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Sample-to-Sample Variation of the Volume Fraction of Superconducting Material in  $\text{YBa}_{2}\text{Cu}_{3}\text{O}_{(7-\delta)}$

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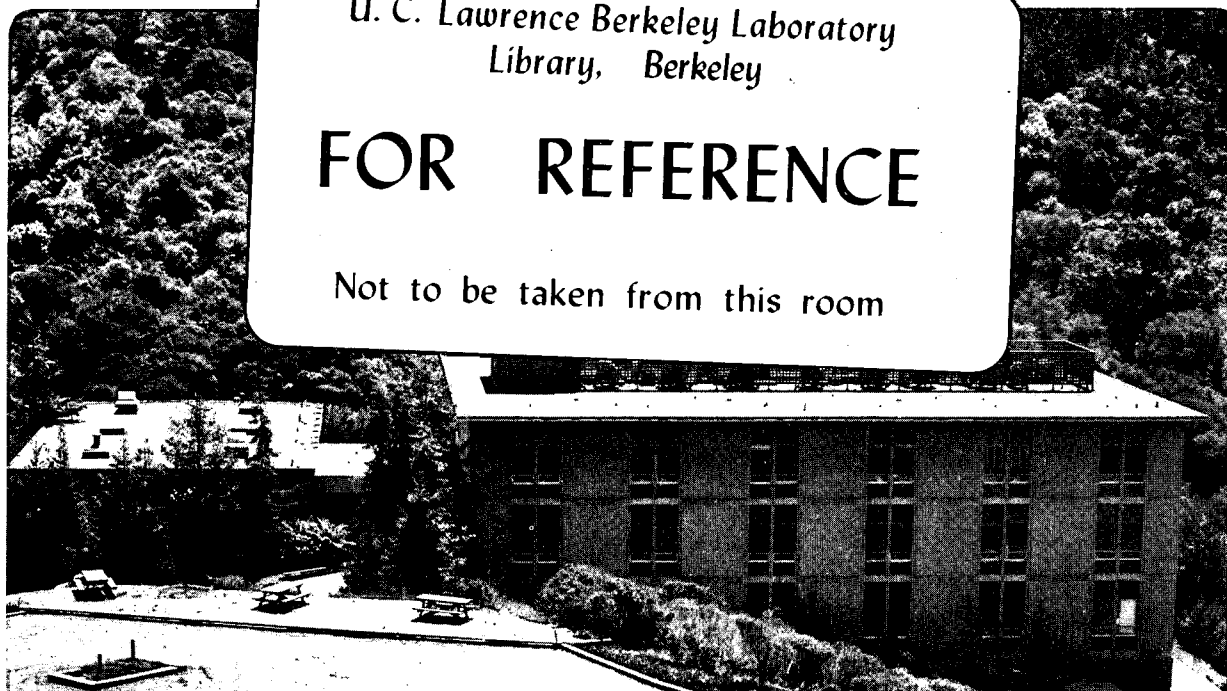
## Materials & Chemical Sciences Division

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### Sample-To-Sample Variation of the Volume Fractions of Superconducting Material in $\text{YBa}_2\text{Cu}_3\text{O}_{7-\delta}$

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December 1991



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**SAMPLE-TO-SAMPLE VARIATION OF THE  
VOLUME FRACTION OF  
SUPERCONDUCTING MATERIAL IN  $\text{YBa}_2\text{Cu}_3\text{O}_{7.6}$**

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**Presented at the MRS Meeting in Boston, MA  
December 2-7, 1991**

