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Inferring the Subsurface Geometry and Strength of Slow-Moving Landslides Using 3-D Velocity Measurements From the NASA/JPL UAVSAR

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Authors

Handwerger, Alexander L
Booth, Adam M
Huang, Mong-Han
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Peer reviewed

1 **Inferring the Subsurface Geometry and Strength of Slow-moving**
2 **Landslides using 3D Velocity Measurements from the NASA/JPL UAVSAR**

3
4 **A. L. Handwerger^{1,2}, A. M. Booth³, M.-H. Huang⁴, and E. J. Fielding²**

5
6 ¹Joint Institute for Regional Earth System Science and Engineering, University of California, Los
7 Angeles, CA, USA

8 ²Jet Propulsion Laboratory, California Institute of Technology, Pasadena, CA, USA

9 ³Department of Geology, Portland State University, Portland, OR, USA

10 ⁴Department of Geology, University of Maryland, College Park, MD, USA

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13 Corresponding author: Alexander L. Handwerger (alexander.handwerger@jpl.nasa.gov)

14 **Key Points:**

- 15 • Landslide thickness can vary by tens of meters within a single landslide
- 16 • The largest landslide complexes get larger by increasing area rather than increasing
- 17 thickness
- 18 • Landslide strength is scale-dependent, such that large landslides tend to be weaker than
- 19 small landslides

26 **Abstract**

27 The hazardous impact and erosive potential of slow-moving landslides depends on
28 landslide properties including velocity, size, and frequency of occurrence. However, constraints
29 on size, in particular, subsurface geometry, are lacking because these types of landslides rarely
30 fully evacuate material to create measurable hillslope scars. Here we use pixel offset tracking
31 with data from the NASA/JPL Uninhabited Aerial Vehicle Synthetic Aperture Radar (UAVSAR)
32 to measure the three-dimensional surface deformation of 134 slow-moving landslides in the
33 northern California Coast Ranges. We apply volume conservation to infer the actively deforming
34 thickness, volume, geometric scaling, and frictional strength of each landslide. These landslides
35 move at average rates between $\sim 0.1\text{--}3$ m/yr and have areas of $\sim 6.1 \times 10^3\text{--}2.35 \times 10^6$ m², inferred
36 mean thicknesses of $\sim 1.1\text{--}25$ m, and volumes of $\sim 7.01 \times 10^3\text{--}9.75 \times 10^6$ m³. The best-fit volume-
37 area geometric scaling exponent is $\gamma \sim 1.2\text{--}1.5$, indicating that these landslides fall between
38 typical soil and bedrock landslide scaling. A rollover in the scaling relationship suggests that the
39 largest landslide complexes in our dataset become large primarily by increasing in area rather
40 than thickness. In addition, the slow-moving landslides display scale-dependent frictional
41 strength, such that large landslides tend to be weaker than small landslides. This decrease in
42 frictional strength with landslide size is likely because larger landslides are composed of higher
43 proportions of weak material. Our work shows how state-of-the-art remote sensing techniques
44 can be used to better understand landslide processes and quantify their contribution to landscape
45 evolution and hazards to human safety.

46 **1 Introduction**

47 Landslides are a major natural hazard and are often the dominant process that erodes
48 mountainous landscapes (Korup et al., 2007; Larsen et al., 2010; Mackey & Roering, 2011;

49 Simoni et al., 2013). Both their hazardous impact and erosive potential depend on landslide
50 properties including the velocity, size, and frequency of occurrence. Measuring these landslide
51 properties is challenging because landslides exhibit a wide range of velocities (mm/yr to m/s),
52 spatial areas ($10^0 - 10^8 \text{ m}^2$), and volumes ($10^{-1} - 10^{10} \text{ m}^3$), and can occur in large numbers
53 (hundreds to tens of thousands) over broad spatiotemporal scales (Cruden & Varnes, 1996;
54 Hungr et al., 2014; Lacroix, Handwerger, et al., 2020; Larsen et al., 2010). Importantly, the
55 landslide failure style also impacts our ability to measure landslide properties, such as thickness
56 and volume, which can strongly influence runout and erosion rate (e.g., Korup et al., 2007;
57 Larsen et al., 2010; Legros, 2002). Some landslides create clear and identifiable scars and
58 deposits by evacuating material from the hillslope, making it possible to directly measure
59 landslide properties from field data, digital elevation models (DEMs), and remote sensing
60 observations (e.g., Bessette-Kirton et al., 2018; Warrick et al., 2019; Wartman et al., 2016).
61 However, for landslides that move slowly for years or centuries (Lacroix, Handwerger et al.,
62 2020; Mackey et al., 2009; Rutter & Green, 2011), referred to as slow-moving landslides, and do
63 not create hillslope scars, it is difficult to constrain their thickness and volume because data are
64 usually limited to isolated point measurements from boreholes (Schulz et al., 2018; Simoni et al.,
65 2013; Travelletti & Malet, 2012), which do not capture the spatial variability exhibited by these
66 landslides. It is therefore advantageous to develop and apply tools and methods that can be used
67 to construct large inventories of slow-moving landslides and quantify their surface and
68 subsurface properties.

69 Modern remote sensing tools, such as synthetic aperture radar (SAR), optical imagery,
70 and lidar, provide high-resolution measurements of topography and ground surface deformation
71 that can be used to identify and monitor landslides with millimeter- to centimeter-scale accuracy

72 at spatial resolutions of a few centimeters to hundreds of meters. Recent work using pixel offset
73 tracking and SAR interferometry with these data has quantified the two-dimensional (2D) and
74 three-dimensional (3D) surface deformation of slow-moving landslides (Aryal et al., 2015;
75 Booth et al., 2020; Hu et al., 2020; Lacroix, Dehecq et al., 2020; Stumpf et al., 2017; Travelletti
76 et al., 2014). These studies, along with numerous ground-based investigations (e.g., Iverson &
77 Major, 1987; Malet et al., 2002; Schulz et al., 2017), have shown that slow-moving landslides
78 exhibit non-uniform spatial and temporal kinematic patterns. In addition, high-resolution 3D
79 surface deformation measurements can be used to infer the thickness and subsurface geometry of
80 the actively moving part of the landslide. Previous studies (Aryal et al., 2015; Booth et al., 2020;
81 Booth, Lamb, et al., 2013; Delbridge et al., 2016; Hu et al., 2020) have suggested that active
82 landslide thickness can vary by tens of meters within a single landslide, and the slip surfaces
83 have an irregular and bumpy morphology that differs considerably from commonly assumed,
84 idealized geometric forms, such as semicircles, ellipsoids, and log spirals (see a detailed review
85 paper by Michel et al., 2020). These large changes in thickness within a single landslide mass
86 have important implications for estimating volume and sediment flux, designing field
87 instrumentation and landslide mitigation strategies, and determining the stresses that control
88 landslide kinematics. Although techniques that invert surface observations for subsurface
89 characteristics are becoming more common, most studies have focused on individual landslides
90 occurring under different and site-specific environmental conditions, making it difficult to
91 identify more generic geometric scaling relations for slow-moving landslides.

92 In this study, we use data from the NASA/JPL Uninhabited Aerial Vehicle Synthetic
93 Aperture Radar (UAVSAR) to construct an inventory of 134 active slow-moving landslides in a
94 ~ 1621 km² area of the northern California Coast Ranges between 2016 and 2019 (Figure 1).

95 These landslides occur in the Eel River catchment, a region well known for its slow-moving
96 landslides, and are driven by high seasonal rainfall (Bennett, Roering, et al., 2016; Booth,
97 Roering, et al., 2013; Handwerger et al., 2013, 2015; Handwerger, Fielding, et al., 2019;
98 Handwerger, Huang, et al., 2019; Kelsey, 1978; Mackey et al., 2009; Mackey & Roering, 2011;
99 Roering et al., 2009, 2015; Schulz et al., 2018). The landslides are underlain by the Central Belt
100 Franciscan mélangé, a mechanically weak and pervasively sheared bedrock with an argillaceous
101 matrix that surrounds blocks of stronger rock types, including sandstone, chert and greenstone
102 (Jayko et al., 1989; Jennings et al., 1977; McLaughlin et al., 1982, 2000). We measure the 3D
103 surface deformation and geometry of each landslide, and use these data in a volume conservation
104 framework to invert for their active thickness, volume, and strength. We derive new geometric
105 scaling relations for slow-moving landslides and make comparisons with a worldwide inventory
106 of soil and bedrock landslides. Our work is the first to use volume conservation methods to
107 invert for the thickness of a large inventory of landslides, and this approach could be applied to
108 other groups of slow-moving landslides around the world. Our work also shows how state-of-
109 the-art remote sensing techniques can be used to better understand landslide processes and
110 quantify their contribution to landscape evolution.

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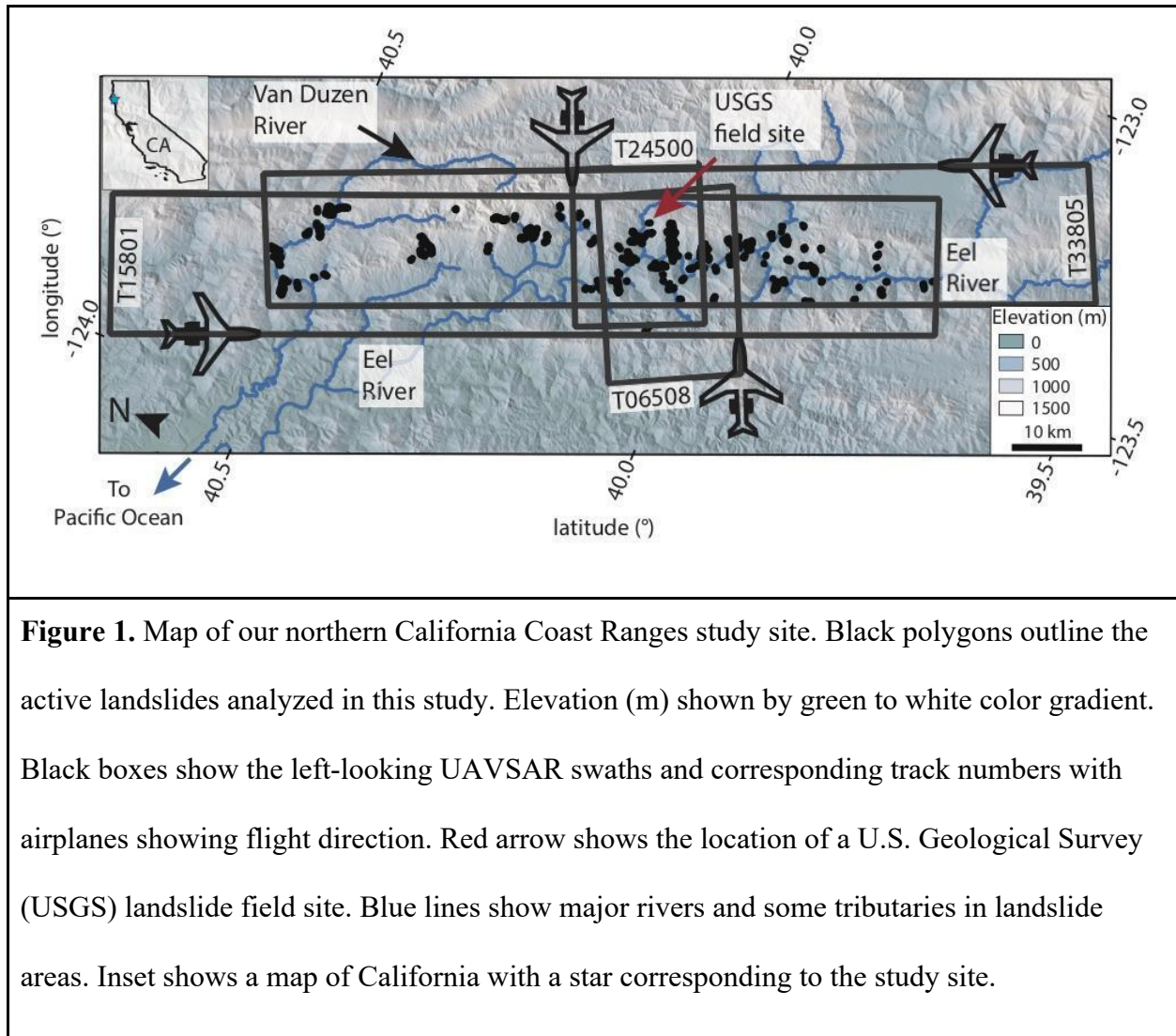
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119 2 Materials and Methods

120 2.1 UAVSAR Data and Processing

121 We use SAR data acquired by the NASA/JPL UAVSAR airborne system for our
 122 landslide investigation. UAVSAR has a left-looking radar attached to a NASA Gulfstream III
 123 airplane that operates with a L-band wavelength (~ 23.8 cm) and a swath width of ~ 20 km. The
 124 NASA Gulfstream III autopilot flies at 13 km above sea level and repeats the flight lines within a
 125 five-meter radius tube, so the spatial baselines are always short and have no impact on

126 deformation measurements. UAVSAR data have a pixel spacing of 1.67 m in the range direction
127 (measured along the line-of-sight, LOS) and 0.6 m in the azimuth direction (measured along the
128 UAVSAR flight direction). We designed the UAVSAR data collection for the northern
129 California Coast Ranges site specifically to monitor a large quantity of slow-moving landslides
130 that were initially identified by several previous studies (e.g., Bennett, Miller, et al., 2016;
131 Handwerger et al., 2015; Kelsey, 1978; Mackey & Roering, 2011; Roering et al., 2009). Some of
132 these UAVSAR data were used in a recent study by Handwerger, Fielding, et al. (2019) to
133 analyze changes in landslide activity due to extreme rainfall. We collected data on 4 partially
134 overlapping flight paths to increase data redundancy and to provide between 4 and 8 independent
135 deformation measurements (Figure 1). There were 12 data acquisitions at our field site between
136 April 2016 and May 2019. The time between data acquisitions ranges between 47 and 237 days,
137 with a mean of 104 days (Table S1). UAVSAR Single-Look Complex (SLC) data are freely
138 available at <https://uavsar.jpl.nasa.gov/>.

