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PALEOCLIMATE

Antarctic surface temperature and elevation during the Last Glacial Maximum

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Water-stable isotopes in polar ice cores are a widely used temperature proxy in paleoclimate reconstruction, yet calibration remains challenging in East Antarctica. Here, we reconstruct the magnitude and spatial pattern of Last Glacial Maximum surface cooling in Antarctica using borehole thermometry and firn properties in seven ice cores. West Antarctic sites cooled $\sim 10^\circ\text{C}$ relative to the preindustrial period. East Antarctic sites show a range from -4° to -7°C cooling, which is consistent with the results of global climate models when the effects of topographic changes indicated with ice core air-content data are included, but less than those indicated with the use of water-stable isotopes calibrated against modern spatial gradients. An altered Antarctic temperature inversion during the glacial reconciles our estimates with water-isotope observations.

Using oxygen and hydrogen isotope ratios in ancient polar ice as records of past site temperature requires a calibration (1). Surface temperature and the isotopic composition of precipitation correlate spatially in Antarctica, with a regression coefficient α_S (spatial slope) of 0.80 per mil per kelvin (‰K^{-1}) for $\delta^{18}\text{O}$ (the ratio of ^{18}O to ^{16}O) (2). Reconstructing past temperatures requires regression over time, and this temporal slope α_T may differ from α_S . In East Antarctica, where the longest continuous ice core records, going back to 800 thousand years before present (ka BP), have been extracted (3), independent temperature estimates are not available, and the spatial slope is commonly used to convert isotopic ratios to temperature (1); this approach gives a surface temperature difference ΔT_S of around -9°C between the Last Glacial Maximum (LGM) (26 to 18 ka BP) and the preindustrial period (1, 4, 5).

Antarctic LGM-preindustrial isotope changes depend on many factors, including hemispheric sea surface temperatures (6), sea ice extent (7), ice sheet elevation (8), vapor origin and transport, precipitation seasonality, and post-

depositional isotopic exchange (9). Isotope-enabled general circulation models seek to capture these physical processes, making them an invaluable tool for studying isotopic variations. Such models simulate LGM-preindustrial α_T ranging from 0.3 to 1.4‰K^{-1} in central East Antarctica (implied ΔT_S of -4° to -20°C), which implies that several aforementioned processes are poorly constrained (8, 10–12).

We distinguish three temperatures: (i) the climatic temperature T_{CLIM} at constant elevation (relative to the present-day geoid); (ii) the surface temperature T_S , which may differ from the climatic temperature because of changing ice sheet topography; and (iii) the vapor condensation temperature T_C , which is warmer than the surface because of the strong Antarctic inversion (2, 13).

In this study, we empirically reconstruct LGM surface temperature across Antarctica (Fig. 1) using two independent methods. We investigated five East Antarctic ice cores—EPICA (European Project for Ice Coring in Antarctica) Dome C (EDC), EPICA Dronning Maud Land (EDML), Dome Fuji (DF), Talos

Dome (TAL), and South Pole (SP)—and two West Antarctic cores—West Antarctic Ice Sheet (WAIS) Divide (WD) and Siple Dome (SDM).

First, we estimated ΔT_S at EDC and DF from the measured borehole temperature profiles (Fig. 2) using a method similar to that used recently at WD (14). Owing to the downward ice flow and low thermal diffusivity, the ice sheet maintains an imprint of its past surface temperature history. The large ice sheet thickness at EDC and DF is favorable for preserving past temperatures, yet the low accumulation rate is not. Consequently, the relative uncertainty in the EDC and DF borehole reconstructions is larger than that at WD. To constrain the problem better, we used downward ice velocities measured by means of phase-sensitive radio-echo sounding (EDC only) and accurate age constraints derived through volcanic synchronization to the layer-counted WD time scale.

