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A local baseline of the black hole mass scaling relations for active galaxies. IV. Correlations between \$M\_{\rm BH}\$ and host galaxy \$σ\$, stellar mass, and luminosity

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#### <span id="page-1-0"></span>A local baseline of the black hole mass scaling relations for active galaxies. IV. Correlations between  $M_{\text{BH}}$  and host galaxy  $\sigma$ , stellar mass, and luminosity.

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#### ABSTRACT

The tight correlations between the mass of supermassive black holes  $(M_{BH})$  and their host-galaxy properties have been of great interest to the astrophysical community, but a clear understanding of their origin and fundamental drivers still eludes us. The local relations for active galaxies are interesting in their own right and form the foundation for any evolutionary study over cosmic time. We present Hubble Space Telescope optical imaging of a sample of 66 local active galactic nuclei (AGNs); for 14 objects, we also obtained Gemini near-infrared images. We use state of the art methods to perform surface photometry of the AGN host galaxies, decomposing them in spheroid, disk and bar (when present) and inferring the luminosity and stellar mass of the components. We combine this information with spatially-resolved kinematics obtained at the Keck Telescopes to study the correlations between  $M_{\rm BH}$  (determined from single-epoch virial estimators) and host galaxy properties. Our sample extends the correlation found for quiescent galaxies down to  $M_{\text{BH}} \sim 10^7$  M<sub>⊙</sub>, along a consistent line. The correlations are uniformly tight for our AGN sample, with intrinsic scatter  $0.2 - 0.4$  dex, smaller or equal to that of quiescent galaxies. We find no difference between pseudo and classical bulges or barred and non-barred galaxies. We show that all the tight correlations can be simultaneously satisfied by AGN hosts in the  $10^7$ - $10^9$   $M_\odot$  regime, with data of sufficient quality. The  $M_{\text{BH}}$ - $\sigma$  relation is also in excellent agreement with that of AGN with  $M_{\rm BH}$  obtained from reverberation mapping, providing an indirect validation of single-epoch virial estimators of  $M_{\rm BH}$ .

Keywords: accretion, accretion disks – black hole physics – galaxies: active – galaxies: evolution – galaxies: Seyfert

#### 1. INTRODUCTION

When growing through accretion, supermassive black holes (BHs) can be seen as bright nuclei in active galaxies (AGNs). The observed relations between the mass of the BH  $(M_{BH})$  and the properties of the host-galaxy spheroid such as luminosity, stellar mass and stellar-

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velocity dispersion  $\sigma$ , are thought to result from the coevolution between BHs and galaxies (for a review see, e.g., [Kormendy & Ho](#page-22-0) [2013;](#page-22-0) [Graham](#page-21-0) [2016\)](#page-21-0). Such a coevolution is either regulated by AGN feedback (e.g., [Di](#page-21-1) [Matteo et al.](#page-21-1) [2005;](#page-21-1) [Croton](#page-21-2) [2006;](#page-21-2) [Dubois et al.](#page-21-3) [2013,](#page-21-3) [2016;](#page-21-4) [DeGraf et al.](#page-21-5) [2015;](#page-21-5) [Hopkins et al.](#page-22-1) [2016\)](#page-22-1), or hierarchical assembly of  $M_{\rm BH}$  and stellar mass through galaxy merging (e.g., [Peng](#page-22-2) [2007;](#page-22-2) [Hirschmann et al.](#page-22-3) [2010;](#page-22-3) [Jahnke & Maccio](#page-22-4) [2011\)](#page-22-4). To shed light on the origin of these relations, recent years have seen an explosion of observational studies both in the local Universe (e.g.,

Ferrarese  $&$  Ford [2005;](#page-21-6) Greene  $&$  Ho [2006;](#page-21-7) Gültekin et [al.](#page-22-5) [2009;](#page-22-5) [Bennert et al.](#page-21-8) [2011a;](#page-21-8) [Kormendy et al.](#page-22-6) [2011;](#page-22-6) [Beifiori et al.](#page-21-9) [2012;](#page-21-9) Läsker et al. [2016;](#page-22-7) [Davis et al.](#page-21-10) [2018;](#page-21-10) [Sahu et al.](#page-22-8) [2019\)](#page-22-8) and as a function of cosmic history (e.g. [Treu et al.](#page-22-9) [2004;](#page-22-9) [Peng et al.](#page-22-10) [2006a,](#page-22-10)[b;](#page-22-11) [Woo et al.](#page-22-12) [2006;](#page-22-12) [Salviander et al.](#page-22-13) [2007;](#page-22-13) [Riechers et al.](#page-22-14) [2009;](#page-22-14) [Jahnke et al.](#page-22-15) [2009;](#page-22-15) [Bennert et al.](#page-21-11) [2010;](#page-21-11) [Decarli et al.](#page-21-12) [2010;](#page-21-12) [Merloni et](#page-22-16) [al.](#page-22-16) [2010;](#page-22-16) [Bennert et al.](#page-21-13) [2011b;](#page-21-13) [Park et al.](#page-22-17) [2015;](#page-22-17) [Sexton](#page-22-18) [et al.](#page-22-18) [2019;](#page-22-18) [Silverman et al.](#page-22-19) [2019;](#page-22-19) [Ding et al.](#page-21-14) [2020\)](#page-21-14).

By necessity, all studies beyond the local Universe focus on broad-line (or type-1) AGNs (BLAGNs). For BLAGNs,  $M_{\text{BH}}$  can be estimated to within a factor of 2-3 using empirically calibrated relations based on a sample of reverberation-mapped AGNs. Reverberation mapping (RM) is a technique that studies the time delay between the variability of the accretion disk and the response of ionized gas in the vicinity of the BH, the broad-line region (BLR) (e.g., [Wandel et al.](#page-22-20) [1999;](#page-22-20) [Woo](#page-22-21) [& Urry](#page-22-21) [2002;](#page-22-21) [Vestergaard](#page-22-22) [2002;](#page-22-22) [Vestergaard & Peterson](#page-22-23) [2006;](#page-22-23) [McGill et al.](#page-22-24) [2008\)](#page-22-24). Using light-travel time arguments, the time delay translates into a size of the BLR. Combining the size with the Doppler-broadened width of the emission lines (e.g., the Hydrogen Balmer series in the optical) results in an estimate of the  $M_{\rm BH}$  up to an unknown factor that depends on the geometry and kinematics of the gas clouds. Traditionally, this factor f (also known as virial factor) has been derived as a sample average by matching the scaling relation between  $M_{\rm BH}$  and (spheroid) stellar-velocity dispersion  $\sigma$  of the RM AGNs with that of quiescent galaxies (e.g., [Onken et](#page-22-25) [al.](#page-22-25) [2004;](#page-22-25) [Park et al.](#page-22-26) [2012;](#page-22-26) [Woo et al.](#page-22-27) [2010,](#page-22-27) [2015\)](#page-23-0). More recently, dynamical modeling of RM data has been used to constrain both geometry and kinematics of the BLR and thus determine  $M_{\rm BH}$  for individual objects, finding consistent results (e.g., [Brewer et al.](#page-21-15) [2011;](#page-21-15) [Pancoast et](#page-22-28) [al.](#page-22-28) [2011;](#page-22-28) [Li et al.](#page-22-29) [2013;](#page-22-29) [Pancoast et al.](#page-22-30) [2018;](#page-22-30) [Williams et](#page-22-31) [al.](#page-22-31) [2018,](#page-22-31) [2020\)](#page-22-32). While RM is time-consuming, the RM AGN sample revealed a relation between BLR size and AGN luminosity that can be used to estimate  $M_{\rm BH}$  for BLAGNs from one spectrum, known as the single-epoch method. In the single-epoch virial estimation, the width of broad emission lines is combined with the AGN luminosity which serves as a proxy for BLR size. As such, the RM AGN sample serves as a  $M_{\rm BH}$  calibrator beyond the local Universe. The single-epoch method has been used for virial mass estimates of hundreds of thousands of AGNs (e.g., [Rakshit et al.](#page-22-33) [2020\)](#page-22-33), across cosmic history (e.g., [Mortlock et al.](#page-22-34) [2011\)](#page-22-34), to study the cosmic evolution of the  $M_{\rm BH}$  scaling relations (e.g., [Treu et al.](#page-22-9) [2004;](#page-22-9) [Peng et al.](#page-22-10) [2006a;](#page-22-10) [Woo et al.](#page-22-12) [2006;](#page-22-12) [Bennert et al.](#page-21-11) [2010;](#page-21-11) [Merloni et al.](#page-22-16) [2010;](#page-22-16) [Park et al.](#page-22-17) [2015;](#page-22-17) [Ding et al.](#page-21-14)

[2020\)](#page-21-14) and distribution of Eddington ratios (e.g., [Shen](#page-22-35) [2013\)](#page-22-35).

Studies of the evolution of the  $M_{\rm BH}$ -host-galaxy scaling relations with redshift constrain theoretical interpretations and shed light onto their origin (e.g., [Croton](#page-21-2) [2006;](#page-21-2) [Hopkins et al.](#page-22-36) [2007\)](#page-22-36); however, they depend on our understanding of the slope and intrinsic scatter of local relations, in particular those for active galaxies. Moreover, studying dependencies of the correlations on bulge structure and other morphological components at highredshifts is difficult if not impossible, especially given the presence of the bright AGN point source in the center. Late-type galaxies are often known to host pseudobulges, characterized by exponential light profiles, ongoing star formation or starbursts, and nuclear bars. It is generally believed that they have evolved secularly through dissipative processes rather than mergers (e.g., [Courteau et al.](#page-21-16) [1996;](#page-21-16) [Kormendy & Kennicutt](#page-22-37) [2004\)](#page-22-37). Classical bulges, in contrast, are thought of as centrally concentrated, mostly red and quiescent, merger-induced systems. Pseudo-bulges and minor mergers provide a valuable test of some hypotheses for the origin of the  $M_{\rm BH}$  scaling relations: If they lie off the relations, major mergers can be considered the fundamental driver [\(Peng](#page-22-2) [2007;](#page-22-2) [Hirschmann et al.](#page-22-3) [2010;](#page-22-3) [Jahnke & Maccio](#page-22-4) [2011;](#page-22-4) [Kormendy & Ho](#page-22-0) [2013\)](#page-22-0); if they lie on the relation, as found by our results based on SDSS images [\(Bennert](#page-21-17) [et al.](#page-21-17) [2015\)](#page-21-17), it could indicate that secular evolution has a synchronizing effect, growing BHs and bulges simultaneously at a small but steady rate for late-type galaxies [\(Cisternas et al.](#page-21-18) [2011a,](#page-21-18)[b\)](#page-21-19).

