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Authors

Sanders, Isabelle C

Oberlander, Katherine A

Spearrin, R Mitchell

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Polytetrafluorethylene (PTFE) burn characteristics and toxicant formation in an oxidizer cross-flow via laser absorption tomography

Isabelle C. Sanders^{a,*}, Katherine A. Oberlander^a, R. Mitchell Spearrin^a

^a*Department of Mechanical and Aerospace Engineering, University of California, Los Angeles (UCLA), Los Angeles, California, 90095*

Abstract

A two-dimensional solid-fuel combustion experiment for fire-resistant polymers under forced convective cross-flow was developed to assess burn characteristics and toxicant formation using advanced laser absorption diagnostics that enable in situ species and temperature measurements near the fuel surface. The method was used to examine the thermochemical flow-field structure near the surface of polytetrafluoroethylene (PTFE) exposed to a well-defined solid-fuel pilot flame burning polymethyl methacrylate (PMMA). Infrared diode and quantum cascade lasers were used to probe rovibrational absorption transitions of hydrogen fluoride (HF) and carbon monoxide (CO), respectively, at the exit plane of a heterogeneous cylindrical fuel grain from which temperature and mole fraction could be inferred. A laser absorption tomography (LAT) technique with a Tikhonov-regularized Abel inversion was applied to reconstruct radially-resolved profiles in the axisymmetric reacting flow which, when compiled across a range of fuel lengths, provided a two-dimensional image of the near-surface reaction layer. Thermochemical data and reaction rate parameters from existing fluorocarbon chemistry models were used to generate high-temperature simulations of product species concentrations and temperature as a function of oxidizer-to-fuel ratio to elucidate observed trends and identify key reactions relevant to the novel dataset. The geometry and scale of the composite fuel experiments is intended to be tractable for reactive multi-physics models, enabling quantitative comparison of combustion characteristics and ultimately improved predictive capability of toxicant quantities resulting from fluoropolymer combustion in fire environments.

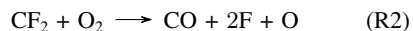
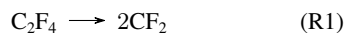
Keywords: fluorocarbon combustion; laser absorption spectroscopy; fire-resistant polymers; hydrogen fluoride

1. Introduction

Fire-resistant polymers are used extensively in modern structures to reduce flammability and the probability of fire-related loss or death. However, gas-phase emissions from these materials are often more toxic than those from conventional fuels. Specifically, halogenated polymers are considered fire resistant due to their chemical stability at high temperatures and the release of halides or halogen acids which can interfere with gas-phase combustion and limit heat release [1]. Unfortunately, hydrogen halides are acutely toxic and incomplete combustion leads to elevated emissions of carbon monoxide and unburned hydrocarbons including aromatics. These toxicants significantly increase the health and safety risk associated with smoke inhalation for first responders and civilians in close proximity to structural fires. While these challenges are well known, the detailed physics (kinetics, transport) associated with toxicant formation are insufficiently understood for many halogenated polymers, precluding accurate model prediction of fire toxicity and inhibiting risk quantification.

Polytetrafluoroethylene (PTFE), or Teflon, is a common fire-resistant fluoropolymer used extensively in construction and industrial applications due to its chemical inertness and high-temperature stability. The stability of PTFE may be attributed to the carbon-fluorine bond associated with the monomer tetrafluoroethylene (C_2F_4). The bond dissociation energy of the C-F single bond is higher than any other carbon-halogen as well as carbon-hydrogen bonds. As such, the pyrolysis of PTFE requires higher temperatures than most other synthetic or natural polymers and the material does not easily burn. However, upon thermal decomposition of PTFE in a fire environment, fluorine and fluorocarbon radicals will form and participate in combustion. Fluorine has the highest electronegativity of any element, and thus strongly attracts shared electrons, more so than oxygen. In the presence of H_2O , which is common to any ambient or fire environment, fluorinated radicals will react to form hydrogen fluoride (HF). HF is an acutely toxic gas with a lethal human exposure limit of approximately 50 ppm for a 30 minute period [2]. Here we aim to examine the incipient formation of HF (as well as CO) and associated burn characteristics of PTFE at fire conditions.

