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# Yet More Evidence That Myelin Protons Can Be Directly Imaged With UTE Sequences on a Clinical 3T Scanner: Bicomponent $T_2^*$ Analysis of Native and Deuterated Ovine Brain Specimens

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**Purpose:** UTE sequences with a minimal nominal TE of  $8\,\mu s$  have shown promise for direct imaging of myelin protons (T<sub>2</sub>, <1 ms). However, there is still debate about the efficiency of 2D slice-selective UTE sequences in exciting myelin protons because the half excitation pulses used in these sequences have a relatively long duration (e.g., 0.3–0.6 ms). Here, we compared UTE and inversion-recovery (IR) UTE sequences used with either hard or half excitation pulses (durations  $32\,\mu s$  or  $472\,\mu s$ , respectively) for imaging myelin in native and deuterated ovine brain at 3T.

**Methods:** Freshly frozen ovine brains were dissected into  $\sim 2 \text{ mm-thick}$  pure white matter and  $\sim 3$  to 8 mm-thick cerebral hemisphere specimens, which were imaged before and/or after different immersion time in deuterium oxide.

**Results:** Bicomponent  $T_2^*$  analysis of UTE signals obtained with hard excitation pulses detected an ultrashort  $T_2$  component (STC) fraction ( $f_S$ ) of 0% to 10% in native specimens, and up to ~86% in heavily deuterated specimens.  $f_S$  values were significantly affected by the TIs used in IR-UTE sequences with either hard or half excitation pulses in native specimens but not in heavily deuterated specimens. The STC  $T_2^*$  was in the range of 150 to 400  $\mu$ s in all UTE and IR-UTE measurements obtained with either hard or half excitation pulses.

**Conclusion:** Our results further support myelin protons as the major source of the ultrashort  $T_2^*$  signals seen on IR-UTE images and demonstrate the potential of IR-UTE sequences with half excitation pulses for directly imaging myelin using clinical scanners. **Magn Reson Med 80:538–547, 2017.** © **2017 International Society for Magnetic Resonance in Medicine.** 

Key words:  $T_2^*$ ; bicomponent; myelin; white matter; UTE; inversion recovery

### INTRODUCTION

Conventional imaging of brain tissue commonly detects three distinct pools: a long  $T_2$  pool associated with free

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water in cerebrospinal fluid (T<sub>2</sub>,  $\sim$ 2 s), an intermediate  $T_2$  pool involving extra- and intracellular water ( $T_2$ ,  $\sim$ 100 ms), and a short T<sub>2</sub> pool composed of water trapped in myelin bilayers and other associated macromolecules (myelin water) (T2, 1-10 ms) (1,2). Studying free water signals is helpful for detecting increased permeability of the blood-brain barrier in stroke and other pathological conditions (3-5). Charactering extra- and intracellular water signals have been useful for probing not only cellular structural characteristics but also functional activation (6-9). Investigating myelin water indirectly provides insights about the integrity of myelin and myelin loss (10-12). These water protons can be characterized using currently available clinical MRI techniques with TEs of 1 ms and longer, and are referred to as long  $T_2$  components (LTCs) in this study.

Unlike most water protons, the semisolid highly constrained lipid protons trapped in myelin sheaths, named as myelin protons hereafter (13), have ultrashort  $T_2$  (< 1 ms) and are not directly detectable with most conventional MRI sequences (14). Direct imaging of myelin protons may improve the specificity of MRI for evaluation of neurological diseases characterized by demyelination and remyelination, such as multiple sclerosis (15,16), and may also be of value in monitoring therapeutic response. In recently developed 2D UTE sequences, nominal TEs as short as 8 µs can be achieved through half pulse excitation, variable rate selective excitation, radial ramp sampling, and fast transmit/receive switching (17,18), making it possible to directly image myelin protons. Furthermore, inversion-recovery (IR) UTE sequences can provide robust suppression of signals from LTCs using an adiabatic IR preparation pulse and therefore have the potential to directly and selectively image myelin with high contrast (19).