This paper is the last of a series aimed at creating a robust local baseline of the  $M<sub>BH</sub>$  scaling relations of BLAGNs for comparison with high redshift studies. We selected a sample of ∼100 Seyfert-1 galaxies from SDSS  $(0.02 \leq z \leq 0.1; M_{\text{BH}} > 10^7 \text{M}_{\odot})$  based on their broad  $H\beta$  emission in the same fashion as high-redshift samples used for evolutionary studies [\(Bennert et al.](#page-21-11) [2010;](#page-21-11) [Park et al.](#page-22-17) [2015;](#page-22-17) [Ding et al.](#page-21-14) [2020\)](#page-21-14), allowing for a direct comparison. The majority of AGNs (∼80%) in our sample reside in galaxies classified as Sa or later [\(Bennert et al.](#page-21-17) [2015\)](#page-21-17), comparable to our high-redshift samples [\(Bennert et al.](#page-21-11) [2010;](#page-21-11) [Park et al.](#page-22-17) [2015\)](#page-22-17), perhaps not surprisingly, given that all studies focus on Seyfert-1 galaxies. In paper I and III [\(Bennert et al.](#page-21-8) [2011a,](#page-21-8) [2015\)](#page-21-17), multi-filter SDSS images yielded photometric parameters such as the spheroid effective radius, the spheroid luminosity, the host-galaxy free  $5100\text{\AA}$  luminosity of the AGN (for an accurate  $M_{\rm BH}$  measurement), and spheroid stellar masses. In paper II [\(Har](#page-22-38)[ris et al.](#page-22-38) [2012\)](#page-22-38), high-quality long-slit Keck/LRIS spectra provided both  $M_{\rm BH}$  estimates as well as accurate spatially-resolved stellar-velocity dispersions  $(\sigma)$  and rotation curves.

Given the wide range of  $M_{\rm BH}$ , host-galaxy morphologies, and stellar masses, our sample is well suited to determine the slope and intrinsic scatter of the local scaling relations and to study dependencies on other parameters such as bulge structure and mergers. However, relying on low-quality optical photometry (such as SDSS) whose insufficient angular resolution and limited sensitivity to dust extinction significantly increase observational scatter in the  $M_{\rm BH}$  scaling relations, ultimately limits conclusive results. High-resolution images are essential to resolve (pseudo-) bulges, given that roughly half of all objects have bulge effective radii smaller than  $\sim$ 1.5" (corresponding to  $\sim$ 1.7 kpc for a typical distance of the galaxies in our sample; [Bennert et al.](#page-21-17) [2015\)](#page-21-17). This is not only crucial for proper morphological classification and determination of bulge luminosity, it also is important for the effective radius measurement. The latter, in turn, is important for a robust measurement of spatiallyresolved stellar-velocity dispersion within the effective spheroid radius [\(Bennert et al.](#page-21-17) [2015\)](#page-21-17). In other words, angular resolution is the key for an accurate determination of all  $M_{\rm BH}$  scaling relations.

In this paper, to overcome these problems and to obtain high-quality host-galaxy images, we took a twopronged approach. A sub-set of the parent sample (15 objects), selected to cover a wide range of morphologies (as based on SDSS images), was observed with the Near InfraRed Imager and spectrograph (NIRI) on Gemini North. Gemini-NIRI was chosen (i) for its high-spatial resolution (instrument plus site seeing) to distinguish between classical and pseudo-bulges in the presence of an AGN point source; (ii) for its large field-of-view  $(2' \times 2'$  at f/6) to measure the surface brightness profile of these nearby galaxies out to large radii; and (iii) because near-infrared observations maximize the contrast between AGN and host and minimize dust extinction, revealing the presence of (pseudo-) bulges, bars and (minor) mergers. The reduced dust extinction also makes NIR luminosities a better tracer of stellar mass. At the same time, the parent sample was part of a Hubble Space Telescope (HST) snapshot (SNAP) program (PI Bennert). HST images provide both a high spatial resolution and a stable point-spread function (PSF). WFC3/UVIS was used with broad-band filter F814W to maximize the contrast between AGN and host, avoiding strong AGN emission lines, while taking full advantage of the high resolution of UVIS. Compared to existing SDSS images, the HST images have a factor of ∼40 increase

in resolution. A total of 66 objects were observed with HST, 14 of which also have Gemini images. Gemini and HST images naturally complement each other and together provide a long wavelength range for stellar-mass determination for overlapping objects. By construction, SDSS images in five filters are available for all objects to further assist in constraining stellar masses. Combining the high-quality spectroscopic data (paper II) with highquality imaging provides a representative spectral and spatial coverage of supermassive BHs and their hosts for a detailed mapping of the local  $M_{\rm BH}$  scaling relations for active galaxies and their underlying drivers.

The paper is organized as follows. Section [2](#page-3-0) summarizes the sample selection, HST and Gemini observations, and data reduction. Section [3](#page-7-0) describes the analysis and derived quantities. Section [4](#page-14-0) presents hostgalaxy morphologies and discusses the resulting  $M_{\rm BH}$ scaling relations. Section [5](#page-16-0) concludes with a summary. Throughout this paper, magnitudes are given in AB magnitudes. For conversion to luminosities, absolute solar magnitudes were taken from [Willmer](#page-22-39) [\(2018\)](#page-22-39) and a Hubble constant of  $H_0 = 70$  km s<sup>-1</sup> and a flat Universe with a cosmological constant of  $\Omega_{\lambda} = 0.7$  are assumed.

#### <span id="page-3-0"></span>2. SAMPLE SELECTION, OBSERVATIONS, AND DATA REDUCTION

#### 2.1. Sample selection

<span id="page-3-1"></span>The parent sample is 102 type-1 Seyfert galaxies selected from the Sloan Digital Sky Survey (SDSS) data release six [\(Adelman-McCarthy et al.](#page-21-20) [2008\)](#page-21-20) based on redshift  $(0.02 \le z \le 0.1)$  and  $M_{\text{BH}}$  (>  $10^7 M_{\odot}$ ), and observed with Keck/LRIS, presented in detail in papers I, II and III in this series [\(Bennert et al.](#page-21-8) [2011a;](#page-21-8) [Harris et al.](#page-22-38) [2012;](#page-22-38) [Bennert et al.](#page-21-17) [2015\)](#page-21-17). Accurate spatially-resolved stellar-velocity dispersions were obtained for 84 objects (paper II) and formed the sample for our HST snapshot (SNAP) program. 15 objects with a wide variety of host-galaxy properties, as determined from SDSS images [\(Bennert et al.](#page-21-17) [2015\)](#page-21-17), were observed with NIRI on Gemini North (PI: Bennert; program ID GN-2016B-Q-33). 68 objects were observed as part of HST SNAP, although we were unable to determine a robust  $M_{\rm BH}$  for 2 of them due to a lack of broad  $H\beta$  in the Keck spectrum (despite it being present in prior SDSS spectra; [Runco et al.](#page-22-40) [2016\)](#page-22-40). Of the remaining 66, 14 overlap with the Gemini sample and so we include here only those 14 objects observed with both HST and Gemini. Fully-reduced SDSS images are available for all objects through the SDSS archive. Sample properties (coordinates and redshift) and derived quantities can be found in Table [1.](#page-3-1)



Table 1. Sample and derived quantities Table 1. Sample and derived quantities

# Table 1 continued Table 1 continued



 ${\bf Table\ 1\ } (continued)$ Table 1 (continued)

Table 1 continued

Table 1 continued



#### 2.2. HST observations and data reduction

84 objects were part of an HST SNAP program (HST GO 15215 and HST GO 16014, PI Bennert; Cycles 25- 27), yielding images for 68 AGN host galaxies, a completion rate of almost 80% (significantly higher than the 30% typical for SNAP programs). To obtain the dynamic range needed for an accurate decomposition of the host galaxy and the AGN, the long exposures (400 seconds) were complemented by short, unsaturated ones (between 20-100 seconds, depending on the object's brightness). To avoid buffer dump (which occurred during the long exposures), a sequence of one short and one long exposure at the same location was followed by another sequence of one long and one short exposure at a dithered location. POS TARG was used to set up a dither pattern manually that corresponds to the default WFC3-UVIS-GAP-LINE (with center UVIS). Fullframe images were read to trace the host-galaxy disks out to the background and to obtain field stars for PSF fitting of the strong AGN point source in the center.

