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# Crossover from Landau Fermi liquid to non-Fermi liquid behavior: Indications from Hall measurements on CeCoIn<sub>5</sub>

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## Abstract

We conducted Hall effect measurements on the heavy-fermion superconductor CeCoIn<sub>5</sub> for temperatures 0.05–5 K and for pressures up to 1.2 GPa. A scaling of the magnetic field  $H$  is introduced for the differential Hall coefficient,  $R_H^d = \partial\rho_{xy}(T, H)/\partial H$  resulting in a single generic curve for  $R_H^d(H)$  curves obtained at different  $T$ . We argue that the peak feature apparent in this generic curve corresponds to the crossover from non-Fermi liquid to Landau Fermi liquid behavior.

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The compound CeCoIn<sub>5</sub> is particularly suited to investigate the interplay of quantum criticality and unconventional superconductivity (SC) in which the pairing might be mediated by magnetic fluctuations. It exhibits the highest superconducting critical temperature,  $T_c$ , among the Ce-based ambient pressure superconductors [1], a magnetic field tuned quantum critical point (QCP) may exist [2] close to the upper critical field of SC,  $H_{c2}$ , and SC may hide an antiferromagnetic (AFM) order.

Hall effect measurements are a well established tool to shed light on the electronic properties of materials close to a QCP. Accordingly, such measurements have early been conducted for  $T \geq 1$  K [3], even for applied pressures  $p$  [4]. In our case, we want to concentrate on the low- $T$  region  $0.05$  K  $\leq T \leq 5$  K and  $p \leq 1.2$  GPa. At these  $T$  well

below the coherence temperature  $T^* \approx 40$  K no anomalous Hall contribution is found. However, interpretation of Hall effect in CeCoIn<sub>5</sub> is complicated since SC inhibits a determination of the initial Hall coefficient and multiple bands at the Fermi level contribute to the Hall signal with a field dependent cyclotron mass [5].

For Hall measurements, isothermal field sweeps were conducted on single crystalline CeCoIn<sub>5</sub> samples with  $H \parallel c$ . Measurements under pressure were carried out in a piston cylinder type pressure cell. The evolution of the Hall resistivity  $\rho_{xy}$  (left) and its differential  $R_H^d = \partial\rho_{xy}/\partial H$  (right) for increasing  $p$  at  $T = 120$  mK is shown in Fig. 1. A changing slope of  $\rho_{xy}(H)$ , as obvious from the  $T = 120$  mK data, is observed for  $0.1 \leq T \leq 0.3$  K at  $p = 0$  and 0.3 GPa resulting in a minimum of  $|R_H^d|$  (arrow). This feature is suppressed with increasing  $p$  and can no longer be resolved at 1.2 GPa.

For further analysis, the  $H$ -values of the ambient pressure isothermal  $R_H^d(T, H)$  vs.  $H$  curves were scaled by  $H_{\min}^d$ . Here,  $H_{\min}^d$  denotes the field value at which  $|R_H^d|$  assumes its minimum for  $70 \leq T \leq 200$  mK. As seen in

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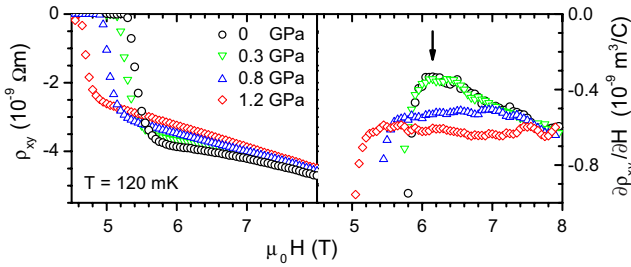


Fig. 1. Evolution of the isothermally measured Hall resistivity  $\rho_{xy}$  (left) and differential Hall coefficient  $\partial\rho_{xy}/\partial H$  (right) of CeCoIn<sub>5</sub> with pressure at 120 mK. The arrow indicates the “peak feature” referred to in the text.