A number of previous theoretical and experimental studies have examined fluorocarbon combustion to advance chemical models towards predictive capability of toxicant formation. The foundational shock tube studies by Modica and LaGraff introduced a simple two-step mechanism to capture high-temperature C_2F_4 pyrolysis and oxidation, respectively, determining formation rates of CO and F atoms [3]:



Subsequent studies have built more detailed reaction mechanisms that account for additional steps

and pathways for pyrolysis, oxidation, and hydrogen interactions, with the latter required to predict HF. Modeling hydrofluorocarbon chemistry has required more sophisticated theoretical approaches to determine the many potential reaction pathways. Burgess et al used a combination of empirical data and ab initio calculations to produce a comprehensive mechanism of C1 and C2 fluorinated hydrocarbon chemistry [4], which has been tested and modified in other works [5–7]. Recent efforts towards understanding high-temperature hydrofluorocarbon chemistry have included shock tube studies to interrogate additional pyrolysis pathways [8, 9] and refined calculations of thermochemical properties [10, 11]. To ultimately predict toxic emissions in a realistic fire environment, additional physics including heat transfer mechanisms and fluid dynamics must be considered, and coupled with the chemistry. State-of-the-art computational studies have recently shown promise to couple such detailed chemistry with multi-dimensional solid-fuel combustion models [12–14]. Complementary and computationally tractable experiments are needed with quantitative in situ measurements of species, temperature, and other spatially-evolving parameters to constrain and anchor such models.

In this work, we present a solid-fuel combustion experiment for fire-resistant polymers subject to forced convective cross-flow to assess burn characteristics and toxicant formation from PTFE using laser absorption diagnostics for in situ species and temperature measurements near the fuel surface. The paper first describes the axi-symmetric solid-fuel experimental configuration and optical diagnostic integration, which involves an oxygen-enriched pre-burner flame using polymethyl methacrylate (PMMA) to create well-defined high-temperature boundary conditions to which the PTFE is exposed. The laser absorption tomography methods are then detailed, highlighting the advantage of the axi-symmetric geometry to reconstruct quantitative two-dimensional images of thermochemical flow structure near the material surface. A novel dataset of quantitative hydrogen fluoride and carbon monoxide species profiles as well as temperature is presented. The results are discussed with an accompanying first-order analysis using available thermochemical models that describe high-temperature chemical reaction pathways for fluorocarbon oxidation in a fire-like environment.

2. Methods

2.1. Solid Fuel Burner

The solid-fuel combustion experiment was adapted from our prior work [15] and is shown in Fig. . The experiment involves forced convection of a gaseous oxidizer across a cylindrical solid fuel grain. The cylindrical grain geometry provides 2D axi-symmetry (r, z) of the gas-solid interface, enabling quantitative tomographic reconstruction of line-of-sight optical measurements from a single view angle (discussed further below) while avoiding the edge ef-

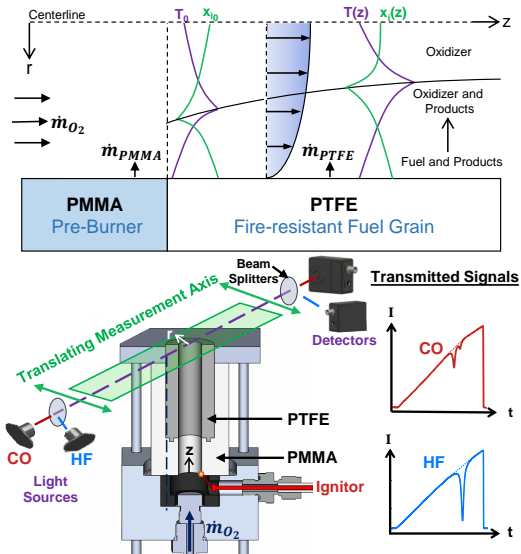


Fig. 1: solid fuel combustion experiment with forced oxidizer convection and laser absorption tomography setup enabling two-dimensional measurements.

fects of planar slab burners. The cylindrical axisymmetry also provides a relatively simple geometry for detailed computational modeling [12]. The test rig is designed to accommodate fuel grains of varying composition and length, the latter of which enables two-dimensional measurements based on exit plane measurements at different axial distances (z). Due to the high-temperature stability and fire resistance of PTFE, a composite fuel grain assembly was used in place of a homogeneous fuel cylinder. Notably, PTFE combustion proved difficult to initiate even in a pure oxygen environment at ambient temperature. As such, a 25.4 mm (1 in.) long polymethyl methacrylate (PMMA) fuel grain was used as a pilot burner to expose PTFE to both sustained high temperatures that initiate pyrolysis as well as hydrocarbon combustion product gases typically present in fires. PTFE fuel grains from 12.7-127mm (0.5-5 in.) in length were examined, each mated to the top surface of the PMMA pilot fuel grain. The boundary condition (temperature, species) created by the pre-burner was characterized, as shown in Fig. 5 and held constant across experiments. Importantly, the PMMA pre-burner combustion facilitates HF formation by supplying hydrogen (via water vapor) to the oxygen-enriched core flow with which the fluorocarbons can react, as would be the case in a fire environment. The test rig, and all associated experimental components, were contained within a high-flow (1150 cfm) laboratory fume hood to exhaust combustion products.