Data processed through the standard WFC3 calibration pipeline were retrieved from the HST archive. L.A. Cosmic (Laplacian Cosmic Ray Identification) [\(van](#page-22-41) [Dokkum](#page-22-41) [2001\)](#page-22-41) was run to remove cosmic rays. All long exposures were carefully checked for saturation, especially of the bright AGN point source. For objects with saturated pixels, the short exposures taken at the same dither location as the long ones were scaled according to exposure time and used to replace the saturated pixels. Pyraf package astrodrizzle was then used to combine the two long exposures. A wide range of combinations of the final drizzle parameters scale and pixfrac were applied. After careful examination of the images and based on resolution, image quality and FWHM of the PSF, the following parameters were adopted: driz sep bits=336, final bits=336, final wcs=yes, final pixfrac=0.9, final scale=0.035, resulting in a final pixel scale of  $0.035''/\text{pix}$ . For objects without any saturated pixels, the long exposures were combined with astrodrizzle directly in the same way. Final images are shown in Figures [1-](#page-8-0)[2.](#page-9-0)

#### 2.3. Gemini observations and data reduction

15 Seyfert-1 galaxies were selected from the parent sample covering a wide range of morphologies (based on SDSS images). All galaxies were observed with NIRI on Gemini North with the largest field-of-view (FOV;  $2' \times 2'$ ) at  $f/6$ ; pixel scale of  $0.117''$  in Ks band. Observations have an average FWHM of  $0.33''$   $(0.24 - 0.44'')$  and were obtained at airmass less than 1.5. This image quality is 3-4 times better than the SDSS images which have a typical seeing of  $1.5$ ". Exposure times range between 210 seconds per image with 3-6 co-adds and 18-24 images on source, resulting in a total exposure time of 144-540 seconds, depending on the brightness of the object. The sky positions were observed with guiding enabled and were carefully selected to include a nearby bright field star. This ensured that we could generate accurate PSF models for every galaxy. Details of the observations are given in Table [2.](#page-10-0)

Data reduction was performed following standard procedures using the Gemini IRAF package customized for NIRI and included dark subtraction and flat fielding using off-target exposures. Flux calibration was obtained by standard IRAF photometry of UKIRT faint standard stars observed directly before or after the science images. Absolute magnitudes take into account extinction [\(Schlafly & Finkbeiner](#page-22-42) [2011\)](#page-22-42); Ks-band luminosities were determined assuming an absolute Ks-band magnitude of the Sun of  $M_{\text{Ks}} = 5.08$  [\(Willmer](#page-22-39) [2018\)](#page-22-39) Resulting images are shown in Figure [3.](#page-10-0)

#### 3. DERIVED QUANTITIES

#### 3.1. Surface photometry

<span id="page-7-2"></span><span id="page-7-0"></span>To perform a detailed 2D host-galaxy fitting, we use the public image analysis software lenstronomy<sup>[1](#page-7-1)</sup> [\(Birrer](#page-21-22) [et al.](#page-21-22) [2018\)](#page-21-22). Lenstronomy supersedes GALFIT [\(Peng et](#page-22-43) [al.](#page-22-43) [2002\)](#page-22-43) by applying an MCMC technique to provide realistic errors and explore the covariance between the various model parameters. It allows for a more general surface brightness reconstruction possible with a large non-parametric basis set; the coefficients are determined through a linear minimization rather than a non-linear parameter fitting [\(Birrer et al.](#page-21-23) [2015\)](#page-21-23). Also, iterative PSF reconstruction is possible and allows one to incorporate residual uncertainties due to PSF mismatch into the analysis. While lenstronomy was originally developed for galaxy-scale strong gravitational lensing, it has a much broader application, including general 2D galaxy decompositions.

A well subtracted background and a matching PSF is important for obtaining reliable host galaxy properties from 2D surface photometry. We estimate and remove the local background light in 2D space based on the SExtractorBackground algorithm built in the photutils package (Python based), which effectively accounts for gradients in the background light distribution. For all objects, PSF stars were created from suitable stars in the field-of-view (FOV) of each object, following the criteria: bright, unsaturated star without any nearby objects, located close to the AGN/center of the FOV and an over-

<span id="page-7-1"></span><sup>1</sup> <https://github.com/sibirrer/lenstronomy>

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<span id="page-8-0"></span>Figure 1. HST UVIS/F814W images of our sample. Position angles and image sizes are listed in table [1.](#page-3-1) Thanks to the high spatial resolution and S/N data, the wide variety of host-galaxy morphologies can clearly be seen.



1137+4826  $1143 + 5941$ 1144+3653 1147+0902 1205+4959 1206+4244 1216+5049 1223+0240 1246+5134 1306+4552 1307+0952 1312+2628 1416+0137 1419+0754 1405-0259 1434+4839 1545+1709 1557+0830 1605+3305 1606+3324 1611+5211 1636+4202 1708+2153 2116+1102 2140+0025 2215-0036 2221-0906 2222-0819 2233+1312 2254+0046 2327+1524 2351+1552

<span id="page-9-0"></span>Figure 2. Continuation of Figure [1.](#page-8-0)



<span id="page-10-0"></span>Figure 3. Gemini NIRI Ks images for 14 galaxies. North is up, East is to the left. Image sizes are listed in table [2.](#page-10-0)

all profile as expected for a PSF star. These individual PSF stars were then recentered and stacked, resulting in a PSF with high signal-to-noise ratio  $(S/N)$ , centered in the image. Each individual PSF star was recentered using AstroObjectAnalyzer<sup>[2](#page-10-1)</sup> through an iterative interpolation algorithm. Finally, masks were created to mask any nearby sources. Three objects have close-by neighboring galaxies which were fitted simultaneously.

The central AGN was fitted by a PSF, the host galaxies with three different models: (1) a spheroid-only component (free Sérsic index n; Sérsic [1963\)](#page-22-44); (2) a spheroid plus disk component (Sérsic index  $= 1$ ) (3) a spheroid plus disk plus bar component (Sérsic index  $= 0.5$ ). Based on pre-defined starting parameters and constraints, lenstronomy determines the maximum likelihood fit adopting a Particle Swarm Optimizer (PSO) [\(Kennedy & Eberhart](#page-22-45) [2001\)](#page-22-45). We use the following limits: effective radius  $r_{\text{eff}}$  for all components:  $3 \times$  pixel size  $\leq r_{\text{eff}} \leq 30 \times \text{pixel size}$ ; spheroid Sérsic index n:  $1 \leq n \leq 5$ . Also, for the spheroid-disk fit, we force the disk to be larger than the spheroid and more elliptical. Likewise, for the spheroid-disk-bar fit, we force the disk to be larger than the bar and the spheroid, and the spheroid component to be the most round one of three components.

After running various trials, the two main challenges that we encountered were (i) determining the best values for the PSO chains, to make sure the code converged and (ii) making sure that the code converged to the true global minimum and not to a local one. The latter may depend on the starting parameters used, especially if more than one component is fitted to the host galaxy, due to degeneracies involved and the high-dimensional parameter volume. After some experimenting, the following procedure was shown to be successful. We chose a PSO chain that guaranteed convergence even for the largest image size. For a spheroid-only fit, we chose a PSO with 200 particles and 70 iterations, for spheroiddisk and spheroid-disk-bar, a PSO with 300 particles and 100 iterations. For all fits, as a diagnostic, we displayed the log (likelihood) of the fit, particle position and parameter velocity for the different parameters as a function of iteration to ensure that the chain indeed converged.

<span id="page-10-1"></span><sup>2</sup> <https://github.com/sibirrer/AstroObjectAnalyser>

Object	Date of Obs.	Exp. time	<b>FWHM</b>	Frame size
	(UT)	(s)	$($ ")	$($ ")
(1)	(2)	(3)	(4)	(5)
$0013 - 0951$	2016 Oct 25	378	0.31	21.1
$0109 + 0059$	2016 Oct 14	360	0.32	18.7
$0121 - 0102$	2016 Sep 06	360	0.35	23.4
$0150 + 0057$	2017 Jan 09	216	0.27	23.4
$0206 - 0017$	2017 Jan 09	144	0.24	46.8
$0301 + 0115$	2016 Nov 06	360	0.36	16.4
$0813 + 4608$	2016 Oct 20	360	0.30	15.2
$0845 + 3409$	2016 Oct 30	324	0.40	28.1
2140+0025	2016 July 16	360	0.29	16.4
$2221 - 0906$	2016 July 16	540	0.32	16.4
2222-0819	2016 Sep 08	378	0.27	21.1
2233+1312	2016 Aug 7	216	0.32	25.7
2327+1524	2016 Aug 8	216	0.44	35.1
$2351 + 1552$	2016 Oct 15	432	0.26	18.7

Table 2. Gemini Observations.

NOTE— Col. (1): Target ID used throughout the text (based on R.A. and declination). Col. (2): Date of observation (UT). Col. (3): Total exposure time in seconds. Col. (4): Full-Width at Half Maximum (FWHM) of PSF star in arcsecond. Col. (5): Frame size of image shown in Fig. [3](#page-10-0) (in arcsecond, in x and y).

First, we ran the spheroid-only fit which was shown to be robust, i.e. giving the same fitting result, regardless of starting parameters used. We then used the results from the spheroid-only fit as starting parameters for the disk in the subsequent spheroid-disk and spheroid-diskbar fit. For objects for which there is at least an indication of a visual bar, the bar parameters size, position angle and ellipticity were carefully determined manually and used as starting parameters for the bar in the spheroid-disk-bar fit. For both the spheroid-disk and spheroid-disk-bar fit, we chose three different starting parameters for the size of the spheroid and the disk, to cover a broader range. For the spheroid effective radius, we chose three different starting parameters: (i) pixel size  $*$  4, (ii) pixel size  $*$  8, and (iii) pixel size  $*$  30. For the disk effective radius we chose these three different starting parameters: (i) spheroid effective radius from spheroid-only fit, (ii) twice the size used in (i), and (iii) half the size used in (i). When combined, this yields nine different starting parameters for both the spheroid-disk and spheroid-disk-bar fits. The nine fits were compared in terms of image residuals and resulting chi-squared values. In most cases, the nine different fits yielded identical results, showing the convergence to a true global minimum. Occasionally, outliers were identified through higher chi-square values and/or from the residual image and excluded. By careful inspection

of the images and final fitting results, we determined the best model and fit for all objects. A disk and bar were included in the host-galaxy fitting only if they were clearly visible in the image and/or if their inclusion significantly improved the fit (as evidenced by chi square and residuals), beyond the typical scatter of values seen for the nine different fits for a given model. We conservatively adopt 0.04 dex as uncertainty on the derived luminosities. Figure [4](#page-13-0) shows example fits by lenstronomy with our chosen procedure. Table [4](#page-1-0) in the appendix gives the results of the fitting.