Although UTE imaging of myelin protons has been investigated by several groups (14,18,20–23), there is still debate about whether UTE sequences, especially when used with half excitation pulses (for which the pulse duration is on the order of hundreds of microseconds), can directly detect myelin proton signals that have a reported  $T_2^*$  of ~300 µs (22) or shorter (14,21). This study aimed to compare UTE and IR-UTE sequences used with either short hard or longer half excitation pulses for probing myelin proton signals at 3T, and meanwhile to probe the specificity and sensitivity of these sequences to myelin proton signals by comparing the ultrashort  $T_2$ 

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FIG. 1. Diagrams of the 2D IR-UTE pulse sequences with TE of 10  $\mu$ s used with (**A**) a short hard-excitation pulse (rectangular shape, duration 32  $\mu$ s, bandwidth = 8.2 kHz) followed by 2D radial ramp sampling, and (**B**) a half excitation pulse (duration 472  $\mu$ s, bandwidth = 2.7 kHz) followed by 2D radial ramp sampling. (**C**) Illustration of the contrast mechanism for imaging white-matter ultrashort T<sub>2</sub> components (WM<sub>s</sub>) using IR-UTE with the TI set for nulling of signals from the white-matter long T<sub>2</sub> components (WM<sub>L</sub>), with the TI termed as WM<sub>L</sub> TI<sub>null</sub>. In magnitude IR-UTE images, the signal in white matter quickly decays to zero with increasing TE, whereas the signal in gray matter is still high at TE = 2.2 ms because it is dominated by long T<sub>2</sub> components (GM<sub>L</sub>). (**D**) M<sub>Z</sub> of WM<sub>L</sub> and WM<sub>S</sub> plotted against different TIs at the time UTE acquisition starts. At TI<sub>c</sub>, signals from WM<sub>L</sub> are nulled. At TI<sub>a</sub>, which is much shorter than TI<sub>c</sub>, signal in the white matter is dominated by WM<sub>L</sub>. At TI<sub>b</sub>, which is only slightly shorter than TI<sub>c</sub>, signal from WM<sub>S</sub> is cancelled out by that from WM<sub>L</sub>. At TI<sub>d</sub> and TI<sub>e</sub>, signals from WM<sub>S</sub> and WM<sub>L</sub> coexist in the image but with different fractions. Note that the time parameters in the diagrams were not proportionally illustrated. IR, inversion recovery; GM, gray matter; WM, white matter.

component (STC) fraction ( $f_S$ ) and  $T_2^*$  ( $T_{2S}^*$ ) seen in fresh and deuterated ovine brain white matter (WM). Deuterons in deuterium oxide  $(D_2O)$  have a MR frequency  ${\sim}6.5$  times lower than that of protons and are not detectable with proton MRI (24). By immersing tissue specimens in highly purified D<sub>2</sub>O, it is possible to replace the majority, if not all, of the long  $T_2$  water protons with deuterons (14). The specimens were expected to have different proportions of <sup>1</sup>H MRI-visible LTCs after they were immersed in D<sub>2</sub>O for different durations, whereas myelin protons would survive the D<sub>2</sub>O exchange (14) and still be detectable with UTE and IR-UTE sequences. If myelin protons are the major source of the ultrashort T<sub>2</sub> signals and the sequences are sensitive to proportional changes of these signals,  $T_{2S}^*$ might not change; however,  $f_S$  would change significantly after the exchange.

#### METHODS

#### Pulse Sequences and Contrast Mechanisms

Figure 1 shows diagrams of the 2D IR-UTE pulse sequences and the associated contrast mechanisms. Either a short hard pulse (rectangular shape, duration  $32 \,\mu$ s, bandwidth = 8.2 kHz) or a half pulse (half-sinc shape, variable rate selective excitation-corrected, duration 472  $\mu$ s, bandwidth = 2.7 kHz) was used for signal excitation (Figs. 1A–1B). Each sequence contained an adiabatic IR preparation pulse (Silver-Hoult pulse, duration 8.64 ms, bandwidth = 1.4 kHz) and a minimal nominal TE of 10  $\mu$ s. The basis for contrast seen on the IR-UTE images for selective myelin proton imaging in WM is illustrated

