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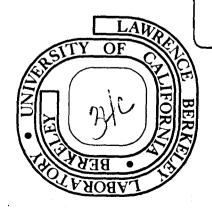
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HIGH-SPIN STATES IN ¹²⁷, ¹²⁹La; A TEST OF BACKBENDING IN THE EVEN Ba AND Ce NUCLEI*

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ABSTRACT: The $h_{11/2}$ band in 127,129 La has been identified to spin $35/2^-$, and it does not backbend. This suggests that rotation-aligned $h_{11/2}$ protons are a major component of high-spin states in neighboring Ba and Ce nuclei which do exhibit backbending.

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The rotation-aligned coupling scheme [1] has had considerable success in describing the collective bands in a number of weakly deformed odd-mass nuclei. It has also been pointed out that the presence of backbending in an odd-mass rare-earth nucleus is an important clue as to the cause of backbending in the adjacent doubly-even nuclei. If a particular rotation-aligned particle delays or prevents the occurrence of backbending in the odd-mass case, it is likely that this particle is involved in the structure of the band whose intersection with the ground band causes the backbending in the doubly-even nucleus [2]. Conversely, if the oddparticle band does backbend, this particle must not be such an important component [3]. Application of this idea has shown the importance of $i_{13/2}$ neutrons in the backbending behavior observed in the lighter deformed nuclei Dy [4], Er [2], and Yb [5]. Very recently the importance of rotation-aligned $h_{q/2}$ protons in the backbending of the heavier deformed Hf, W, and Os nuclei has been demonstrated by studying 167 Lu [6] and 181 Re, 181,183 Os [7]. It may be that in some of these cases both the $i_{13/2}$ neutrons and the $h_{0/2}$ protons are important, presumably as a mixed band, but this is not yet established.

In this letter we report a study of the high-spin decoupled states of ^{127,129}La from which we can draw conclusions concerning the causes of the backbending observed in the neighboring Ba and Ce nuclei. To our knowledge, this is the first time this method has been applied outside the rare-earth deformed region.

Previous studies of high-spin states in 126,128 Ba [8] and 130 Ce [9] have shown that the ground bands backbend around spin 10^+ . This

corresponds to a spin of 31/2 in the 11/2 band in La, but the earlier work on ^{127,129}La [10] only assigned levels as high as 27/2. In the present work we have identified levels to spin 35/2, just beyond the critical region. To obtain as complete a comparison as possible, we have also studied the ground band in ¹²⁸Ce, since it was previously known [11] only to spin 8⁺.

The experiments were performed at the LBL 88" Cyclotron, using the reactions $^{117}\mathrm{Sn}(^{14}\mathrm{N},4\mathrm{n}\gamma)^{127}\mathrm{La}$ with 75 MeV $^{14}\mathrm{N}$, $^{119}\mathrm{Sn}(^{14}\mathrm{N},4\mathrm{n}\gamma)^{129}\mathrm{La}$ with 67 and 75 MeV $^{14}\mathrm{N}$, and $^{116}\mathrm{Sn}(^{16}\mathrm{O},4\mathrm{n}\gamma)^{128}\mathrm{Ce}$ with 88 MeV $^{16}\mathrm{O}$. The targets were made by evaporating the tin isotope to $^{\circ}$ 1 mg/cm² thickness on backings of $^{\circ}$ 50 mg/cm² of lead. Gamma rays were detected both in singles and in coincidence mode using two 45 cc coaxial Ge(Li) detectors. The 0° to 90° anisotropy of the γ -rays was measured to test that the main cascade sequences were indeed stretched E2 transitions.

To reduce the number of magnetic tapes collected in the γ-γ coincidence runs, whilst retaining high statistics, only events were stored which corresponded to restricted regions of the pulse height spectrum from one of the γ-detectors. Four S.C.A.'s were used to define the regions of interest. For example, in ¹²⁹La, windows about 20 keV wide were set on the three lowest transitions of the decoupled band and the last window was set from just below the fourth transition (786 keV) up to about 1200 keV. It was estimated that this procedure increases the useful data on a magnetic tape by about a factor of eight compared with storing all coincident events. Furthermore, the effective count-rate could be increased considerably, since the dead time involved in writing events on magnetic tape was the limiting

factor with our apparatus.