Galactic foreground extinction was subtracted based on dust reddening measurements from [Schlafly &](#page-22-42) [Finkbeiner](#page-22-42) [\(2011\)](#page-22-42), assuming F814W = 0.61  $A_V$ . Magnitudes were converted to luminosities, applying a kcorrection using Astrolib  $PySymbol<sup>3</sup>$  $PySymbol<sup>3</sup>$  $PySymbol<sup>3</sup>$  and a [Kinney et](#page-22-46) [al.](#page-22-46) [\(1996\)](#page-22-46) Sa galaxy template. Note that given the low redshift of our galaxies, using a different template does not significantly change our results. Also, pysynphot gave (V-I) colors less than 1.2 magnitude for all galaxy templates (elliptical, S0, Sa, or Sb) and thus, the F814W filter magnitudes can be considered identical to I-band magnitudes [\(Harris](#page-22-47) [2018\)](#page-22-47).

<span id="page-11-0"></span><sup>3</sup> https://pysynphot.readthedocs.io/en/latest/

To derive colors, we fitted the Gemini and SDSS images (in the filters  $g'$ ,  $r'$ ,  $i'$  and  $z'$ ) in a similar way using lenstronomy. We first used the Gemini images to independently determine the host galaxy morphology of each galaxy, based on both visual inspection of the images and lenstronomy fitting results. They agree well with the conclusions reached from the HST fitting. We then adopted the same host-galaxy parameters derived from the fitting of the HST images and used lenstronomy to fit the Gemini and SDSS images, leaving only the magnitudes of PSF and host-galaxy components as free parameters. This gives us magnitudes in 5 or 6 different filters for the different host-galaxy components spheroid, (pseudo-) bulge, and disk, if present. Dust extinction and k correction were applied.

As many literature studies rely on GALFIT, we also ran GALFIT on the same background-subtracted HST images for comparison, using the same PSFs and overall procedure as for lenstronomy. Overall, the results agree, especially (and not surprisingly) for a single-component fit (spheroid only), as well as the total host-galaxy magnitude (similar conclusions were also reached by [Yang](#page-23-1) [et al.](#page-23-1) [2020\)](#page-23-1). For more complicated models, the effective radii for individual objects scatter, but the biggest difference is seen in the spheroid Sérsic index (since disk and bar have fixed Sérsic indices). While the mean of the ratio between Sérsic index  $n$  as determined from lenstronomy and GALFIT is still around 1, it scatters greatly  $(0.97\pm0.6$  for spheroid-disk fit and  $0.97\pm0.87$  for spheriod-disk-bar fit). This cautions the usage of  $n$  alone as an indication of the nature of the spheroid (classical vs. pseudo-bulge) and in this paper, we use a conservative approach (see discussion in Section [4.1\)](#page-14-1). However, we want to stress that GALFIT tends to need more user interaction and visual inspections of results to ensure a true global minimum was reached which we did not do here. Lenstronomy's design of semi-linear inversion and PSO resulted in a significant improvement in automation and reduction of labor-intensive work in the fitting process relative to GALFIT.

#### 3.2. Stellar-velocity dispersion and black hole mass

In the literature, the observed correlation between  $M_{\rm BH}$  and  $\sigma$  is generally considered the tightest and thus most fundamental of the  $M_{\rm BH}$ -host-galaxy scaling relations [\(Tremaine et al.](#page-22-48) [2002;](#page-22-48) [Beifiori et al.](#page-21-9) [2012;](#page-21-9) [Saglia](#page-22-49) [et al.](#page-22-49) [2016;](#page-22-49) [Shankar et al.](#page-22-50) [2016;](#page-22-50) [van den Bosch](#page-22-51) [2016;](#page-22-51) [de](#page-21-24) [Nicola et al.](#page-21-24) [2019\)](#page-21-24). Moreover, it is used to calibrate  $M_{\text{BH}}$  by matching the  $M_{\text{BH}}$ - $\sigma$  relation of RM AGNs to that of quiescent galaxies. Thus, robust measurements of  $\sigma$  are essential. There are several definitions of  $\sigma$  used in the literature, resulting in widely differing measurements depending on aperture size used and host-galaxy morphology (see paper III of this series [Ben](#page-21-17)[nert et al.](#page-21-17) [2015\)](#page-21-17), with the most robust being spatiallyresolved stellar-velocity dispersions within the effective spheroid radius. Spatially-resolved stellar-velocity dispersions were presented in paper II [\(Harris et al.](#page-22-38) [2012\)](#page-22-38) based on our Keck spectra. In paper III [\(Bennert et al.](#page-21-17) [2015\)](#page-21-17), we determined  $\sigma$  from spatially-resolved  $\sigma$  measurements integrated within the effective spheroid radius (see equation (1) in paper III; [Bennert et al.](#page-21-17) [2015\)](#page-21-17). However, the effective spheroid radius in paper III was based on surface-photometry of SDSS images. We here repeat the same calculation, now using robust effective spheroid radii from the HST surface photometry. When compared, on average, the  $\sigma$  values are similar (the ones based on HST radii are larger by 1%), but with a large scatter of 10%, and a couple of individual objects having changed by as much as 50%. For 16 objects, the lack of sufficient spatially-resolved  $\sigma$  measurements hindered a robust determination of  $\sigma$  within the effective spheroid radius and they were excluded here. Thus, the  $M_{\rm BH}$ - $\sigma$  relation presented in section [4.2](#page-14-2) includes 50 objects.

 $M_{\rm BH}$  was determined for the entire sample in [Bennert](#page-21-17) [et al.](#page-21-17) [\(2015\)](#page-21-17), based on the second moment of the broad  $H\beta$  emission line determined from Keck spectroscopy. The 5100A AGN luminosity was used as a proxy for BLR size and combined with the width of  $H\beta$  to estimate  $M_{\text{BH}}$  as in equation (2) in [Bennert et al.](#page-21-17) [\(2015\)](#page-21-17). In [Bennert et al.](#page-21-17) [\(2015\)](#page-21-17), a virial factor of  $\log f = 0.71$ was assumed [\(Park et al.](#page-22-26) [2012;](#page-22-26) [Woo et al.](#page-23-0) [2015\)](#page-23-0). However, since this virial factor is based on matching the RM AGN sample to a sample of quiescent galaxies from [Mc-](#page-22-52)[Connell et al.](#page-22-52)  $(2011)$ , we here derived f independently by matching the  $M_{\rm BH}$ - $\sigma$  relation to that of [Kormendy &](#page-22-0) [Ho](#page-22-0) [\(2013\)](#page-22-0). To do so, we first fix the slope of the  $M_{\rm BH}$ - $\sigma$  relation to the one from [Kormendy & Ho](#page-22-0) [\(2013\)](#page-22-0) and then adjust  $\log f$  to match the intercept, resulting in  $\log f = 0.97$ . A wide spread in virial factors ( $\log f$  ranging between 0.5 and 1.2) has also been found in previous studies, depending on the choice of different quiescent samples, fitting methods and  $M_{\rm BH}$  range (e.g., [Park et](#page-22-26) [al.](#page-22-26) [2012;](#page-22-26) [Shankar et al.](#page-22-53) [2019\)](#page-22-53), possibly due to selection effects in the local sample of quiescent black holes.

#### 3.3. Stellar and dynamical masses

From our surface photometry (Section [3.1\)](#page-7-2), we have magnitudes for five to six different bands (HST UVIS/F814W, SDSS  $g'$ ,  $r'$ ,  $i'$ ,  $z'$  for all objects plus Gemini NIRI/Ks for 14 objects) for the different hostgalaxy components, (pseudo-) bulge, disk, and bar, if present. To estimate stellar masses from colors, we use a Bayesian stellar-mass estimation code with priors on



<span id="page-13-0"></span>Figure 4. Example fits from lenstronomy. From left to right: HST image ("data"), best fit model derived from lenstronomy ("model"), PSF-subtracted image ("data - Point Source"), residual image after subtraction of best fit model from data ("normalized residual") and surface-brightness profile showing the data as black circles, the model as a blue line, the PSF fitting the central AGN as an orange line and the host-galaxy total fit as a green line. From top to bottom, five different objects are shown, representing a variety of host-galaxy morphologies:  $1043+1105$  is an elliptical galaxy fit with a single Sérsic;  $1416+0137$ is a disk galaxy with classical bulge; 1419+0754 is a disk galaxy with pseudo-bulge; 0845+3409 is a barred disk galaxy with classical bulge; and 1206+4244 is a barred disk galaxy with pseudo-bulge.

age, metallicity and dust content of the galaxy and error bars on the different magnitudes [\(Auger et al.](#page-21-25) [2009\)](#page-21-25). To explore the full parameter space and quantify degeneracies, it uses an MCMC sampler. A Chabrier initial mass function (IMF) was assumed, but later, for comparison with literature, converted to a Kroupa IMF (by adding  $0.075$  to  $log M$ ). This gives us stellar masses for the different host-galaxy components of 63 objects; for 3 objects, no robust stellar masses could be determined. Thus, the  $M_{\rm BH}$ -stellar-mass relations presented in section [4.2](#page-14-2) include 63 objects.