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SUMMARY

Specific heat ( $C$ ) measurements on high- $T_c$  superconductors show dramatic sample-to-sample variations. Three quantities derived from  $C$  for YBCO that show such a sample dependence are mutually proportional and suggest that the sample dependence reflects a corresponding variation in the volume fraction of superconductivity ( $f_s$ ). They are (1)  $d\gamma^*/dH$ , the field dependence of the coefficient of the "linear term," which is associated with the creation of vortices in the superconducting regions (see Fig.1); (2)  $\Delta C(T_c)$  the mean-field discontinuity in  $C$  at  $T_c$  (see Fig. 2); (3)  $\Delta S$ , an entropy change in the superconducting regions produced by the application of a magnetic field (see Fig. 3). These quantities have been used to determine relative values of  $f_s$  (see Fig. 4). Their correlation with a concentration of  $\text{Cu}^{2+}$  magnetic moments ( $n_2$ ) (see Fig. 1) suggests that the transition to the superconducting state is suppressed either by these moments themselves or by some other defect that is correlated with these moments. The limit  $n_2=0$  is taken to correspond to  $f_s=1$ , thus establishing absolute values of  $f_s$  (see Fig. 5). The conclusion is that substantial fractions of these materials, even of samples believed to be "good" superconducting material, can be non superconducting.

Fig. 1. Analysis of the low-temperature  $C$  for a typical YBCO sample into its four components: (1)  $C_l$  = lattice specific heat. (2)  $C_e(H) \equiv \gamma^*(H)T$ , the field-dependent linear contribution, one term of which is  $Hd\gamma^*/dH$ . (3)  $C_m(H)$ , associated with  $\text{Cu}^{2+}$  moments. For  $H=0$  it appears here as the well-known "upturn" in  $C/T$ ; for  $H=7T$  it is a Schottky anomaly that determines the concentration of  $\text{Cu}^{2+}$  moments ( $n_2$ ). The points for  $C_m(7T)$  represent experimental data from which the other contributions have been subtracted. (4)  $C_h(H)$ , the hyperfine component.

Fig. 2. The determination of  $\Delta C(T_c)$  for a sample that shows no obvious fluctuation effects. The simple entropy-conserving construction, represented by the dashed lines, gives essentially the same value of  $\Delta C(T_c)$  as a fit with a gaussian distribution of  $T_c$ , represented by the solid curve. The dash-dot curve, a harmonic-lattice extrapolation from high temperatures, represents the background used for the gaussian fit.

Fig. 3.  $\Delta S$ , a measure of the effect of a 7-T magnetic field on the entropy associated with the anomaly at  $T_c$ . Thermodynamically,  $\Delta S$  is approximately related to  $d\gamma^*/dH$  by  $\Delta\gamma^* = Hd\gamma^*/dH - \Delta S/T_c$ , but experimentally it is an independent measure of  $f_s$ .