139 We perform pixel offset tracking on the coregistered UAVSAR stack SLC data using the
140 *Ampcor* module, which is part of the JPL InSAR Scientific Computing Environment (ISCE)
141 version 2 software package (Rosen et al., 2012). Pixel offset tracking (sometimes referred to as
142 subpixel correlation) uses cross-correlation between SAR amplitude images to quantify image
143 offsets (i.e., displacement) due to ground surface motion in two dimensions; 1) the range or look
144 direction, and 2) the azimuth or along-track direction (e.g., Fialko et al., 2001; Fielding et al.,
145 2020; Pathier et al., 2006). We use the terms range/look direction and azimuth/along-track
146 direction, interchangeably. Pixel tracking has a precision up to $\sim 1/10$ of the pixel size, which
147 corresponds to ~ 6 cm in the along-track direction and ~ 17 cm in the range direction for a pair of
148 UAVSAR images. Although this technique is less precise than conventional InSAR, it does not

149 involve phase unwrapping and thus is better suited for measuring the decimeter- to meter-scale
150 displacements commonly displayed by many slow-moving landslides (Lacroix, Handwerger, et
151 al., 2020). To account for the differences in the range and along-track pixel size, we use a cross-
152 correlation window length of 128 pixels with a skip size of 32 pixels (distance between matching
153 window calculations) in the along-track direction and a cross-correlation window width of 64
154 pixels with a skip size of 16 pixels in the range direction, resulting in a window size of 77 m by
155 107 m. This cross-correlation window size was found to provide the best landslide deformation
156 signal from UAVSAR pixel offset tracking by Handwerger, Fielding et al. (2019). We geocode
157 the pixel offset measurements to a 0.4 arcsecond (~ 12 m) pixel using the TanDEM-X DEM
158 provided by the German Aerospace Center (DLR). We process all possible combinations of pixel
159 offset tracking pairs, which results in 66 pixel offset tracking maps on each track (264 in total)
160 with single pair time spans ranging from 47 to 1148 days (Table S1). We exclude 35 poor-
161 quality pixel offset tracking maps from our analysis that included a large number of pixels with
162 physically incorrect displacements (e.g., upslope motion or unusually large values) and
163 significant noise that obscured the landslide signals. We found these poor-quality data tend to
164 result from long duration pairs that exceed ~ 2 years, which are subject to numerous changes in
165 the ground surface (e.g., vegetation changes, anthropogenic changes) that can deteriorate the
166 cross-correlation result (Table S1). We convert all of the displacement offset maps to velocities
167 and then take the temporal average of the 31 remaining pixel offset velocity maps to make a
168 mean velocity map for our thickness inversions.

169

170 2.2 Three-dimensional Ground Surface Deformation

171 To solve for 3D deformation from SAR requires at least three independent measurements
 172 of surface deformation. Each UAVSAR flight path provides two independent measurements of
 173 surface motion from pixel offset tracking (i.e., along-track and range). Therefore, using pixel
 174 offset tracking velocity maps, data from at least two flights is required for 3D inversions.
 175 Because UAVSAR acquires data on four different flight paths in our field area (Figure 1), we
 176 have a maximum of eight deformation measurements in the central region of our field area where
 177 all four flight paths overlap and a maximum of two deformation measurements in the northern
 178 and southern extents where only two flight paths overlap. Thus, we are always able to achieve an
 179 overdetermined 3D inversion.

180 Each deformation measurement from pixel tracking is composed of the true displacement
 181 vector projected onto the along-track or range direction of the UAVSAR. We use a least-squares
 182 inversion to isolate the east, north, and vertical components of deformation defined in the form
 183 $\mathbf{d} = \mathbf{Gm}$,

$$184 \begin{bmatrix} v_{rng1} \\ v_{azi1} \\ \vdots \\ v_{rng,M} \\ v_{azi,M} \end{bmatrix} = \begin{bmatrix} \cos \xi_1 \sin \theta_1 & \sin \xi_1 \sin \theta_1 & -\cos \theta_1 \\ \cos \xi_1 & \sin \xi_1 & 0 \\ \vdots & \vdots & \vdots \\ \cos \xi_M \sin \theta_M & \sin \xi_M \sin \theta_M & -\cos \theta_M \\ \cos \xi_M & \sin \xi_M & 0 \end{bmatrix} \begin{bmatrix} v_{ew} \\ v_{ns} \\ v_{ud} \end{bmatrix}, \quad (1)$$

185 where $v_{rng,M}$ is the range (or look direction) velocity, $v_{azi,M}$ is the azimuth (or along-track
 186 direction) velocity, M is the flight path number (minimum of two needed for pixel offset
 187 tracking), ξ is the UAVSAR heading direction (i.e., along track direction) with counterclockwise
 188 as positive, θ is the UAVSAR look angle, and v_{ew} , v_{ns} , v_{ud} are the east-west, north-south, and
 189 vertical components of velocity, respectively.

190 The overdetermination of the 3D inversion allows us to constrain the uncertainty from the
 191 inversion (e.g., Delbridge et al., 2016). To constrain the inversion uncertainty, we repeat the 3D

192 inversion multiple times using different combinations of v_{rng} and v_{azi} . For instance, for landslides
193 with eight deformation measurements (i.e., four range and four azimuth measurements), we
194 perform the 3D inversion 198 times using between three and eight deformation measurements.
195 We then take the mean and standard deviation of all of the inversions and use these values as the
196 3D velocities and inversion uncertainty, respectively. We further constrain the uncertainty in our
197 velocity measurements by examining the apparent deformation rate of stable hillslopes. To
198 reduce noise and error (i.e., unrealistically large displacements), we apply velocity thresholds
199 and mask out pixels with apparent velocities > 50 m/yr, which is much faster than the typical
200 velocity range displayed by the northern California Coast Ranges landslides (Bennett, Roering,
201 et al., 2016; Handwerger, Fielding, et al., 2019; Roering et al., 2015). We also mask out pixels
202 that have mean velocities less than their inversion uncertainty and use nearest neighbor
203 interpolation with a five pixel maximum radius to fill in these masked pixels.

204

205 2.3 Landslide Thickness Inversion

206 We use 3D surface velocity measurements from pixel offset tracking to infer the
207 thickness, volume, and shear zone geometry of the active parts of each landslide using a
208 conservation of volume approach. We apply the method originally described by Booth, Lamb, et
209 al., (2013) and more recently by Booth et al. (2020), which assumes that during our ~ 3 year
210 study period, the measured surface velocity is representative of the depth-averaged velocity, the
211 sliding surface does not change in time, there is minimal direct erosion or deposition of the
212 landslide surface, and the landslide material density is uniform and constant. While landslides
213 may violate these assumptions in general, they are reasonable for our study area for the following
214 reasons: (1) at the Two Towers landslide, a U.S. Geological Survey (USGS) instrumented

215 landslide in our study site (Schulz et al., 2018), the measured surface velocity was approximately
 216 equal to the depth-averaged velocity, and a narrow shear zone was identified (Figure S1); (2) the
 217 landslides were continuously active with fixed spatial boundaries over the time periods that 3D
 218 displacements were measured, suggesting movement on the same slip surface; (3) minor
 219 amounts of direct surface erosion or deposition were likely confined to gully systems on the
 220 landslides' surfaces, which occupy a small percentage of the landslides' surface area (~1%) and
 221 therefore have a minimal effect on the inversion; and (4) dilation/compaction or
 222 shrinking/swelling that would cause changes in density is likely on the order of centimeters or
 223 less (Booth et al., 2020; Delbridge et al., 2016; Iverson, 2005; Schulz et al., 2018), which is
 224 typically small compared to surface velocity gradients, thus having limited influence of the
 225 measured 3D surface velocity. Therefore, for a landslide of constant density with no erosion or
 226 deposition, conservation of volume implies that

$$227 \quad v_{ud} = -\nabla \cdot (\bar{u}h) + u_{surf} \cdot \nabla z_{surf}, \quad (2)$$

228 where v_{ud} is the vertical component of the 3D landslide surface velocity vector, h is the active
 229 landslide thickness, u_{surf} is the vector of horizontal components of landslide surface velocity, \bar{u} is
 230 the depth-averaged vector of horizontal components of landslide velocity, and z_{surf} is the surface
 231 elevation measured from the ~12 m TanDEM-X DEM. The first term on the right-hand side of
 232 equation 2 is the contribution of flux divergence to the vertical component of the surface
 233 velocity, and the second term is the contribution due to advection of the sloped land surface.
 234 Because UAVSAR measures the velocity of the ground surface, u_{surf} , we assume that
 235 $\bar{u} = fu_{surf}$, where f is a constant that characterizes the thickness of the shear zone at the base of
 236 the landslide relative to the total landslide thickness. We constrain f using borehole inclinometer
 237 data from two boreholes at the USGS field station on the Two Towers landslide (supporting

238 information and Figure S1). Unfortunately, the Two Towers landslide is not detectable with pixel
 239 tracking from UAVSAR data because the landslide is small (250 m long and 40 m wide) and
 240 moving too slowly (maximum speed ~ 6 cm/yr) (Schulz et al., 2018). Using these data, we find
 241 that $f \sim 0.96$, which indicates that the landslide moves along a narrow shear zone with the
 242 material above translating essentially as a rigid block. For simplicity, we assume that $f = 1$ and
 243 that the landslides move as a rigid block. Other studies in California (e.g., Keefer & Johnson,
 244 1983, Swanston et al., 1995) and around the world (e.g., van Asch & van Genuchten, 1990;
 245 Simoni et al., 2013) have also found that similar type slow-moving landslides move as a rigid
 246 plug above a narrow shear zone such that $f \sim 1$ is a reasonable approximation, however more
 247 ground-based investigations are required to better constrain the f parameter for multiple
 248 landslides. Although f generically represents the ratio of depth-averaged to surface velocity, it
 249 can be related to specific rheologies if desired (Booth, Lamb, et al., 2013; Delbridge et al., 2016)
 250 and we discuss the implications of different f values in Section 4.2.

251 Incorporating f into equation 2 gives

$$252 \quad v_{ud} = -\nabla \cdot (f u_{surf} h) + u_{surf} \cdot \nabla z_{surf}, \quad (3)$$

253 which is a statement of conservation of volume in a Lagrangian reference frame (Booth et al.,
 254 2020; Delbridge et al., 2016). We discretize equation 3 using centered finite differences,
 255 rearrange it as a system of linear equations, and then solve for thickness by minimizing the value
 256 of

$$257 \quad |Xh - b|^2 + \alpha^2 |\nabla^2 h|^2, \quad (4)$$

258 subject to non-negative constraints,

259 where X is a diagonally dominant matrix that contains the depth-averaged horizontal velocity
 260 data, b is a vector defined as $u_{surf} \cdot \nabla z_{surf} - v_{ud}$, and α is a damping parameter to regularize

261 the ill-posed inverse problem. Since both the matrix X and the vector b contain data with
262 uncertainties, and the damping parameter necessarily introduces bias, estimating total uncertainty
263 of the resulting thickness model is not straightforward. However, we make a minimum estimate
264 following standard techniques from inverse theory, which reflects uncertainty in b only
265 (supporting information). We explore a wide range of α from 10^{-3} to 10^1 and determine the best
266 level of regularization using the Generalized Cross-Validation method (supporting information
267 and Figure S2). We resample our ~ 12 m pixel spacing grid to square 10 m x 10 m pixel and
268 perform the thickness inversion in the MATLAB software package using the CVX program, a
269 package for specifying and solving convex programs (Grant & Boyd, 2014). For the largest
270 landslide in our inventory (i.e., Boulder Creek landslide complex) we had to downsample the
271 grid to a 20 m x 20 m pixel due to computational limitations. The inferred thickness values
272 represent the best solution that does not violate conservation of volume and assumes that the
273 surface velocity is equal to the depth-averaged velocity.

274 It is important to further emphasize that the thickness inversions are only relevant to the
275 active parts of landslides such that there needs to be detectable surface deformation to invert for
276 the landslide thickness. Specifically, the values of b (equation 4) need to differ from background
277 values on known stable ground to infer non-zero thicknesses. Landslides or areas and kinematic
278 zones within landslides that are not moving are therefore considered to have zero depth.
279 Landslide thickness in this study therefore specifically means the “active thickness” during our
280 study period.

281

282 2.4 Landslide Inventory and Geometric Scaling

283 To select landslides for 3D surface velocity and thickness inversions, we assemble a new
284 inventory of active landslides in our ~1621 km² study area in the northern California Coast
285 Ranges that includes only those landslides that show a significant deformation signal using the
286 pixel offset tracking method. This limits our analysis to the faster-moving landslides that exhibit
287 rates of decimeters to meters per year. Our landslide inventory was guided by a number of pre-
288 existing landslide inventories for the northern California Coast Ranges (Bennett, Miller, et al.,
289 2016; Handwerger, Fielding, et al., 2019; Kelsey, 1978; Mackey & Roering, 2011). We map the
290 landslide boundaries in QGIS using the 3D velocity maps, hillshade maps constructed from 1 m
291 pixel spacing lidar provided by OpenTopography (Roering, 2012), the ~12 m pixel spacing
292 TanDEM-X DEM, and Google Earth imagery. Because slow-moving landslides display non-
293 uniform spatial kinematic zones and complex kinematic histories (e.g., Nereson & Finnegan,
294 2019; Schulz et al., 2017; Stumpf et al., 2017), there are often differences between the landslide
295 boundaries mapped with kinematic data and those mapped based on geomorphic interpretation of
296 hillshades or aerial photos. These differences in mapping are especially important for our
297 thickness inversions because including the parts of landslides that are not currently moving can
298 cause the thickness inversion to produce unreliable results. Therefore, we use the temporally
299 averaged landslide velocity and only map areas of each landslide that are moving during our
300 study period. For larger landslides with multiple kinematic zones, we perform separate thickness
301 inversions for any isolated, faster-moving areas of the landslide, as well as for the entire
302 landslide complex as a whole. If results had substantially different spatial patterns of thickness,
303 we adopt the more reliable results for the smaller isolated landslides. We use QGIS to quantify
304 the spatial metrics of each landslide, including length, average width (defined as area divided by

305 length), area, and slope angle. We also report the mean, median, 75th percentile, and maximum
 306 horizontal velocity, 3D velocity magnitude, and 3D inversion velocity errors for each landslide.