Given  $\sigma$  within the effective radius as described in the previous section, we can also calculate a dynamical mass:

$$
M_{\rm sph, dyn} = c r_{\rm eff, sph} \sigma_{\rm ap, reff}^2 / G \tag{1}
$$

with  $c = 3$  for comparison with literature [\(Courteau et](#page-21-26) [al.](#page-21-26) [2014\)](#page-21-26). Since robust  $\sigma$  measurements within the effective spheroid radius were only obtained for 50 objects, the  $M_{\rm BH}$ - $M_{\rm sph,dyn}$  relation in section [4.2](#page-14-2) includes 50 objects.

#### 3.4. Comparison samples

<span id="page-14-3"></span>To compare the resulting scaling relations of  $M_{\rm BH}$  and  $\sigma$ , luminosity and mass with the literature, we use the sample presented by Kormendy  $\&$  Ho [\(2013\)](#page-22-0) as a quiescent galaxy comparison sample, 85 local galaxies with  $M_{\rm BH}$  based on dynamical modeling of spatially-resolved kinematics. Their sample consist of 44 elliptical galaxies, 20 spiral and S0 galaxies with classical bulge and 21 spiral and S0 galaxies with pseudo-bulge. Five of the elliptical galaxies are mergers in progress. Pseudo-bulges and mergers are significant outliers in [Kormendy & Ho](#page-22-0)  $(2013)$  and ignored here. For 11 objects, the  $M_{\rm BH}$  is considered uncertain and these objects are also ignored. We are thus left with 51 objects total, 32 elliptical galaxies and 19 spiral and S0 galaxies with classical bulges.

The stellar-velocity dispersions are adopted in most cases from Gültekin et al. [\(2009\)](#page-22-5) and represent effective velocity dispersions within  $r_{\text{eff}}/2$  as average of  $V^2(r) + \sigma^2(r)$  weighted by  $I(r)dr$ , thus consistent with the way we derived stellar-velocity dispersions, since the difference between averaging inside  $r_{\text{eff}}$  and  $r_{\text{eff}}/2$ is small [\(Kormendy & Ho](#page-22-0) [2013\)](#page-22-0).

Kormendy  $&$  Ho [\(2013\)](#page-22-0) list spheroid magnitudes in Ks and V, and (V-Ks) and (B-V) colors. We use a variety of elliptical and spiral spectral templates from [Bruzual](#page-21-27) [& Charlot](#page-21-27) [\(2003\)](#page-21-27) and [Kinney et al.](#page-22-46) [\(1996\)](#page-22-46) and derive a linear least-square fit of the form  $(V-I) = \alpha * (B-V) + \beta$ with  $\alpha = 0.72$  and  $\beta = 0.41$ , for conversion to I-band magnitudes.

The stellar spheroid masses given by [Kormendy & Ho](#page-22-0) [\(2013\)](#page-22-0) are derived from a mean of mass-to-light ratios

based on  $\sigma$  and  $(B-V)_0$  (their equations 8 and 9) and Kband magnitude. The mass-to-light ratio based on color is derived from [Into & Portinari](#page-22-54) [\(2013\)](#page-22-54), who assume a Kroupa (2001) IMF, but [Kormendy & Ho](#page-22-0) [\(2013\)](#page-22-0) adjust to the dynamical zeropoint.

For the  $M_{\text{BH}}$ - $\sigma$  relation, we also show the 29 RM AGN sample presented by [Woo et al.](#page-23-0) [\(2015\)](#page-23-0). We adjust their  $M_{\text{BH}}$  to match the virial factor of  $\log f = 0.97$  adopted here.

#### 4. RESULTS AND DISCUSSION

#### 4.1. Host-galaxy morphology

<span id="page-14-1"></span><span id="page-14-0"></span>The host-galaxy morphology was determined based on visual inspection of images and the results of the surfacebrightness fitting (Section [3.1\)](#page-7-2). Of the full sample of 66 AGNs with HST images, we conclude that 3 are hosted by bona-fide elliptical galaxies and 63 by spiral or S0 galaxies. Out of the latter, 26 galaxies are found to have a bar. In order for a spheroidal component to be classified as pseudo-bulge, we conservatively require that at least three of the following four criteria are met (fol-lowing [Kormendy & Ho](#page-22-0) [2013\)](#page-22-0): (i) Sérsic index  $\langle 2;$  (ii) bulge-to-total luminosity ratio  $< 0.5$ ; (iii) rotation dominated, i.e., ratio between maximum rotational velocity at effective spheroid radius and central stellar-velocity dispersion  $> 1$ ; (iv) for face-on galaxies, the presence of a bar is considered an indicator for the existence of a pseudo-bulge. In this way, of the 63 spiral or S0 galaxies, 21 spheroids are classified as pseudo-bulges, the majority of which (19) are in barred spiral galaxies. Table [1](#page-3-1) gives the host galaxy classification for all objects. Four objects show signs of interaction and/or merger activity (0206−0017, 0904+5536, 1708+2153, 2254+0046). The distribution of host-galaxy morphologies is typical for Seyfert-type AGNs. Given the wide range of host-galaxy morphologies (with mass-to-light ratios ranging between 0.7 and 2.2) and the high quality of the imaging, our sample is ideal to study dependency of the  $M_{\rm BH}$ -hostgalaxy scaling relations with other parameters such as (pseudo) bulges and bars.

#### 4.2. Scaling relations

<span id="page-14-2"></span>We present scaling relations between  $M_{\rm BH}$  and stellar velocity dispersion  $\sigma$  (within effective radius of spheroid), dynamical spheroid mass, stellar mass, and I-band luminosity (Figure [5\)](#page-18-0). We choose [Kormendy &](#page-22-0) [Ho](#page-22-0) [\(2013\)](#page-22-0) for a consistent comparison for all these different scaling relations, even though there are more recent studies with a compilation of larger samples. However, in a review by [Greene et al.](#page-21-28) [\(2020\)](#page-21-28), the authors note that their results on the  $M_{\text{BH}}\text{-}\sigma$  relation would not have changed if, instead of using a recent literature compi-

lation, they had used exclusively the [Kormendy & Ho](#page-22-0) [\(2013\)](#page-22-0) sample. For the  $M_{\text{BH}}$ - $\sigma$  relation, we also show 29 RM AGNs measured in a similar way [\(Woo et al.](#page-23-0) [2015,](#page-23-0) see Section [3.4](#page-14-3) for details).

Following common practices, we fit the different scaling relations as the linear relation with  $\alpha$  and  $\beta$  as slope and intercept values (Table [3\)](#page-17-0). The error bars of the measurements (in both x and y) are taken into account to perform the inference; in this step we adopt the EMCEE package (Python based) to derive the bestfit value and the uncertainties of  $\alpha$  and  $\beta$ . The intrinsic scatter is estimated so that when the squares of the observed uncertainties are summed up, the best-fit reduced chi-square value is close to unity. For all of the  $M_{\text{BH}}$  scaling relations, our sample of 66 local AGNs naturally extends the correlations for quiescent galaxies down to  $M_{\rm BH} \sim 10^7$  M<sub>⊙</sub> along the same line, with the same slope and normalization. However, by itself, the dynamic range in  $M_{BH}$  covered by our sample is too small to determine the slope. Thus, when deriving fits to the different scaling relations, we either fit both samples (AGNs and quiescent galaxies) together or when fitting our AGN sample alone, we fix the slope to that of [Kormendy & Ho](#page-22-0) [\(2013\)](#page-22-0).

The  $M_{\text{BH}}$ - $\sigma$  relation of our local AGN sample with  $M_{\rm BH}$  determined using the single-epoch method and  $\sigma$ based on spatially-resolved kinematics agrees well with that of AGNs with  $M_{BH}$  obtained from reverberation mapping [\(Park et al.](#page-22-26) [2012;](#page-22-26) [Woo et al.](#page-23-0) [2015\)](#page-23-0). The importance of the RM AGN sample cannot be overstated since it serves as the  $M_{\rm BH}$  calibrator beyond the local Universe. Given that slope and scatter of the  $M_{\rm BH}$ - $\sigma$  relation of our local AGNs, selected based solely on broad  $H\beta$  emission line width, agree with that of RM AGNs provides an independent confirmation that the selection of the RM AGN sample based on variability (not on well defined galaxy/black-hole mass properties) does not introduce biases. Moreover, the close agreement between both samples provides an indirect validation of the single-epoch method for the estimation of  $M_{\text{BH}}$ . Note that these conclusions are independent of the fact that the  $M_{\rm BH}$ - $\sigma$  scaling relation of RM AGNs is matched to that of inactive galaxies, since that only affects the normalization, but not slope and scatter.

To illustrate the effect of un-identified bars, we also include scaling relations for stellar mass and luminosity with spheroid+bar component added. This may help in comparison with literature data, especially given the difficulties and potential ambiguities involved in decomposition of images with poor data quality. Since the spheroids in the majority of barred spiral galaxies in our sample are classified as pseudo-bulges (19/21), this affects the location of pseudo-bulges the most. It moves the pseudo-bulges further to the right in the  $M_{\rm BH}$ -stellar mass and  $M_{\rm BH}$ -luminosity relations which tends to move them into better agreement with the scaling relations of quiescent galaxies. However, within the uncertainty, the difference is small. For none of the scaling relations do we find a significant difference between pseudo- and classical bulges in terms of correlations with  $M_{\rm BH}$ . This is in line with some studies (e.g., [Davis et al.](#page-21-10) [2018\)](#page-21-10), but contrary to many others (e.g., [Hu](#page-22-55) [2008;](#page-22-55) [Greene et al.](#page-21-29) [2010;](#page-21-29) [Sani et al.](#page-22-56) [2011;](#page-22-56) [L¨asker et al.](#page-22-7) [2016;](#page-22-7) [Saglia et al.](#page-22-49) [2016;](#page-22-49) [Menci et al.](#page-22-57) [2016;](#page-22-57) [de Nicola et al.](#page-21-24) [2019\)](#page-21-24). For example, [Kormendy & Ho](#page-22-0) [\(2013\)](#page-22-0) went so far to conclude that "any  $M_{\rm BH}$  correlations with the properties of diskgrown pseudobulges [...] are weak enough to imply no close coevolution" (see also, [Kormendy et al.](#page-22-6) [2011\)](#page-22-6).