We forced a one-dimensional heat transport-ice flow model at the surface boundary with a temperature history that is based on the $\delta^{18}\text{O}$ record scaled with a constant α_T value (10). Applying traditional isotope scaling ($\alpha_T \approx 0.7 \text{‰K}^{-1}$, yielding $\Delta T_S = -9^\circ\text{C}$ at EDC and -7.5°C at DF) simulates temperature profiles that do not fit the borehole observations at either site (Fig. 2). At EDC, the model-data fit is optimized for $\alpha_T = 1.14 \text{‰K}^{-1}$, which is consistent with $\Delta T_S = -5.5^\circ\text{C}$ (95% confidence range is -6.9° to -3.1°C). At DF, the optimal ΔT_S is in the -2.0° to -5.4°C range; we provide a range without a best estimate because, at DF, there are no direct constraints on the downward ice velocity. In Fig. 1, the WD, EDC, and DF borehole estimates are marked “BH.”

Second, we reconstructed past climate at all seven sites using the dependence of firn densification, the gradual transformation of polar snow to ice, on T_S and accumulation rate (A). Air bubbles are isolated from the atmosphere at the lock-in depth (50 to 120 m below the surface), an event preserved in two ice core signals (15): $\delta^{15}\text{N}$ of N_2 which records past firn column thickness by means of gravitational enrichment, and the gas age-ice age difference, or Δage . The $\delta^{15}\text{N}$ and Δage -isopleths are perpendicular in T_S -A space (Fig. 3A), meaning that if $\delta^{15}\text{N}$ and Δage are independently known, a distinctive climatic (T_S , A)

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solution exists (subject to the uncertainties of the firn model).

Synchronization using both volcanic deposits and globally synchronous abrupt atmospheric methane variations allowed us to estimate Δ_{age} empirically for the Antarctic ice cores (10, 16). We used an inverse dynamical firn densification–heat transport model (17, 18) to reconstruct T_S and A histories that optimize the fit to Δ_{age} and $\delta^{15}\text{N}$ data (Fig. 3, B and C). Reconstructed accumulation rates agree (within uncertainty) with independent estimates (fig. S8). Methodological biases and uncertainties are estimated by using a Monte Carlo approach (10). The histograms in Fig. 1 give the ΔT_S distribution of the Δ_{age} -based reconstruction.

In East Antarctica, ΔT_S ranges from $-3.8^\circ \pm 2.0^\circ\text{C}$ (DF) to $-7.1^\circ \pm 1.7^\circ\text{C}$ (TAL); at DF, EDC, and EDML, ΔT_S is substantially lower than estimates from isotope scaling that use α_S . The two West Antarctic sites have similar ΔT_S of $-10.2^\circ \pm 2.4^\circ\text{C}$ (SDM) to $-10.3^\circ \pm 1.3^\circ\text{C}$ (WD). The Δ_{age} - and borehole-based reconstruction methods agree within uncertainty at all sites (Fig. 1). Allowing for more flank-like ice flow at EDC during the glacial period (which would occur if the divide position were different from that at present) improves the agreement by changing the borehole estimate to around -4.5°C (10); we choose to report the -5.5°C value to keep both methods independent. PMIP4 (Paleoclimate Modeling Intercomparison Project phase 4) simulations (19) find a seven-site-mean ΔT_S magnitude that is $1.2^\circ \pm 4.6^\circ\text{C}$ larger than our Δ_{age} -based reconstructions (mean and spread of 10 climate models; Fig. 1).

We emphasize that the firn method is primarily constrained by the empirical Δ_{age} estimates. Because T_S and A broadly covary via the saturation vapor pressure, the deglacial climatic changes run parallel to the $\delta^{15}\text{N}$ -isopleths (Fig. 3A). Therefore, $\delta^{15}\text{N}$ data alone do not constrain the magnitude of climate change meaningfully. The effects of T_S and A are additive in Δ_{age} , however, making Δ_{age} a sensitive proxy for climate change (Fig. 3D), as first noted by Schwander *et al.* (20). The empirical Δ_{age} at 24 ka BP is larger than at 18 ka BP for all five cores where both are available, and coldest conditions in Antarctica occur ~ 27 to ~ 24 ka BP in our reconstructions (fig. S8h); this follows expectations from local insolation (21).