307 We then derive empirical geometric scaling relations for landslide thickness (h) and
 308 volume V from the measured landslide area A . Geometric scaling relations are commonly used to
 309 quantify erosion rates of large inventories of landslides and are important for understanding
 310 landslide mechanics (e.g., Bunn et al., 2020a; Guzzetti et al., 2009; Larsen et al., 2010; Milledge
 311 et al., 2014). Larsen et al. (2010) showed that these scaling relations hold over 9 orders of
 312 magnitude in area and 12 orders of magnitude in volume. Landslide scaling relations take the
 313 form of a power function where

$$314 \quad V = c_V A^\gamma \text{ and } h = c_h A^\zeta, \quad (5a \text{ and } 5b)$$

315 where γ and ζ are scaling exponents and c_V and c_h are the intercepts. We constrain the
 316 coefficients of these power functions by log-transforming our data and finding the best-fit
 317 parameters with 95% confidence intervals using a linear least square inversion in MATLAB.

318

319 2.5 Frictional Strength

320 We estimate the frictional strength of each landslide by following the 3D Simplified
 321 Janbu method (Bunn et al., 2020b; Hungr, 1987; Hungr et al., 1989; Leshchinsky, 2019). This
 322 method assumes that the vertical intercolumn shear forces are negligible. Each landslide is
 323 discretized into 3D columns with a surface area S_{basal} and total weight W . The basal surface area
 324 is defined by

$$325 \quad S_{basal} = \Delta x \Delta y \frac{(1 - \sin^2 \beta_y \sin^2 \beta_x)^{1/2}}{\cos \beta_y \cos \beta_x}, \quad (6)$$

326 where Δx and Δy are the grid spacing in the x and y direction, respectively, β_x is the local dip
 327 angle perpendicular to the direction of motion and β_y is the local dip in the direction of motion.

328 The normal force N at the base of each column is defined by

$$329 \quad N = \frac{W - C S_{basal} \sin \beta_x / F + p S_{basal} \tan \phi \sin \beta_x / F}{\cos \Delta_z \left(1 + \frac{\sin \beta_x \tan \phi}{F \cos \Delta_z} \right)}, \quad (7)$$

330 where p is the mean pore pressure acting at the base of each column, C is the cohesion, ϕ is the
 331 residual friction angle, F is the factor of safety, and Δ_z is the local dip angle defined in terms of
 332 the motion-parallel and motion-perpendicular dips by

$$333 \quad \cos \Delta_z = \left(\sqrt{\frac{1}{1 + \tan^2 \beta_y + \tan^2 \beta_x}} \right). \quad (8)$$

334 Finally, F is defined by

$$335 \quad F = \frac{\sum C S_{basal} \cos \beta_x + (N - p S_{basal}) \tan \phi \cos \beta_x}{\sum N \cos \Delta_z \tan \beta_x}, \quad (9)$$

336 where the summation is over all columns. The numerator is the resisting force, with the term in
 337 the parentheses defining the effective normal force, and $\tan \phi$ is the friction coefficient, and the
 338 denominator is the shear force. We assume that cohesion is negligible since these landslides are
 339 moving, some of which have been moving for decades (Mackey and Roering, 2011). We set $F =$
 340 1 (i.e., balanced forces at failure) and solve for friction angle under both dry and fully saturated
 341 (hydrostatic conditions) end members to produce a minimum and maximum estimate. Table S2
 342 shows the dry and wet landslide density values used for our calculations. Recent work by Bunn
 343 et al., (2020b) used a similar approach to infer the strength of several hundred landslides in
 344 Oregon, USA.

345

346 **3 Results**

347 **3.1 Landslide Inventory and 3D Velocity**

348 We identified 134 active landslides in our northern California Coast Ranges field site
349 (Figure 1), 19 of which were unmapped by previous studies (Bennett, Miller, et al., 2016;
350 Handwerger, Fielding, et al., 2019; Mackey & Roering, 2011). These landslides have average
351 widths from 66 to 556 m, lengths from 68 to 4727 m, areas from 7.8×10^3 to 2.63×10^6 m², and
352 mean slope angles from 10 to 29 degrees (Table S3). Each landslide exhibited a non-uniform
353 spatial velocity pattern (see examples in Figure 2). The spatial kinematic patterns remain fixed
354 during our study period and are similar to those mapped in previous studies (see Bennett,
355 Roering, et al., 2016; Handwerger, Fielding, et al., 2019; Mackey & Roering, 2011). The
356 maximum 3D velocity magnitude of the individual landslides, calculated as $v_{3D} =$
357 $(v_{ns}^2 + v_{ew}^2 + v_{ud}^2)^{1/2}$, ranged from 0.198 to 8.58 m/yr. The average 3D velocity magnitude of the
358 individual landslides ranged from 0.123 to 2.11 m/yr. The landslide motion was always primarily
359 in the downslope direction (see example in Figures 2e and 2f), but at different locations we do
360 measure areas of both uplift and subsidence within a single landslide (see example in Figure 2d).
361 We note that local surface uplift occurs when the vertical component of the velocity vector dips
362 less steeply than the topographic surface at a given point. As a result, the vertical velocity is
363 often still negative even in areas where the topographic surface is locally being uplifted, and only
364 when the vertical motion is upwards relative to horizontal do we observe positive vertical
365 velocities. The mean 3D velocity uncertainty from the 3D inversion (equation 1) for the
366 individual landslides ranged from 0.0179 to 1.91 m/yr. We report the full uncertainty statistics
367 for each individual landslide in Table S3. The 3D velocity magnitude uncertainty from
368 examining the apparent velocity of stable hillslopes was ≤ 0.1 m/yr.

369 We classified the slow-moving landslides into three subgroups based on their geometry
370 and kinematic patterns (Table S3). Figure 2 shows three example landslides which we define as

371 slumps, earthflows, and landslide complexes. The landslide complex shown in Figure 2 is the
372 largest landslide in our dataset and is also known as the Boulder Creek landslide in several other
373 studies (e.g., Bennett, Miller, et al., 2016; Bennett, Roering, et al., 2016; Handwerger et al.,
374 2015, 2015; Handwerger, Fielding, et al., 2019; Handwerger, Huang, et al., 2019; Mackey &
375 Roering, 2011; Roering et al., 2009). We defined slumps as landslides with lower length/width
376 aspect ratios (median = 1.57 ± 1.00 , ± 1 standard deviation), a strong signal of positive vertical
377 velocity components in the toe and negative vertical velocity components in the source area, and
378 one primary kinematic zone (Figure 2a). We defined earthflows as those with medium aspect
379 ratios (median = 3.56 ± 1.88 , ± 1 standard deviation), one primary kinematic zone, and small
380 magnitude, but mostly negative, vertical velocity components (Figure 2b). And we defined
381 landslide complexes as those with higher aspect ratios (median = 5.13 ± 2.34 , ± 1 standard
382 deviation), that are composed of multiple kinematic zones or even multiple landslides that
383 coalesce into a single landslide mass (Figure 2c). Landslide complexes are relatively common in
384 areas with slow-moving landslides (e.g., Cerovski-Darriau & Roering, 2016; Keefer & Johnson,
385 1983; Simoni et al., 2013). 33% of our inventory were classified as slumps, 31% as earthflows,
386 and 36% as landslide complexes. The mean 3D velocity magnitude was 0.585, 0.606, and 0.670
387 m/yr for slumps, earthflows, and landslide complexes, respectively.

388

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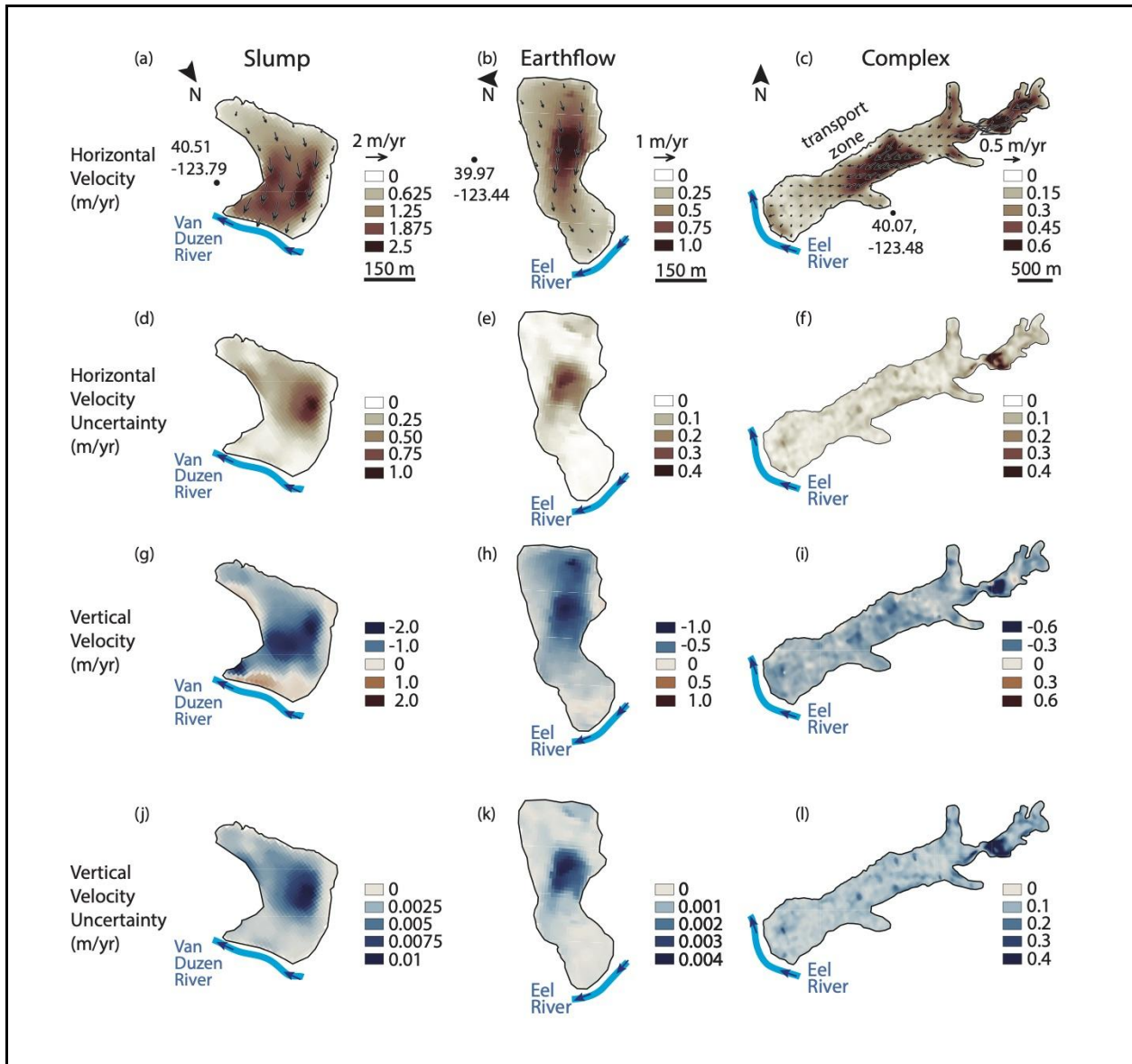


Figure 2. 3D velocity maps for example slump, earthflow, and landslide complex. (a–c) Horizontal velocity maps. Black arrows show horizontal vectors. Black circle shows latitude and longitude coordinates. (d–f) Horizontal velocity inversion uncertainty maps. (g–i) Vertical velocity maps for the three landslides. (j–l) Vertical velocity inversion uncertainty maps. Negative values correspond to vertically downward motion. Thick blue lines show the approximate location of the river channel at the toe of each landslide with dark blue arrows showing water flow direction.

390 3.2 Thickness, Volume, and Geometric Scaling Relations

391 The non-uniform kinematic patterns exhibited by these landslides are also reflected in
392 their inferred subsurface geometry (Figure 3). We find that the thickness of each landslide varies
393 spatially and can vary by tens of meters within the landslide boundaries. The slip surfaces are
394 generally concave-up, but are rough and irregular in places, especially for landslide complexes.
395 The mean active thickness of the individual landslides ranged from 0.4 to 22.4 m, and the
396 maximum active thickness ranged from 2.25 to 89.6 m. The mean, median, minimum, maximum,
397 and standard deviation active thickness for each landslide are reported in Table S3.

398 We calculated the minimum thickness uncertainty from uncertainties in the data in vector
399 b following standard inverse theory for a sample of seven landslides representing the variety of
400 style, size, and shape found in the study population (supporting information). We found that
401 minimum thickness uncertainty increased with landslide size (Figure S3), ranging from ± 1.5 to
402 ± 3.8 m from the smallest to largest landslide sampled. To reduce computation time, we estimated
403 the minimum thickness uncertainty for each landslide using a power function (Figure S3d) and
404 propagated these uncertainties into the landslide volume calculations (Table S3).

405 Next we describe our thickness inversion results for the three example types of landslides
406 shown in Figure 2. We note again that these landslides represent their subgroups to first order.
407 The example slump has one primary deep zone and the slip surface has a concave-up profile
408 (Figure 3a). The slope of the slip surface deviates from the ground surface and is steeper near the
409 headscarp and gentler near the toe. Some areas within the head of the landslide are inferred to
410 have no active thickness because the values of b (equation 4) are slightly negative near the
411 headscarp (Figure S4). For b to be negative, the divergence of the horizontal landslide flux (first
412 term on the right-hand side of equation 3) must also be negative, which requires the landslide

413 thickness to decrease in the direction of movement. This is not physically possible because the
414 landslide thickness is by definition zero at the headscarp, so an inferred thickness of zero
415 minimizes the misfit there.

416 The example earthflow generally has a concave-up slip surface with some irregular
417 bumps (Figure 3b). The slip surface more closely mimics the ground surface in the main
418 transport zone, however there are some low thickness zones near the headscarp and landslide
419 margins that result from negative b values (Figure S4). Lastly, the example landslide complex
420 (Boulder Creek landslide complex) has several different active zones, each with an alternating
421 concave-up and convex-up slip surface profile (Figure 3c). The landslide slip surface is rough
422 and irregular over the length of the entire landslide, but each deep zone generally corresponds to
423 the different kinematic units that comprise the landslide complex (Figure 2c). This large
424 landslide has several areas that do not have a resolvable active thickness. These patches with low
425 active thickness result from low velocity zones (i.e., the landslide toe) and the same
426 characteristics of the velocity field described for the example slump and earthflow (Figures S4).
427 Patches with negative b values must have negative flux divergence, which tends to force the
428 inferred thickness to decrease in the direction of movement at those locations.