Pseudo-bulges, considered to have evolved secularly through dissipative processes rather than through galaxy mergers (e.g., [Courteau et al.](#page-21-16) [1996;](#page-21-16) [Kormendy &](#page-22-37) [Kennicutt](#page-22-37) [2004\)](#page-22-37), play an important role for understanding the origin of the  $M_{\rm BH}$  scaling relations. If major mergers are the fundamental drivers of the  $M_{\rm BH}$  scaling relations, only classical bulges, centrally concentrated, mostly red and quiescent, merger-induced systems, should follow these tight correlations. On the basis of high-quality HST imaging, a careful analysis and a conservative classification of bulges as pseudo-bulges, our results clearly show that pseudo-bulges follow the same relations as classical bulges, confirming findings of an earlier study of ours based on SDSS images [\(Bennert](#page-21-8) [et al.](#page-21-8) [2011a\)](#page-21-8). This rules out hierarchical assembly as the sole origin of the  $M_{\rm BH}$ -host-galaxy scaling relations [\(Peng](#page-22-2) [2007;](#page-22-2) [Hirschmann et al.](#page-22-3) [2010;](#page-22-3) [Jahnke & Mac](#page-22-4)[cio](#page-22-4) [2011;](#page-22-4) [Kormendy & Ho](#page-22-0) [2013\)](#page-22-0), a conclusion reached by others, based on different arguments (e.g., Anglés-Alcázar et al. [2013;](#page-21-30) [Graham](#page-21-0) [2016\)](#page-21-0).

In fact, our study shows that there are no significant outliers that could be attributed to any specific category, whether it be galaxies with pseudo-bulges, bars or signs of interactions/mergers. For example, the four objects with signs of mergers/interaction do not tend to lie off the relations. Likewise, barred galaxies (26 out of 63 disk galaxies in our sample) do not form outliers, in line with some literature [\(Beifiori et al.](#page-21-9) [2012;](#page-21-9) [Sahu et al.](#page-22-8) [2019\)](#page-22-8). The location of barred galaxies on the scaling relations is not only important since over half of the disk galaxy population is barred (e.g., [Weinzirl et al.](#page-22-58) [2009\)](#page-22-58), but also because of the relevance of bars in secular evolution and potentially fueling of BHs. Moreover, it is much easier to identify a bar than a pseudo-bulge (for a discussion, see [Graham](#page-21-0) [2016\)](#page-21-0), a reason why some studies choose to distinguish between barred and non-barred galaxies

rather than classical vs. pseudo-bulges (e.g., [Graham](#page-21-31) [& Scott](#page-21-31) [2013\)](#page-21-31). Most previous literature studies found barred galaxies to lie off the  $M_{\rm BH}$ -scaling relations, in particular in the case of  $M_{\text{BH}}-\sigma$  (e.g., [Graham](#page-21-32) [2008;](#page-21-32) [Graham & Li](#page-21-33) [2009;](#page-21-33) [Hu](#page-22-55) [2008,](#page-22-55) for conflicting results, see also [Beifiori et al.](#page-21-9) [2012\)](#page-21-9). This is not surprising, given that the stellar dynamics in galactic sub-structures such as bars and pseudo-bulges is very different from that of elliptical galaxies or classical bulges. Moreover,  $\sigma$  measurements can depend significantly on, e.g., size of the fiber (as is the case, for example, for SDSS), orientation of the slit (in case of long-slit observations) and aperture size used (for details and comparisons, see [Bennert et al.](#page-21-17) [2015\)](#page-21-17). Integral-field spectroscopy and spatially-resolved spectroscopy is an obvious step forward and has been obtained for a sub-sample of the RM AGN sample [\(Batiste](#page-21-34) [et al.](#page-21-34) [2017\)](#page-21-34). While our  $\sigma$  measurements were obtained using long-slit spectroscopy, we mitigate these effects by using spatially-resolved measurements integrated within the spheroid effective radius which is a robust way to determine  $\sigma$  (see also [Bennert et al.](#page-21-17) [2015\)](#page-21-17).

The majority of AGNs in our sample reside in host galaxies of S0 or late-type morphology (63/66), out of which almost half of the galaxies are barred and a third of spheroids are classified as pseudo-bulges. The fact that all of them are obeying the same tight  $M_{\rm BH}$ -scaling relations, highlights the importance of secular evolution for the growth of BHs and bulges. Secular evolution may have a synchronizing effect, growing BHs and bulges simultaneously at a small but steady rate for late-type galaxies, and keeping them on tight relations over time. Comparison with semi-analytical models for galaxy formation including secular evolution (such as e.g., [Menci](#page-22-57) [et al.](#page-22-57) [2016,](#page-22-57) who, however, find little or no correlation of pseudo-bulge mass with  $M_{\rm BH}$ ) can further shed light on such a scenario, but is beyond the scope of this paper.

Interestingly, we do not find significant differences in the tightness of the different correlations. The scatter in the relations ranges between 0.2-0.4 dex, smaller or equal to that of quiescent galaxies (Gültekin et al. [2009;](#page-22-5) [McConnell & Ma](#page-22-59) [2013;](#page-22-59) [Kormendy & Ho](#page-22-0) [2013\)](#page-22-0). This is contrary to most previous studies that have often concluded that  $M_{\text{BH}}-\sigma$  is the tightest (0.3 dex in log  $M_{\text{BH}}$ ) and thus the most fundamental of the scaling relations [\(Tremaine et al.](#page-22-48) [2002;](#page-22-48) [Beifiori et al.](#page-21-9) [2012;](#page-21-9) [Saglia](#page-22-49) [et al.](#page-22-49) [2016;](#page-22-49) [Shankar et al.](#page-22-50) [2016;](#page-22-50) [van den Bosch](#page-22-51) [2016;](#page-22-51) [de](#page-21-24) [Nicola et al.](#page-21-24) [2019\)](#page-21-24), at least for late-type spiral galaxies (Gültekin et al. [2009;](#page-22-5) [Greene et al.](#page-21-29) [2010;](#page-21-29) Läsker et al. [2016,](#page-22-7) see, however, [Davis et al.](#page-21-35) [2019\)](#page-21-35). We attribute this difference to (i) our homogeneous sample selection, (ii) high-quality data, for both imaging and spectroscopy, and (iii) reliable surface photometry for a detailed structural decomposition of the host galaxy components and spatially-resolved kinematics. Given the fact that the biggest uncertainty in host-galaxy surface-brightness fitting is the classification and identification of individual structures, a combination of (ii) and (iii) is essential if one wants to determine the role of host-galaxy substructures on the correlation with  $M_{\rm BH}$ .

#### 5. SUMMARY

<span id="page-16-0"></span>This paper presents a study of 66 local  $(0.02 \leq z \leq$ 0.1) active galactic nuclei (AGNs) homogeneously selected based on the presence of a broad  $H\beta$  emission line in SDSS spectra. High-quality HST optical (66 objects) and Gemini NIR imaging (14 of 66 objects) are complemented by spatially-resolved kinematics from spectra obtained at the Keck Telescopes.  $M_{\rm BH}$  is determined based on the single-epoch method with broad  $H\beta$  emission-line width measured from Keck spectra. Surface photometry is performed using state of the art methods, providing a structural decomposition of the AGN host galaxies into spheroid, disk and bar (when present), with the spheroid component conservatively being classified as either classical or pseudo-bulge. Scaling relations between  $M_{\text{BH}}$ -and host galaxy properties — spatially-resolved stellar-velocity dispersion, dynamical spheroid mass, stellar spheroid mass and spheroid luminosity — are presented, in comparison with quiescent galaxies and RM AGNs taken from the literature. Our findings can be summarized as follows:

- 1. The majority of AGNs (63/66) are hosted by galaxies classified as spiral or S0 with a high fraction of bars  $(26/63)$  and pseudo-bulges  $(21/63)$ , typical for Seyfert-type galaxies. The wide variety of host-galaxy morphologies makes our sample ideally suited to study the dependency of the  $M_{\text{BH}}$ -host-galaxy scaling relations with other parameters such as (pseudo) bulges and bars.
- 2. Tight correlations are found between  $M_{\rm BH}$  and spatially-resolved stellar-velocity dispersion, dynamical spheroid mass, stellar spheroid mass and spheroid luminosity, without significant differences in the scatter. This is contrary to the widely accepted paradigm that the  $M_{\rm BH}$ - $\sigma$  relation is the most fundamental of all scaling relations.
- 3. The intrinsic scatter of 0.2-0.4 dex is smaller than or comparable to that of quiescent galaxies, showing that spiral galaxies hosting AGNs follow the same tight  $M_{\rm BH}$ -scaling relations, contrary to many literature studies.
- 4. We do not find any particular outliers: objects with bars, pseudo-bulges or signs of merger activity all fall within the intrinsic scatter of the relations. Our results rule out hierarchical assembly as the sole origin of the  $M_{\rm BH}$ -host-galaxy scaling relations and highlight the importance of secular evolution for growing both  $M_{\rm BH}$ and spheroid.
- <span id="page-17-0"></span>5. The  $M_{\rm BH}$ - $\sigma$  relation of our AGNs is indistinguishable from the relation of AGNs with  $M_{\rm BH}$  obtained through reverberation mapping. This indirectly validates single-epoch virial estimators of  $M_{\rm BH}$  and is consistent with no significant selection bias for RM AGNs.

