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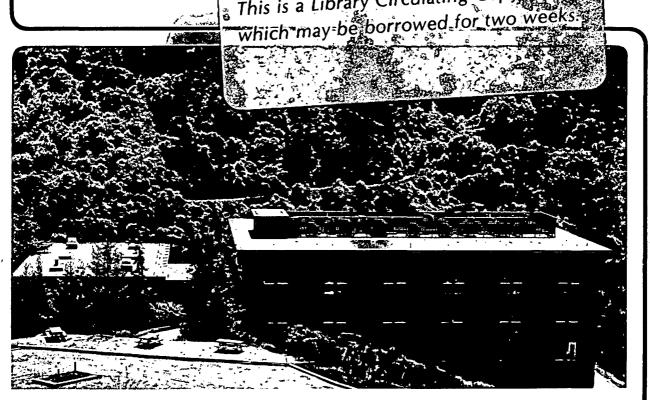
VIBRATIONAL FREQUENCIES OF SMALL METAL CLUSTERS. THE BERYLLIUM TETRAMER

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VIBRATIONAL FREQUENCIES OF SMALL METAL CLUSTERS.
THE BERYLLIUM TETRAMER

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ABSTRACT. The structure and harmonic vibrational frequencies of Be $_{\parallel}$ have been predicted using <u>ab initio</u> molecular electronic structure theory. A better than double zeta plus polarization (DZ+P) basis set was used in conjunction with self-consistent-field (SCF) and configuration interaction (CI) methods. The predicted frequencies (SCF followed by CI in parentheses) are a₁ 651 (680), t₂ 576 (589), and e 489 (487) cm⁻¹, respectively.

INTRODUCTION

During the past two years, the generation of well-defined metal clusters in the gas phase has progressed from dream to reality [1-6]. It is now becoming possible to study both the spectroscopic properties [1-3] and reactivity [4-6] of such small metal clusters. One of the primary motivations for the preparation of such "naked" clusters is to attempt to make analogies with metal surfaces and with metal particles involved in heterogeneous catalysis [7].

One of the simplest metal clusters to be well-characterized theoretically is the beryllium tetramer, Be $_{\mu}$. It was established [8] in 1975 that Be $_{\mu}$ has a tetrahedral equilibrium geometry and is much more strongly bound than Be $_2$ or Be $_3$. Subsequent theoretical studies [9-19] have confirmed this view and provided quantitatively reliable predictions. For example, Bauschlicher, Bagus, and Cox [17] have

addressed the problem using a triple zeta plus double polarization (TZ+2P) basis in conjunction with configuration interaction (CI) including all single and double excitations [20]. After appendage of the Davidson correction for quadruple excitations [21], the Be-Be bond distance was predicted to be 2.07 Å and the dissociation energy $D_{\rm e}$ to be 64 kcal/mole.

Vibrational spectroscopy has in recent years become a powerful tool for the characterization of metal surfaces and of chemisorption on metal surfaces [22]. In light of the analogy between surfaces and gas-phase metal clusters [7], it is apparent that the vibrational spectroscopy of the latter will in time become an important area of research. Since experimental studies of the infrared and Raman spectra of naked metal clusters are essentially nonexistent, theory is in a good position to provide some potentially helpful guidance to laboratory research. In this vein, we report here a theoretical study of the vibrational frequencies of Be_{ij} , a simple, relatively stable metal cluster.

THEORETICAL APPROACH

Previous theoretical research has shown [14,17,19] that the reliable prediction of the dissociation energy of Be $_{\mu}$, namely $\Delta E(Be_{\mu} \rightarrow 4Be)$, requires a high level of theory. Specifically, d orbitals must be included in the basis set, and the effects of electron correlation must be carefully considered.

Two basis sets of contracted gaussian functions were used in the present research and are designated A and B in Table 1. The Be s primitive funcions are those of van Duijneveldt [23], truncated to four significant figures in the orbital exponents. This s basis is contracted to (9s/5s) so as to maintain maximum flexibility in the valence shell. The Be p functions are those optimized by Yarkony [24] for the lowest ³P state (electron configuration 1s²2s2p) of the atom. As seen in Table 1, the p basis functions are contracted

(4p/2p) in Basis A and (4p/3p) in Basis B. Both basis sets include a set of six d-like functions (x^2 , y^2 , z^2 , xy, xz, and yz multiplied by $e^{-\alpha r^2}$), with orbital exponent of α = 0.5. Thus the final Basis Set A may be designated Be(9s4p1d/5s2p1d), while the same notation describes Basis Set B as (9s4p1d/5s3p1d).

All theoretical studies on Be_{μ} began with self-consistent-field (SCF) wave functions for the ground electronic configuration

At the SCF level of theory, the equilibrium geometry of Be_{ij} was determined using analytic gradient methods [25] with both basis sets. Harmonic vibrational frequencies were evaluated at the SCF level using analytic second derivative methods [26].