429 Landslide zones with approximately zero inferred thickness should correspond to parts of
430 landslides that are not currently active, however, as shown in Figure 3, we also observed low
431 thickness zones in areas with detectable landslide motion. These low thickness areas in our
432 inversions are likely a consequence of issues related to our landslide mapping, noisy velocity or
433 slope data, or violations of the conservation of volume assumptions (e.g., non-uniform landslide
434 density), and are better interpreted as zones where thickness is undefined, rather than where
435 thickness is low. Because it is not possible to independently identify the exact cause of the

436 negative b values that result in low thickness zones with our dataset, we exclude these low
 437 thickness zones (< 0.1 m) from our analyses since the thickness is not determined there. We
 438 selected this threshold because it characterizes the typical thin soil depth in the Central Belt
 439 Franciscan mélange (Hahm et al., 2019). We find these areas typically correspond to regions
 440 near the landslide margins for slumps and earthflows, but are scattered throughout the body of
 441 larger landslide complexes, downflow from regions with negative b values (Figure 3). After
 442 excluding the low thickness zones, the mean active thickness of the individual landslides ranged
 443 from 1.06 to 25.4 m, which, as expected, is higher than the mean thickness range including the
 444 low thickness zones (0.4 to 22.4 m). For the remainder of the paper, we will report landslide
 445 metrics with these low thickness zones excluded and will report metrics including the low
 446 thickness zones in Table S3.

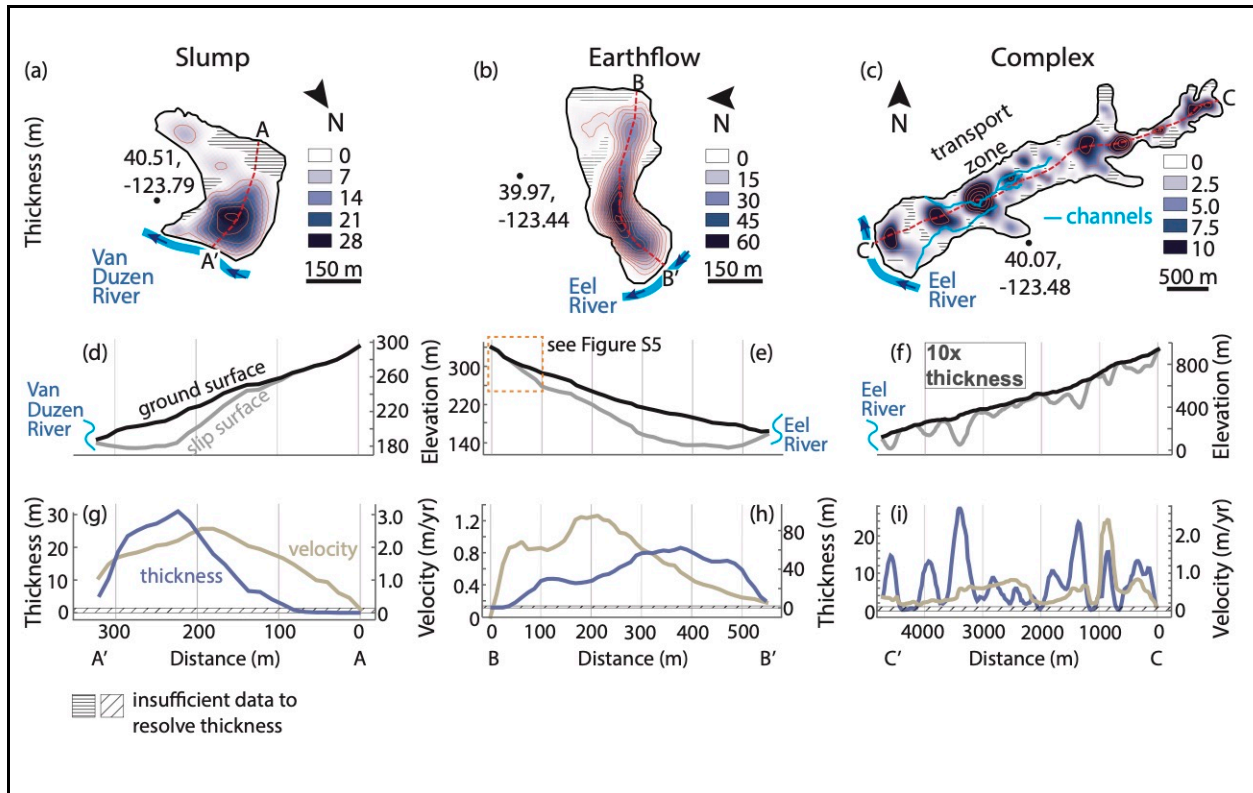


Figure 3. Landslide thickness inversions for example slump, earthflow, and landslide complex. (a–c) Landslide thickness maps. Thin orange lines show 5-meter thickness contours. Red dashed line shows profiles plotted in (d–i). Black dots show latitude and longitude coordinates. Thick blue lines show rivers and thin blue lines show deep channels incised into the landslide body. (d–f) Ground surface and slip surface elevation profiles. Dashed orange rectangle in (e) shows location of landslide headscarp in Figure S5. In subplot (f) the results of thickness inversion are vertically exaggerated by a factor of 10 relative to the elevation profile. (g–i) Landslide thickness and 3D velocity magnitude profiles. Hachures (a–c) and (g–i) identify areas with insufficient data to resolve thickness.

447

448 Although we do not have borehole data to confirm our thickness estimates, we used the
449 topography to verify the inferred slip surface elevation in several cases. Figure S5 shows the
450 example earthflow has a clear headscarp that can be used to trace the sliding surface underneath
451 the ground surface. The extension of the headscarp slip surface under the landslide provides
452 confirmation that the inversion is approximating the slip surface elevation correctly. Figure S6
453 shows another slow-moving landslide that has filled into a pre-existing valley. Transects across
454 this landslide show the ground surface of the filled-in valley and that the slip surface has the
455 shape of the pre-existing valley, providing additional confirmation that our inversions are
456 approximating the slip surfaces correctly. In addition, we compared our thickness inversions to
457 thickness estimates from lidar. Mackey and Roering (2011) used lidar to measure the toe height
458 at the channel interface for dozens of landslides in the Eel River catchment, which is assumed to
459 be minimum thickness estimates at those locations. Of those landslides, 10 (including slumps,
460 earthflows, and complexes) can be used to make comparisons with our dataset. We found overall

461 good agreement between the landslides toe thickness estimated from lidar and from our
 462 inversions (Figure 4).

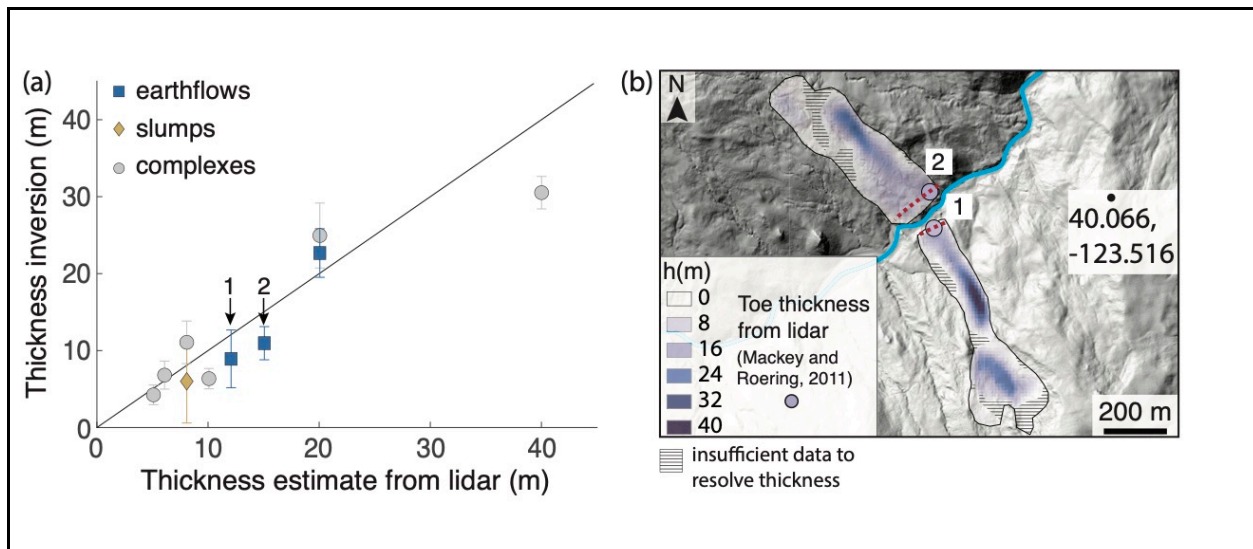


Figure 4. Landslide toe thickness estimates from lidar compared to inverted toe thickness estimates. (a) Colored symbols show inverted thickness compared to lidar estimates for 10 slumps, earthflows, and landslide complexes. We calculated the inverted toe thickness from profiles extracted across the toes and show the mean and standard deviation with error bars. Black line shows 1-to-1 line. (b) Two example landslide thickness maps draped on a lidar hillshade. Labels 1 and 2 correspond to the data points shown in (a). Hachures show zones with insufficient data to resolve thickness. Blue line shows the river channel. Red dashed line shows the profile line used to calculate toe thickness from the inversion. Toe thickness from Mackey and Roering (2011) is shown by the colored circles with black outline.

463

464 Using our thickness inversions for each landslide, we estimate that the individual
 465 landslide volumes range from 7.012×10^3 to 9.747×10^6 m³ (Figure 5 and Table S3). Figure 5
 466 also shows the distribution of mean thickness, area, and volume for each landslide type. Slumps
 467 are the smallest landslide type with a median thickness of 5.49 ± 2.99 m (± 1 standard deviation),

468 median area of $2.71 \pm 2.05 \times 10^4 \text{ m}^2$, and median volume of $1.53 \pm 1.88 \times 10^5 \text{ m}^3$. Earthflows are
469 medium sized with an inventory median thickness of $6.99 \pm 5.33 \text{ m}$, median area of $4.99 \pm 3.26 \times$
470 10^4 m^2 , and median volume of $2.87 \pm 5.36 \times 10^5 \text{ m}^3$. And landslide complexes are the largest
471 landslides, with a median thickness of $8.05 \pm 4.34 \text{ m}$, median area of $1.58 \pm 3.46 \times 10^5 \text{ m}^2$, and
472 median volume of $1.22 \pm 2.19 \times 10^6 \text{ m}^3$.

473 We fit a power function to the volume-area to characterize the geometric scaling relations
474 (equation 5a) for these slow-moving landslides. We also compared our inventory to a worldwide
475 inventory of soil, undifferentiated, and bedrock landslides compiled by Larsen et al. (2010). We
476 find that the slow-moving landslides in the northern California Coast Ranges are larger in both
477 area and volume than most soil landslides, but smaller than the largest bedrock landslides around
478 the world (Figure 5). The best fit volume-area power function exponent (with 95% confidence)
479 for our inventory was $\gamma = 1.306$ (1.213, 1.399) (Figure 5). We observed an apparent break in the
480 slope of the volume-area relation for the largest landslides in our inventory with area $> 10^5 \text{ m}^2$.
481 To further investigate this break in slope, we also fit volume-area scaling as a function of
482 landslide type and find that the break in slope is primarily associated with the landslide
483 complexes. By fitting a power function to each landslide type, we find slumps $\gamma_S = 1.493$ (1.224,
484 1.762), earthflows $\gamma_{Ef} = 1.535$ (1.273, 1.796), and complexes $\gamma_C = 1.172$ (0.9858, 1.357).
485 Although these parameters are not statistically distinct at the 95% confidence level, the fact that
486 γ_S and γ_{Ef} overlap more with each other than with γ_C supports the argument that the break in slope
487 is likely related to landslide type. We report all of the geometric scaling parameters in Table S4.

488 In addition, we calculated the thickness-area scaling relations using the mean thickness
489 (equation 5b) to represent each landslide (Figure 5). We compared these scaling relations to
490 point based estimates (lidar) and measurements (boreholes) of landslide thickness for slow-

491 moving landslides in the northern California Coast Ranges (Mackey and Roering, 2011) and the
 492 Reno River catchment, Apennines, Italy (Simoni et al., 2013). The best fit thickness-area power
 493 function exponent (with 95% confidence) for the inventory $\zeta = 0.3058$ (0.2129, 0.3987),
 494 indicating a weak increase in mean thickness with area for the inventory as a whole. We also fit
 495 thickness-area scaling as a function of landslide type and find slumps $\zeta_S = 0.4926$ (0.2236,
 496 0.7615), earthflows $\zeta_{Ef} = 0.5348$ (0.2734, 0.7963), and for landslide complexes $\zeta_C = 0.1716$ (-
 497 0.0142, 0.3573). Therefore, landslide thickness significantly increases with area for slumps and
 498 earthflows (p-value = 0.0002 and 0.0006, respectively), but does not significantly vary with area
 499 for landslide complexes (p-value = 0.0694).

