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Aysto, J.

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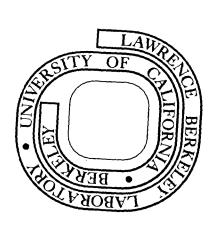
J. Äystö, D. M. Moltz, M. D. Cable, R. D. von Dincklage, R. F. Parry, J. M. Wouters and Joseph Cerny

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DECAY OF A NEW ISOTOPE, 24si:

A TEST OF THE ISOBARIC MULTIPLET MASS EQUATION.*

J. Aysto, D. M. Moltz, M. D. Cable, R. D. von Dincklage, R. F. Parry, J. M. Wouters and Joseph Cerny

Department of Chemistry and Lawrence Berkeley Laboratory University of California Berkeley, California 94720

ABSTRACT

Observation of beta-delayed protons from the decay of the mass separated T $_{\rm Z}$ = -2 nuclide $^{24}{\rm Si}$ (t $_{1/2}$ $^{\sim}$ 100 ms) establishes the mass excess of the lowest T = 2 (0 $^{+}$) state in $^{24}{\rm Al}$ (5.903 ± 0.009 MeV). An excellent fit to the quadratic form of the isobaric multiplet mass equation is obtained.

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[†] On leave from: University of Jyväskylä, Finland.

On leave from: University of Göttingen, West Germany.

Recent developments in nuclear instrumentation techniques have made possible the detection of exotic light nuclei up to the proton drip line. Such measurements have also led to the completion of the mass 8 and 20 isospin quintets [1,2] as well as to the observation of four members of the mass 32 and 36 quintets [3,4], thereby extending the tests of the isobaric multiplet mass equation (IMME) beyond the very well-studied isospin quartets [5]. The masses of analog states of an isospin multiplet are given in first order by the IMME as:

$$M(A,T,T_z) = a(A,T) + b(A,T) \cdot T_z + c(A,T) \cdot T_z^2$$
,

where the coefficients a, b and c are related to diagonal reduced matrix elements of the charge dependent part of the Hamiltonian.

Although this quadratic form of the IMME fits the vast majority of the data on known isospin quartets, a persistent deviation has been reported in the mass 9 quartet. In addition, the mass 8 isospin quintet also shows a deviation from the simple IMME, so that tests at higher masses have become imperative to establish comparable systematics for quintets. Such deviations from the quadratic form are generally represented by additional terms $d(A,T) \cdot T_Z^3$ and $e(A,T) \cdot T_Z^4$, in which the d and e coefficients can be derived from second order perturbation theory. Deviations from the quadratic form in cases of the A = 9 quartet and the A = 8 quintet have been explained as resulting from an expansion of the radial wavefunctions due to strong Coulomb repulsion as well as to isospin mixing [1,5].

Since, in general, the most experimentally inaccessible mass in the A = 4n isospin quintets is that of the 0^+ , T=2 state in the $T_z=-1$ nuclide, we have initiated a program to characterize these states by observing the beta-delayed proton decay of the mass-separated $T_z=-2$ nuclei [2]. Mass analysis is of critical importance due to the presence of other, copiously produced, strong delayed particle emitters; it is accomplished with a heliumjet fed, on-line mass separator system, RAMA, described in ref. 6. In particular, we wish herein to report the observation of the decay of 24 Si (whose accurate mass has not yet been determined). The lowest T=2 state in 24 Al is expected to be fed via a pure Fermi (superallowed) transition from the ground state of 24 Si. Accurate measurement of the energy of the subsequent isospinforbidden proton decay from this 24 Al state yields a precise value for its mass. Inasmuch as all members of the mass 24 multiplet are bound to isospinallowed particle decays, an exacting test of the isobaric multiplet mass equation using four members of the mass 24 quintet is now possible.

Beams of 70 MeV ³He ions of intensity 3-5 µA from the Lawrence

Berkeley Laboratory 88-inch cyclotron were used to produce ²⁴Si nuclei via

the ²⁴Mg(³He,3n) reaction. A special target system, consisting of three

1.2 mg/cm² Mg targets, was constructed to optimize the yield of the low

cross section and short-lived ²⁴Si. Reaction products, recoiling from

each target and stopping in 1.3 atm of helium, were collected by four adjacent

capillaries spaced evenly over a distance of the maximum recoil range (~ 3cm).

This 12-unit multiple capillary system fed a single 6 m stainless steel

capillary (1.27 mm i.d.), which transported activity to the skimmer-ion

source region. The average total transport time was approximately 250 ms.