in Figures 1C and 1D. The inversion time (TI) is chosen to null the signals from LTCs in WM ( $WM_L$ , and this TI is abbreviated as  $WM_L$   $TI_{null}$  is this paper). At the time the UTE acquisition starts (TE =  $10 \,\mu$ s), LTCs in gray matter (GM<sub>L</sub>) have nonzero negative magnetization (because  $GM_{I}$  has a longer  $T_1$  than  $WM_{I}$ , whereas STCs in WM and GM (i.e., WM<sub>S</sub> and GM<sub>S</sub>) have nonzero-positive magnetizations (Fig. 1C). Therefore, on the IR-UTE images obtained at  $WM_L$   $TI_{null}$  and  $TE = 10 \,\mu s$ , the WM signal comes from STCs, whereas the GM signal contains mixed contributions from both LTCs and STCs. When the selected TI is incrementally offset from  $WM_L$  TI<sub>null</sub>, WM signals in the image will be increasingly contaminated by LTCs, leading to inaccurate measurement of STC signals (presumably mainly from myelin protons) (Fig. 1D). WM<sub>L</sub> TI<sub>null</sub> can be estimated from IR-UTE images obtained with varying TIs at a later yet still short TE (e.g., TE = 2.2 ms). At this TE, the magnetization of STCs is zero or near zero due to fast signal decay, whereas that of LTCs is largely unchanged, allowing selective measurement of the signals from LTCs for the estimation of their  $T_1$  and  $TI_{null}$ .

Hard excitation pulses have higher power and shorter duration than half excitation pulses, and thus should be more efficient and reliable in exciting extremely short  $T_2$ protons (25). However, a half pulse has to be used for 2D slice-selective imaging. In this study, the slice-selective gradients were turned off in the half pulse excitation experiments based on the following three considerations: First, the specimens were carefully prepared into pure WM specimens (~2 mm thick) or thin coronal hemisphere slabs (~3–8 mm thick); thus, a region of interest (ROI)



FIG. 2. Representative magnitude images (acquired with hard excitation pulses) of one native whitematter short block. (A) Images obtained with varying TIs at TE = 2.2 ms. (B)  $T_1$  map calculated from the images in (A) showing a relatively homogeneous center and a rim with a slightly longer T<sub>1</sub>. The black box shows the region of interest used for quantitative  $T_2^*$  analyses. (C) Images obtained with different TEs w/o IR prep. (D-F) Images obtained with different TEs and varying TIs. There was an obvious signal intensity decrease with the increase of TE in the images at TIs = 270 ms and 275 ms but not at TI = 295 ms. IR, inversion recovery; w/o IR prep, without inversion recovery preparation.

could be defined in WM without contamination from neighboring GM when the specimen was excited without slice selection. Second, errors induced by eddy currents in regular half pulse excitation experiments could be eliminated by turning off the slice-selective gradients, thus providing more accurate estimations of  $T_{2S}^*$  and  $f_S$ , as well

as allowing a more direct comparison between the hard pulse and half pulse excitations. Third, regular 2D UTE and IR-UTE imaging using half excitation pulses would require two excitations per scan to achieve slice selection. The total scan time could therefore be halved by turning off the slice-selective gradients.



FIG. 3. Representative bicomponent  $T_2^*$  fitting curves of one native white-matter short block. (A) Measurement with hard excitation pulses and no inversion-recovery preparation. (**B–D**) Measurements with hard excitation pulses and varying TIs. All measurements provided different  $f_S$  but similar  $T_{2S}^*$  values.  $f_S$ , ultrashort  $T_2$  component fraction;  $T_{2S}^*$ , ultrashort  $T_2$  component  $T_2^*$ ; w/o IR prep, without inversion recovery preparation.