We propose that elevation changes explain the spatial differences in ΔT_S (8). Let Δz be the LGM elevation anomaly relative to the present. We present WD and DF total air content data (fig. S12) and interpret them in terms of elevation change (22). These data suggest a 420-m (range, 280 to 590 m) contrast in Δz between WD and central East Antarctica (here, DF and EDC)—for example,

$\Delta z = +300$ m at WAIS and $\Delta z = -120$ m in central East Antarctica (Fig. 4B). Our estimate is broadly in agreement with LGM ice sheet reconstructions that suggest a West-East Δz contrast between 160 and 560 m (10). Although the implied Δz at WAIS exceeds the observed highstand at ice margin nunataks (23), such data do not strongly constrain the elevation at WD more than 500 km away. The corresponding ΔT_S contrast (WD ΔT_S minus the average ΔT_S at DF and EDC) is $-6.2^\circ \pm 2.3^\circ\text{C}$ in the Δ_{age} -based reconstructions, $-6.0^\circ \pm 2.0^\circ\text{C}$ in the borehole reconstructions, and $-5.9^\circ \pm 2.7^\circ\text{C}$ in the PMIP4 model ensemble; the level of agreement suggests this is a robust feature of Antarctic LGM climate. This temperature contrast is thus plausibly linked to Δz through the (spatial) lapse rate in the interior of Antarctica of around $-12^\circ\text{C km}^{-1}$ (2, 24).

To further assess the elevation impact on ΔT_S , we perform an atmosphere-ocean general circulation model (AOGCM) sensitivity study of Antarctic LGM climate using the MIROC (Model for Interdisciplinary Research on Climate) and HadCM3 (Hadley Centre Coupled Model version 3) models and a series of LGM topographic reconstructions (10). We first estimated climatic LGM cooling using full LGM boundary conditions (including LGM albedo) but preindustrial Antarctic topography; this yielded a seven-site average ΔT_{CLIM} of -4.7° and -7.0°C in the MIROC and HadCM3 models, respectively, but stronger albedo-driven cooling is found over the Ross and Weddell Seas due to ice growth onto the continental

shelf (Fig. 4A). Simulated climatic ΔT_{CLIM} is similar in interior West and East Antarctica in the absence of topographic change.

Next, we performed climate simulations with five Antarctic LGM topographic reconstructions. These reconstructions suggest Δz of $+100$ to $+600$ m in interior WAIS and down to -250 m in interior East Antarctica (Fig. 4B). These changes result in greater ΔT_S in West than in central East Antarctica (Fig. 4C), in agreement with our reconstructions. By comparing the various topographic reconstructions, we find that ΔT_S is closely linked to Δz in both models through the dry adiabatic lapse rate of $-9.8^\circ\text{C km}^{-1}$ (Fig. 4D). Also, a fraction of the variance is due to the topography altering the atmospheric circulation around Antarctica, rather than the direct lapse-rate effect. We find a correlation $r = 0.96$ between the reconstructed and the simulated site ΔT_S pattern (averaged across the five topographic reconstructions and both models); for the PMIP4 multimodel mean, this correlation is $r = 0.95$. We conclude that changes in LGM ice sheet topography plausibly explain the ΔT_S spatial variability in our reconstruction (8).

Our findings have implications for the interpretation of water isotopes in Antarctic ice cores. We found α_T in the range of 0.9 to 1.4‰K^{-1} in East Antarctica and, therefore, $\alpha_T > \alpha_S$, opposite to Greenland, where $\alpha_T < \alpha_S$ (17, 25). We compared our α_T with those from LGM and preindustrial simulations using the latest-generation isotope-enabled Community Earth System Model (iCESM) (Fig. 4E). The

