not been in recent contact with the continent based on low radon values. The aerosol surface area during this period (day of year (DOY) 331–333) was always $>84 \mu\text{m}^2 \text{cm}^{-3}$, and thus was not conducive to new particle production [Clarke *et al.*, this issue]. However, UF particles with diameters between 10 and 20 nm appeared in this air mass after the frontal passage (Figures 6 and 7a). During DOY 332, measurements aboard the NCAR-130 revealed layers of high concentrations (2000 to 4000cm^{-3}) of small UF particles (narrow mode centered at about 12 nm) in the broken convective cloud region behind the front [Clarke *et al.*, this issue]. The layers occurred between 2 and 4 km and were associated with cloud outflows. In

the cloud-free region of postfrontal subsidence, elevated concentrations of UF particles in the 10 to 20 nm size range extended down to the MBL inversion [Clarke *et al.*, this issue]. On the basis of these data the most likely source of the UF particles measured aboard *Discoverer* in the lower MBL was entrainment from aloft. Although not all the frontal passages and UF events are as clear as this example, many show similar features that support the hypothesis that UF particles are produced in the FT or upper MBL and mixed down into the lower MBL.

Other UF events (5, 9, 10, 11, and 13, Figure 5), often with higher concentrations of particles <10 nm diameter and more variable concentrations of UF mode particles (Figure 7b), appeared to coincide with regions of active cumulus clouds [Cotton *et al.*, 1995]. The degree of convective activity was estimated from the lifted index (LI) which was calculated from the shipboard radiosonde data [Galway, 1956]. Mathematically, $LI = T_{500} - T_L$, where T_{500} is the in-situ temperature at

500 mbar and T_L is the resultant temperature after a moist adiabatic lifting of an air parcel within the MBL to 500 mbar. Lower values of LI, indicating greater instability, coincided with UF events 9, 10, 11, and 13 (Figure 5). Event 13, for example, occurred in a field of broken cumulus clouds. The temperature/dew point sounding from *Discoverer* on DOY 344 (Figure 8b) showed a capping inversion at 350 mbar and a relatively deep boundary layer extending to 700 mbar (-20°C). Advanced very high resolution radiometer (AVHRR) satellite data showed that cloud tops frequently extended above the boundary layer (cloud top temperatures $< -20^{\circ}\text{C}$). These active cumulus clouds vent gases from the MBL to the free troposphere where conditions favor the formation of new particles [Perry and Hobbs, 1994; Raes, 1995; Clarke *et al.*, this issue]. The photochemical production of new particles in the cloud processes air was observed in this region on DOY 334 [Clarke *et al.*, this issue]. To balance the cumulus updrafts, air can descend in the cloud-free air between clouds [Cotton *et al.*, 1995], providing a transport mechanism to bring UF particles to the MBL. A contrasting case occurred on DOY 323 where stratus clouds with cloud-top temperatures of -5°C capped the well-defined boundary layer at 850 mbar (Figure 8a). Here there was little exchange with the FT and no evidence of UF particles (Figure 7c).

MBL Aitken mode particles originate from the condensational growth and coagulation of UF mode particles [Hoppel *et al.*, 1990; Weber and McMurry, 1996] and from entrainment from aloft [Clarke *et al.*, 1996]. As Aitken mode particles are mixed within the MBL, they are "processed" by clouds and grow by liquid phase conversion of trace gases to nonvolatile solute compounds and by coagulation with other Aitken and UF mode particles [Hoppel and Frick, 1990]. On the basis of the water supersaturation level in the cloud, some particles will eventually grow sufficiently large to activate and form cloud droplets. When the cloud dissipates, the interstitial (unactivated) particles and residual (from evaporated cloud droplets) particles remain. This process is thought to differentiate the Aitken and accumulation modes with the MBL cloud supersaturation and aerosol chemistry defining the diameter of the intermode minimum [Hoppel *et al.*, 1986, 1994a, b]. On average, particles will be cycled through many non-precipitating clouds before being removed by rain [Hoppel and Frick, 1990]. Thus cloud processing will occur multiple times to create and magnify this bimodal feature.