Ethylene glycol was employed as an additive in the helium to build up high

molecular weight aerosols; nuclides attached to these aerosols possess

excellent transport and skimming properties. After skimming, the reaction products were ionized in a hollow cathode ion source operated at about 1800 °C. Ionized species were subsequently extracted and accelerated to 18 keV for magnetic mass separation. Total efficiency of the system (target to focal plane) for Si was about 0.01%.

The detection system for delayed protons on the focal plane of the separator consisted of a counter telescope subtending a solid angle of $\sim 30\%$ of 4π sr. This telescope consisted of a 42 μm ΔE detector and a 300 μm E detector, each 1.5 cm in diameter. Mass separated ion beams of interest were implanted on a 50 $\mu gm/cm^2$ carbon foil placed just before the ΔE counter. Given this high detection geometry and the large area detectors, a decreasing detection efficiency is obtained with decreasing proton energy; the minimum observable energy was approximately 1.5 MeV.

Having $J^T = 0^+$ and T = 2, 24 Si is expected to undergo superallowed β^+ decay to the 0^+ (T=2) analog state in 24 Al in addition to strong allowed transitions to lower-lying 1^+ states, predicted by shell model calculations [7]. The quadratic IMME prediction and the proton binding energy in 24 Al lead one to expect ~ 3.9 MeV laboratory energy for the proton transition from the analog state to the ground state of 23 Mg. Calibration of the telescope was accomplished by detecting the well-known beta-delayed proton emitter 25 Si [8], produced in much higher yield in the 24 Mg(3 He,2n) reaction. A proton spectrum arising from the decay of 220 ms 25 Si is shown in fig. 1(b); the groups at $^{4089\pm2}$ keV and $^{5403\pm2}$ keV provided convenient calibration points. These calibration energies were derived from the most recent value of the excitation energy of the lowest T = 3/2 state in 25 Al [9]. Events in the shaded areas in both spectra are due to the simultaneous detection of a proton and its preceding positron in this high geometry telescope.

The proton spectrum arising from the decay of 24 Si after bombardment for 560 mC is shown in fig. 1(a). Only one peak is evident in the spectrum; it occurs at a laboratory energy of 3.914±0.009 MeV. Possible lower energy groups arising from positron decay to 1 states were not observed, a result partly due to the low detection efficiency of the telescope below 3 MeV. A half-life of 100^{+90}_{-40} ms was estimated for the observed peak by comparing the 24 Si focal plane yield to the yields of 20 Na(t $_{1/2}$ = 446 ms), 24 Al(2.07s), 24m Al(129ms) and 25 Si(220ms). Relative cross sections for these nuclides were estimated by using the overlaid ALICE code as well as by the relative experimental cross section value for the previously studied 20 Mg - 21 Mg case [2]. The weak beta-delayed alpha emitter 24 Al [10] also served as a very convenient monitor during the 24 Si experiment. It should be noted that the peak attributed to the decay of 24 Si can not arise from the possible beta-delayed proton decay of 24 Al or 24m Al, since the maximum available proton energies in these latter decays are 2.1 and 2.5 MeV, respectively.

The proposed decay scheme for ²⁴Si is shown in fig. 2, in which the 3.914 MeV proton group [fig. 1(a)] is attributed to the isospin-forbidden proton decay of the lowest 0⁺, T=2 state in ²⁴Al. Since the superallowed transition to this ²⁴Al state has a calculated log ft-value of 3.18, allowed beta decay to other states near this excitation energy of 5.955 MeV would lead to considerably lower intensities in the proton spectrum.

Taking the center of mass proton energy together with the ²³Mg mass excess [11], one obtains a mass excess of 5.903±0.009 MeV for the 0⁺(T=2) analog state in ²⁴Al. Since the ²⁴Si ground state mass has not yet been measured, a mass excess of 10.750±0.016 MeV was predicted using the quadratic IMME. The decay energy and estimated half-life combined with the calculated log ft-value yield a branching ratio of 7±⁶₄% for the superallowed beta transition.

Experimental mass values for the four known members of the A = 24 isospin quintet are given in Table I [12-15]. These four accurate masses provide a good test of the validity of the quadratic isobaric multiplet mass equation as shown also in Table I. In addition, calculations with a cubic form for this mass equation are presented. Both sets of IMME coefficients show clearly the insignificance of the cubic term. The value of a possible d-coefficient, -0.9±2.5 keV, is consistent with zero, reflecting the absence of higher order charge dependent effects. This result is in good agreement with the results for the other quintets, in which four or more members are known, the only exception being mass 8 as noted earlier. In the mass 8 system nonzero d and e coefficients are obtained, a result attributed to the strong Coulomb repulsion associated with particle unbound members and isospin mixing effects, especially in the $T_2 = 0$ member of the quintet. Since isospin quintets with A = 24, 32 and 36 now have four members known in each, accurate measurements of the fifth members of these multiplets, although a difficult experimental task, would clearly be of substantial value.