The effects of electron correlation were considered via the method of configuration interaction (CI) [20]. The CI included all single and double excitations, with the restriction that the four core is-like molecular orbitals were held doubly occupied in all configurations. Basis Set A only was used for the CI treatment. The Be $_{\downarrow}$ structure was optimized using analytic CI gradient methods [27] and the harmonic vibrational frequencies were evaluated from finite differences of analytic gradients. To determine the quadratic force constants of Be $_{\downarrow}$ in internal coordinates, only two gradient calculations are required, both in C $_{2v}$ symmetry. In point group C $_{2v}$ the CI included 7,591 configurations.

RESULTS AND DISCUSSION

Reported in Table 2 are equilibrium geometries, total energies, and harmonic vibrational frequencies for Be_{ij} predicted at three levels of theory. The bond distances and total energies are consistent with earlier theoretical studies [8-19]. The true Be-Be bond distance is probably 2.07 \pm 0.02 Å. This distance is noticeably shorter than the

Be-Be distances in metallic beryllium [28], namely 2.29 Å and 2.23 Å. Another point worth noting is that correlation effects decrease the bond distances in Be $_{\mu}$, contrary to what is typically found. However, this result may be explained when one appreciates that the lowest unoccupied molecular orbitals (LUMO) of Be $_{\mu}$ are bonding (rather than nonbonding or antibonding, as is usually the case) in nature. Thus double excitations into these bonding orbitals qualitatively strengthen and therefore shorten the six Be-Be bonds.

SCF vibrational frequencies predicted with Basis Sets A and B are in close agreement. The largest difference occurs for the e vibrational mode, for which Basis Set B gives a result 4 cm⁻¹ larger than Basis Set A. For this reason only the smaller Basis Set A was used at the CI level of theory.

The predicted CI vibrational frequencies are "nonintuitive" in the sense that two of the three lie above the analogous SCF predictions. For ordinary closed shell molecules, the opposite is usually true. For example, for the set of molecules CH_4 , H_2O , H_2CO , and HCN, the DZ+P CI vibrational frequencies were found [29] to be an average 4.8% lower than the analogous DZ+P SCF predictions. To speculate on whether this result is peculiar to Be_4 or characteristic of other small metal clusters would be premature.

For the sample set of four molecules studied by Yamaguchi and Schaefer [29], the remaining average error in the DZ+P CI harmonic vibrational frequencies was 3.5%. There is of course an additional error relative to the fundamentals due to the neglect of anharmonicity, and for normal closed shell molecules this error is in the same direction [29]. Thus it would not be unreasonable to expect the predicted Be $_{\mu}$ frequencies to be ~ 5% higher than the true frequencies. With this guidance from theory, we would hope to see some intrepid spectroscopic group take up the Be $_{\mu}$ problem in the laboratory, perhaps using infrared matrix isolation techniques.

Herzberg's discussion [30] of the isostructural P_{ij} molecule shows that for such a tetrahedral molecule, only the triply-degenerate t_2 vibrational mode is infrared active, i.e., both a_1 and e modes are

forbidden in the IR. Using analytic IR intensity methods [31], the t_2 mode is predicted to have intensity (per degenerate component) 0.25 (Basis Set A) and 0.27 (Basis Set B) km/mole. For comparative purposes, this is about the intensity of the 0-H symmetric stretching frequency in water. All three fundamental vibrations are Ramanallowed.

As Herzberg notes in his book [30] a molecule with as much symmetry as P_{μ} (or Be_{μ}) gives rise to a remarkably orderly pattern of vibrational frequencies if the central force approximation (CFA) is adopted. In particular Herzberg shows that in this approximation the vibrational frequencies of a tetrahedral M_{μ} molecule display the relationship

$$v(a_1):v(t_2):v(e) = 2:\sqrt{2:1}$$

The theoretical frequencies in Table 2 do not fit this formula well — a result which is not surprising when one realizes that the CFA amounts to neglecting all off-diagonal elements in the internal coordinate force constant matrix.

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Table 1. Be atom basis sets for the theoretical description of Be_{μ} .

Basis Set A:	Be(9s4p1d/5s2p1d) Function Type	Orbital Exponent α	Contraction Coefficient
	s	2732.	0.001916
	S	410.3	0.014720
	S	93.67	0.074290
	3	26.59	0.275402
	3	8.63	0.720340
	3	3.056	1.0
•	S	1.132	1.0
	S	0.1817	1.0
	S	0.05917	1.0
	p	3.202	0.052912
	р	0.6923	0.267659
	р	0.2016	0.792085
	p	0.06331	1.0
	d	0.5	1 .0

Basis Set B: Be(9s4p1d/5s3p1d)
Same as Basis Set A except for p function contraction

Function Type	Orbital Exponent α	Contraction Coefficient	
p	3.202	0.177439	
р	0.6923	0.897586	
p	0.2016	1.0	
р	0.06331	1.0	

Table 2. Structures, total energies, and harmonic vibrational frequencies (in $\text{cm}^{-1})$ for the Be_{μ} molecule.

·		Be[5s2p1d] SCF	Be[5s3p1d] SCF	Be[5s2p1d] CI
Bond distance (A)		2.083	2.073	2.062
Total energy (hartrees)		-5 8.35128	-58.35340	-58.52608
	a ₁	651	652	680
Vibrational	t ₂	576	577	589
Frequencies	e	489	493	487

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