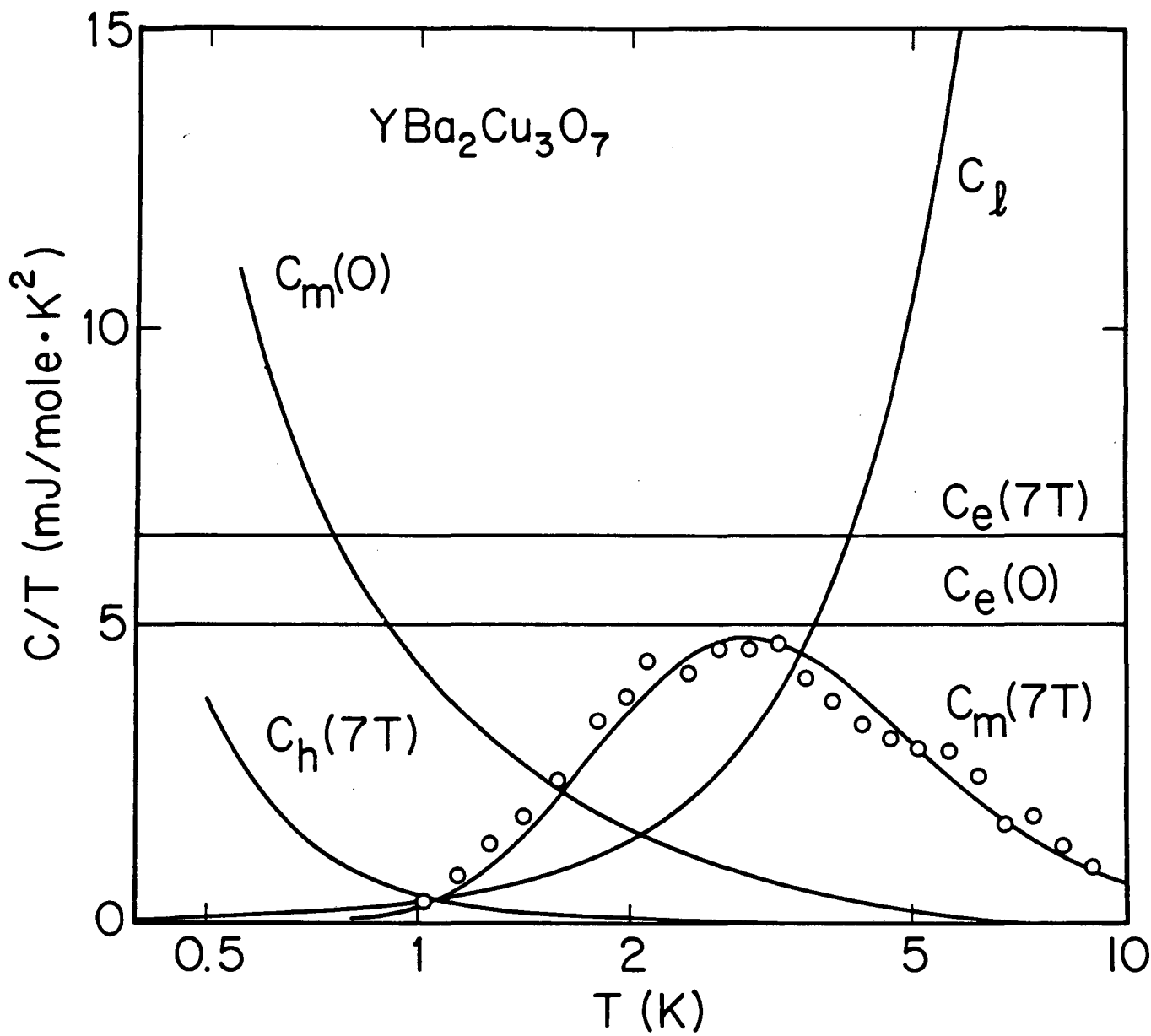
Fig. 4. The correlations among  $d\gamma^*/dH$ ,  $\Delta C(T_c)$ , and  $\Delta S$ , scaled to fall on the same line for a number of polycrystalline YBCO samples. For each sample the three quantities are plotted at the relative value of  $f_s$  that best represents all three.

Fig. 5. The correlation of the relative values of  $f_s$  with  $n_2$ , which suggests that  $n_2$  measures the concentration of a defect that suppresses the transition to the superconducting state. The association  $n_2=0$  with  $f_s=1$  puts the values of  $f_s$  on an absolute scale.

## POSSIBLE EFFECT OF OXYGEN STOICHIOMETRY

Oxygen deficiency,  $\delta > 0$ , reduces  $\Delta C(T_c)$  and the density of electron states, and, presumably,  $d\gamma^*/dH$  and  $\Delta S$ . However, it is unlikely that variations in oxygen stoichiometry play a significant role in the effects represented in Fig. 3 because:

- (1) The Pauli susceptibility for these samples is essentially constant, ruling out the variation in electron density of states that would be required to explain the results.
- (2) For these samples the 0-K Debye temperature and  $\gamma^*(0)$  vary strongly with  $f_s$ , but other work [Y. Nakazawa and M. Ishihara, *Physica C* **458**, 381 (1989)] shows that they are nearly independent of  $\delta$  in the relevant range of  $\delta$ .
- (3) Similar correlations have been observed for  $(\text{La}_{1.85}\text{Sr}_{0.15})\text{CuO}_4$  for which there is no reason to expect variations in oxygen stoichiometry.