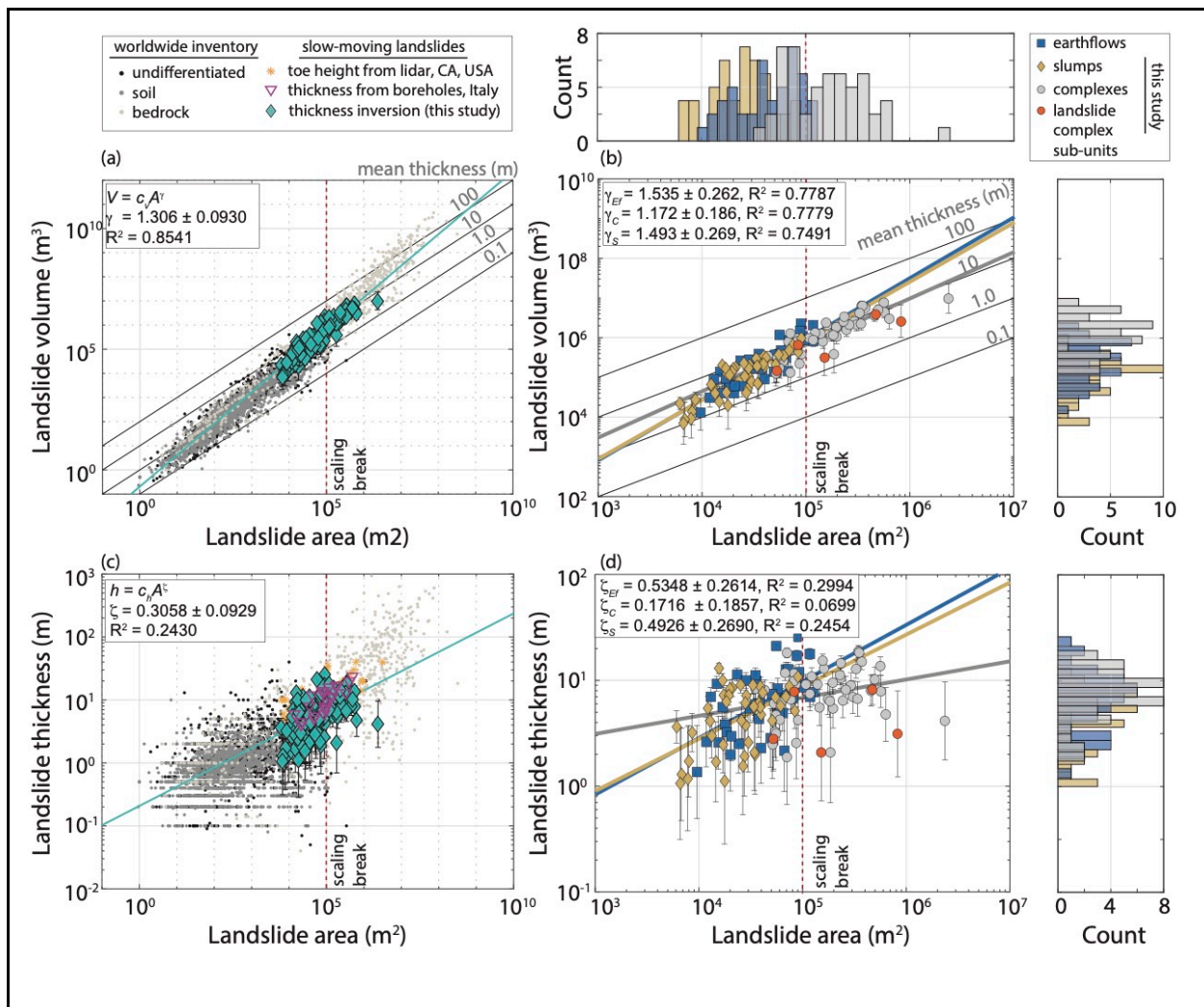


Figure 5. Landslide thickness, volume, and area geometric scaling relations. (a) Volume-area relations for our inventory and a worldwide inventory of soil, undifferentiated, and bedrock landslides (Larsen et al. 2010). (b) Volume-area relations for slumps, earthflows, and landslide complexes. (a, b) Thin diagonal black lines show volume-area for various constant mean thicknesses. (c) Thickness-area relations for our inventory (mean thickness), the worldwide inventory (Larsen et al., 2010), and slow-moving landslides in the northern California Coast Ranges (Mackey and Roering, 2011) and the Apennine mountains, Italy (Simoni et al., 2013). (d) Landslide thickness-area relations by landslide type. Orange circles in (b, d) correspond to the Boulder Creek landslide complex split into 5 smaller landslides (see Figure S7). Error bars show estimated minimum uncertainty estimates (supporting information). Red dashed vertical line shows an apparent break in scaling for the largest landslide complexes in our dataset. Histograms of landslide thickness, area, and volume show the size distributions for each landslide type. All fit parameter values are in Table S4.

500

501 3.3 Frictional Strength

502 Using equation 9, we back-calculated the landslide friction angle ϕ under dry and
 503 saturated conditions end members assuming nil cohesion. Additional landslide properties used in
 504 computations are listed in Table S2. The inferred friction angle ranged from $\sim 6.8^\circ$ to $\sim 28^\circ$ for
 505 dry conditions and $\sim 13^\circ$ to $\sim 54^\circ$ for saturated conditions (Table S3). Our inferred friction angles
 506 encompass friction angle values measured in the laboratory for Franciscan *mélange* rocks and
 507 landslide material (Figure 6). We also analyzed the friction angle as a function of landslide size
 508 and mean slope angle (Figure 6). We found a weak decreasing power-function relationship with
 509 increasing size and a linear increasing relationship with mean slope angle. The negative trend

510 with length indicates that the largest landslides are weaker, on average, than smaller landslides,
 511 while the positive trend with mean slope angle indicates that landslides with gentle slopes are
 512 weaker on average. Figure 6 also shows that the weakest landslides are the large landslide
 513 complexes that have relatively gentle slope angles while slumps are the strongest and steepest
 514 landslides in our inventory.

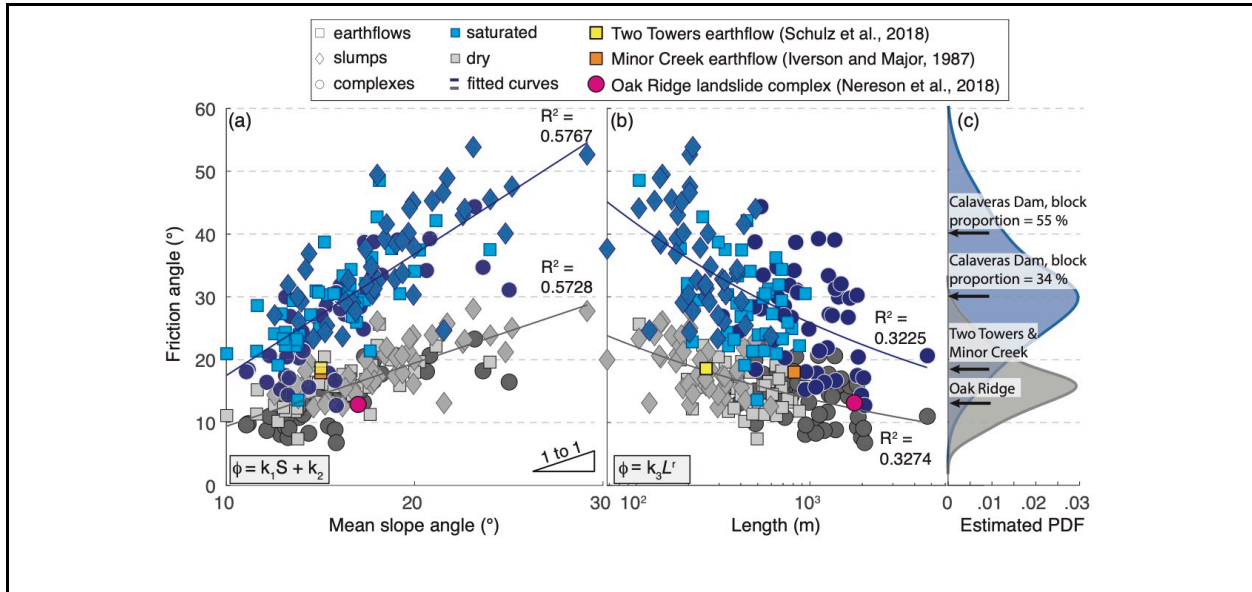


Figure 6. Inferred friction angle for dry and saturated end-members. Friction angle compared to mean hillslope angle (a) and landslide length (b). Solid line lines in (a,b) correspond to best-fit linear and power function curves. For dry conditions, best-fit parameters (with 95% confidence) $k_1 = 1.009$ (0.8586, 1.158), $k_2 = -0.7137$ (-3.279, 1.852), $k_3 = 57.1$ (39.66, 74.55), and $r = -0.2069$ (-0.2582, -0.1556). For wet conditions, $k_1 = 1.935$ (1.649, 2.22), $k_2 = -1.816$ (-6.699, 3.067), $k_3 = 108.4$ (74.81, 142), and $r = -0.2076$ (-0.2597, -0.1555). (c) Estimated probability density function for the full inventory. Black arrows and colored symbols show lab-based and back-calculated friction angle values for the Franciscan mélangé hosted Oak Ridge (Nereson et al., 2018), Two Towers (Schulz et al., 2018), and Minor Creek landslides (Iverson & Major, 1987; Iverson 2000) and the Calaveras Dam, which is founded on

Franciscan mélange (Roadifer et al., 2009). The Calaveras Dam samples are plotted for two different block-in-matrix proportions, which are reported as percentages.

515

516 **4 Discussion**

517 4.1 Landslide Kinematics

518 Our 3D UAVSAR velocity measurements reveal 134 active slow-moving landslides in
519 the northern California Coast Ranges moving at average rates from cm/yr to m/yr between 2016
520 and 2019. The 3D velocity data confirm that the motion of these landslides is generally in the
521 downslope direction. Many of the landslides had relatively low vertical velocities compared to
522 their horizontal velocities that are due to the gradual slope angle (inventory mean $\sim 17^\circ$)
523 exhibited by these slow-moving landslides. However, we did observe segments with vertical
524 uplift that tended to be at the landslide toe due to the concave-up slip surface geometry, and the
525 tendency for longitudinal shortening in the direction of motion to occur at the toe. It is possible
526 that a component of uplift of landslide surfaces could also result from dilation or swelling
527 (volumetric expansion), but the magnitude is small, likely on the order of a few centimeters at
528 most (Booth et al., 2020; Delbridge et al., 2016; Iverson, 2005; Schulz et al., 2018). Including
529 volume changes such as this in the thickness inversion may help reduce uncertainty and improve
530 our results, especially in the zones of low thickness found in many of the landslides, but the
531 amount of dilation or compaction occurring throughout an entire landslide and its variation is
532 generally unknown.

533 Our findings agree with previous work in this region that shows that these landslides
534 exhibit slow, spatially non-uniform downslope motion. Several of the landslides in our study
535 area (e.g., Boulder Creek) have been moving in this manner since at least 1944 (Bennett,

536 Roering, et al., 2016; Mackey & Roering, 2011). Our findings also show that pixel offset
537 tracking with very high resolution UAVSAR data is well-suited for monitoring landslides
538 moving at rates > 10 cm/yr. Some satellites acquire very high resolution SAR with Spotlight
539 modes, including the German TerraSAR-X and Italian COSMO-SkyMed that could provide
540 similar measurements (e.g., Madson et al., 2019), however these data are not open-access.
541 Lastly, we note that there are likely active landslides or landslide zones moving below the
542 precision of our pixel offset tracking technique (< 10 cm/yr) and therefore cannot be observed
543 with our approach. Landslides in our inventory that contain very slow-moving zones may result
544 in unreliable thickness estimates.

545

546 4.2 Landslide Geometry

547 Our study is the first (to our knowledge) to apply the conservation of volume approach to
548 invert for the thickness of multiple landslides in a given region. Previous work (Booth et al.,
549 2020; Booth, Lamb, et al., 2013; Delbridge et al., 2016) has used the same approach to analyze
550 individual landslides, but these landslides occur in different regions and environmental
551 conditions. Like these previous studies, however, we found that the active landslide thickness is
552 variable and that the slip surfaces are rough and irregular in places. The non-uniform thickness
553 and velocity of each landslide results in a non-uniform sediment flux, which has implications for
554 understanding sediment motion along hillslopes (Booth et al., 2020; Guerriero et al., 2017). The
555 shape of the slip surface likely also impacts the landslide kinematics and groundwater flow (Coe
556 et al., 2009; Guerriero et al., 2014; Iverson & Major, 1987; Keefer & Johnson, 1983). Slip
557 surfaces that are bumpy and rough may create additional resisting stresses that act to prevent
558 runaway acceleration and permit long periods of slow landslide motion (Baum & Johnson, 1993;

559 Booth et al., 2018; Leshchinsky, 2019). Investigation of tectonic faults and glaciers also shows
560 that slip surface roughness is an important parameter that controls frictional strength (Brodsky et
561 al., 2016; Fang & Dunham, 2013; Meyer et al., 2018).

562 For our thickness inversions we assumed that the depth-averaged velocity was equal to
563 the surface velocity (i.e., $f = 1$) for all landslides. This block on slope approximation was made
564 to simplify our regional scale analyses. Yet the borehole data from the Two Towers landslide
565 shows that $f \sim 0.96$. While changing f uniformly for each landslide does not alter the spatial
566 pattern of thickness or scaling exponents, it does impact the magnitude of the thickness and
567 therefore the volume. Setting $f = 0.96$ would cause a 4% increase in the inferred thickness and
568 volume of each landslide ($h \sim 1/f$) (Table S3). More work is needed to better constrain the depth-
569 averaged velocity for individual landslides in our field area, particularly to see if f differs with
570 landslide type. Nonetheless, our findings indicate that most of the sliding surfaces are deep-
571 seated (mean thickness for inventory ~ 7.2 m) and thus are expected to lie within the
572 unweathered Central Belt Franciscan mélangé bedrock (Hahm et al., 2019). Therefore, the slow-
573 moving landslides in the northern California Coast Ranges can be classified as bedrock
574 landslides.

575 Using our landslide inventory, we developed new volume-area and thickness-area
576 geometric scaling relations for slow-moving landslides. Geometric scaling relations are
577 particularly useful for slow-moving landslides because these landslides rarely (if ever) evacuate
578 hillslopes, or create clear scars or deposits that can be easily measured. As a result, most
579 measurements of landslide thickness come from isolated boreholes, which are logistically
580 challenging and expensive to install, and are difficult to extrapolate over an entire landslide. Our
581 results provide best-fit volume-area power function exponents ($\gamma \sim 1.2 - 1.5$) that are comparable

582 to power function exponents for bedrock and soil landslides (Guzzetti et al., 2009; Larsen et al.,
583 2010; Bunn et al., 2020a). Recent work by Bunn et al. (2020a) found that deep-seated bedrock
584 landslides in Oregon, USA had $\gamma_{bedrock} \sim 1.4 - 1.6$. Analysis of a worldwide landslide inventory
585 by Larsen et al (2010) showed that soil landslides had a $\gamma_{soil} \sim 1.1 - 1.3$, while bedrock landslides
586 had $\gamma_{bedrock} \sim 1.3 - 1.6$.

587 In addition, our best-fit thickness-area scaling power exponents ($\zeta \sim 0.17 - 0.53$) are also
588 comparable (with a wide range) to previously published values for deep-seated landslides (Figure
589 4c). Bunn et al. (2020a) found $\zeta \sim 0.41 - 0.58$ for deep-seated bedrock landslides. Simoni et al.
590 (2013) reported $\zeta = 0.44$ from borehole inclinometer data from 23 slow-moving landslides in the
591 Apennine Mountains, Italy. Handwerger et al. (2013) reported $\zeta = 0.29$ derived from lidar-based
592 estimates of landslide toe thickness from 69 landslides in the Eel River catchment, several of
593 which are also analyzed in this study (e.g., Figure 4). Importantly, neither Simoni et al. (2013) or
594 Handwerger et al. (2013) used large inventories (> 100) or spatially extensive measurements of
595 landslide thickness, which are especially important for slow-moving landslides with variable
596 thicknesses. Therefore, our new scaling relationships provide the most appropriate values for
597 deep-seated slow-moving landslides, like earthflows, and could be used to help estimate
598 sediment flux and landslide stresses in similar areas around the world. Yet, we note that the large
599 range of scaling exponents suggests that scaling relations should be used with caution. Applying
600 an incorrect scaling exponent to estimate volume for landslides with unknown thickness can lead
601 to large errors in volume calculations (Larsen et al., 2010).

