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## ELASTIC NEUTRON SCATTERING IN UCd<sub>11</sub>

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We report results of neutron powder diffraction experiments on the heavy-electron magnet UCd<sub>11</sub> at temperatures above and below its ordering temperature. These measurements set an upper limit on the ordered moment of  $1.5\mu_B/U$ .

Specific heat measurements on UCd<sub>11</sub> show a large electronic contribution ( $\gamma \approx 840 \text{ mJ/mol K}^2$ ) above 5 K, at which temperature occurs an anomaly indicative of a magnetic phase transition in the bulk [1]. Associated with this anomaly is a very weak feature in the magnetic susceptibility. More recent muon-spin rotation/relaxation measurements [2] suggest antiferromagnetic order at 5 K but neither the ordered moment nor magnetic structure could be determined from these experiments. For the two other U-based heavy-electron magnets U<sub>2</sub>Z<sub>17</sub> [3] and UCu<sub>5</sub> [4], neutron scattering experiments have established that the magnetic order is simple antiferromagnetic, with ordered moments of  $0.8 \pm 0.1$  and  $(0.9 \pm 0.1)\mu_B/U$ , respectively, values close to 1/3 of their corresponding paramagnetic moments. Interestingly, the factor 1/3 also appears in the ratio of their electronic specific heat  $\gamma$  below  $T_N$  to that above  $T_N$  [5]. UCd<sub>11</sub> also exhibits the ratio  $\gamma(T=0)/\gamma(T \geq T_N) = 0.3$  [1] from which, if the analogy holds, we would imply an ordered moment of  $\geq 1\mu_B/U$ .

To minimize the problem associated with the extremely large neutron absorption cross section of naturally abundant Cd, we have performed neutron diffraction on samples enriched to 98.55% <sup>114</sup>Cd, which has a much smaller absorption cross

section. An approximately 5 g sample was prepared by sealing under vacuum fine U powder in a quartz tube with distilled <sup>114</sup>Cd and heating to 470°C for 7 days. X-ray analysis of the reacted powder showed the material to be  $\geq 97\%$  UCd<sub>11</sub>, having the Pm3m crystal structure with a lattice parameter  $a_0 = 9.290 \text{ \AA}$ . Magnetic susceptibility measured on 85 mg of the sample showed Curie-Weiss behavior for  $20 \text{ K} < T < 350 \text{ K}$ , with  $\mu_{\text{eff}} = 3.86\mu_B/U$  and  $\theta = -46.6 \text{ K}$ , and a plateau near 5 K, characteristic of the magnetic transition in UCd<sub>11</sub>.

Neutron scattering was performed on the General Purpose Powder Diffractometer at Argonne's Intense Pulsed Neutron Source. For these experiments, the 5 g sample was encapsulated in a thin-walled aluminum holder that was mounted on the cold finger of a <sup>4</sup>He cryostat whose temperature could be controlled down to 2 K. Parts of the sample holder exposed to the neutron beam and not containing sample were masked by a boron-nitride shield. Approximately 20 hours of data were collected at three temperatures 8, 4 and 2 K. Each data set was analyzed by Rietveld refinement using the codes developed by Larson and Von Dreele [6].

Results of the measurements at 2 K are shown in fig. 1. In fitting these data, it was necessary to

Table 1  
 Crystal Data for UCd<sub>11</sub> at 2 K

Atom	Atomic positions			$u(\text{\AA}^2)$	Fraction
	<i>x</i>	<i>y</i>	<i>z</i>		
U1	0	1/2	0	0.0072(5)	1
Cd1	1/2	1/2	1/2	0.0033(13)	0.81(3)
Cd2	0.1567(2)	0.1567(2)	0.1567(2)	0.0063(6)	0.84(1)
Cd3	0	0.3441(2)	0.3441(2)	0.0082(5)	0.87(1)
Cd4	1/2	0.2664(1)	0.2664(1)	0.0078(4)	0.88(1)