Clouds were a dominant meteorological feature in the ACE 1 study area. Frontal passages and subsidence around venting clouds introduced Aitken mode particles and their UF mode precursors to the MBL. However, it is not possible to directly attribute any increase in Aitken mode number to the frequent high concentrations of UF mode particles (Figure 4) since the shipboard observations are not Lagrangian. The frequent

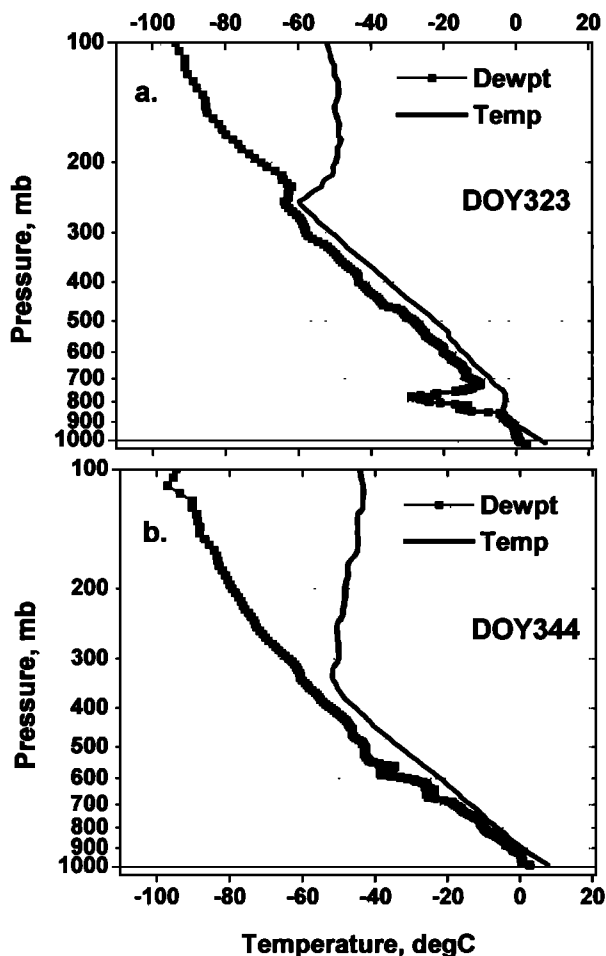


Figure 8. Vertical soundings of dry bulb and dew point temperatures from radiosondes launched from *Discoverer* on (a) DOY 323 at 1700 where *Discoverer* was located under a stratus cloud deck and (b) DOY 344 at 1100 where *Discoverer* was located in a region of cumulus clouds.