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FIG. 1

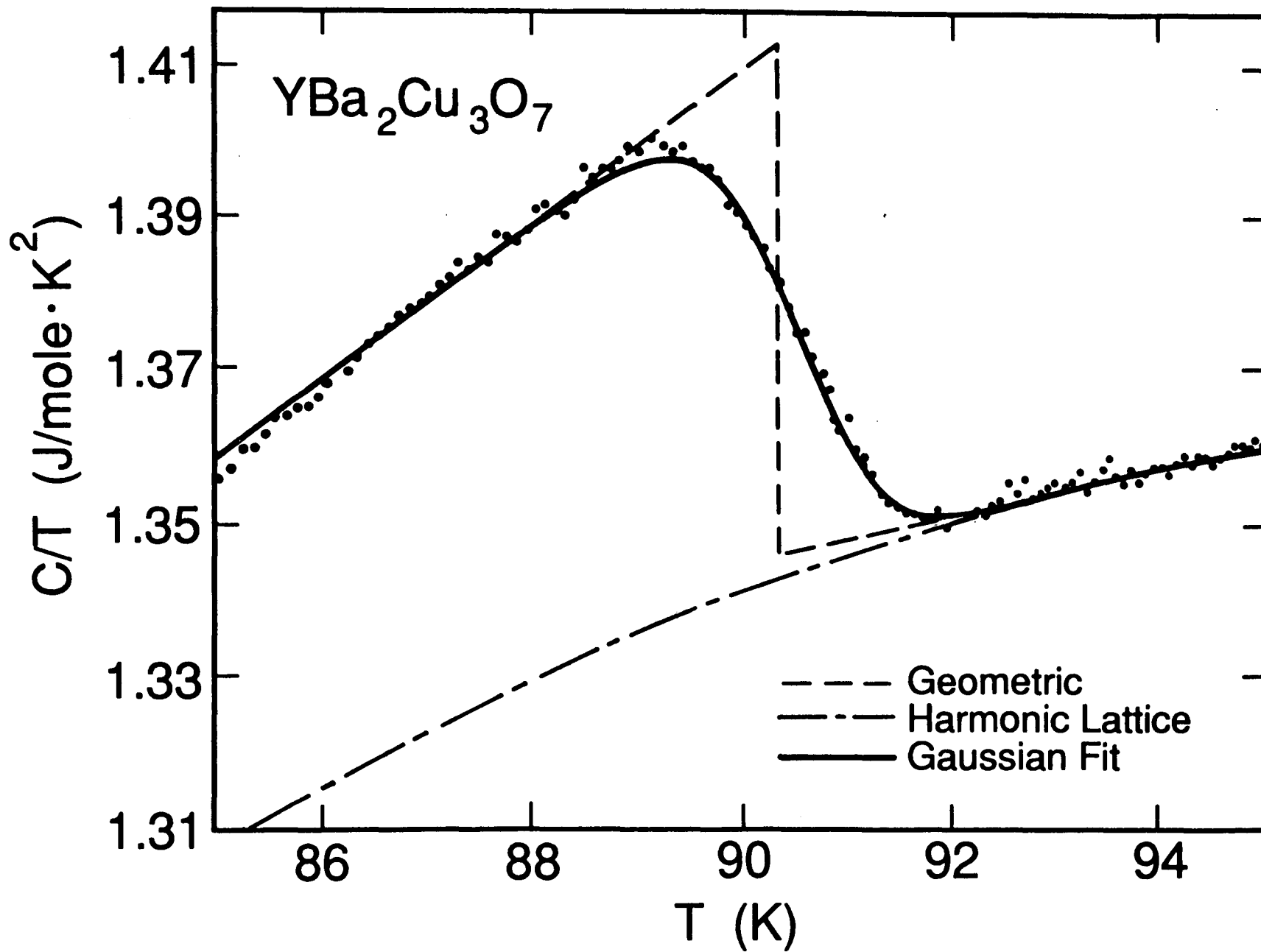
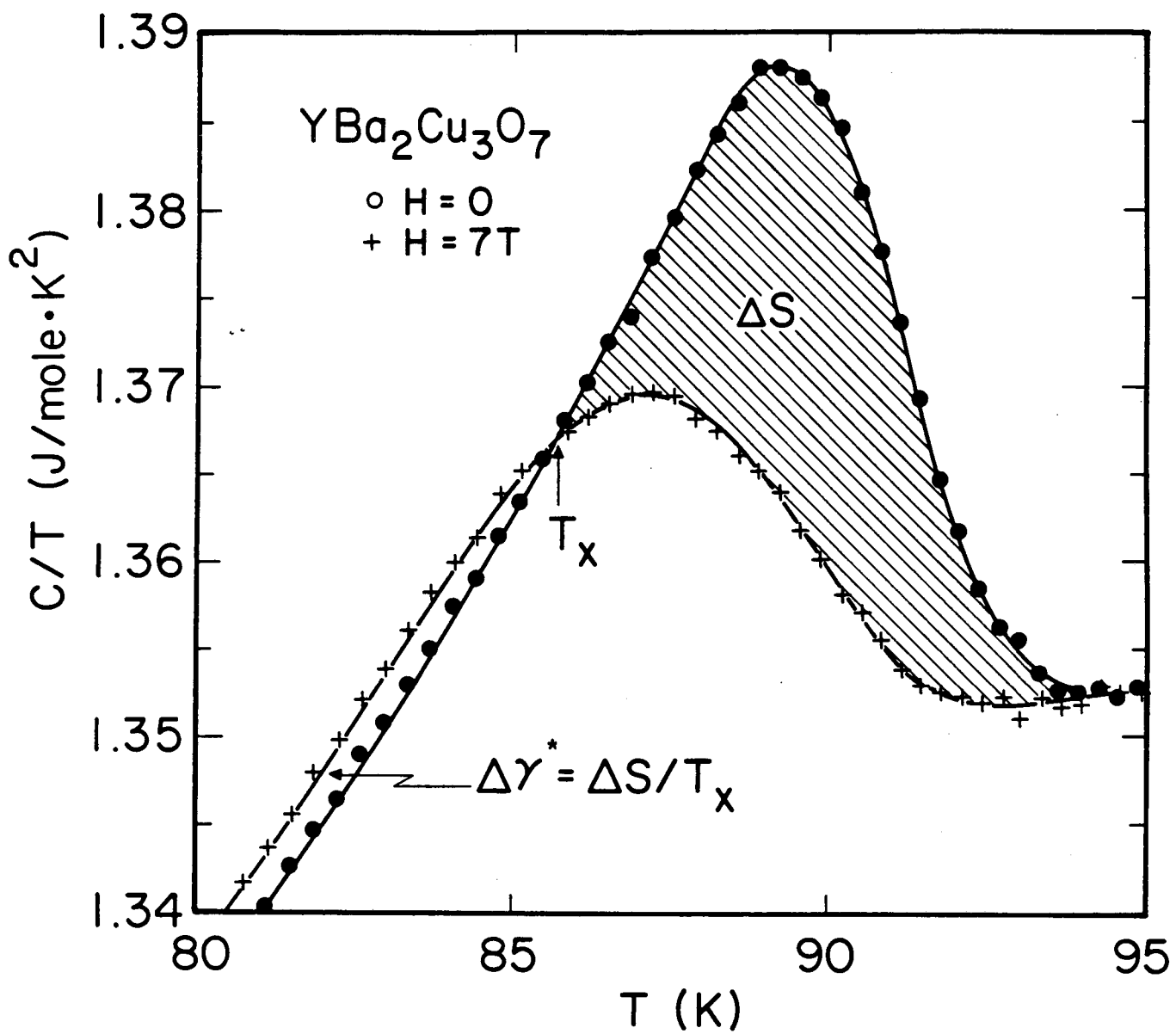