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Table I.	Properties	of the	A=24	isobaric	quintet	and	coefficents	of	the .	IMME.

Nucleus	T _z	Mass Excess M [MeV]	E _X [MeV]	Reference
²⁴ si	-2	10.750(16)	0	IMME-prediction (quadratic)
24 _{Al}	-1	5.903(9)	5.955(10)	This work
24 _{Mg}	0	1.5016(16)	15.432(2)	[13],[14],[15]*
24 Na	1	-2.4473(12)	5.9702(9)	[12]
24 _{Ne}	2	-5.949(10)	0	[11]
24 Na	. 1	-2.4473(12)	5.9702(9)	[12]

Predicted coefficients for the IMME: M= $a+bT_z+cT_z^2+dT_z^3$

a	b	С	đ	χ ²
1.5018 (15)	-4.1742 (37)	0.2250 (34)	-	0.13
1.5016 (16)	-4.1743 (37)	0.2263 (48)	-0.0009 (25)	-

 $^{^{\}star}$ The quoted value is the weighted mean of the results given in the references.

FIGURE CAPTIONS

- Fig. 1. Spectra of beta-delayed protons from (a) ²⁴Si and (b) ²⁵Si.

 Arrows at low and high energy indicate telescope cutoffs. Events in the shaded areas in both spectra are a pileup effect as discussed in the text.
- Fig. 2. Proposed decay scheme for 24 Si. The 24 Si 24 Al mass difference is taken from the quadratic IMME prediction.

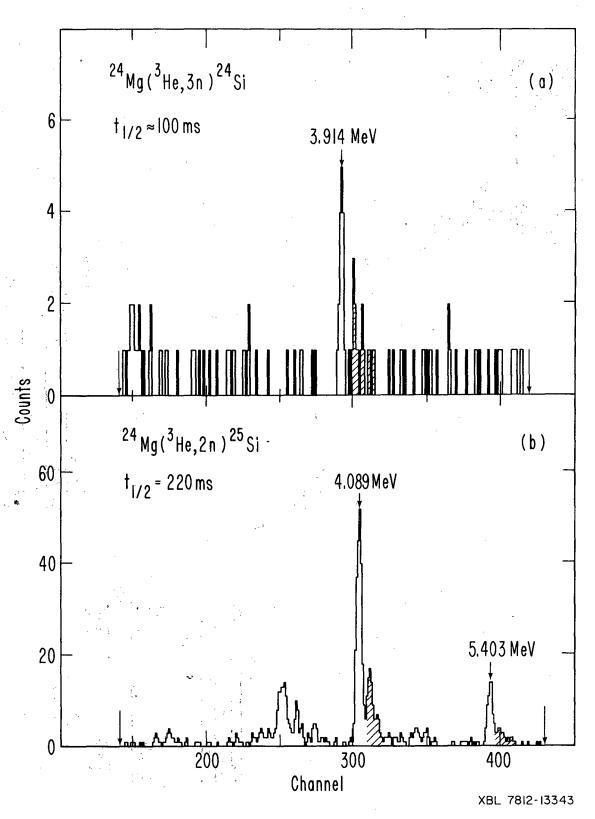
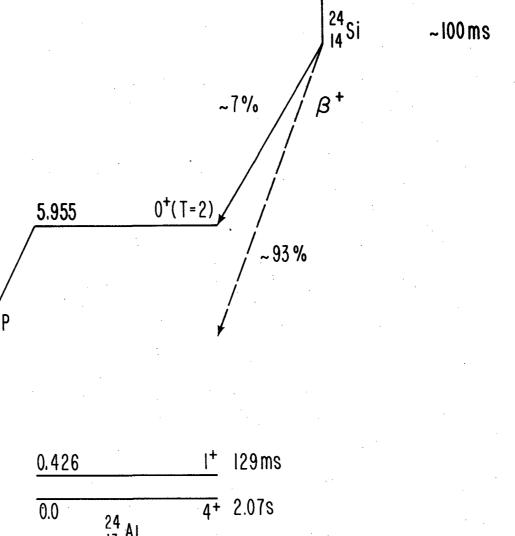


Fig. 1



10.802

0.0 24 AI

5/2+

3/2+

2.321

1.870

²³ M g

Fig. 2

LBL-8578

-12-

0⁺(T=2)

XBL 7812-13344

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