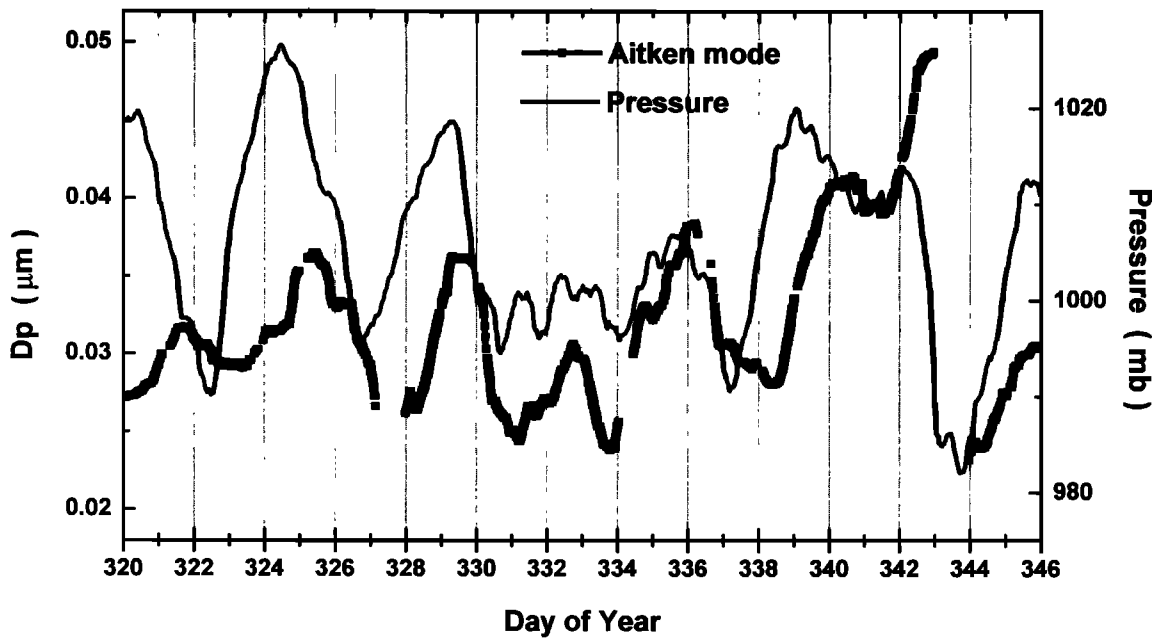


Figure 9. Twenty-four hour geometric, number mean diameter, D_{gn} , of the Aitken mode and atmospheric pressure as a function of DOY.

abrupt changes in the aerosol size distribution (Figure 5) are also evidence that the system was rarely in steady state. Still, there are several re-occurring features in the data. In general, the diameter of the Aitken mode increased with time in high-pressure systems (Figure 9), which is consistent with increased condensational growth in the MBL and reduced vertical entrainment through a well-capped boundary layer. Calculated isokinetic back trajectories suggest that the air mass stayed in the MBL (RH 60–70%) for several days. In contrast, the Aitken mode diameter either decreased or was low during periods of frontal passages and convective subsidence, presumably due to mixing of smaller, less aged particles from the FT. Elevated concentrations of UF particles and lower concentrations of Aitken mode particles existed just above the MBL in these regions [Clarke *et al.*, this issue].

The period from DOY 337 to DOY 341 is shown in more detail in Figure 10. The frontal passage with associated rainfall occurred at the beginning of DOY 337 (Figure 10a). The trough following the front generated a large area of cumuliform clouds which brought UF particles to the MBL through DOY 338 (Figure 4). The photochemically generated SO_2 [De Bruyn *et al.*, this issue] provided a source of sulfur for both condensational particle growth and within-cloud liquid phase growth (Figure 10b). During this period the average Aitken and accumulation mode diameters remained constant (Figure 10c), presumably due to a balance between the effects of mixing from the FT and growth within the MBL. As the high-pressure system passed, a field of broken stratocumulus clouds capped the boundary layer at 880 mbar. During DOY 339 and 340 the Aitken mode

diameter appeared to grow during the day, while the mean diameter of the accumulation mode constantly increased. A rain event (Figure 10a) may have been responsible for the decrease in mean modal diameters at DOY 340.5. Again, these observations did not follow a specific air mass, and thus the observed changes could be simply due to advection. Still, the data are consistent with the cloud processing described above.