#### Specimen Preparation and Experimental Design

Five freshly frozen ovine brains were thawed and dissected into six  $\sim 2 \text{ mm-thick} \sim 2 \times 5 \text{ mm}^2$  pure WM short blocks, three  ${\sim}2\,\text{mm-thick}\,{\sim}4{\times}\,10\,\text{mm}^2$  pure WM long segments from three consecutive slabs of the same brain, one  $\sim$ 8 mmthick cerebral hemisphere slab, and three  $\sim$ 3 mm-thick cerebral hemisphere slabs. Two WM short blocks, all three WM long segments, and all three  $\sim$ 3 mm-thick cerebral hemisphere slabs were subject to D<sub>2</sub>O exchange. Each WM short block was immersed in 5 mL D<sub>2</sub>O (99.9%, Sigma-Aldrich, St. Louis, Missouri) for 8h with D<sub>2</sub>O refreshed four times (8 h five-pass D<sub>2</sub>O exchange). Each WM long segment was immersed in 5 mL D<sub>2</sub>O for 1 h, 2 h (two 1 h consecutive periods with  $D_2O$  refreshed once), or 27 h (1 h, 1 h, 2 h, and 23 h consecutive periods with D<sub>2</sub>O refreshed three times, that is, 27 h four-pass  $D_2O$  exchange), respectively. Each cerebral hemisphere slab was immersed in 10 mL D<sub>2</sub>O and subject to 27 h four-pass  $D_2O$  exchange.

All of the pure WM short blocks were imaged using a 5 mm-diameter transmit/receive solenoid coil to achieve high signal-to-noise ratio (SNR). All the pure WM long segments and cerebral hemisphere slabs were imaged with a 7.6-cm receive-only surface coil that can be used clinically. First, three native WM short blocks were imaged with UTE sequences and IR-UTE sequences with varying TIs using hard excitation pulses. Then one deuterated WM short block was imaged with the UTE sequence using hard excitation pulse; and one native and one deuterated WM short block were each subject to free induction decay (FID) acquisition to evaluate water signal loss after  $D_2O$  exchange. Next, the ~8 mm-thick cerebral hemisphere slab was imaged in the native condition with UTE and IR-UTE sequences using both hard and half excitation pulses. Lastly, the three WM long segments were imaged after progressive D<sub>2</sub>O exchange, and the three  $\sim$ 3 mm-thick cerebral hemisphere slabs were all imaged before and after 27 h D<sub>2</sub>O exchange with UTE and IR-UTE sequences using half excitation pulses.

#### Data Acquisition

Both single-slice 2D UTE and IR-UTE sequences were implemented on a 3T Signa TwinSpeed scanner (GE Healthcare Technologies, Milwaukee, Wisconsin), which had a maximum gradient strength of 40 mT/m and a maximum slew rate of 150 mT/m/ms. The UTE sequence was performed with TR = 1,000 ms and a series of TEs ranging from 10 µs to up to 30 ms. At each imaging time point, the IR-UTE sequence was first used with TR/TE = 1,000/2.2 ms and a series of TIs (50, 100, 300, 500, and 800 ms) to measure the  $WM_L T_1$  and determine  $WM_L TI_{null}$ . The same IR-UTE sequence was then used with TR = 1,000 ms,  $TI = WM_L TI_{null}$  and/or TIs of 5 to 130 ms longer than  $WM_L~TI_{\rm null},$  and a series of TEs ranging from  $10\,\mu s$  to 30ms to measure  $f_S$  and  $T^*_{2S}$  in WM. The acquisition matrix was  $96 \times 96$ , and the scan time was 97s per acquisition. T<sub>2</sub>-weighted multislice fast spin echo images were acquired with TR/TE = 3,000/40 ms, slice thickness = 1 mm, and no slice gap to guide the ROI definition in WM on the UTE and IR-UTE images. FID signals were acquired on an 11.7T small bore preclinical scanner (Bruker BioSpec 117/11 USR/R, Rheinstetten, Germany). Briefly, the specimen was



FIG. 4. Bicomponent  $T_2^*$  analysis of images acquired with hard excitation pulse and no inversion recovery preparation from a WM short block after a 27 h four-pass D<sub>2</sub>O exchange, showing a much higher  $f_S$  with similar  $T_{2S}^*$  compared with a native WM short block as shown in Figure 3A. Inserts are the images of the WM short block after 27 h D<sub>2</sub>O exchange, which showed fast signal decay with the increase in TE. WM, white matter; D<sub>2</sub>O, deuterium oxide;  $f_S$ , ultrashort T<sub>2</sub> component fraction;  $T_{2S}^*$ , ultrashort T<sub>2</sub> component T<sub>2</sub>\*.

placed inside a 5-mm NMR proton-free tube and examined using a custom-made proton-free coil with a block excitation pulse (pulse width =  $20 \,\mu$ s), TR =  $20 \,s$ , 4,096 points, and sweep bandwidth =  $50 \,\text{kHz}$ .