The composite fuel grain, including both the PTFE and PMMA components, had an inner diameter of 12.7 mm (0.5 in.). All tests were performed with a consistent oxidizer mass flow rate of $\dot{m}_{ox} = 1.5$ g/s using a single central injection port previously characterized [16]. Ignition was controlled remotely with an

electrically-initiated pyrotechnic charge directed into a precombustion chamber directly beneath the bottom surface of PMMA [15]. The mass of the heterogeneous fuel grain was measured before and after each test to estimate a test time-integrated mass loss, equivalence ratio and regression rate. A graphite gasket (~ 3 mm thick) was moulded to the top of the PTFE component to mitigate top surface pyrolysis. Burn durations for tests, including start up time and the quasi steady-state measurement period, were approximately 14 seconds. At the end of the test time, nitrogen gas was immediately flushed through the oxygen line to extinguish the flame and purge the plumbing.

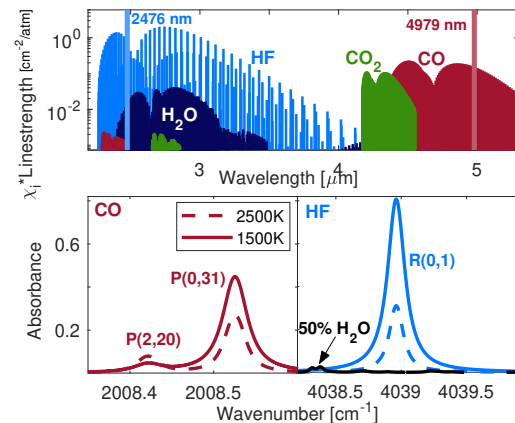


Fig. 2: *Top*. Absorption line strengths of HF, CO, and other relevant combustion species simulated at 2500 K [17]. *Bottom*. Simulated target transitions of CO and HF ($X = 15\%$) with exaggerated water interference.

2.2. Laser Absorption Spectroscopy

In situ species and temperature measurements in the near-surface reaction layer were performed using infrared laser absorption spectroscopy (LAS). LAS relates the attenuation of light through a gas medium to thermophysical properties of the absorbing molecules. Species-specific absorption spectra can be obtained by tuning lasers in spectral frequency or wavelength across individual lines associated with quantum energy transitions of the species of interest. In this study, we probe the mid-infrared fundamental vibrational bands of HF and CO in order to quantitatively recover species mole fraction and temperature in the reacting flow. The targeted wavelengths or frequencies are highlighted in Fig. 2. Specifically, we probe the hydrogen fluoride spectra in the R-branch of the fundamental band near 2475 nm. The R(1) rovibrational line at 4038.965 cm^{-1} , simulated in Fig. 2, was selected due to its high strength and lack of spectral interference from other atmospheric and combustion gases. It should be noted that non-optical ex situ detection and quantification of HF is particularly difficult due to its propensity to irreversibly interact with solid surfaces associated with sensors and sampling [18]. For carbon monoxide, we probe the P(0,31) and P(2,20) transitions near 2008.5 cm^{-1}

1 as also shown in Fig. 2. This line pair can be measured
 2 simultaneously within the typical injection current
 3 tuning range of a semi-conductor laser. The CO
 4 lines were selected in part due to the large difference
 5 in lower state energies, $\Delta E'' = 3151\text{cm}^{-1}$, which
 6 enhances temperature sensitivity [15]. We note the
 7 spectroscopic parameters of all target transitions are
 8 compiled in the HITRAN database [17].

9 The laser spectroscopy optical setup was coupled
 10 with the solid fuel combustion test rig as depicted
 11 in Fig. . A distributed feedback (DFB) diode laser
 12 (Norcada) with ~ 5 mW output centered around 4040
 13 cm^{-1} was utilized to target the hydrogen fluoride line.
 14 For carbon monoxide, a DFB quantum cascade laser
 15 (Alpes) with ~ 75 mW output centered around 2008
 16 cm^{-1} was used. The output beams of both lasers were
 17 combined via a CaF_2 beam splitter to make them co-
 18 linear before being focused across the exit plane of
 19 the fuel grain with a beam waste of ~ 0.5 mm. On
 20 the catch side, a second beam splitter was used to
 21 spatially de-multiplex the colinear beams, producing
 22 two perpendicular beams. Narrow bandpass spectral
 23 filters were placed along the separate beam paths to
 24 isolate each wavelength in front of distinct photode-
 25 tectors (VIGO PVI-4TE-5 and Thorlabs PDA10D)
 26 used for CO and HF measurements. Both lasers were
 27 injection-current scanned in intensity and wavelength
 28 at a frequency of 2 kHz using a sawtooth waveform,
 29 while the transmission signals from the photodetec-
 30 tors were recorded at a 2.5 MHz sample rate. The
 31 multi-plexed optical line of sight was translated at 2.5
 32 mm/s across the exhaust plane. The recorded trans-
 33 mission signals were time-averaged (every 200 scans)
 34 to yield a measurement radial resolution of half the
 35 beam diameter, such that the beam diameter (0.5 mm)
 36 effectively determines the radial spatial resolution.