A key question in this study is whether the emissions of DMS from the ocean can quantitatively account for the nss sulfate in marine aerosols. The frequent frontal passages and mixing between the MBL and FT complicate a regionally averaged assessment of this question. Given the average atmospheric DMS concentration (Table 2), an average 1 km well-mixed boundary layer, and the calculated ocean-atmosphere DMS fluxes (Table 2), the lifetime of DMS in the MBL in this region is approximately 1 day. Photochemical model calculations indicate that 30–50% of this DMS is converted to SO_2 [De Bruyn *et al.*, this issue]. The diurnal cycles of SO_2 during DOY 337 to DOY 341 (Figure 10b) are evidence of a daytime photolytic production coupled with a continuous loss to aerosol and surface dry deposition. Using a regional average OH concentration of $6 \times 10^{-5} \text{ mol cm}^{-3}$ [Spivakovsky *et al.*, 1990] and an OH + SO_2 reaction rate constant of 9.5×10^{-13} [DeMore *et al.*, 1992], we can estimate the mean photochemical production rate of H_2SO_4 from observed SO_2 concentrations. The average photochemical production rate of H_2SO_4 for the period DOY 339 to DOY 341 is $260 \text{ mol cm}^{-3} \text{ s}^{-1}$. Using the observed aerosol surface areas, we can then estimate the average steady state H_2SO_4 concentration and the flux of H_2SO_4 to each aerosol mode.

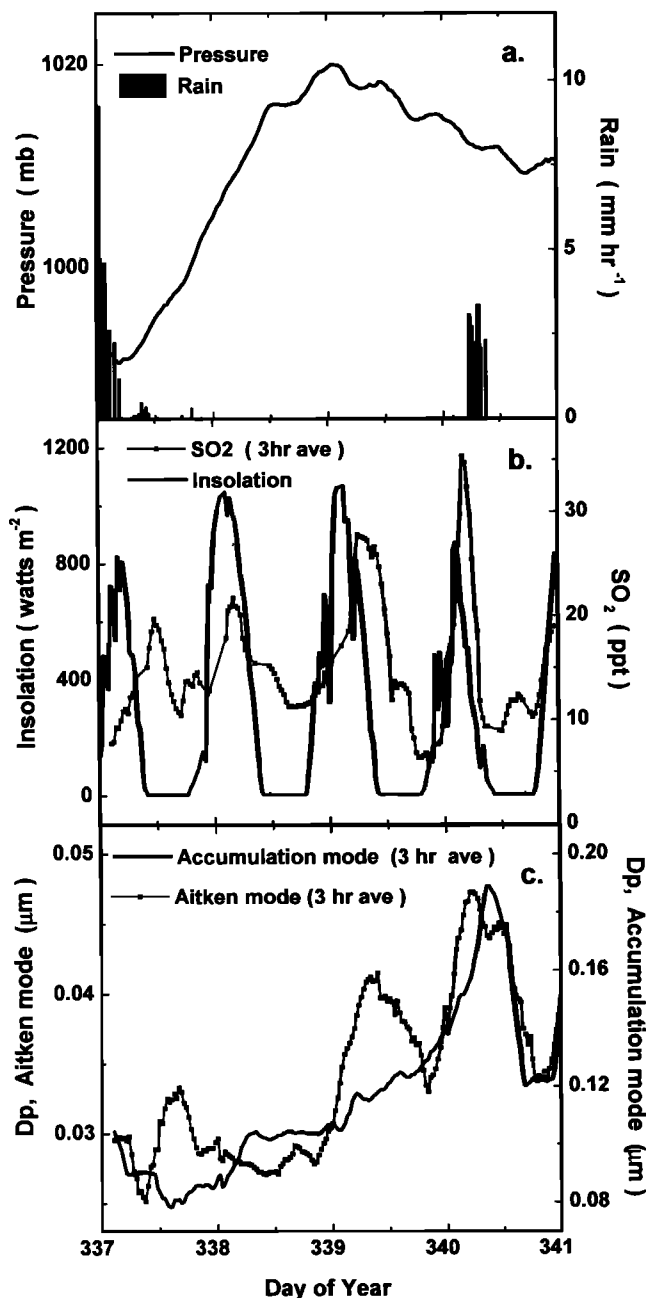


Figure 10. (a) Atmospheric pressure (millibars) and rainfall rate (mm hr^{-1}), (b) solar insolation (W m^{-2}) and 3-hour average SO_2 concentrations (ppt), and (c) 3-hour average diameters of the Aitken and accumulation modes (micrometers) as a function of DOY.