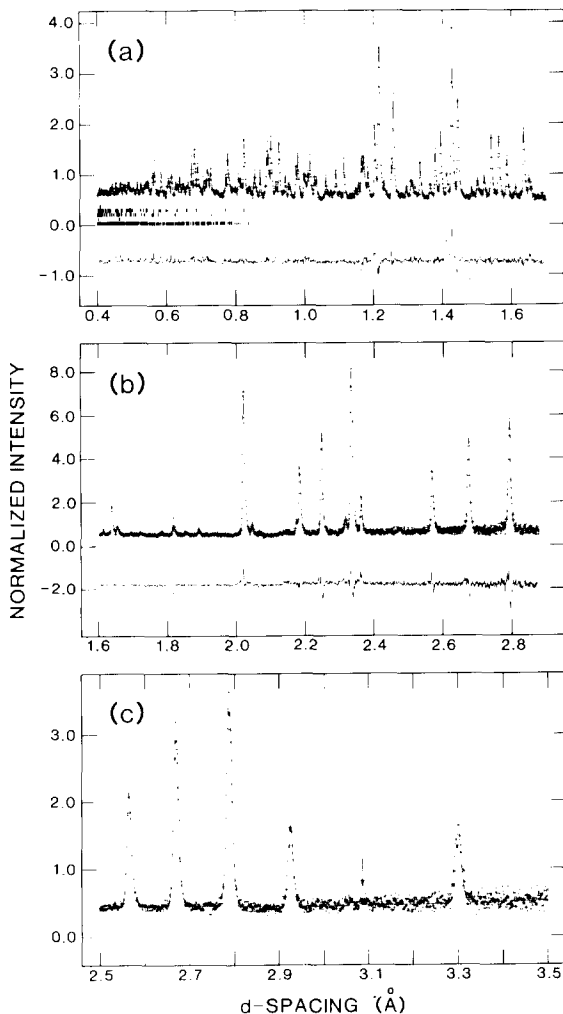


Fig. 1. Time-of-flight neutron diffraction spectra for  $UCd_{11}$  at 2 K. Data in panels (a) and (b) were collected with detectors at  $2\theta = 148^\circ$  and those in (c) with  $2\theta = 90^\circ$ . Crosses are data, the solid line is the Rietveld fit and the lower curve in (a) and (b) is the residual. Vertical tick marks indicate  $d$ -spacings of Bragg reflections for  $UCd_{11}$ , Al, Cd, BN and displaced Al, respectively, from the bottom set upwards.

include scattering from the Al sample holder, BN mask, and “displaced” Al arising from scattering off the Al tail of the cryostat. In addition to the majority phase  $UCd_{11}$ , we also found evidence for a small amount ( $\sim 2\%$ ) of free Cd, in agreement with the X-ray analysis. By taking these scattering sources into account, all peaks (except one with  $d = 1.409 \text{ \AA}$ ) could be indexed and the refined parameters for  $UCd_{11}$  are given in table 1. The

unidentified peak also appeared in the 4 and 8 K spectra. Additional structural details of  $UCd_{11}$  will be given elsewhere [7] but we note here that assuming unity occupation of the U-site gives a stoichiometry of  $UCd_{9.5}$  and that no structural change could be detected to within  $\Delta a_0 = \pm 0.00008 \text{ \AA}$ . Further, within the resolution of these experiments, no additional lines could be found below 5 K that might be expected to appear because of antiferromagnetic order.

A minimum detectable moment was estimated assuming a simple ferromagnetic structure, which would provide an upper limit. We assume that a magnetic peak would be observed if its intensity were greater than the  $(3,0,0)/(2,1,1)$  nuclear peak indicated by the arrow in fig. 1(c). Because the structure factor of this peak is known, we can calculate a limit of detectability to be  $1.5\mu_B/U$ . Therefore, we conclude that the ordered moment of “ $UCd_{11}$ ” must be less than  $1.5\mu_B/U$ , that the magnetic structure is not simple, or possibly both. The first conclusion is implicated on the basis of analogy to other U-based heavy-electron magnets.

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## References

- [1] Z. Fisk, G.R. Stewart, J.O. Willis, H.R. Ott and F. Hulliger, Phys. Rev. B 30 (1984) 6360.
- [2] S. Barth, H.R. Ott, F. Hulliger, F.N. Gyax, A. Schenck and T.M. Rice, Hyperfine Interactions 31 (1986) 403.
- [3] D.E. Cox, G. Shirane, S.M. Shapiro, G. Aeppli, Z. Fisk, J.L. Smith, J. Kjems and H.R. Ott, Phys. Rev. B 33 (1986) 3614.
- [4] A. Murasik, S. Ligenza and A. Zygmunt, Phys. Stat. Sol. (a) 23 (1974) K163.
- [5] H.R. Ott and Z. Fisk, in: Handbook on the Physics and Chemistry of the Actinides, eds. A.J. Freeman and G.H. Lander (Elsevier, New York, 1987) p. 85.
- [6] A.C. Larson and R.B. Von Dreele, Generalized Crystal Structure Analysis System. LA-UR-86-748 (1986).
- [7] A.C. Lawson, J.D. Thompson, M.W. McElfresh, A.P. Sattelberger and Z. Fisk (unpublished).