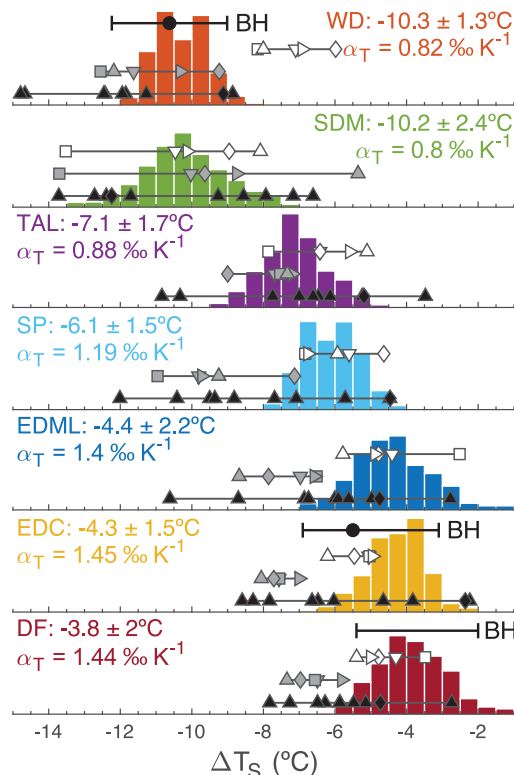


Fig. 1. Summary of Antarctic LGM cooling estimates. Black markers with horizontal error bars marked “BH” give borehole estimates; WD results are from (14). Histograms give distribution of Δ_{age} -based temperature reconstructions from a Monte Carlo sampling ($N = 1000$) of model parameters; listed are mean and 2σ standard deviation of the distribution, as well as the implied temporal isotope slope α_T . ΔT_S is the LGM (18 to 21.4 ka BP) minus preindustrial (0.5 to 2.5 ka BP) condition. White (MIROC), gray (HadCM3), and black (PMIP4) show AOGCM-simulated ΔT_S , with symbols denoting different LGM topography reconstructions (10): Pollard and Deconto (downward triangle) (32); Whitehouse *et al.* (square) (33); Glac-1D (diamond) (29); Gollidge *et al.* (rightward triangle) (34); and Ice-6G (upward triangle).

good agreement ($r = 0.91$; 0.06 ‰K^{-1} mean offset) demonstrates that our reconstructed α_T are consistent with isotope physics, yet the large intermodel spread in simulated α_T [see section S3.5 in (10) for a review] prevents us from claiming that it validates our results. Although the α_T agree well, iCESM simulated a ΔT_S and LGM-preindustrial $\delta^{18}\text{O}$ change that are both too large (compared with our reconstructions and ice core data, respectively).

Last, we investigated changes to the strong surface-based inversion in the Antarctic boundary layer (Fig. 4F). The condensation temperature T_C is higher than T_S , and they correlate spatially with a slope dT_C/dT_S in the 0.63 to 0.67 range (2, 13, 26). T_C controls precipitation $\delta^{18}\text{O}$, with a present-day spatial sensitivity of $d\delta^{18}\text{O}/dT_C = d\delta^{18}\text{O}/dT_S \times dT_S/dT_C \approx 0.80/0.65 = 1.23 \text{ ‰K}^{-1}$. We now assume that, unlike ΔT_S , the LGM-preindustrial change ΔT_C

can be estimated by using this spatial slope via $\Delta T_C = \Delta\delta^{18}\text{O}/1.23$ (Fig. 4F). At WD and SDM, the $\alpha_T \approx \alpha_S$ assumption holds, suggesting that the ratio $\Delta T_C/\Delta T_S$ is close to the present-day ratio of 0.65; in central East Antarctica, the ratio $\Delta T_C/\Delta T_S$ exceeds 0.65, which is consistent with $\alpha_T > \alpha_S$. We plotted simulated ΔT_S versus ΔT_C across interior Antarctica from a wide range of AOGCMs and topographies; we found that the ratio $\Delta T_C/\Delta T_S$ ranges from 0.48 to 1.3 (95% interval, gray lines), with our empirical reconstructions falling within the model data cloud (Fig. 4F). In aggregate, these simulations find that $\Delta T_C/\Delta T_S$ tends to exceed the present-day ratio of 0.65 (~79% of model data points); such a change to the inversion structure would result in $\alpha_T > \alpha_S$ for ΔT_S . In the iCESM simulations, the $\Delta T_C/\Delta T_S$ and α_T fields look similar, with the $\Delta T_C/\Delta T_S = 0.65$ contour line broadly aligning with the $\alpha_T = 0.8 \text{ ‰K}^{-1}$ contour line (fig. S11). We conclude that physically plausible changes to the inversion (27, 28) may reconcile our reconstructions with previous work on Antarctic LGM water isotopes.