The average steady state H_2SO_4 concentration for the period DOY 337 to DOY 341 is 0.014 pptv. The flux of H_2SO_4 to accumulation mode aerosols is estimated to be $1.7 \times 10^{-17} \text{ mol S cm}^{-3} \text{ d}^{-1}$. The observed increase in the accumulation mode diameter (Figure 10c) for the $100 \text{ particles cm}^{-3}$ present corresponds to a mass increase of $2 \times 10^{-17} \text{ mol S cm}^{-3} \text{ d}^{-1}$. While again, these observations did not follow a specific air mass and thus the observed changes could be due to advection, the ocean-atmosphere flux of DMS, its conversion to

SO_2 , its oxidation to H_2SO_4 , and the subsequent flux of H_2SO_4 to existing particles is sufficient to explain the observed growth in this mode.

The coarse mode, which consisted almost entirely of sea salt (Table 1), had an average number concentration 18 m above the sea surface of $15 \pm 9.6 \text{ particles cm}^{-3}$ (Figure 11a). The number mean diameter of this mode was relatively constant at $540 \pm 73 \text{ nm}$ (at 10% RH). Particles in this size range have a residence time of approximately 2 days [Gong *et al.*, 1997], and thus the number concentration is determined by the balance of source and sink processes integrated over this time period. Sea-salt aerosols in this size range have been shown to have a strong exponential dependence on wind speed [O'Dowd and Smith, 1993]. A similar analysis of the ACE 1 coarse mode data shows that the instantaneous wind speed can account for approximately one third ($R^2 = 0.32$) of the variance (Figure 11b). This low correlation is not surprising given the highly variable wind speed over the region and regular frontal passages during the month-long study [Hainsworth *et al.*, this issue] in comparison to the 2-day residence time of the aerosol. Wind speeds measured at the ship varied from 0 to 22 m s^{-1} with a mean of $9.5 \pm 3.8 \text{ m s}^{-1}$ (Figure 11a). However, the winds, measured at the ship, do not necessarily reflect the wind/wave history experienced by the air parcel in the 2-day period prior to arriving at the ship. Sink processes, that is, wet and dry deposition and vertical turbulent transport, also vary with wind speed and frontal passages. The time and space variability in the processes controlling the sea-salt number concentrations require the use of regional aerosol transport models driven by observation-derived meteorology to better describe the fluctuations in the sea-salt number concentration. Local instantaneous wind speed is not a good indicator of sea-salt number concentrations in regions of rapidly varying wind speeds.

4. Conclusions

The MBL aerosol in the region south of Australia was composed of four distinct modes with a mean total particle number concentration of 500 cm^{-3} . The dominant process affecting the three smaller modes was the dynamical mixing between the FT and MBL which brought UF mode particles from the FT to the MBL and diluted the Aitken and accumulation mode populations. During the 29-day experiment an UF mode was observed approximately 50% of the time and coincided with the passage of cold fronts, postfrontal subsidence, or cloud pumping. There was no evidence of major new particle production in the MBL. The mean diameters of Aitken and accumulation mode particles increased in air masses that had been in the MBL for longer periods of time suggesting growth from the DMS sulfur that was emitted from the ocean to the MBL. Photochemical model calculations showed that the emissions of DMS could quantitatively account for the observed particle