2.3. Tomographic Imaging

37 Attaining the spatial profiles of temperature and
 38 species involves a tomographic inversion of the path-
 39 integrated line-of-sight laser absorption signals. Each
 40 of the three target lines for HF and CO are spectrally-
 41 resolved via the laser scans. The intensity signals can
 42 be converted to spectral absorbance, $\alpha(\nu)$, and spec-
 43 trally integrated to yield projected absorbance areas,
 44 A_{proj} , for each target line, spatially-resolved across
 45 the measurement plane (see figure 3, top). These
 46 spatially-resolved path-integrated line areas, A_{proj} ,
 47 can be directly related to the line-of-sight distribution
 48 of the integrated spectral absorption coefficient, $K(l)$:

$$\begin{aligned}
 A_{proj}(y) &= \int_{-\infty}^{\infty} \alpha(\nu) d\nu = \int_0^{L(y)} K(l) dl \\
 &= \int_0^{L(y)} PS_j(T(l)) X_i(l) dl
 \end{aligned}
 \quad (1)$$

50 To invert the line-of-sight absorption measure-
 51 ments, we exploit the axisymmetry of the flow-field
 52 to employ a Tikhonov-regularized onion peeling de-
 53 convolution [15, 16, 19], commonly referred to as 1D

54 tomography, to convert the projected absorbance ar-
 55 eas to line-integrated spectral absorption coefficients,
 56 $K(r)$, in the radial dimension (see figure 3, bottom).
 57 The Tikhonov regularization (or smoothing) param-
 58 eter was tuned through assessment of the L-curve to
 59 balance over-fitting versus over-smoothing. The
 60 radially-resolved ratio of the integrated absorption co-
 61 efficients of the two CO lines is used to obtain temper-
 62 ature, $T(r)$, where the ratio reduces to a linestrength
 63 ratio that is only a function of temperature: $R(r) =$
 64 $K_1(r)/K_2(r) = S_1(T(r))/S_2(T(r))$. With the gas
 65 temperature distribution known, the radially-resolved
 66 mole fraction can be determined from the radially-
 67 resolved spectral absorption coefficient of a single
 68 line via: $X(r) = K_j(r)/PS_jT(r)$. To produce
 69 two-dimensional images of the reaction layer, the ra-
 70 dial absorption profiles were acquired at different ax-
 71 ial lengths of the PTFE fuel grain, while holding the
 72 PMMA pre-burner length and oxidizer flow rate con-
 73 stant.

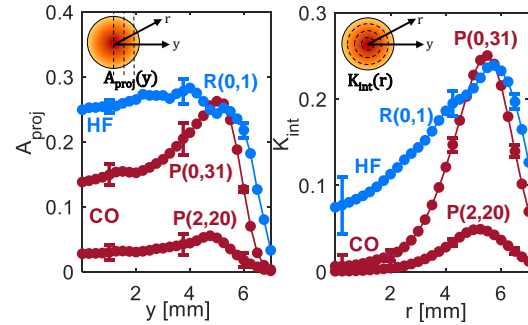


Fig. 3: *Left.* Projected integrated absorbance area versus distance y from the center of the reacting flow. *Right.* Tomographically-inverted radially-resolved integrated absorption coefficient.

3. Results

74 Experiments were conducted for PTFE fuel grain
 75 lengths of 25.4mm–127mm (1–5 in.) with 25.4 mm
 76 (1 in.) interval spacing, with an additional test of a
 77 12.7 mm (0.5 inch) fuel grain to add spatial resolu-
 78 tion in a region where a higher gradient was expected
 79 near the PMMA boundary condition. The forced con-
 80 vection oxidizer flow rate was held constant at $\dot{m}_{ox} =$
 81 1.5 g/s. Several tests were conducted at a single fuel
 82 grain length to verify experimental repeatability and
 83 to assess the quasi-steady state and axisymmetric as-
 84 sumptions over the measurement time.

85 Global fuel regression behavior was assessed based
 86 on post-burn mass loss and geometric measurements
 87 of the fuel grain; key results of this analysis are pre-
 88 sented in Table. 1. The total consumed fuel mass was
 89 determined for each of the PMMA pilot burner and
 90 the PTFE fuel grain by measuring the masses before
 91 and after each hot-fire test. This mass loss is used
 92 along with the total burn time to determine the space-
 93 time averaged parameters discussed. The space-time
 94 averaged regression rate, \bar{r} , is observed to decrease
 95 with fuel grain length due to the declining availability