#### Data Analysis

The WM<sub>L</sub> T<sub>1</sub> and TI<sub>null</sub> were calculated offline in each manually defined ROI using a single-component threeparameter fitting model in MatLab (MathWorks, Natick, Massachusetts).  $T_{2S}^*$  and  $f_S$  were calculated using a previously published bicomponent fitting model in MatLab (MathWorks) (26):

$$S(TE) = A_S \times e^{-\frac{TE}{T2S^*}} + A_L \times e^{-\frac{TE}{T2L^*}} + noise, \quad [1]$$

where S(TE) is the UTE or IR-UTE MR signal;  $A_S$  and  $A_L$  are the signal amplitudes of the STCs and LTCs;  $T_{2S}^*$  is the STC  $T_2^*$ ; and  $T_{2L}^*$  is the LTC  $T_2^*$ .  $f_S$  was defined as  $A_S/(A_S + A_L)$ . Initially, all datasets were to bicomponent fitting. However, in some cases with exceptionally high fractions of LTCs or STCs, the two components were not clearly separated (i.e., both  $T_2^*$  values < 1 ms or both > 20 ms) using Equation [1]. Thus, to improve algorithmic stability, these cases were reanalyzed using a single-component fitting model:

$$S(TE) = S_0 \times e^{-TE/T2^*} + noise, \qquad [2]$$

where  $S_o$  is the signal intensity at the minimum TE. The key fitting parameters are provided in the online



FIG. 5.  $T_2^*$  fitting results from the ~8 mm-thick hemisphere slab imaged using both hard and half excitation pulse sequences (**A–C**) and the average results from three pure WM short blocks imaged using hard excitation pulse sequence (**D–F**). IR-UTE images were acquired with TI = WM<sub>L</sub> TI<sub>null</sub> (labeled as TI<sub>c</sub>, equivalent to that in Fig. 1D, on the *x*-axes of the figures) and 5 to 70 ms longer. The ultrashort T<sub>2</sub> component fraction ( $f_S$ ) was very low when measured with UTE (<10%), much higher when measured with IR-UTE at WM<sub>L</sub> TI<sub>null</sub>, and decreased as TI increased from WM<sub>L</sub> TI<sub>null</sub>. The UTE sequence used with half pulse excitation did not detect the ultrashort T<sub>2</sub> component in this native specimen.  $f_S$  was consistently lower when measured with half pulse excitation than with hard pulse excitation at various TIs greater than WM<sub>L</sub> TI<sub>null</sub> (**A** and **D**). The ultrashort T<sub>2</sub> component T<sup>\*</sup><sub>2</sub> ( $T_{2S}^*$ ) values were all in the range of 150 to 250 µs (**B** and **E**), and the long T<sub>2</sub> component T<sup>\*</sup><sub>2</sub> ( $T_{2L}^*$ ) changed significantly with the increase in TI when hard pulse excitation was used (**C** and **F**). At WM<sub>L</sub> TI<sub>null</sub>, the T<sup>\*</sup><sub>2</sub> signal followed a single-component decay (R<sup>2</sup> > 0.99) (F). These results demonstrated the efficiency of the half-pulse UTE and IR-UTE sequences in measuring  $T_{2S}^*$  and detecting  $f_S$  changes in the presence of different proportions of signals from long T<sub>2</sub> components. WM, white matter; WM<sub>L</sub> TI<sub>null</sub>, inversion time for nulling signals from white-matter long T<sub>2</sub> components.

Supporting Information. All images used for analysis and presentation are magnitude images. Each FID was subject to Fourier transform as well as automatic phase and baseline corrections in iNMR (http://www.inmr.net) and displayed as a spectrum.