FIG. 2

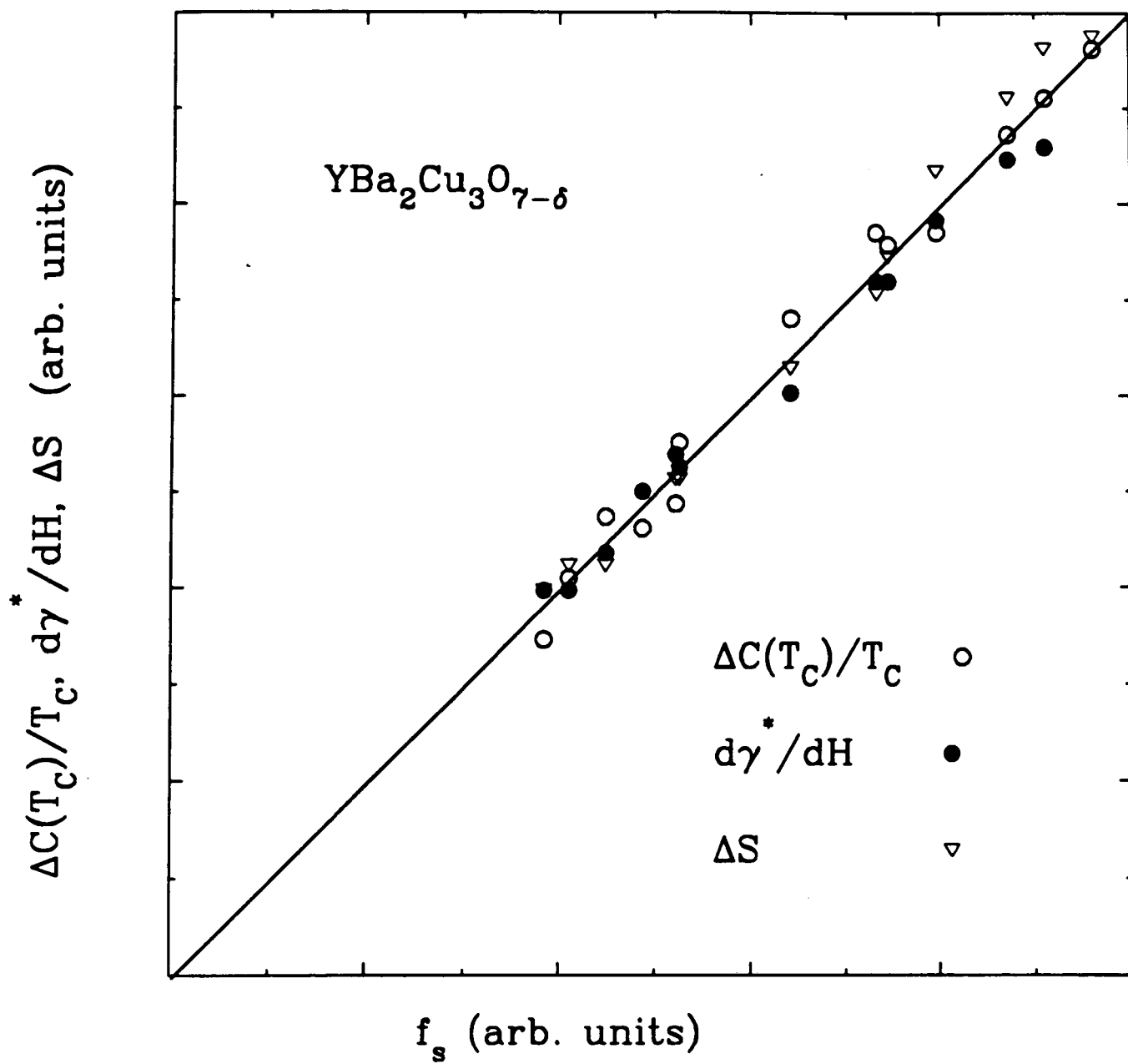
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FIG. 3