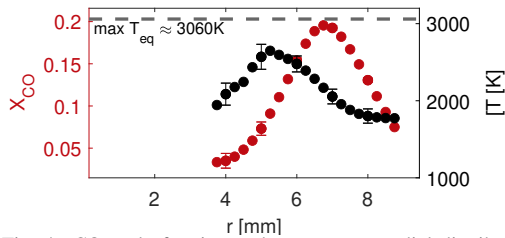


Fig. 4: CO mole fraction and temperature radial distribution measured from the 25.4 mm long PMMA pilot burner in GOx cross-flow.

of oxygen from ongoing reactions and formation of oxygenated product species. The averaged mass burn rate of regressing PTFE with fuel grain length are also shown. The rates were corrected for burn time variation by accounting for variation in mass burn rates measured from the 25.4 mm (1 in.) PMMA pilot burner fuel grain which ranged from 0.22 +/- 0.03 g/s. Over axial length, we observe a diminishing marginal increase in \dot{m} , consistent with the declining regression rate. Compared to a study with PMMA as the sole fuel, the rates measured for PTFE regression are over 50% lower [16]. The oxidizer-to-fuel (O/F) ratio decreases with fuel grain length as fuel mass is added to the constant oxidizer mass flow, trending from fuel-lean towards fuel-rich conditions. These global burn characteristics, although space-time averaged, provide an important reference for more granular species and temperature measurements.

Table 1: Global fuel regression parameters for PTFE with PMMA/GOx pre-burner and $\dot{m}_{ox} = 1.5$ g/s

L_f in.	\bar{r} [mm/s]	\bar{G} [kg/m ² s]	\bar{m}_{PTFE} [g/s]	O/F _{t_b}
0.5	0.069	0.010	0.17	9.8
1	0.041	0.011	0.23	7.6
2	0.039	0.012	0.35	5.1
3	0.024	0.014	0.44	3.0
4	0.021	0.014	0.55	3.1
5	0.017	0.015	0.58	2.2

3.1. PMMA pre-burner boundary condition

The pre-burner boundary condition was measured using the laser absorption measurements with the 1D tomography deconvolution method to determine the initial temperature and gas composition to which the PTFE was exposed. The first subplot in Fig. 5 shows the measured temperature and CO species as a function of radial position from the core. Given the PMMA mass burn rate, the global equivalence ratio at the boundary was approximately $\phi = 0.28$. Despite this very fuel-lean global condition, a peak temperature (~ 2950 K) slightly lower than the stoichiometric value is reached within 4 mm of the fuel surface due to the radial distribution of O/F ratio moving from fuel rich at the surface to the core which is primary oxygen. The carbon monoxide mole fraction is measured

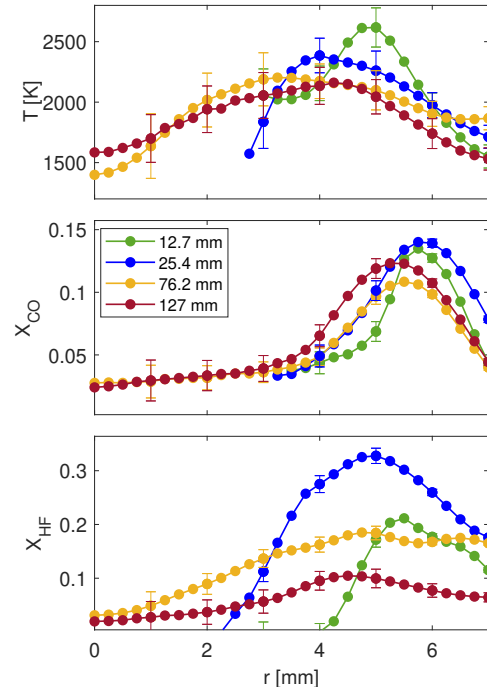


Fig. 5: Planar species and temperature measurements for PTFE combustion along the radial dimension at different axial measurement planes.

to peak at nearly 20% within 2 mm of the fuel surface while reducing to nearly zero within 5 mm of the PMMA surface. Overall, based on the global O/F, it is estimated that approximately 50-60% of the bulk cross-flow remains oxygen.

3.2. PTFE reaction layer measurements

Laser absorption tomography results from the aforementioned combustion tests of PTFE with the hot vitiated oxygen cross-flow are shown for select axial distances in Fig. 5 and compiled across the full experimental domain in two-dimensional images in Fig. 6. Quantitative spatially-resolved temperature and species profiles illustrate the thermochemical evolution in the reaction layer. Select trends are noted here and discussed further in the following section. It is notable that due to the inherently faster fuel regression of PMMA compared to PTFE, the respective port diameters and thus radial domains differ, rendering a slight flow convergence at the interface. Temperature and species uncertainty, as represented by error bar magnitudes, reflects variation in axisymmetry across the flow centerline, uncertainty associated to tomographic deconvolution, and uncertainty in spectral parameters from HITRAN [17]. Further details on this uncertainty analysis can be found in previous works from our group [16, 19].