602 Our findings show that the slow-moving landslides located in the northern California
603 Coast Ranges have geometric scaling exponents that lie in between the soil and bedrock type
604 landslides. However, examining the best-fit power function exponents by landslide type suggests

605 that slumps and earthflows display close to self-similar scaling ($\gamma_{self-similar}=1.5$), which is
606 characteristic of bedrock landslides, while landslide complexes display scaling that is
607 characteristic of soil landslides. Figure 5b shows that the landslide complexes with the largest
608 areas display a scaling that tends to follow a constant mean thickness. We propose that landslide
609 complexes have scaling relations that are close to soil landslides because: 1) the mean landslide
610 thickness is limited by a strong layer in the *mélange* (and thus are similar to soil landslides that
611 are limited by the soil thickness), or 2) that landslide complexes are an amalgamation of multiple
612 smaller and shallower landslides. The second explanation provides a reason for why large
613 landslide complexes tend to have multiple kinematic units (e.g., Aryal et al., 2012; Hu et al.,
614 2020) and further emphasizes the importance of having detailed landslide maps, especially when
615 applying geometric scaling relations (e.g., Marc and Hovius, 2015).

616

617 4.3. Thickness Inversion Challenges

618 The inferred thickness of many of the slow-moving landslides, in particular the landslide
619 complexes, can be highly variable with deeper active zones and thinner or zero thickness areas
620 that are not currently moving (Figure 3). In addition, some patches with an inferred thickness of
621 zero occurred in areas where $b < 0$, such that a negative divergence was required to match the
622 observations (Figure S4). These negative b values typically arose when the product of the
623 horizontal velocity and the topographic gradient was more negative than the vertical component
624 of the surface velocity vector (equations 3 and 4). This situation could result from artifacts in the
625 velocity or topographic data or from actual physical processes occurring in the landslide that
626 would tend to increase the magnitude of the horizontal velocity, increase the magnitude of the
627 topographic gradient, or decrease the magnitude of the vertical velocity relative to their true

628 values, assuming conservation of volume. In particular, one plausible physical mechanism that
629 would decrease the magnitude of v_{ud} relative to that of u_{surf} or ∇Z_{surf} is dilation of landslide
630 material as it deforms. That increase in volume would cause an additional positive vertical
631 component to v_{ud} . Although we cannot determine whether errors in the velocity and topographic
632 data, or actual physical mechanisms are responsible for the low inferred thickness zones, we find
633 dilation a plausible explanation, especially near landslide headscarps, or in other zones of
634 extension, indicating macro-scale decreases in density.

635 Additionally, it is important to note that the irregular thickness patterns observed in some
636 landslides may not align with inferred thickness based on geomorphic or structural
637 interpretations. This discrepancy is likely related to the long-lasting geomorphic imprint that
638 slow-moving landslides have made on the landscape. Landslide surface morphology may last for
639 decades or longer after a landslide completely stops moving (e.g., Booth et al., 2017), which can
640 make it challenging to infer the active landslide thickness without kinematic data. Although our
641 approach is useful for identifying the currently active portions of landslides and inferring their
642 thickness based on volume conservation (with assumptions), it does not allow us to infer the
643 subsurface geometry of the often larger inactive landslide body. As a result, we emphasize the
644 need for more comparisons between ground- and remote sensing-based investigation of landslide
645 geometry. In particular, direct comparison between numerous ground-based measurements from
646 boreholes and structural mapping are needed to widely test the results of our remote sensing
647 approach. Nonetheless, we find our thickness inversions are producing reasonable estimates of
648 landslide thickness in the cases we were able to test (Figures 4, S5, and S6).

649

650 4.4 The Boulder Creek Landslide Complex

651 We found that the inferred active thickness for the Boulder Creek landslide complex was
652 particularly irregular and challenging to explain based on a priori assumptions of landslide
653 geometry. While we expect areas that are not currently active to thin, and even have zero
654 thickness in places (e.g., parts of the landslide toe), the active transport zone on Boulder Creek
655 also contains thin and thick patches (Figures 3c, 3f, and 3i). One possible explanation for this
656 variability is related to patches of local density changes (e.g., dilation) that could result in
657 negative b values. In addition to these potential artifacts, another possible cause of these low
658 thickness zones is related to the large channel network incised into the landslide (Figure 3). In
659 some places the channel reaches depths of 15-20 meters (Figure S8). Since the thickness is
660 measured as the vertical distance from the ground surface to the inferred basal sliding surface,
661 the predicted thickness is expected to be low in places surrounding the channel if the channel
662 depth is similar to the landslide thickness. Our findings indicate that the channel has incised to
663 depths that approach the predicted sliding surface in several places (Figure S8). However, the
664 channel has not incised deeper than the landslide base because we find the channel is moving
665 with similar velocity to the surrounding regions (Figure 2c).

666 The distinct kinematic zones within Boulder Creek landslide complex also indicate that
667 smaller, faster, and possibly shallower features are superimposed on a larger, slower, and
668 possibly deeper-seated failure (Figure 2c). If multiple failure planes are indeed present, that
669 would violate the assumption of a constant f throughout the landslide and cause unreliable
670 thickness estimates. Specifically, the surface velocity would be much greater than the depth-
671 averaged velocity (i.e., f would be much smaller) within the superimposed landslide. This would
672 systematically cause the inferred thickness to be too large near the headscarp of the
673 superimposed landslide and too shallow near its toe, since the divergence of the surface velocity

674 field would be much greater than the divergence of the depth-averaged velocity field at those
675 locations. To further explore the hypothesis that the Boulder Creek landslide complex is
676 composed of multiple smaller landslides, we delineated Boulder Creek into 5 smaller sub-
677 landslides and performed a thickness inversion for each sub-landslide (Figure S7). While the
678 thickness patterns are similar to the thickness inversion for the full landslide complex, the
679 magnitude of the inferred thickness differs in some places, and the area of each landslide is
680 smaller, which places them into the space mostly populated by earthflows on the thickness-area
681 and volume-area plots (orange circles in Figures 5b and 5d). Some of these differences in the
682 magnitude of the thickness estimate are due to differences in the pixel resolution of the sub-
683 landslides (10 m pixel) and the full landslide (20 m pixel). Nonetheless, mapping landslide
684 complexes as one large landslide results in a lower mean thickness relative to the landslide area
685 which affects the geometric scaling relations. While more investigation is warranted, our
686 thickness inversions have caused us to reevaluate how we think about large landslide complexes.

687

688 4.5 Landslide Strength

689 Our back-analysis of landslide strength suggests that there is a weak decreasing
690 relationship between landslide size and strength and an increasing relationship between mean
691 slope angle and strength (Figure 6; Figure S9). The increasing relationship between mean slope
692 angle and friction angle was expected because steeper landslides must be stronger to maintain
693 force balance (equation 9). The decreasing relationship between landslide size and friction angle
694 is notable and intriguing. We hypothesize that larger landslides are weaker than smaller
695 landslides because of strength heterogeneity in the Franciscan *mélange* bedrock and the
696 increased likelihood of incorporating weak material within larger volumes. Laboratory

697 measurements of the strength of the Franciscan *mélange* rocks have shown that the proportion of
698 the blocks hosted in the argillaceous matrix controls the overall rock strength (Roadifer et al.,
699 2009) (Figure 6). This implies that larger landslides may have a decreased proportion of blocks,
700 which are not uniformly distributed, and are therefore controlled by the weak argillaceous
701 matrix.

702 Scale-dependent strength has also been observed along other landslides and faults.
703 Brodsky et al. (2016) suggested that faults are weaker at large spatial scales because they
704 encompass larger weak zones. A recent study by Bunn et al., (2020b) found that the inferred
705 shear strength of landslides decreases with increasing landslide size. They proposed that smaller
706 landslides were stronger because they occur in cemented cohesive materials and larger landslides
707 were in a residual state. Although we assumed nil cohesion to back-calculate the residual
708 frictional strength of the active landslides, it is likely that cohesion is important in controlling the
709 initial landslide failure due to the high-clay content of the Central Belt Franciscan *mélange* (e.g.,
710 Milledge et al., 2014).

711 Our inferred friction angles also depend on wetness conditions. Due to the high seasonal
712 rainfall in the northern California Coast Ranges, these slow-moving landslides are typically
713 saturated (or nearly saturated) during the wet season and partially saturated or dry during the dry
714 season (Hahm et al., 2019; Iverson & Major, 1987; Schulz et al., 2018). Direct comparison with
715 friction angle values measured in the laboratory and back-calculated for Franciscan *mélange*
716 rocks and landslide materials provides some insight into our findings. For saturated conditions
717 we find that the inferred friction angles for medium to large earthflows and landslide complexes
718 overlap the measured friction values from the Two Towers earthflow (Schulz et al., 2018), Minor
719 Creek earthflow (Iverson and Major, 1987), and Oakridge landslide complex (Nereson et al.,

720 2018). The majority of the smaller slumps have saturated friction angles that are significantly
721 higher than these three landslides. The saturated friction values for smaller slumps, earthflows,
722 and landslide complexes have more overlap with measured rock friction values that depend on
723 the block-in-matrix proportion (Roadifer et al., 2009). Interestingly, the dry friction angles for all
724 landslide types have more overlap with lab-based friction measurements for the landslides. Yet,
725 it is unlikely that most of these landslides, especially the larger landslides, become completely
726 dry. Instead, the true landslide-scale friction angle values likely lie somewhere between our
727 inferred values for saturated and dry conditions.

728 We suggest that some of these differences between lab-based and inferred friction angles
729 may be attributed to commonly observed differences in laboratory- and field-scale measurements
730 that are often related to large scale spatial heterogeneity in the field (e.g., Marone, 1998; Van
731 Asch et al., 2007). In addition, our assumption of nil cohesion can partially explain the higher
732 friction values for saturated conditions (Bunn et al., 2020b). The additional strength imparted by
733 cohesion would act to reduce the inferred friction angle values to maintain equilibrium (equation
734 9). We assumed nil cohesion because the landslides have moved significantly over the study
735 period (and likely much longer), but it is likely that cohesion is important for the clay-rich
736 landslide material and future work needs to better account for temporal changes in cohesion,
737 which may be especially important for landslides that completely stop moving during dry
738 periods. While the large spread of inferred friction values makes it difficult to identify a single
739 representative value for slow-moving landslides in the northern California Coast Ranges, our
740 results further highlight the heterogeneous nature of the Central Belt Franciscan *mélange*
741 lithologic unit. Similar to the recent findings of Bunn et al., (2020b), our findings also suggest
742 that landslide type, mean slope angle, and wetness conditions may provide some first-order

743 information on relative landslide strength at the regional scale. Furthermore, our findings have
744 implications for understanding landscape evolution and agree with previous work that shows that
745 over geomorphic timescales, we generally expect to find steeper hillslopes where hillslope
746 materials are stronger (e.g., Korup et al., 2007; Roering et al., 2015). More work is needed to
747 understand our findings in the context of landscape evolution because the currently active
748 landslides are just the most recent snapshot of the landscape, and the slopes they occur on have
749 probably been shaped by numerous previous generations of similar landslides (e.g., Mackey and
750 Roering, 2011; Roering et al., 2015).

751

752 4.6 What Controls the Size of Slow-moving Landslides?

753 Landslide size is set by the landslide mechanical properties, slope geometry, and
754 environmental conditions. For most landslides, the maximum size is typically limited to the
755 maximum hillslope size, such that the landslide length does not exceed the hillslope length. The
756 landslide thickness is typically set by the location of a weak layer beneath the ground surface, or
757 at a depth where there are changes in strength and permeability, such as the soil to bedrock
758 transition or the bottom of the critical zone (i.e., the zone that extends from the ground surface
759 down to unweathered bedrock) (Booth, Roering, et al., 2013; Larsen et al., 2010; Milledge et al.,
760 2014). Using a 3D slope stability model for shallow soil landslides that accounts for the forces
761 acting on the landslide basal slip surface, lateral margins, and passive/active wedges at the
762 toe/head, Milledge et al. (2014) found that the critical area and depth that can fail as a landslide
763 depends on the topography, pore-water pressure, and landslide material properties, including
764 density, cohesion, and friction angle. We note that their modeled landslides have less complex
765 geometries than the landslides in our inventory. In their model the pore-water pressure plays a

766 fundamental role in determining the critical landslide size and failure depth, such that higher
767 pore-water pressures decrease the critical size required for failure. Large landslides therefore
768 occur when high pore pressures are reached over a correspondingly large spatial area. At our
769 northern California Coast Range study site, the relatively thin, but laterally extensive critical
770 zone that is often saturated during the wet season (Hahm et al., 2019), may promote laterally
771 extensive landslides by elevating the water table height simultaneously over large areas.

772 Milledge et al. (2014)'s model also predicts that landslide thickness should increase as
773 the square root of the landslide area and that the failure depth sets the minimum landslide area.
774 Our best-fit thickness-area scaling exponents for slumps and earthflows are close to a square root
775 scaling (exponents ~ 0.5 with large 95th confidence intervals). Our results also suggest that the
776 landslide thickness controls the minimum area, but does not bound its maximum size. Instead,
777 slow-moving landslides can continue to grow in area by becoming a landslide complex
778 consisting of multiple, connected, sub-landslides without becoming significantly deeper on
779 average. Large landslide complexes can occupy multiple hillslopes, and fill valleys and
780 catchments such that their size may exceed the typical hillslope size, in contrast to landslides that
781 fully evacuate their hillslopes (e.g., Jeandet et al., 2019). Thus, it seems that the catchment size
782 sets the maximum area for slow-moving landslides. Our thickness inversion results also indicate
783 that large landslides are weaker than small landslides. This finding may indicate that large
784 landslides become large by incorporating weak material. It is possible that the largest landslides
785 grow over time and take decades to develop (e.g., Mackey & Roering, 2011). As many of our
786 landslide complexes seem to be composed of several smaller sub-landslides or kinematic zones,
787 it is possible that these features have connected through time as slip surfaces propagate along the
788 slope.