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FIG. 4

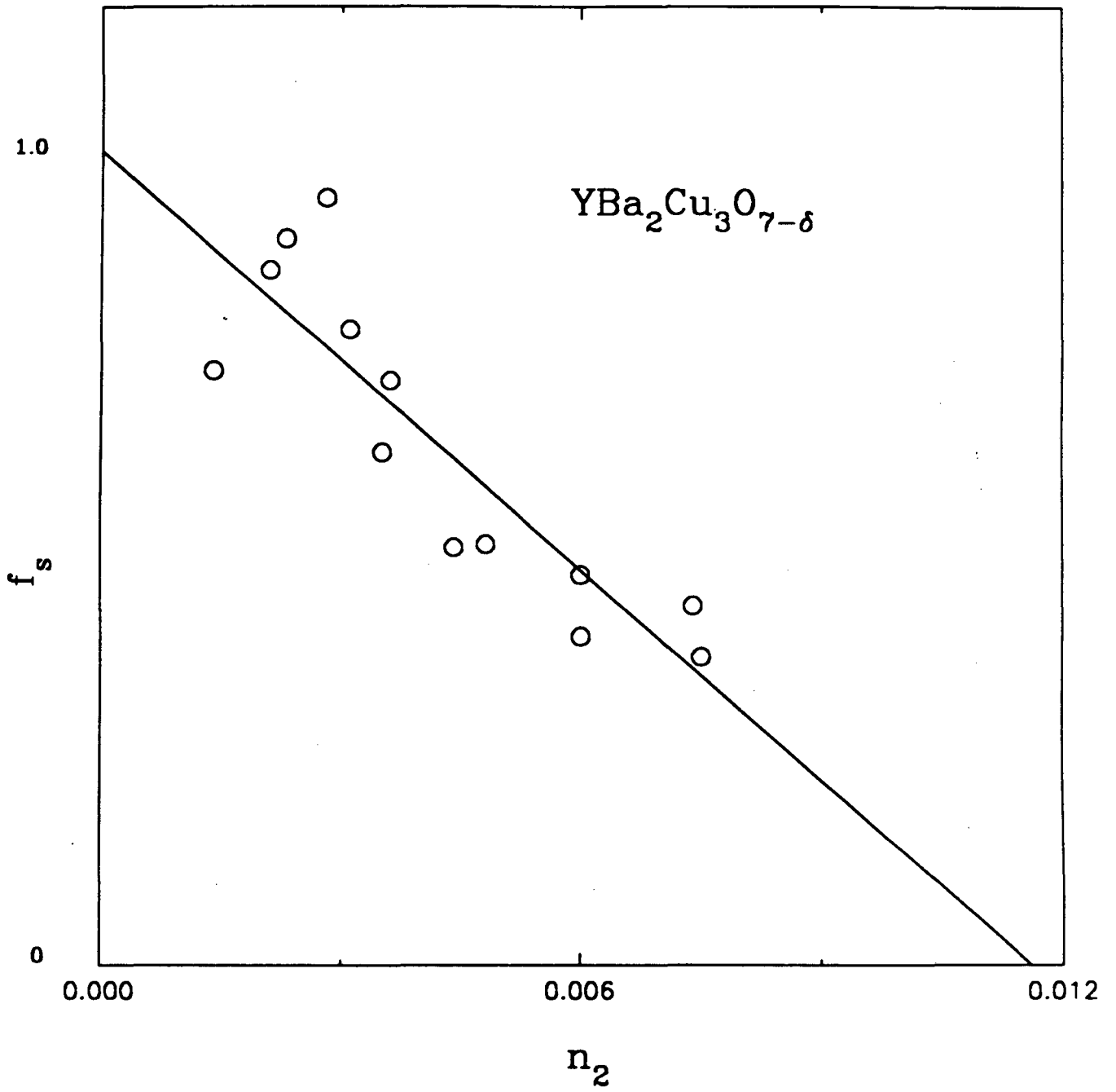


FIG. 5

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