Radial temperature distributions indicate a decrease in peak temperature and broadening with axial distance reflecting radial diffusion and mixing. The radial location of the peak temperature shifts towards

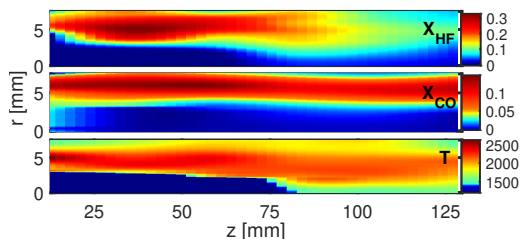


Fig. 6: Two-dimensional thermochemical structure in the near-surface reaction layer of PTFE in an oxidizer cross-flow

the oxygen-enriched core flow with increased axial distance as the O/F ratio decreases. The first measurement downstream of the fuel interface, at 12.7 mm (0.5 in.), most closely resembles the temperature profile measured at the PMMA surface initial condition. For the first few axial distances, there is insufficient concentration of carbon monoxide in the core flow (near $r = 0$) to obtain a reliable temperature measurement, thus the profiles are not fully-resolved. Moving downstream, more CO has been produced and has diffused into the core region such that local temperature can be calculated, albeit with greater uncertainty than in the more fuel and product rich regions near the fuel surface. The axial temperature decrease can also be attributed to increased distance from the PMMA pilot burner, and lower expected heat release of PTFE relative to PMMA.

The species measurements reflect an abundance of toxicant formation in the reaction layer, with HF reaching localized mole fraction values above 30%, and CO mole fraction peaking near 15%. In the bottom-right plot of Fig. 5, we observe a trend in X_{HF} with two distinct temporal phases. From the first to second measurement locations beyond the PMMA-PTFE interface where the gas is hottest, spanning 12.7 - 25.4 mm, the peak concentration of HF increases significantly and shifts about 0.5 mm towards the core. Beyond this point, however, we observe a decrease in peak HF concentration coupled with a rapid broadening across the radial domain. The more concentrated peaks seen at shorter lengths appear to spread out from their maxima a few millimeters inward of the fuel grain wall and distribute into the oxidizer core flow, though the concentrations in the core remain relatively low ($X < 0.05$) throughout the range of fuel lengths tested. Conversely, the high CO concentrations are largely maintained with increasing axial distances, likely reflecting a balance of radial diffusion with increasing CO formation with lower O/F. Throughout the full axial domain, encompassing 127 mm, the peak CO location moves gradually towards the core for a total variation of approximately 1 mm.

4. Discussion/Analysis

The measured species and temperature results discussed in the prior section reflect radial and axial variation of O/F as well as chemical reaction kinetics that compete with the transport timescales. To assess the observed trends, a thermochemical equilibrium

model and chemical kinetics mechanism were used to simulate relevant conditions. Simulations were performed using 0-D reactors in the Cantera toolbox, with hydrofluorocarbon thermodynamic data and reaction rate parameters from Burgess et al. [6, 7].

4.1. Chemical equilibrium analysis

To estimate chemical equilibrium for PTFE combustion with the pilot PMMA burner, a two-step simulation was performed. First a chemical equilibrium simulation of PMMA and oxygen was used to obtain product gas composition over a range of O/F reflecting the radial variation of O/F at the PMMA-PTFE boundary. Second, PMMA-O₂ product gas was used to prescribe the vitiated oxidizer initial conditions for PTFE combustion, using the TFE monomer as the primary fuel. The final simulated product composition is presented here in Figure 7 as a function of the experimental range of mass burn rates of PTFE, reflecting an increase in axial distance and corresponding reduction in O/F, which is similarly expected in radial variation moving towards the PTFE surface. The subplots show resulting composition for different initial equivalence ratios of MMA/GO_x, representative of the radially-varied initial conditions [20]. The final subplot of figure 7 indicates the maximum temperature (for any O/F) expected at equilibrium based on the ratio of PTFE and PMMA mass consumption rates at different axial distances, where as $\dot{m}_{PTFE} \rightarrow 0$, the temperature would be the maximum temperature for pure PMMA. The simulations presented here are intended to provide a reference for the range of species concentration and temperatures measured in the PTFE reaction layer.

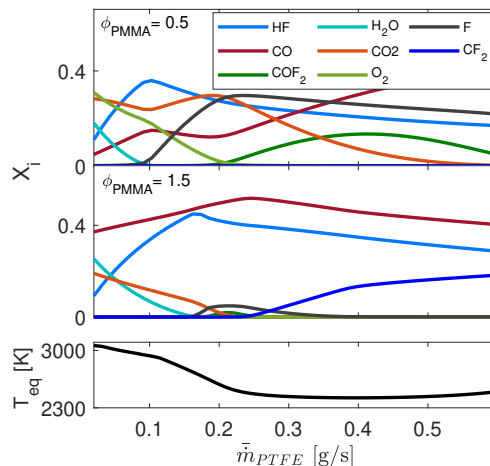


Fig. 7: Species mole fractions (X_i) and max temperature at chemical equilibrium, T_{eq} , for PTFE combustion over a range of mass flow rates of PTFE, with oxidizer vitiated with MMA/GO_x products at equivalence ratios of 0.5 and 1.5.