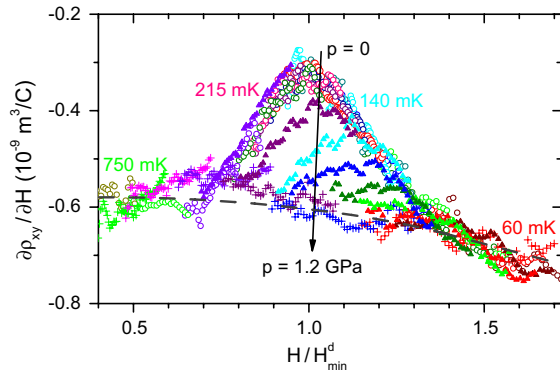


Fig. 2. Summarized differential Hall coefficient for  $p = 0$  (○), 0.8 GPa (▲) and 1.2 GPa (+) illustrating the increasing suppression of the Hall “peak feature” (arrow in Fig. 1). Magnetic fields are scaled with respect to  $H_{\min}^d$  and different temperatures are presented by different colors.

Fig. 2(○), all scaled  $R_H^d$  curves collapse onto a single, generic curve. For  $T$  below/above the given range, the  $H$ -values were scaled such that this generic curve is further completed towards larger/smaller values  $H/H_{\min}^d$ . Here,  $\mu_{\text{eff}} \propto 1/H_{\min}^d$  can be viewed as effective mobility, averaged over all Fermi surfaces (FS) contributing.

The ambient pressure  $R_H^d$  data between  $0.7 \leq H/H_{\min}^d \leq 1.3$  mark a distinct “peak feature”, whereas smaller and larger scaled fields appear to form an underlying “base line” of weak  $H$ -dependence (dashed line, Fig. 2). This “peak feature” is likely related with the AFM spin fluctuations (SF), based on the following observations:

- (i)  $p$  dependence: Applying pressure to CeCoIn<sub>5</sub> drives the system towards a heavy Landau Fermi liquid (LFL) state [6] by gradually suppressing the AFM SF [7]. This is likely related to the increase of  $T_c$  with  $p$  for small  $p$ , leading to maximum  $T_c$  at  $p^* \sim 1.3$  GPa. Our  $R_H^d$  data at  $p = 1.2$  GPa (+,

Fig. 2), i.e. slightly below  $p^*$ , approach the base line, with minor deviations at  $H/H_{\min}^d \sim 0.7$ . For intermediate  $p = 0.8$  GPa, the  $R_H^d$  values appear to be reduced for lower fields only. Note that for  $H$  scaling at  $p > 0$  the  $H_{\min}^d$  values obtained at  $p = 0$  were used.

- (ii)  $R_H^d$  values: At low  $H$  ( $< 0.7H_{\min}^d$ ) we obtain  $-R_H^d \approx 6 \times 10^{-10} \text{ m}^3/\text{C}$  with a slight  $H$  dependence ( $7 \times 10^{-10} \text{ m}^3/\text{C}$  at  $1.5H_{\min}^d$ ). This value agrees remarkably well with the one reported [4] for the non-magnetic analogue LaCoIn<sub>5</sub> ( $-5.5 \times 10^{-10} \text{ m}^3/\text{C}$ ). Generally, pressure drives Ce from a  $4f^1$  towards a non-magnetic  $4f^0$  configuration. Moreover,  $R_H^d$  of the Ce- and the La-based compound agree well at  $\mu_0 H = 7$  T, i.e. in the LFL regime.
- (iii) The  $T$  dependence of  $H_{\min}^d$  as obtained from the scaling (Fig. 2) tracks the crossover [2] from non-Fermi liquid to LFL behavior (not shown).

The “peak feature” might be related to AFM SF or to the opening of an AFM gap at the FS if a spin density wave is formed. The latter may also cause a discontinuity in  $R_H^d$  [8]. However, pressure suppresses the “peak feature” while changing the FS only little [5]. Note that Hall measurements (unlike thermodynamic ones) are sensitive to even weak fluctuations. Hence, the anisotropic AFM SF might be considered as a precursor of a gap opening.

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