789

790 **5 Conclusions**

791 We measured the 3D surface velocity of more than one hundred slow-moving landslides
792 in the northern California Coast Ranges with data from the NASA/JPL UAVSAR. We used
793 volume conservation techniques to infer the active thickness, volume, and strength of each
794 landslide. The thickness of each landslide is variable and can vary by tens of meters sometimes
795 resulting in an irregular slip surface geometry. Volume-area geometric scaling relations suggest
796 that these landslides have similarities to both soil and bedrock landslides around the world.
797 Although their failure planes are likely hosted in unweathered bedrock, their thickness seems to
798 be limited, producing a scaling similar to soil landslides for the largest landslide complexes. The
799 inferred residual friction angles are also scale-dependent, like faults, such that large landslides
800 complexes tend to be weaker than small landslides such as slumps. This decrease in inferred
801 friction angle with landslide size is likely because larger landslides are composed of larger
802 proportions of weak material. Our study represents the first to use the conservation of volume
803 approach for numerous landslides occurring under the same environmental conditions. Our
804 results provide key insights into the subsurface geometry and strength that control the behavior
805 of slow-moving landslides. Our work shows how state-of-the-art remote sensing techniques can
806 be used to better understand landslide processes for hazards and to quantify their contribution to
807 landscape evolution.

808

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819

820 **Data Availability**

821 Landslide geometry data used in this study are listed in the references: Larsen et al.,
822 (2010), Mackey and Roering, (2011), Simoni et al., (2013) and are included in the figures.
823 Borehole thickness data at the Two Towers landslide is in reference: Schulz et al. (2018). Lidar
824 digital elevation models are provided by OpenTopography and may be downloaded online
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830 To acquire these data, proposals may be submitted to the DLR online ([https://tandemx-
831 science.dlr.de/](https://tandemx-science.dlr.de/)). NASA/JPL UAVSAR data used in this study are freely available and may be
832 downloaded through their website (<https://uavsar.jpl.nasa.gov/>).

833

834 **References**

835 Aryal, A., Brooks, B. A., Reid, M. E., Bawden, G. W., & Pawlak, G. R. (2012). Displacement fields from

- 836 point cloud data: Application of particle imaging velocimetry to landslide geodesy. *Journal of*
837 *Geophysical Research: Earth Surface* (2003–2012), 117(F1).
838 <https://doi.org/10.1029/2011JF002161>
- 839 Aryal, A., Brooks, B., & Reid, M. E. (2015). Landslide subsurface slip geometry inferred from 3D surface
840 displacement fields. *Geophysical Research Letters*. 42(5), 1411-1417.
841 <https://doi.org/10.1002/2014GL062688>
- 842 Aster, R. C., Borchers, B., & Thurber, C. H. (2013). Chapter Four - Tikhonov Regularization. In R. C.
843 Aster, B. Borchers, & C. H. Thurber (Eds.), *Parameter Estimation and Inverse Problems (Second*
844 *Edition)* (Second Edition, pp. 93–127). Boston: Academic Press. [https://doi.org/10.1016/B978-0-](https://doi.org/10.1016/B978-0-12-385048-5.00004-5)
845 [12-385048-5.00004-5](https://doi.org/10.1016/B978-0-12-385048-5.00004-5)
- 846 Baum, R. L., & Johnson, A. M. (1993). Steady movement of landslides in fine-grained soils: A model for
847 sliding over an irregular slip surface. *US Geological Survey Bulletin (USA)*.
- 848 Bennett, G. L., Roering, J. J., Mackey, B. H., Handwerger, A. L., Schmidt, D. A., & Guillod, B. P.
849 (2016). Historic drought puts the brakes on earthflows in Northern California. *Geophysical*
850 *Research Letters*, 43(11), 5725–5731. <https://doi.org/10.1002/2016GL068378>
- 851 Bennett, G. L., Miller, S. R., Roering, J. J., & Schmidt, D. A. (2016). Landslides, threshold slopes, and
852 the survival of relict terrain in the wake of the Mendocino Triple Junction. *Geology*, 44(5), 363–
853 366. <https://doi.org/10.1130/G37530.1>
- 854 Bessette-Kirton, E. K., Coe, J. A., & Zhou, W. (2018). Using stereo satellite imagery to account for
855 ablation, entrainment, and compaction in volume calculations for rock avalanches on glaciers:
856 Application to the 2016 Lamplugh rock avalanche in Glacier Bay National Park, Alaska. *Journal*
857 *of Geophysical Research: Earth Surface*, 123(4), 622–641. <https://doi.org/10.1002/2017JF004512>
- 858 Booth, A. M., Lamb, M. P., Avouac, J.-P., & Delacourt, C. (2013). Landslide velocity, thickness, and
859 rheology from remote sensing: La Clapière landslide, France. *Geophysical Research Letters*,
860 40(16), 4299–4304. <https://doi.org/10.1002/grl.50828>
- 861 Booth, A. M., Roering, J. J., & Rempel, A. W. (2013). Topographic signatures and a general transport law

- 862 for deep-seated landslides in a landscape evolution model. *Journal of Geophysical Research:*
863 *Earth Surface*, 118(2), 603–624. <https://doi.org/10.1002/jgrf.20051>
- 864 Booth, A. M., LaHusen, S. R., Duvall, A. R., & Montgomery, D. R. (2017). Holocene history of deep-
865 seated landsliding in the North Fork Stillaguamish River valley from surface roughness analysis,
866 radiocarbon dating, and numerical landscape evolution modeling. *Journal of Geophysical*
867 *Research: Earth Surface*, 122(2), 456–472. <https://doi.org/10.1002/2016JF003934>
- 868 Booth, A. M., McCarley, J., Hinkle, J., Shaw, S., Ampuero, J.-P., & Lamb, M. P. (2018). Transient
869 Reactivation of a Deep-Seated Landslide by Undrained Loading Captured With Repeat Airborne
870 and Terrestrial Lidar. *Geophysical Research Letters*, 45(10), 4841–4850.
871 <https://doi.org/10.1029/2018GL077812>
- 872 Booth, A. M., McCarley, J. C., & Nelson, J. (2020). Multi-year, three-dimensional landslide surface
873 deformation from repeat lidar and response to precipitation: Mill Gulch earthflow, California.
874 *Landslides*, 1–14. <https://doi.org/10.1007/s10346-020-01364-z>
- 875 Brodsky, E. E., Kirkpatrick, J. D., & Candela, T. (2016). Constraints from fault roughness on the scale-
876 dependent strength of rocks. *Geology*, 44(1), 19–22. <https://doi.org/10.1130/G37206.1>
- 877 Bunn, M., Leshchinsky, B., & Olsen, M. J. (2020a). Estimates of three-dimensional rupture surface
878 geometry of deep-seated landslides using landslide inventories and high-resolution topographic
879 data. *Geomorphology*, 367, 107332. <https://doi.org/10.1016/j.geomorph.2020.107332>
- 880 Bunn, M., Leshchinsky, B., & Olsen, M. J. (2020b). Geologic Trends in Shear Strength Properties
881 Inferred through Three-Dimensional Back-Analysis of Landslide Inventories. *Journal of*
882 *Geophysical Research: Earth Surface*, e2019JF005461. <https://doi.org/10.1029/2019JF005461>
- 883 Cerovski-Darriau, C., & Roering, J. J. (2016). Influence of anthropogenic land-use change on hillslope
884 erosion in the Waipaoa River Basin, New Zealand. *Earth Surface Processes and Landforms*,
885 41(15), 2167–2176. <https://doi.org/10.1002/esp.3969>
- 886 Coe, J. A., McKenna, J. P., Godt, J. W., & Baum, R. L. (2009). Basal-topographic control of stationary
887 ponds on a continuously moving landslide. *Earth Surface Processes and Landforms*, 34(2), 264–

- 888 279. <https://doi.org/10.1002/esp.1721>
- 889 Cruden, D. M., & Varnes, D. J. (1996). Landslides: Investigation and Mitigation. Chapter 3-Landslide
890 types and processes. *Transportation Research Board Special Report*, (247).
- 891 Delbridge, B. G., Bürgmann, R., Fielding, E., Hensley, S., & Schulz, W. H. (2016). Three-dimensional
892 surface deformation derived from airborne interferometric UAVSAR: Application to the
893 Slungullion Landslide. *Journal of Geophysical Research: Solid Earth*, 121(5), 3951–3977.
894 <https://doi.org/10.1002/2015JB012559>
- 895 Fang, Z., & Dunham, E. M. (2013). Additional shear resistance from fault roughness and stress levels on
896 geometrically complex faults. *Journal of Geophysical Research: Solid Earth*, 118(7), 3642–3654.
897 <https://doi.org/10.1002/jgrb.50262>
- 898 Fialko, Y., Simons, M., & Agnew, D. (2001). The complete (3-D) surface displacement field in the
899 epicentral area of the 1999 Mw7. 1 Hector Mine earthquake, California, from space geodetic
900 observations. *Geophysical Research Letters*, 28(16), 3063–3066.
901 <https://doi.org/10.1029/2001GL013174>
- 902 Fielding, E. J., Liu, Z., Stephenson, O. L., Zhong, M., Liang, C., Moore, A., et al. (2020). Surface
903 Deformation Related to the 2019 M w 7.1 and 6.4 Ridgecrest Earthquakes in California from
904 GPS, SAR Interferometry, and SAR Pixel Offsets. *Seismological Research Letters*.
905 <https://doi.org/10.1785/0220190302>
- 906 Grant, M., & Boyd, S. (2014). *CVX: Matlab Software for Disciplined Convex Programming, version 2.1*.
907 Retrieved from <http://cvxr.com/cvx>
- 908 Guerriero, L., Coe, J. A., Revellino, P., Grelle, G., Pinto, F., & Guadagno, F. M. (2014). Influence of slip-
909 surface geometry on earth-flow deformation, Montaguto earth flow, southern Italy.
910 *Geomorphology*, 219(0), 285 – 305. <http://dx.doi.org/10.1016/j.geomorph.2014.04.039>
- 911 Guerriero, L., Bertello, L., Cardozo, N., Berti, M., Grelle, G., & Revellino, P. (2017). Unsteady sediment
912 discharge in earth flows: A case study from the Mount Pizzuto earth flow, southern Italy.
913 *Geomorphology*, 295, 260–284. <https://doi.org/10.1016/j.geomorph.2017.07.011>

- 914 Guzzetti, F., Ardizzone, F., Cardinali, M., Rossi, M., & Valigi, D. (2009). Landslide volumes and
915 landslide mobilization rates in Umbria, central Italy. *Earth and Planetary Science Letters*, 279(3–
916 4), 222–229. <https://doi.org/10.1016/j.epsl.2009.01.005>
- 917 Hahm, W. J., Rempe, D. M., Dralle, D. N., Dawson, T. E., Lovill, S. M., Bryk, A. B., et al. (2019).
918 Lithologically controlled subsurface critical zone thickness and water storage capacity determine
919 regional plant community composition. *Water Resources Research*, 55(4), 3028–3055.
920 <https://doi.org/10.1029/2018WR023760>
- 921 Handwerger, A. L., Roering, J. J., & Schmidt, D. A. (2013). Controls on the seasonal deformation of
922 slow-moving landslides. *Earth and Planetary Science Letters*, 377, 239–247.
923 <https://doi.org/10.1016/j.epsl.2013.06.047>
- 924 Handwerger, A. L., Roering, J. J., Schmidt, D. A., & Rempel, A. W. (2015). Kinematics of earthflows in
925 the Northern California Coast Ranges using satellite interferometry. *Geomorphology*, 246, 321–
926 333. <https://doi.org/10.1016/j.geomorph.2015.06.003>
- 927 Handwerger, A. L., Huang, M.-H., Fielding, E. J., Booth, A. M., & Bürgmann, R. (2019). A shift from
928 drought to extreme rainfall drives a stable landslide to catastrophic failure. *Scientific Reports*,
929 9(1), 1569. <https://doi.org/10.1038/s41598-018-38300-0>
- 930 Handwerger, A. L., Fielding, E. J., Huang, M.-H., Bennett, G. L., Liang, C., & Schulz, W. H. (2019).
931 Widespread initiation, reactivation, and acceleration of landslides in the northern California Coast
932 Ranges due to extreme rainfall. *Journal of Geophysical Research: Earth Surface*, 124(7), 1782–
933 1797. <https://doi.org/10.1029/2019JF005035>
- 934 Hu, X., Bürgmann, R., Schulz, W. H., & Fielding, E. J. (2020). Four-dimensional surface motions of the
935 Slumgullion landslide and quantification of hydrometeorological forcing. *Nature*
936 *Communications*, 11(1), 1–9. <https://doi.org/10.1038/s41467-020-16617-7>
- 937 Hungr, O. (1987). An extension of Bishop’s simplified method of slope stability analysis to three
938 dimensions. *Geotechnique*, 37(1), 113–117. <https://doi.org/10.1680/geot.1987.37.1.113>
- 939 Hungr, O., Salgado, F., & Byrne, P. (1989). Evaluation of a three-dimensional method of slope stability

- 940 analysis. *Canadian Geotechnical Journal*, 26(4), 679–686. <https://doi.org/10.1139/t89-079>
- 941 Hungr, O, Leroueil, S., & Picarelli, L. (2014). The Varnes classification of landslide types, an update.
942 *Landslides*, 11(2), 167–194. <https://doi.org/10.1007/s10346-013-0436-y>
- 943 Intrieri, E., Raspini, F., Fumagalli, A., Lu, P., Del Conte, S., Farina, P., et al. (2017). The Maoxian
944 landslide as seen from space: detecting precursors of failure with Sentinel-1 data. *Landslides*, 1–
945 11. <https://doi.org/10.1007/s10346-017-0915-7>
- 946 Iverson, R. M. (2000). Landslide triggering by rain infiltration. *Water Resources Research*, 36(7), 1897–
947 1910. <https://doi.org/10.1029/2000WR900090>
- 948 Iverson, R. M. (2005). Regulation of landslide motion by dilatancy and pore pressure feedback. *Journal*
949 *of Geophysical Research: Earth Surface*, 110(F2). <https://doi.org/10.1029/2004JF000268>
- 950 Iverson, R. M., & Major, J. J. (1987). Rainfall, ground-water flow, and seasonal movement at Minor
951 Creek landslide, northwestern California: Physical interpretation of empirical relations.
952 *Geological Society of America Bulletin*, 99(4), 579–594. [https://doi.org/10.1130/0016-](https://doi.org/10.1130/0016-7606(1987)99<579:RGFASM>2.0.CO;2)
953 [7606\(1987\)99<579:RGFASM>2.0.CO;2](https://doi.org/10.1130/0016-7606(1987)99<579:RGFASM>2.0.CO;2)
- 954 Jayko, A., Blake, M., McLaughlin, R., Ohlin, H., Ellen, S., & Kelsey, H. (1989). Reconnaissance
955 Geologic Map of the Covelo 30-by 60-Minute Quadrangle. *Northern California: US Geological*
956 *Survey Miscellaneous Field Investigation Map MF-2001, Scale, 1(100), 000.*
957 <https://doi.org/10.3133/mf2001>
- 958 Jeandet, L., Steer, P., Lague, D., & Davy, P. (2019). Coulomb mechanics and relief constraints explain
959 landslide size distribution. *Geophysical Research Letters*, 46, 4258–4266.
960 <https://doi.org/10.1029/2019GL082351>
- 961 Jennings, C. W., Strand, R. G., & Rogers, T. H. (1977). Geologic map of California: California Division
962 of Mines and Geology, scale 1:750,000.
- 963 Jung, J., & Yun, S.-H. (2020). Evaluation of Coherent and Incoherent Landslide Detection Methods
964 Based on Synthetic Aperture Radar for Rapid Response: A Case Study for the 2018 Hokkaido
965 Landslides. *Remote Sensing*, 12(2), 265. <https://doi.org/10.3390/rs12020265>