Despite the relatively simplicity of the 0-D chemical equilibrium model, many of the observed trends in measured temperature and species are consistent. The maximum magnitude of HF mole fraction is simulated to be approximately 40% across the various

1 initial conditions, just higher than the maximum measured
 2 values. Notably, an initial spike in HF, peak-
 3 ing at fuel mass flow rates between 0.1 and 0.2 g/s,
 4 is followed by a gradual decay. This trend is re-
 5 flected in the experimental data wherein the peak HF
 6 concentration increases between 12.7 mm and 25.4
 7 mm, followed by a steady decline over the sequen-
 8 tial axial positions. The initial rise in HF corresponds
 9 with a decline in H₂O, which is presumed the domi-
 10 nant hydrogen-containing species in the vitiated oxidizer
 11 exiting the PMMA-O₂ pilot burner. After HF peaks,
 12 the subsequent decline corresponds to a rise in
 13 atomic fluorine and carbonyl fluoride, COF₂, pro-
 14 vided sufficient oxygen is present. The atomic fluo-
 15 rine and COF₂ concentrations decline at more fuel
 16 rich conditions, which become dominated by CO and
 17 CF₂, along with diluted HF, pronounced in the most
 18 oxygen-deficient initial condition. CO is predicted
 19 to form in large concentrations at most conditions.
 20 In comparison to the experimental measurements, the
 21 equilibrium model estimates notably higher peak CO
 22 concentrations (~40%) in the fuel-rich regions by
 23 about a factor of 2. Some of the difference can cer-
 24 tainly be related to over-simplification of the model,
 25 but carbon predicted to form carbon monoxide and
 26 other gaseous products may also disproportionately
 27 exist in solid carbon char or particles.

28 The equilibrium temperature shown is the maxi-
 29 mum value across the range of radial fuel-lean to fuel-
 30 rich equivalence ratios for a given global balance of
 31 PTFE and PMMA mass burn rates at each axial posi-
 32 tion. As $\dot{m}_{PTFE} \rightarrow 0$, the maximum equilibrium
 33 temperature approaches that of purely MMA/GOx
 34 combustion, approximately 3060K. Downstream, fur-
 35 ther away from the PMMA fuel grain, the tempera-
 36 ture decreases to a plateau of around 2350K which
 37 is slightly higher, but within uncertainty, compared
 38 to the experimentally measured temperatures down-
 39 stream. Beyond the first 25.4 mm downstream, cor-
 40 responding in the model to \bar{m} of approximately 0.2, the
 41 peak temperature stabilizes. These trends are consist-
 42 ent with experimental data. The formation of HF is
 43 exothermic, however, the formation, by way of fluo-
 44 rocarbon decomposition, of the F radical is endother-
 45 mic. The effects of the trade off in these fluorinated
 46 end product species is seen in the temperature trend.
 47 As less hydrogen is available for F to abstract to form
 48 HF, F radical product concentration increases and the
 49 equilibrium temperature decreases. The formation of
 50 COF₂ is similarly endothermic, thus its concentration
 51 peaks also contribute to local lower temperatures in
 52 fuel lean conditions where it is most present.

53 4.2. Reaction kinetics analysis

54 To better understand the reactions responsible for
 55 the observed trends in species and temperature for
 56 PTFE combustion, a high-temperature chemical kin-
 57 etic reaction pathway analysis was performed. Key
 58 pathways from PTFE to CO and HF formation at
 59 2500 K are shown in figure 8. PTFE thermally de-
 60 composes at high temperatures (> 400°C) primarily

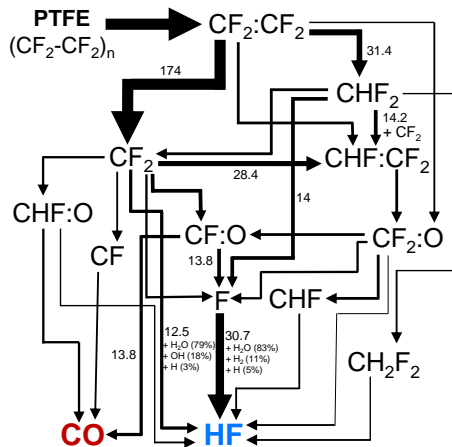
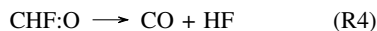
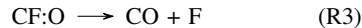