### RESULTS

Figure 2 and Figure 3 show results from one native WM short block imaged using hard pulse excitation. Figure 2A shows obvious signal intensity variation in the magnitude images acquired with different TIs at TE = 2.2 ms, reflecting varying relative contributions from STCs and LTCs. The signal intensity was lowest at TI = 300 ms, suggesting significant suppression of the signals from LTCs (Fig. 2A). The  $T_1$  map (Fig. 2B) calculated from the images in Figure 2A showed a near uniform structure inside the specimen, with marginal regions showing slightly longer  $T_1s$  (the map scale was 400–460 ms), validating that the specimen was pure WM ( $T_1$  of GM matter would be in the range of  $\sim$ 800 to 1,000 ms when measured using the same IR-UTE sequences (27)), although its margins might have been contaminated or injured during handling. Figure 2C shows magnitude images acquired with different TEs and no IR preparation. Figures 2D through 2F show magnitude images acquired with different TEs and varying TIs. At TI = 270 ms (i.e., WM<sub>L</sub> TI<sub>null</sub>), the signal intensities dropped significantly with increasing TE (Fig. 2D). At TIs of 275 ms and 295

ms (Figs. 2E–2F), the signal decay against TE was slower but still faster than that on the images acquired without IR preparation (Fig. 2C). Figure 3 shows bicomponent analyses of the magnitude images acquired without and with IR preparation using different TIs (ROI was defined in the center of the specimen, as illustrated in Fig. 2B). The  $f_S$  was low (7.9%) in the measurement without IR preparation (Fig. 3A), slightly higher (32.9%) at TI=295 ms (Fig. 3D), and much higher (93.0%) at TI=270 ms in the measurements with IR preparation (Fig. 3B), whereas the  $T_{2S}^*$  was all in the range of 150 to 400 µs (Figs. 3A– 3D).

The Supporting Figure S1 (available online) shows enlarged spectra from one native (Supporting Fig. S1A) and one deuterated (Supporting Fig. S1B) WM short block. In the native specimen (Supporting Fig. S1A), the area under the water peak was much larger than that of the lipid peak and contributed  $\sim 88.5\%$  of the total FID signal. In the deuterated specimen (Supporting Fig. S1B), the water peak was extremely narrow and contributed only  $\sim$ 24.6% to the total FID signal. Figure 4 shows the bicomponent fitting results of the magnitude images obtained with hard excitation pulse and no IR preparation from another deuterated WM short block that was subject to an identical deuteration procedure, revealing a  $f_S$  of ~86%. This  $f_S$  was much higher than that in a native specimen ( $\sim$ 7.9%) (Fig. 3A). This difference in  $f_S$ between the native and the deuterated specimens detected with imaging was consistent with the water-



FIG. 6.  $T_2^*$  fitting curves of UTE and IR-UTE signals (all half excitation pulses) from three pure WM long segments that were subject to 1 h (left column, no. 1), 2 h (middle column, no. 2), and 27 h (right column, no. 3) exchange with D<sub>2</sub>O, respectively. (**A**) The UTE sequence detected an increase of  $f_S$  from ~9.5% to ~47.1% with D<sub>2</sub>O exchange time in specimens no. 1 to no. 3. (**B**) The IR-UTE images obtained at WM<sub>L</sub> TI<sub>null</sub> showed single-component  $T_2^*$  signal decay, which changed little with D<sub>2</sub>O exchange time. (**C**) The IR-UTE images obtained at 85 ms longer than WM<sub>L</sub> TI<sub>null</sub> (WM<sub>L</sub> TI<sub>null</sub> +85 ms in figure) showed bicomponent  $T_2^*$  signal decays in specimens after 1 h and 2 h D<sub>2</sub>O exchange, and single-component  $T_2^*$  signal decay in the specimen after 27 h D<sub>2</sub>O exchange. D<sub>2</sub>O, deuterium oxide;  $f_S$ , ultrashort T<sub>2</sub> component fraction; IR, inversion recovery;  $T_{2S}^*$ , ultrashort T<sub>2</sub> component  $T_2^*$ , WM, white matter.

fraction reduction after  $D_2O$  exchange revealed by the FID measurements. The magnitude images (inserts in Fig. 4) from the deuterated specimen showed much faster signal decay with increasing TE than those of the native specimen (Fig. 2C), supporting unexchangeable STCs (presumably mainly myelin protons) as the major source of the UTE signals.