Our results show that all the tight correlations can be simultaneously satisfied by AGN hosts in the  $10^{7}$ - $10^{9}$  $M_{\odot}$  regime if data of sufficient quality are in hand and great care is taken when deriving host-galaxy properties. A simple explanation of the difference between our uniformly tight relations and the larger scatter found in the literature is that  $\sigma$  is generally measured more accurately than the other host galaxy parameters. The sample presented in this paper is meant to serve as a local reference point for studies of the cosmic evolution of the correlations between host galaxy properties and  $M_{\rm BH}$ .

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Facilities: HST (WFC3), Keck:I (LRIS), Gemini:Gillett (NIRI), Sloan



<span id="page-18-0"></span>Figure 5. M<sub>BH</sub> scaling relations. In all panels, black data points correspond to the local quiescent comparison sample from [Kormendy & Ho](#page-22-0) [\(2013\)](#page-22-0), only including elliptical galaxies and spiral galaxies with classical bulge. For our sample, pseudo-bulges are shown in blue and classical bulges in red. To reduce confusion of data points, error bars on  $M_{BH}$  for our sample are omitted and shown instead in the bottom right corner. Top left panel:  $M_{\rm BH}$ - $\sigma$  relation. Cyan data points show 29 RM AGNs from [Woo et al.](#page-23-0) [\(2015\)](#page-23-0). Top right panel:  $M_{\text{BH}}-M_{\text{sph},dyn}$  relation. Middle left panel:  $M_{\text{BH}}-M_{\text{sph}}$  relation. Middle right panel:  $M_{\text{BH}}-M_{\text{sph+bar}}$  relation. Bottom left panel:  $M_{\text{BH}}-L_{\text{sph,I}}$  relation. Bottom right panel:  $M_{\text{BH}}-L_{\text{sph+bar,I}}$  relation.

$X$ in relation	Sample	$\alpha$	β	Scatter	Reference
(1)	(2)	(3)	(4)	(5)	(6)
$\sigma/200 \rm km\,s^{-1}$	AGNs $(50)$ & Quiescent galaxies $(51)$	$8.52 \pm 0.56$	$4.00 \pm 0.25$	$0.35 \pm 0.04$	this paper
$\sigma/200 \rm km\,s^{-1}$	Quiescent galaxies $(51)$	$8.49 \pm 0.035$	$4.38 \pm 0.29$	$0.29 \pm 0.03$	KH13
$\sigma/200 \rm km\,s^{-1}$	AGNs(50)	$8.49 \pm 0.08$	$4.38$ (fixed KH13)	$0.40 \pm 0.09$	this paper
$M_{\rm sph, dyn}/10^{11}M_{\odot}$	AGNs $(50)$ & Quiescent galaxies $(52)$	$8.76 \pm 0.69$	$0.97 \pm 0.06$	$0.37 \pm 0.04$	this paper
$M_{\rm sph, dyn}/10^{11}M_{\odot}$	AGNs(50)	$8.87 \pm 0.07$	$1.16$ (fixed KH13)	$0.21 \pm 0.11$	this paper
$M_{\rm sph}/10^{11}M_{\odot}$	Quiescent galaxies $(52)$	$8.69 \pm 0.06$	$1.16 \pm 0.08$	0.29	KH13
$M_{\rm sph, stellar}/10^{11} M_{\odot}$	AGNs $(63)$ & Quiescent galaxies $(52)$	$8.72 \pm 0.7$	$0.97 \pm 0.07$	$0.39 \pm 0.04$	this paper
$M_{\rm sph+bar, stellar}/10^{11}M_{\odot}$	AGNs (sph+bar) $(63)$ & Quiescent galaxies $(52)$	$8.71 \pm 0.74$	$1.05 \pm 0.07$	$0.38 \pm 0.04$	this paper
$M_{\rm sph, stellar}/10^{11}M_{\odot}$	AGNs(63)	$8.78 \pm 0.07$	$1.16$ (fixed KH13)	$0.33 \pm 0.08$	this paper
$L_{\rm sph,I}/10^{11}L_{\odot}$	AGNs $(66)$ & Quiescent galaxies $(51)$	$9.06 \pm 0.73$	$1.03 \pm 0.07$	$0.39 \pm 0.04$	this paper
$L_{\rm sph,I}/10^{11}L_{\odot}$	AGNs (sph+bar) (66) & Quiescent galaxies (51)	$9.05 \pm 0.76$	$1.11 \pm 0.07$	$0.40 \pm 0.04$	this paper
$L_{\rm sph,I}/10^{11}L_{\odot}$	Quiescent galaxies (51)	$9.11 \pm 1.15$	$1.02 \pm 0.11$	$0.47 \pm 0.06$	this paper
$L_{\rm sph,I}/10^{11}L_{\odot}$	AGNs(66)	$8.88 \pm 0.05$	$1.02$ (fixed KH13)	$0.08 + 0.07$	this paper

Table 3. Fits to the Local Scaling Relations.

NOTE— The relations plotted as dashed lines in Fig. [5](#page-18-0) correspond to the ones given in sample "AGN & Quiescent galaxies." Col. (1): Scaling relation of the form  $\log(M_{\rm BH}/M_{\odot}) = \alpha + \beta \log X$  with X given in the table. Col. (2): Sample used for fitting. In parantheses, number of galaxies are given. Col. (3): Mean and uncertainty of the best fit intercept. Col. (4): Mean and uncertainty of the best fit slope. Col. (5): Mean and uncertainty of the best fit intrinsic scatter. Col. (6): References for fit. [Kormendy & Ho](#page-22-0) [\(2013\)](#page-22-0) is listed as KH13.

	HST I-band					Gemini Ks-band	<b>HST</b>					
Object	<b>AGN</b>	Spheroid	Disk	Bar	<b>AGN</b>	Spheroid	Disk	Bar	$n_{\rm sph}$	$R_{\rm sph}$	$R_{\rm disk}$	$R_{\rm bar}$
	(mag)	(mag)	(mag)	(mag)	(mag)	(mag)	(mag)	(mag)		(11)	$($ $\prime\prime)$	(11)
(1)	(2)	(3)	(4)	(5)	(6)	(7)	(8)	(9)	(10)	(11)	(12)	(13)
0013-0951	18.3	18.0	15.9	$\cdots$	16.0	17.5	15.0	$\cdots$	1.2	0.36	5.19	$\cdots$
$0038 + 0034$	17.9	16.3	18.3	$\ddotsc$	.	$\ldots$	$\ldots$	.	4.5	2.39	2.39	$\cdots$
$0109 + 0059$	18.9	18.3	16.9	18.2	17.8	17.1	15.9	17.2	1.6	0.2	3.2	1.11
0121-0102	16.7	18.4	14.8	16.6	15.1	17.0	14.1	15.6	1.0	0.42	5.35	2.71
$0150 + 0057$	19.1	17.6	15.5	17.2	17.0	16.8	14.6	16.2	1.0	0.39	3.92	1.3
0206-0017	19.0	14.1	14.2	$\ddotsc$	16.0	13.1	13.5	$\ddotsc$	3.8	3.27	9.42	$\ddotsc$
$0212 + 1406$	18.4	16.8	15.6	18.4	.	.	$\ldots$	.	1.0	0.75	4.22	0.37
$0301 + 0110$	18.1	16.9	17.7	$\cdots$	$\cdots$	$\cdots$	$\cdots$	$\cdots$	4.0	1.2	3.35	$\cdots$
$0301 + 0115$	18.0	18.1	16.8	18.1	15.9	16.4	15.7	17.3	1.0	0.23	2.95	1.46
0336-0706	20.7	17.2	16.4	$\ddotsc$	.	.	$\cdots$	.	1.0	0.77	5.4	$\cdots$
0353-0623	18.4	17.3	16.9	18.4	$\ldots$	.	$\cdots$	.	1.0	0.94	4.75	0.27
0737+4244	19.7	17.5	16.6	$\ddotsc$	$\ldots$	.	$\cdots$	.	2.7	0.11	2.85	$\cdots$
$0802 + 3104$	17.2	17.4	15.6	17.7	$\cdots$	.	$\cdots$	.	1.1	0.28	3.27	1.01
$0811 + 1739$	19.1	17.7	16.2	17.9	$\cdots$	.	$\cdots$	.	1.4	0.38	4.48	2.88
$0813 + 4608$	20.1	16.6	16.6	17.0	17.8	15.7	15.7	16.1	2.6	0.67	5.14	3.58
$0845 + 3409$	20.1	16.9	15.9	18.1	17.4	15.7	15.2	17.2	3.9	1.09	7.0	1.66
$0857 + 0528$	18.0	17.7	15.6	$\ddotsc$	.		$\cdots$	.	1.2	0.38	3.56	$\cdots$
$0904 + 5536$	16.7	16.4	16.6	$\ddotsc$	.	.	$\cdots$	.	1.8	1.35	7.42	$\cdots$
$0909 + 1330$	20.6	17.7	15.4	17.3	.	.	$\cdots$		1.5	0.64	7.61	4.61

Table 4. Surface-Photometry Fitting Results.