Fig. 8: Reaction pathway diagram tracking the fluorine atom through PTFE pyrolysis and oxidation with water along pathways that result in hydrogen fluoride (and carbon monoxide) production (T = 2500K)

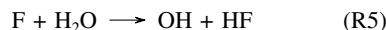
61 into the monomer, tetrafluoroethylene (TFE), from
 62 which numerous reaction pathways exist to form HF
 63 and CO as products by way of intermediate fluorocar-
 64 bons, hydrofluorocarbons (HFCs), and carbon oxoflu-
 65 orides (COFs). As mentioned, the PMMA combustion
 66 products serve as the primary source of hydrogen,
 67 largely via water molecules, which is required in the
 68 formation of HF. The case shown reflects the relative
 69 reaction contributions simulated at an early time step
 70 of 2.5 μs for equal mass contributions from PTFE and
 71 PMMA (the pilot burner fuel) analogous to approxi-
 72 mately 25.4 mm downstream of the PMMA/PTFE in-
 73 terface. Here, the PMMA contribution is again rep-
 74 resented by the product species resulting from a prior
 75 equilibrium simulation of MMA/GOx where $\phi = 1$.
 76 While the relative contributions towards the forma-
 77 tion of HF and CO are dependent on the initial fuel-
 78 oxidizer mixture and temperature, the case in Figure 8
 79 was selected as it is representative of the many major
 80 contributing pathways found across initial conditions
 81 and reactor progression at combustion temperatures.

82 Beyond the initial depolymerization, TFE primar-
 83 ily decomposes into two difluoromethylene (CF₂)
 84 molecules, although some reacts to form difluo-
 85 romethylide (CHF₂) and carbonyl fluoride (COF₂).
 86 The magnitude of these alternate pathways is directly
 87 impacted by the initial concentrations of H₂O and O₂
 88 available as dictated by the preliminary MMA combus-
 89 tion equilibrium. The predominant pathways we
 90 observe here towards CO formation involve the oxida-
 91 tion of CF₂ to produce the carbon oxofluorides, CFO
 92 and CH:FO, followed by decomposition (see R3 and
 93 R4 below); the reaction following the latter species
 94 also produces HF.



1 Although the reaction rates (shown in $kmol/m^3s$) re- 55
2 sulting in CO are shown to be relatively low compared 56
3 to the pathways producing HF, this is largely due to
4 the time step as the formation of HF is relatively rapid
5 compared to that of CO.

6 The fluorocarbons, HFCs, and COFs nearly all 58
7 have direct reactions that produce HF; however, 59
8 the favored pathways consistently involve disso- 60
9 ciation to produce the fluorine atom from these 61
10 molecules. Thereafter, the highly reactive fluorine 62
11 radicals rapidly abstract hydrogen from water and re- 63
12 act with the hydrogen molecules and radicals to pro- 64
13 duce HF, per R5. 65



14 Previous works have found, by way of modeling and 69
15 experimental efforts, hydrogen fluoride to be formed 70
16 preferably over water as a final endpoint product for 71
17 hydrogen [5]. This is due to the relatively high elec- 72
18 tronegativity of fluorine and H-F bond strength, mak- 73
19 ing it a more stable product of the two. In fact, pro- 74
20 vided there are sufficient fluorine atoms available, the 75
21 chemical model predicts that any water present will 76
22 decompose and the hydrogen atoms will react to form, 77
23 as follows from the experimental HF trend and was 78
24 noted in the equilibrium simulation results above. 79

5. Conclusion

25 In this work, the burn characteristics and toxicant 82
26 formation of polytetrafluoroethylene (PTFE) combus- 83
27 tion were examined using laser absorption tomogra- 84
28 phy. Experiments were conducted in a solid fuel com- 85
29 bustion facility with gaseous oxygen cross-flow and a 86
30 PMMA pilot burner to emulate fire-like conditions. 87
31 Radially-resolved temperature and species (CO and 88
32 HF) concentration profiles were obtained by invert- 89
33 ing line-of-sight absorption measurements across the 90
34 fuel grain exit planes. Quantitative, two dimensional 91
35 thermochemical images were produced by compil- 92
36 ing planar results from tests with varied fuel grain 93
37 lengths. The data were compared to chemical mod- 94
38 els for fluorocarbon thermal decomposition and com- 95
39 bustion, revealing similar trends to the experimental 96
40 results and highlighting the underlying reaction path- 97
41 ways responsible for some observations. The geome- 98
42 tries, operating conditions, and experimental mea- 99
43 surements of these experiments are detailed so that 100
44 the data can serve as validation targets for higher- 101
45 fidelity models of toxicant formation and destruction 102
46 in PTFE combustion. 103

Declaration of competing interest

48 The authors declare that they have no known com- 109
49 peting financial interests or personal relationships that 110
50 could influence the work reported in this paper. 111

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