Typical coincidence spectra for 129 La are shown in fig. 1, and the cascade sequences deduced from individual gates set on the known ground-band transitions are summarized in table 1. The A_2 values derived from the anisotropy measurements assuming $A_4 = -0.05 \pm 0.05$ are consistent with stretched E2 assignments, (table 1). However, in several cases the γ -ray lines in the direct (singles) spectra were not pure and it was rather difficult to extract reliable values for their A_2 coefficients.

The coincidence data also enabled us to identify side-bands in both 127 La and 129 La. These are shown in fig. 2. The angular anisotropy of the γ -cascades within the La side-bands indicates that they are probably stretched E2 transitions. Spin and parity assignments can be proposed for these levels, but they are rather uncertain and not relevant to the main point of this letter.

To compare the systematics of the core states in 127,129 La with those of the corresponding Ba and Ce nuclei we have calculated the moments of inertia, \mathcal{F} and the rotational frequencies, ω , according to the formulae given in [3]. Plots of \mathcal{F} vs. ω^2 are shown in fig. 3. It can be seen that the $11/2^-$ decoupled band in La does not backbend, even for rotational frequencies appreciably higher than those for which the Ba and Ce nuclei do backbend. This fact suggests that a rotation alignment of two such $h_{11/2}$ protons is involved in the backbending in the even barium and cerium nuclei. In fact it seems likely that the backbending is caused by an intersection of the ground band with a band composed of two $h_{11/2}$ protons maximally aligned with the rotation axis. It is of interest not only that the backbending is caused by protons rather than neutrons in

this region, but also that it is most likely due to a rotation-alignment effect like that observed in the rare-earth region, rather than to a shape change as is the case for the neutron-deficient Hg nuclei. It appears that the technique of studying an even-even case by observing its behavior with various particle orbits blocked, is applicable to this Ba-Ce region, as well as to the deformed rare-earth region, and can probably be applied rather generally.

ACKNOWLEDGMENTS

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Table 1. Summary of the rotational cascades in 127,129 La and 128 Ce observed in this experiment. Assignments in La nuclei to spin 27/2 (8⁺) have been made previously [10]. Values for 126,128 Ba and 130 Ce are from previous work; 128 Ce assignments [8,9] to spin 8⁺ have been made previously [11]. Uncertainties in the relative intensities and A₂ coefficients are shown in parentheses.

Spin	1	128 _{Ba} [8] E _γ (keV)		129 La (75 MeV)			
I (or I - 11/2)	1			E _Y Int		A ₂	
2		284	269.3	100	0.	35(3)	254.1
4		479	474.8	81(8)	0.	32(3)	456.7
6		644	641.9	55 * (16	0.	.35(3)	614.2
8		782	785.6	29(3)	0.	31(4)	729.6
10		894	910.5	11(2)	0.	26(15)	757.0
12		906	1013.5	8 (2)	0.	.02(30)	503.3
14							549.0
16							693.5
Spin	126 _{Ba[8]}	12	⁷ La (76 Me	v)	128 _{Ce} (88 MeV)		
(or I - 11/2)	1	}	Int	A ₂	E _Y	Int	A ₂
2	256	252.4	100	0.19(3)	207.3	100	0.26(3)
4	455	458.6	73(7)	0.28(3)	399.6	88(9)	0.23(3)
6	621	630.8	48(5)	0.23(3)	550.6	88(9)	0.32(3)
8	757	780.3	3 * (4)	0.13(4)	662.6	50(5)	0.26(4)
10	852	908	8(2)	0.10(15)	711.1	33(3)	0.31(6)
12	805	1003	* (1)	0.16(30)	576.5	28(3)	0.25(6)
14	672				560.6	22(2)	0.24(6)
16	825				689.9	14(2)	0.38(10

^{*}Not single lines in the direct spectra. Relative intensities taken from spectra in coincidence with the $2^+ \rightarrow 0^+$ (15/2 \rightarrow 11/2) transitions.