- 966 Kasper, van W., John A, S., William, N., & Luis, T. (2002). Data and model uncertainty estimation for
967 linear inversion. *Geophysical Journal International*, 149(3), 625–632.
968 <https://doi.org/10.1046/j.1365-246X.2002.01660.x>
- 969 Keefer, D. K., & Johnson, A. M. (1983). Earth flows: Morphology, mobilization, and movement.
970 Washington: United States Government Printing Office. No 1264. <https://doi.org/10.3133/pp1264>
- 971 Kelsey, H. M. (1978). Earthflows in Franciscan mélangé, Van Duzen River basin, California. *Geology*,
972 6(6), 361–364. [https://doi.org/10.1130/0091-7613\(1978\)6<361:EIFMVD>2.0.CO;2](https://doi.org/10.1130/0091-7613(1978)6<361:EIFMVD>2.0.CO;2)
- 973 Korup, O., Clague, J. J., Hermanns, R. L., Hewitt, K., Strom, A. L., & Weidinger, J. T. (2007). Giant
974 landslides, topography, and erosion. *Earth and Planetary Science Letters*, 261(3–4), 578–589.
975 <https://doi.org/10.1016/j.epsl.2007.07.025>
- 976 Lacroix, P., Dehecq, A., & Taipe, E. (2020). Irrigation-triggered landslides in a Peruvian desert caused by
977 modern intensive farming. *Nature Geoscience*, 13(1), 56–60. [https://doi.org/10.1038/s41561-019-](https://doi.org/10.1038/s41561-019-0500-x)
978 0500-x
- 979 Lacroix, P., Handwerker, A. L., & Bièvre, G. (2020). Life and death of slow-moving landslides. *Nature*
980 *Reviews Earth & Environment*, 1–16. <https://doi.org/10.1038/s43017-020-0072-8>
- 981 Larsen, I. J., Montgomery, D. R., & Korup, O. (2010). Landslide erosion controlled by hillslope material.
982 *Nature Geoscience*, 3(4), 247–251. <https://doi.org/10.1038/ngeo776>
- 983 Legros, F. (2002). The mobility of long-runout landslides. *Engineering Geology*, 63(3–4), 301–331.
984 [https://doi.org/10.1016/S0013-7952\(01\)00090-4](https://doi.org/10.1016/S0013-7952(01)00090-4)
- 985 Leshchinsky, B. (2019). Quantifying the influence of failure surface asperities on the basal shear
986 resistance of translational landslides. *Landslides*, 16(7), 1375–1383.
987 <https://doi.org/10.1007/s10346-019-01185-9>
- 988 Mackey, B. H., & Roering, J. J. (2011). Sediment yield, spatial characteristics, and the long-term
989 evolution of active earthflows determined from airborne LiDAR and historical aerial
990 photographs, Eel River, California. *Geological Society of America Bulletin*, 123(7–8), 1560–
991 1576. <https://doi.org/10.1130/B30306.1>

- 992 Mackey, B. H., Roering, J. J., & McKean, J. (2009). Long-term kinematics and sediment flux of an active
993 earthflow, Eel River, California. *Geology*, 37(9), 803–806. <https://doi.org/10.1130/G30136A.1>
- 994 Madson, A., Fielding, E., Sheng, Y., & Cavanaugh, K. (2019). High-resolution spaceborne, airborne and
995 in situ landslide kinematic measurements of the slumgullion landslide in Southwest Colorado.
996 *Remote Sensing*, 11(3), 265. <https://doi.org/10.3390/rs11030265>
- 997 Malet, J.-P., Maquaire, O., & Calais, E. (2002). The use of Global Positioning System techniques for the
998 continuous monitoring of landslides: application to the Super-Sauze earthflow (Alpes-de-Haute-
999 Provence, France). *Geomorphology*, 43(1–2), 33–54. [https://doi.org/10.1016/S0169-
1000 555X\(01\)00098-8](https://doi.org/10.1016/S0169-555X(01)00098-8)
- 1001 Marone, C. (1998). Laboratory-derived friction laws and their application to seismic faulting. *Annual*
1002 *Review of Earth and Planetary Sciences*, 26(1), 643–696.
1003 <https://doi.org/10.1146/annurev.earth.26.1.643>
- 1004 Marc, O., & Hovius, N. (2015). Amalgamation in landslide maps: effects and automatic detection.
1005 *Natural Hazards & Earth System Sciences*, 15, 723–733. [https://doi.org/10.5194/nhess-15-723-
1006 2015](https://doi.org/10.5194/nhess-15-723-2015)
- 1007 McLaughlin, R. J., Kling, S. A., Poore, R. Z., McDougall, K., & Beutner, E. C. (1982). Post–middle
1008 Miocene accretion of Franciscan rocks, northwestern California. *Geological Society of America*
1009 *Bulletin*, 93(7), 595–605. [https://doi.org/10.1130/0016-7606\(1982\)93<595:PMAOFR>2.0.CO;2](https://doi.org/10.1130/0016-7606(1982)93<595:PMAOFR>2.0.CO;2)
- 1010 McLaughlin, R. J., Blake, S., Jayko, M., Irwin, A., Aalto, W., Carver, K., et al. (2000). Geologic map of
1011 the Cape Mendocino, Eureka, Garberville, and southwestern part of the Hayfork 30 X 60
1012 Quadrangles and Adjacent Offshore Area, Northern California.
- 1013 Meyer, C. R., Downey, A. S., & Rempel, A. W. (2018). Freeze-on limits bed strength beneath sliding
1014 glaciers. *Nature Communications*, 9(1), 1–6. <https://doi.org/10.1038/s41467-018-05716-1>
- 1015 Michel, J., Dario, C., Marc-Henri, D., Thierry, O., Marina, P. I., & Benjamin, R. (2020). A review of
1016 methods used to estimate initial landslide failure surface depths and volumes. *Engineering*
1017 *Geology*, 267, 105478. <https://doi.org/10.1016/j.enggeo.2020.105478>

- 1018 Milledge, D. G., Bellugi, D., McKean, J. A., Densmore, A. L., & Dietrich, W. E. (2014). A
1019 multidimensional stability model for predicting shallow landslide size and shape across
1020 landscapes. *Journal of Geophysical Research: Earth Surface*, *119*(11), 2481–2504.
1021 <https://doi.org/10.1002/2014JF003135>
- 1022 Nereson, A. L., Davila Olivera, S., & Finnegan, N. J. (2018). Field and Remote-Sensing Evidence for
1023 Hydro-mechanical Isolation of a Long-Lived Earthflow in Central California. *Geophysical
1024 Research Letters*, *45*(18), 9672–9680. <https://doi.org/10.1029/2018GL079430>
- 1025 Nereson, A. L., & Finnegan, N. J. (2019). Drivers of earthflow motion revealed by an 80 yr record of
1026 displacement from Oak Ridge earthflow, Diablo Range, California, USA. *Geological Society of
1027 America Bulletin*. *131*(3-4), 389-402. <https://doi.org/10.1130/B32020.1>
- 1028 Pathier, E., Fielding, E. J., Wright, T. J., Walker, R., Parsons, B. E., & Hensley, S. (2006). Displacement
1029 field and slip distribution of the 2005 Kashmir earthquake from SAR imagery. *Geophysical
1030 Research Letters*, *33*(20). <https://doi.org/10.1029/2006GL027193>
- 1031 Roadifer, J. W., Forrest, M. P., & Lindquist, E. S. (2009). Evaluation of shear strength of mélangé
1032 foundation at Calaveras Dam. Proceedings of U. S. Society for Dams, Annual Meeting and
1033 Conference, 29th, on "Managing Our Water Retention Systems, ", 507–521.
- 1034 Roering, J.J (2012). Eel River, CA: Landsliding and the Evolution of Mountainous Landscapes in
1035 collaboration with National Center for Airborne Laser Mapping (NCALM), distributed by
1036 OpenTopography. <https://doi.org/10.5069/G9XS5S9P>
- 1037 Roering, J. J., Stimely, L. L., Mackey, B. H., & Schmidt, D. A. (2009). Using DInSAR, airborne LiDAR,
1038 and archival air photos to quantify landsliding and sediment transport. *Geophysical Research
1039 Letters*, *36*(19). <https://doi.org/10.1029/2009GL040374>
- 1040 Roering, J. J., Mackey, B. H., Handwerger, A. L., Booth, A. M., Schmidt, D. A., Bennett, G. L., &
1041 Cerovski-Darriau, C. (2015). Beyond the angle of repose: A review and synthesis of landslide
1042 processes in response to rapid uplift, Eel River, Northern California. *Geomorphology*, *236*, 109–
1043 131. <https://doi.org/10.1016/j.geomorph.2015.02.013>

- 1044 Rosen, P. A., Gurrola, E., Sacco, G. F., & Zebker, H. (2012). The InSAR scientific computing
1045 environment. In *Synthetic Aperture Radar, 2012. EUSAR. 9th European Conference on* (pp. 730–
1046 733). VDE.
- 1047 Rutter, E., & Green, S. (2011). Quantifying creep behaviour of clay-bearing rocks below the critical stress
1048 state for rapid failure: Mam Tor landslide, Derbyshire, England. *Journal of the Geological*
1049 *Society*, 168(2), 359–372. <https://doi.org/10.1144/0016-76492010-133>
- 1050 Schulz, W. H., Coe, J. A., Ricci, P. P., Smoczyk, G. M., Shurtleff, B. L., & Panosky, J. (2017). Landslide
1051 kinematics and their potential controls from hourly to decadal timescales: Insights from
1052 integrating ground-based InSAR measurements with structural maps and long-term monitoring
1053 data. *Geomorphology*, 285, 121–136. <https://doi.org/10.1016/j.geomorph.2017.02.011>
- 1054 Schulz, W. H., Smith, J. B., Wang, G., Jiang, Y., & Roering, J. J. (2018). Clayey landslide initiation and
1055 acceleration strongly modulated by soil swelling. *Geophysical Research Letters*, 45(4), 1888–
1056 1896. <https://doi.org/10.1002/2017GL076807>
- 1057 Simoni, A., Ponza, A., Picotti, V., Berti, M., & Dinelli, E. (2013). Earthflow sediment production and
1058 Holocene sediment record in a large Apennine catchment. *Geomorphology*, 188, 42–53.
1059 <https://doi.org/10.1016/j.geomorph.2012.12.006>
- 1060 Stumpf, A., Malet, J.-P., & Delacourt, C. (2017). Correlation of satellite image time-series for the
1061 detection and monitoring of slow-moving landslides. *Remote Sensing of Environment*, 189, 40–
1062 55. <https://doi.org/10.1016/j.rse.2016.11.007>
- 1063 Swanston D. N., R. R. Ziemer, & R. J. Janda (1995), Rate and mechanisms of progressive hillslope
1064 failure in the Redwood Creek basin, northwestern California, in *Geomorphic Processes and*
1065 *Aquatic Habitat in the Redwood Creek Basin, Northwestern California*, edited by K. M. Nolan,
1066 H. M. Kelsey, and D. C. Marron, U.S. Geological Survey Professional Paper 1454, pp. E1–E16,
1067 United States Government Printing Office, Washington, D.C.
- 1068 Travelletti, J., & Malet, J.-P. (2012). Characterization of the 3D geometry of flow-like landslides: A
1069 methodology based on the integration of heterogeneous multi-source data. *Engineering Geology*,

- 1070 128, 30–48. <https://doi.org/10.1016/j.enggeo.2011.05.003>
- 1071 Travelletti, J., Malet, J.-P., & Delacourt, C. (2014). Image-based correlation of Laser Scanning point
1072 cloud time series for landslide monitoring. *International Journal of Applied Earth Observation
1073 and Geoinformation*, 32, 1–18. <https://doi.org/10.1016/j.jag.2014.03.022>
- 1074 Van Asch, T. J., & Van Genuchten, P. M. B. (1990). A comparison between theoretical and measured
1075 creep profiles of landslides. *Geomorphology*, 3(1), 45-55. [https://doi.org/10.1016/0169-
1076 555X\(90\)90031-K](https://doi.org/10.1016/0169-555X(90)90031-K)
- 1077 Van Asch, T. W., Van Beek, L., & Bogaard, T. (2007). Problems in predicting the mobility of slow-
1078 moving landslides. *Engineering Geology*, 91(1), 46–55.
1079 <https://doi.org/10.1016/j.enggeo.2006.12.012>
- 1080 Warrick, J. A., Ritchie, A. C., Schmidt, K. M., Reid, M. E., & Logan, J. (2019). Characterizing the
1081 catastrophic 2017 Mud Creek landslide, California, using repeat structure-from-motion (SfM)
1082 photogrammetry. *Landslides*, 1–19. <https://doi.org/10.1007/s10346-019-01160-4>
- 1083 Wartman, J., Montgomery, D. R., Anderson, S. A., Keaton, J. R., Benoît, J., dela Chapelle, J., & Gilbert,
1084 R. (2016). The 22 March 2014 Oso landslide, Washington, USA. *Geomorphology*, 253, 275–288.
1085 <https://doi.org/10.1016/j.geomorph.2015.10.022>