Our reconstructions improve the LGM Antarctic temperature estimation and provide a benchmark for testing the ability of (isotope-enabled) climate models to simulate climate states radically different from the late

Fig. 2. Borehole temperature reconstruction for EDC and DF.

(Left) Site borehole temperature observations at EDC (yellow) and DF (red). At both sites, the ice-bedrock interface is at the pressure melting point (-2.2°C). (Right) Model-data mismatch at EDC (yellow) and DF (red) for an ice flow-heat transport model forced by the optimized temperature histories (solid lines, ΔT_S of -5.5°C at EDC and -3.2°C at DF) and forced with water-isotope scaling of 0.7 ‰K^{-1} (dashed lines, ΔT_S of -9.0°C at EDC and -7.5°C at DF).

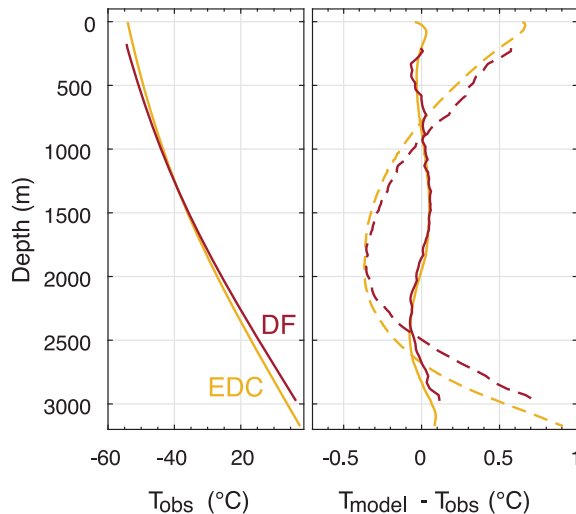
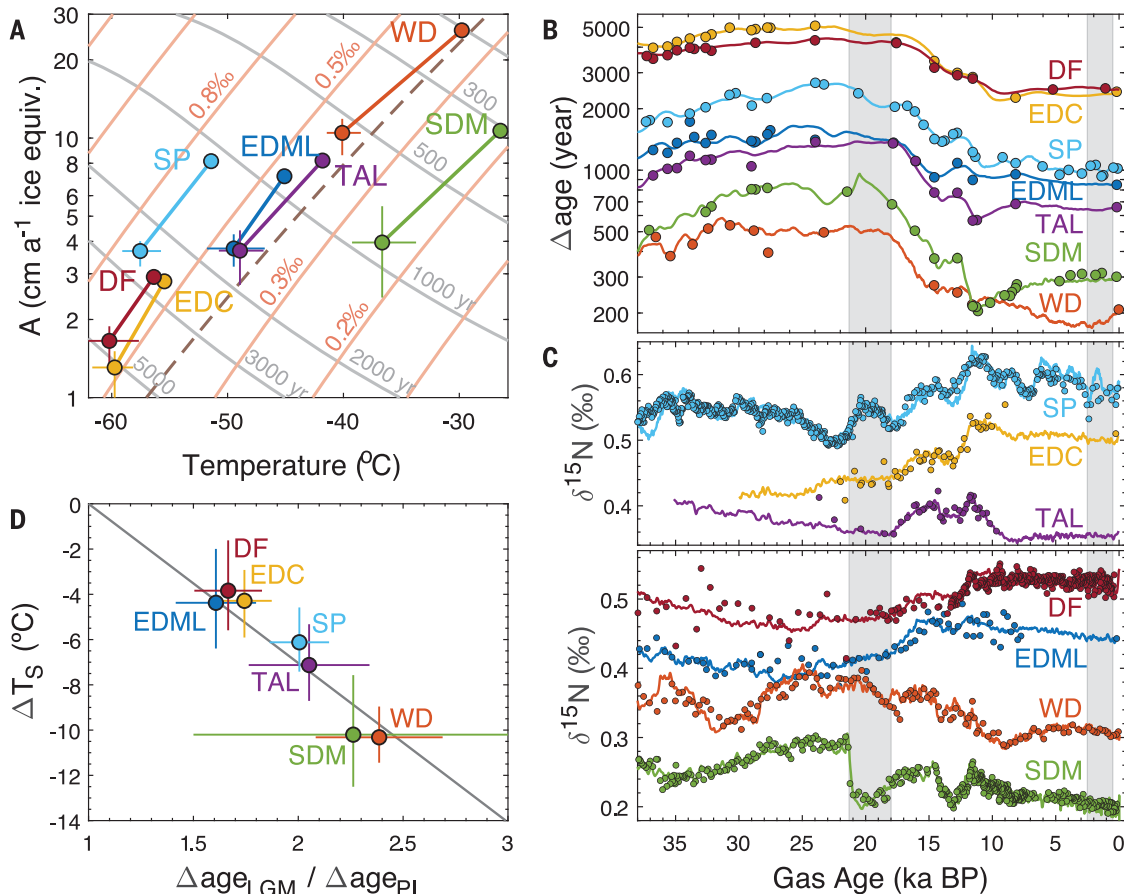