FIGURE CAPTIONS

- Figure 1. Coincidence spectra for 129 La, spectrum (a) is for the sum of gates set on the lowest two rotational transitions, (b) is for the $31/2 \rightarrow 27/2$ transition. Spectra corresponding to background events have been subtracted in an appropriate way.
- Figure 2. Level schemes for 127,129 La and 128 Ce.
- Figure 3. Comparison of backbending plots for *core states* in ^{127,129}La with adjacent even nuclei.



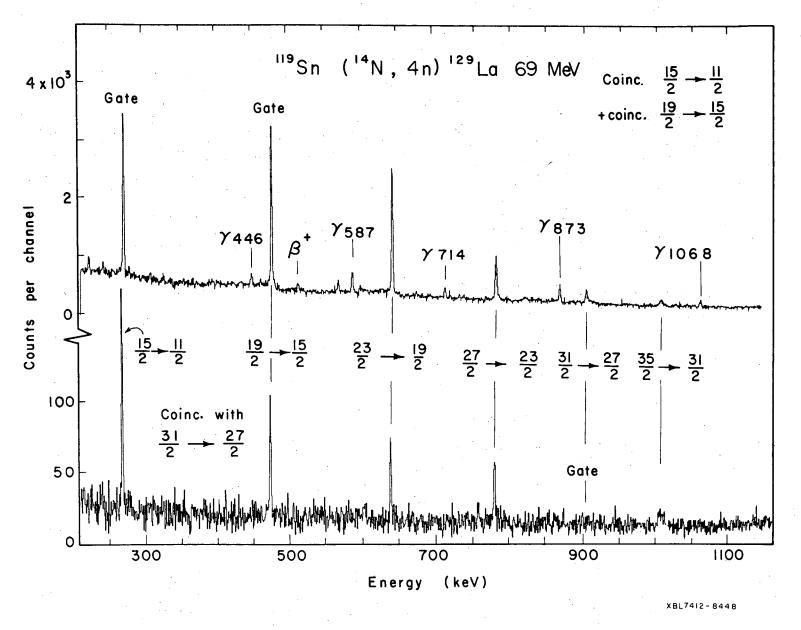
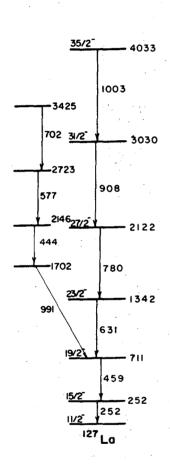
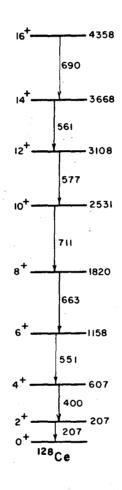
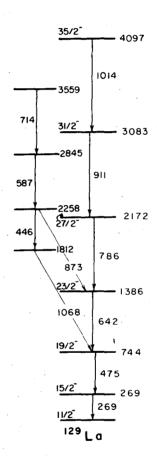


Fig. 1







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Fig. 2

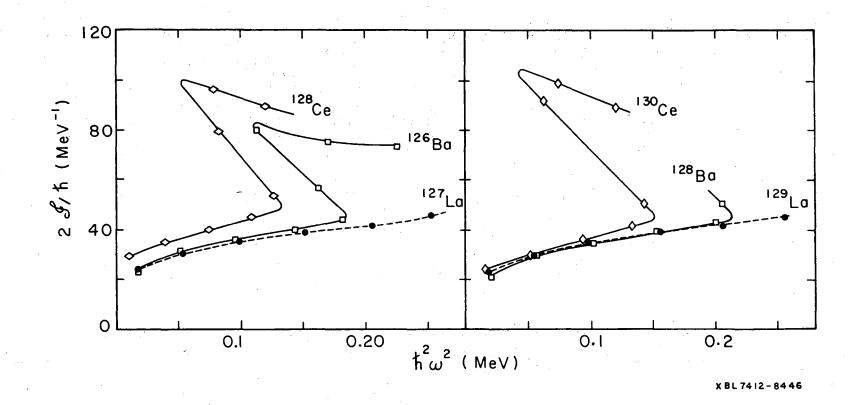


Fig. 3

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