Table 4 continued

 ${\small \bf B}{\small \bf E}{\small \bf N}{\small \bf N}{\small \bf E}{\small \bf T}{\small \bf F}{\small \bf T}{\small \bf A}{\small \bf L}.$ 

Table 4 (continued)

		HST <i>I</i> -band			Gemini Ks-band				<b>HST</b>			
Object	AGN	Spheroid	Disk	Bar	$\rm{AGN}$	Spheroid	Disk	Bar	$n_{\rm sph}$	$R_{\rm sph}$	$R_{\rm disk}$	$R_{\rm bar}$
	(mag)	(mag)	(mag)	(mag)	(mag)	(mag)	(mag)	(mag)		(n)	(n)	(11)
(1)	(2)	(3)	(4)	(5)	(6)	(7)	(8)	(9)	(10)	(11)	(12)	(13)
$0921 + 1017$	18.7	$15.7\,$	14.7	$\cdots$		$\ldots$	$\ldots$	$\cdots$	$5.0\,$	3.88	3.97	$\ddotsc$
$0923 + 2254$	16.4	15.5	14.2	16.4	.	$\ldots$	$\ldots$	$\ldots$	1.4	$\rm 0.95$	10.9	6.4
$0923 + 2946$	21.1	15.5	$\cdots$	$\ldots$	.		$\ldots$		4.8	2.37	$\cdots$	$\ldots$
$0927 + 2301$	17.7	14.8	$13.5\,$	$\ddotsc$	.		.	$\ddotsc$	1.4	1.24	7.46	$\cdots$
$0932 + 0233$	18.3	16.9	16.2	$\ldots$	.		.	.	$1.1\,$	0.68	4.37	$\ldots$
$0936 + 1014$	16.9	$18.0\,$	15.2	$\ldots$	.		.		1.4	$0.65\,$	6.15	$\cdots$
$1029 + 1408$	18.3	$15.4\,$	16.6	$\ldots$		$\ldots$			$4.3\,$	2.42	5.87	$\ldots$
$1029 + 2728$	19.8	$16.2\,$	16.5	$\cdots$	.	$\cdots$	$\ldots$		$3.1\,$	0.78	3.21	$\cdots$
$1029 + 4019$	17.4	17.7	16.2	$\cdots$		$\ldots$	$\ldots$		1.1	0.42	$2.5\,$	$\ldots$
$1042 + 0414$	18.8	17.6	16.5	18.1		$\ldots$			1.6	0.25	3.21	2.02
$1043 + 1105$	16.9	16.8	$\cdots$	$\ldots$	.		$\ldots$		$3.1\,$	$2.3\,$	$\ldots$	$\ldots$
$1058 + 5259$	18.3	16.9	16.7	17.3	.	$\ldots$	$\ldots$	$\cdots$	$1.9\,$	0.63	5.94	2.58
$1101 + 1102$	19.1	15.4	15.7	$\ldots$	.	$\ldots$	$\ldots$	.	$5.0\,$	3.0	4.59	$\cdots$
1104+4334	20.6	15.8	18.7	18.6	.	$\ldots$		.	4.5	2.79	2.79	1.46
$1137 + 4826$	22.9	17.5	17.5	$\ldots$		$\ldots$			2.6	0.24	1.29	$\ldots$
$1143 + 5941$	18.7	17.5	16.5	17.5	.	$\ldots$	.		$2.4\,$	$0.5\,$	6.72	4.26
$1144 + 3653$	16.3	15.8	15.0	$\ldots$		$\ldots$			1.4	$1.2\,$	7.87	$\ldots$
$1145 + 5547$	19.3	$18.2\,$	15.3	18.0	.		$\ldots$		1.0	0.52	7.0	2.1
1147+0902	16.6	15.9	18.2	$\ldots$		$\ldots$	$\ldots$		3.8	2.01	2.01	$\ldots$
$1205 + 4959$	17.2	16.6	15.7	$\cdots$		$\ldots$			$2.0\,$	0.68	4.41	$\cdots$
1206+4244	17.4	17.2	15.5	17.1	.		.	$\cdots$	1.0	0.64	7.0	2.57
$1216 + 5049$	18.4	15.0	14.9	$\cdots$	.	$\ldots$	$\ldots$		$3.3\,$	2.75	8.03	$\cdots$
$1223 + 0240$	16.6	15.3	15.0	$\ldots$	.	$\cdots$	$\ldots$	.	$5.0\,$	3.51	4.3	$\cdots$
$1246 + 5134$	19.1	17.9	17.1	$\ldots$	.	$\ldots$	.		$2.0\,$	0.4	2.88	$\ldots$
$1306 + 4552$	$20.9\,$	$18.3\,$	15.6	17.1	.		.		$1.0\,$	$\rm 0.3$	$5.05\,$	2.88
1307+0952	19.1	17.4	15.3	19.4				$\cdots$	1.0	0.56	4.59	0.27
$1312 + 2628$	17.5	17.9	15.5	17.9		$\ldots$			1.0	0.45	6.03	2.7
1405-0259	18.8	17.4	15.6		.	$\ldots$	$\ldots$	$\cdots$	1.0	0.77	5.01	
$1416 + 0137$	18.5	15.9	15.2	$\ldots$	.		$\ldots$		2.4	1.57	6.75	$\ldots$
$1419 + 0754$	17.6	16.1	14.5	$\cdots$		$\ldots$		$\ldots$	$1.9\,$	0.98	6.45	$\ldots$
		16.3	15.0	$\ldots$		$\ldots$		.	$1.5\,$	0.88	6.29	$\ldots$
1434+4839	17.7			16.4			$\cdots$	$\cdots$				3.98
$1545 + 1709$	18.7	$16.2\,$	16.9	$\cdots$		$\ldots$		$\cdots$	$5.0\,$	1.49	2.42	
1557+0830	17.9	16.8	$\ldots$	$\ldots$	.	$\ddots$			$2.5\,$	1.16	$\ldots$	$\cdots$
$1605 + 3305$	17.9	$16.9\,$	16.5	$\cdots$	.	$\cdots$		$\ldots$	1.2	0.68	3.82	$\cdots$
$1606 + 3324$	18.6	$16.3\,$	$16.5\,$	$\cdots$	.	$\ldots$			$3.3\,$	1.16	4.29	$\cdots$
$1611 + 5211$	18.8	$15.8\,$	$16.2\,$	$\ldots$		$\cdots$	$\cdots$	$\ldots$	3.1	0.75	4.62	$\ldots$
$1636 + 4202$	18.4	$15.8\,$	17.1	18.2		$\ldots$	$\cdots$	$\ldots$	5.0	3.0	3.0	0.33
$1708 + 2153$	16.6	$16.5\,$	16.3	$\ldots$		$\cdots$	$\ldots$	$\ldots$	1.9	1.54	6.95	$\cdots$
2116+1102	18.1	17.6	$16.0\,$	19.0		$\ldots$	$\ldots$	$\ldots$	$1.2\,$	$0.5\,$	5.95	1.96
$2140 + 0025$	$16.5\,$	17.3	$16.8\,$	$\ldots$	15.0	17.0	15.7	$\ldots$	$1.0\,$	0.47	2.24	$\ldots$
2215-0036	17.0	18.1	16.7	$\ldots$	$\ldots$	$\ldots$	$\ldots$	$\ldots$	$1.0\,$	$0.5\,$	3.45	$\ldots$
2221-0906	17.6	17.6	17.3	$\ldots$	16.6	17.6	15.9	$\ddotsc$	$2.7\,$	0.86	2.98	$\ldots$
2222-0819	17.4	$18.7\,$	15.6	17.8	15.0	25.9	14.7	16.6	$1.1\,$	0.33	3.98	1.48
$2233 + 1312$	17.9	$17.9\,$	$15.9\,$	17.5	16.2	16.4	14.8	16.2	1.0	0.38	7.71	1.82

Table 4 continued

	HST <i>I</i> -band					HST						
Object	AGN	Spheroid	Disk	Bar	<b>AGN</b>	Spheroid	Disk	Bar	$n_{\rm sph}$	$R_{\rm sph}$	$R_{\rm disk}$	$R_{\rm bar}$
	(mag)	(mag)	$(\text{mag})$	(mag)	$(\text{mag})$	(mag)	(mag)	(mag)		$\left( \frac{\eta}{\eta} \right)$	$\left( n\right)$	(11)
(1)	(2)	(3)	(4)	(5)	(6)	(7)	(8)	(9)	10)	(11)	(12)	(13)
$2254 + 0046$	17.0	17.4	17.6	$\cdots$	$\cdots$	$\cdots$	$\cdots$	$\cdots$	1.1	0.26	1.91	$\cdots$
$2327 + 1524$	17.9	14.7	14.8	$\cdots$	$15.1\,$	13.5	14.1	$\cdots$	2.4	1.71	7.86	$\cdots$
$2351 + 1552$	17.7	17.0	16.9	$\cdots$	16.5	16.1	15.8	$\cdots$	2.0	0.77	2.69	$\cdots$

Table 4 (continued)

Note—Surface-photometry fitting results using lenstronomy on HST and Gemini images. Col. (1): Target ID used throughout the text (based on R.A. and declination). Col. (2): Point-source (AGN) magnitude in HST I-band (uncertainty 0.1 mag). Col. (3): Spheroid magnitude in HST I-band (uncertainty 0.1 mag). Col. (4): Disk magnitude in HST I-band (if present; uncertainty 0.1 mag). Col. (5): Bar magnitude in HST I-band (if present; uncertainty 0.1 mag). Col. (6): Point-source (AGN) magnitude in Gemini Ks-band (uncertainty 0.2 mag). Col. (7): Spheroid magnitude in Gemini Ks-band (uncertainty 0.2 mag). Col. (8): Disk magnitude in Gemini Ks-band (if present; uncertainty 0.2 mag). Col. (9): Bar magnitude in Gemini Ks-band (if present; uncertainty 0.2 mag). Col. (10): Spheroid Sérsic index  $n$  (5% uncertainty). Col. (11): Spheroid radius in arcseconds (10%) uncertainty). Col. (12): Disk radius in arcseconds (10% uncertainty). Col. (13): Bar radius in arcseconds (10% uncertainty).

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