Fig. 3. Δ age-based temperature reconstructions.

(A) Δ age and $\delta^{15}\text{N}$ -isopleths (gray and salmon, respectively) in the steady-state Herron-Langway firn densification model as a function of T_S and A . The dashed line shows accumulation scaling by means of the saturation vapor pressure at the site (ignoring the atmospheric inversion). Reconstructed preindustrial and LGM conditions at the seven sites are indicated. (B) Model fit to empirical Δ age constraints. Gray vertical bars denote the LGM (21.4 to 18 ka BP) and preindustrial (2.5 to 0.5 ka BP) periods. (C) Model fit to $\delta^{15}\text{N}$ data, divided over two panels to prevent overlapping curves. Data are shown on the WD2014 time scale (30, 31). (D) Reconstructed ΔT_S versus ratio of LGM Δ age over preindustrial Δ age (with linear fit) (Δ age_{LGM} / Δ age_{PI}), showing the utility of Δ age as a climate proxy.



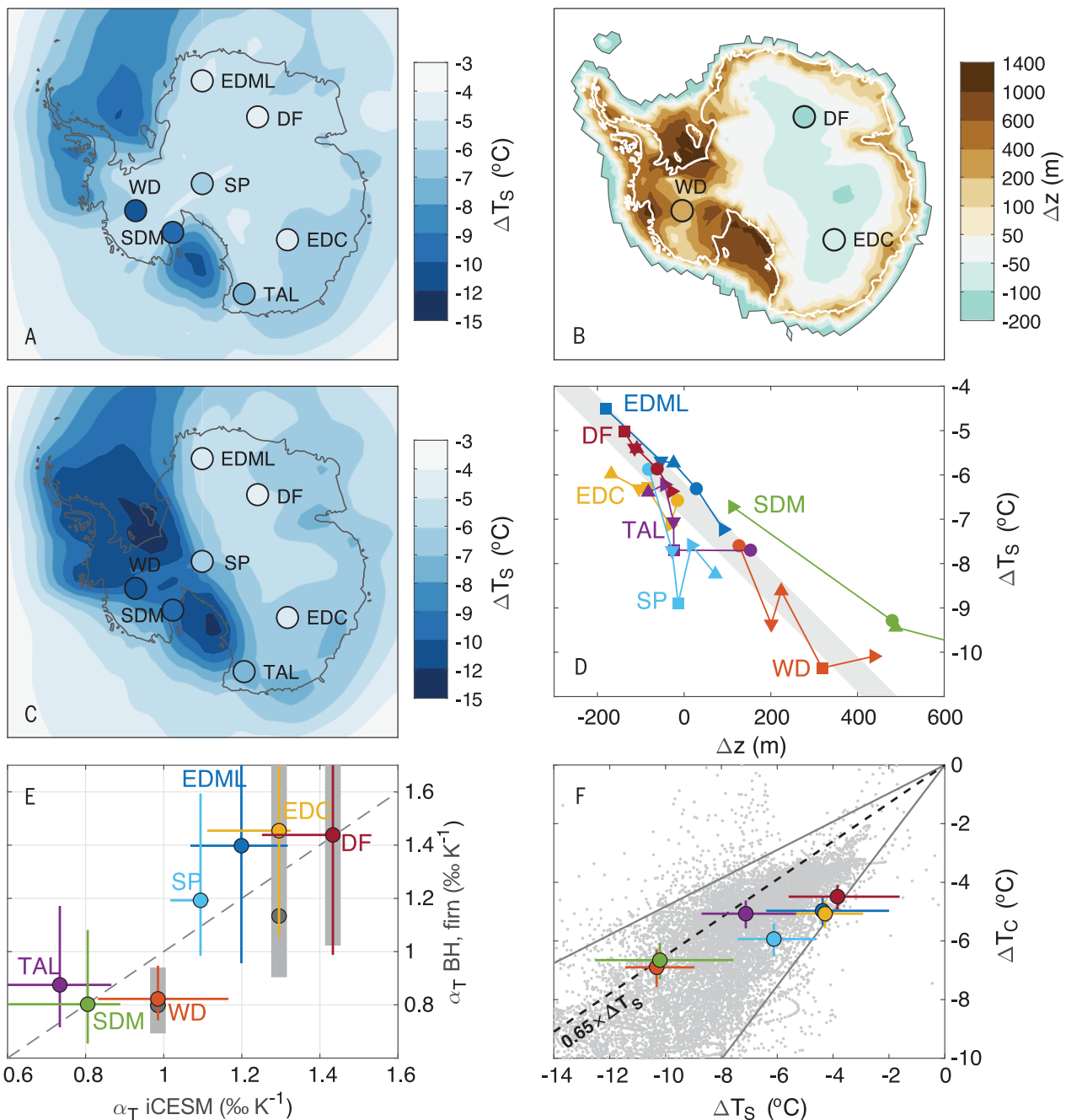


Fig. 4. Climate models and Antarctic topography. (A) AOGCM simulations of ΔT_S using preindustrial ice topography in Antarctica (average of MIROC and HadCM models), with Δ age-based ΔT_S reconstructions for the seven sites. (B) Simulated LGM elevation anomaly (shaded, average of five topographies) with LGM elevation anomaly of +310, -80, and -140 m at WD, EDC, and DF (10). (C) As in (A), but using LGM ice topography in Antarctica (average of five LGM topographies and both MIROC and HadCM models). (D) Elevation change versus ΔT_S in the AOGCM simulations (average of MIROC and HadCM models); symbols denote the different LGM topographic reconstructions (see Fig. 1 caption for legend). The gray bar shows the dry adiabatic lapse rate. (E) Temporal isotope slope α_T from the iCESM model against

Holocene. For surface temperature, the spatial isotopic slope is not always a good approximation of the temporal slope, challenging the prevalent interpretation of ice core water isotopes in Antarctica.

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SUPPLEMENTARY MATERIALS

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Materials and Methods
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Data S1

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Antarctic surface temperature and elevation during the Last Glacial Maximum

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Antarctic paleotemperatures

It has been widely thought that East Antarctica was -9°C cooler during the Last Glacial Maximum, close to the $\sim 10^{\circ}\text{C}$ difference between then and now determined independently for West Antarctica. Buizert *et al.* used borehole thermometry, firn density reconstructions, and climate modeling to show that the temperature in East Antarctica was actually only -4° to 7°C cooler during the Last Glacial Maximum. This result has important consequences for our understanding of Antarctic climate, polar amplification, and global climate change.

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