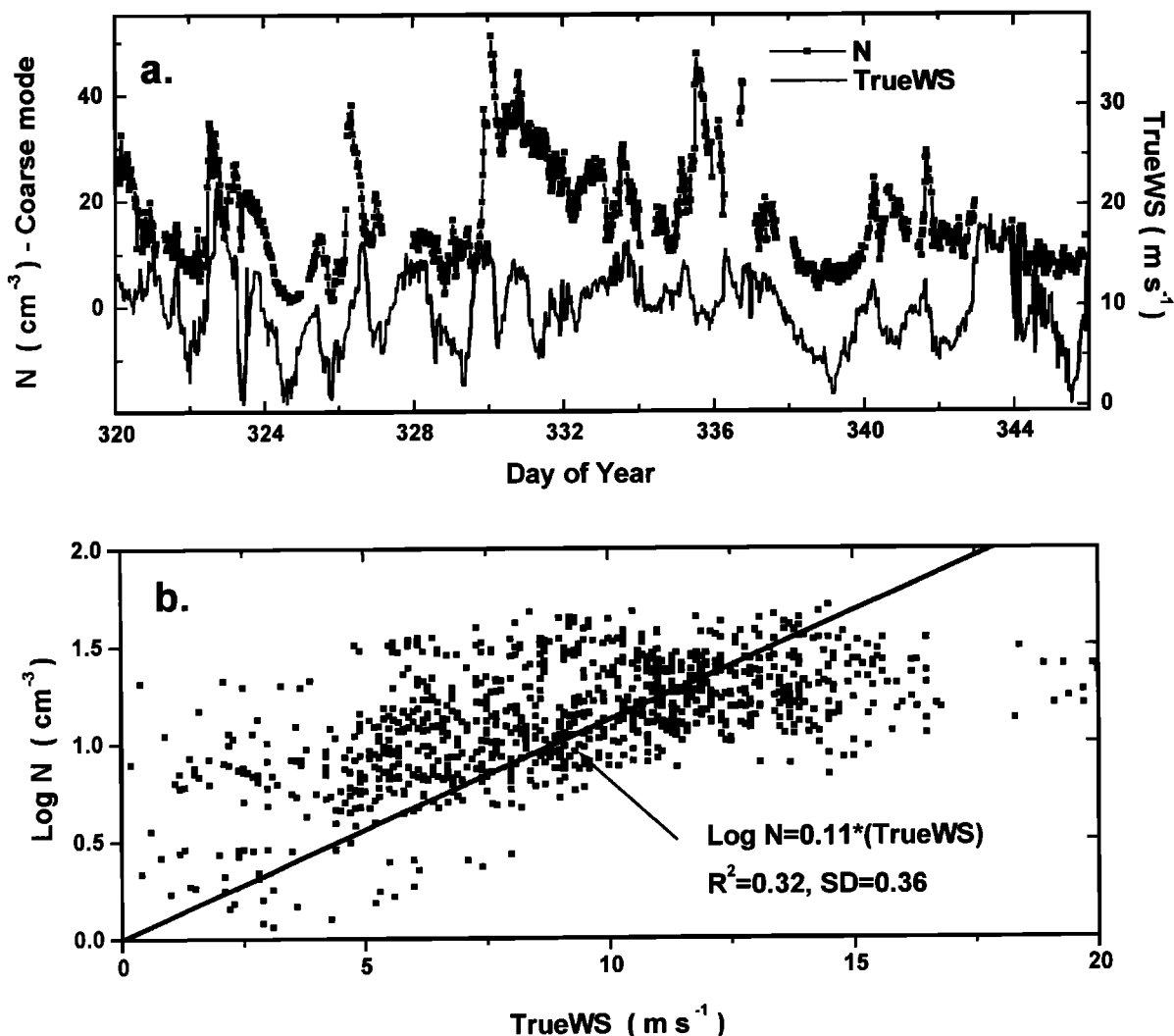


Figure 11. Total number concentration in the coarse mode (particles per cubic centimeters) and true wind speed (WS) (in m s^{-1}) measured aboard *Discoverer* plotted (a) as a function of day of year and (b) as a logarithmic regression of half-hourly data.

growth. The coarse mode was dominated by sea salt. However, in this region of rapidly changing wind speed, the local instantaneous wind speed accounted for only one third of the variance in the number concentration of the coarse mode.

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