The Supporting Figure S2 (available online) illustrates the procedure of ROI definition in the ~8 mm-thick native cerebral hemisphere slabs. Figure 5 summarizes the bicomponent  $T_2^*$  fitting results from the ~8 mm-thick native cerebral hemisphere specimen (Figs. 5A-5C) shown in Supporting Figure S2 imaged using both hard and half excitation pulses, as well as the average results from three native WM short-block specimens imaged using hard excitation pulses (Figs. 5D–5F).  $f_S$  and  $T_{2S}^*$ were identified in all specimens when measured with the UTE sequences using hard excitation pulses. However, UTE signals of the hemisphere specimen imaged using half excitation pulses showed single-component decay, and the  $T_2^*$  value was similar to  $T_{2L}^*$  in the bicomponent measurements with hard excitation pulses (Figs. 5A–5C). In the IR-UTE experiments,  $f_S$  decreased with TI when measured either with hard or half excitation pulses (Figs. 5A and 5D), and was consistently lower when measured with half pulse excitation than with hard pulse excitation at various TIs greater than WM<sub>L</sub> TI<sub>null</sub> (Fig. 5A). In all measurements,  $T_{2S}^*$  was in the range of 150 to 250 µs (Figs. 5B and 5E), and  $T_{2L}^*$  changed significantly with the increase in TI when hard pulse

	Native			Deuterated		
	UTE	IR-UTE (TI = 295 ms)	IR-UTE (TI = 365 ms)	UTE	IR-UTE (TI = 250 ms)	IR-UTE (TI = 380 ms)
T <sup>*</sup> <sub>2S</sub>						
1	n.a. <sup>a</sup>	$339\pm34\mu\text{s}$	$210\pm69\mu s$	$239\pm46\mu\text{s}$	$297\pm21\mu s$	$229\pm39\mu\text{s}$
2	n.a. <sup>a</sup>	$246\pm28\mu\text{s}$	$202\pm22\mu s$	$241\pm48\mu s$	$245\pm10\mu s$	$262\pm35\mu s$
3	n.a. <sup>a</sup>	$215\pm50\mu s$	$205\pm64\mu s$	$215\pm50\mu s$	$247\pm21\mu s$	$262 \pm 42 \mu s$
f <sub>S</sub>						
1	n.a. <sup>a</sup>	pprox100% <sup>b</sup>	$18.0\pm2.1\%$	$42.4 \pm 0.4$ %	pprox100% <sup>b</sup>	pprox100% <sup>b</sup>
2	n.a. <sup>a</sup>	≈100% <sup>b</sup>	$17.5 \pm 1.7\%$	43.1 ± 2.2 %	≈100% <sup>b</sup>	pprox100% <sup>b</sup>
3	n.a. <sup>a</sup>	pprox100% <sup>b</sup>	$16.0\pm1.7\%$	$39.1\pm1.3~\%$	pprox100% <sup>b</sup>	pprox100% <sup>b</sup>

Comparison of UTE and IR-UTE Results Obtained With Half Excitation Pulses in Three Cerebral Hemisphere Slabs Before and After These Specimens Were Subjected to 27 H Four-Pass D<sub>2</sub>O Exchange

<sup>a</sup>Signals on the UTE images showed single component decay with long  $T_2^*$  (22 ± 2 ms), so the  $T_{2S}^*$  and  $f_S$  readings were not available (n.a.). <sup>b</sup>Signals on the IR-UTE images showed single-component decay with ultrashort  $T_2^*$ .

D<sub>2</sub>O, deuterium oxide;  $f_S$ , ultrashort T<sub>2</sub> component fraction; IR, inversion recovery;  $T_{2S}^*$ , ultrashort T<sub>2</sub> component T<sub>2</sub>^\*.

excitation was used (Figs. 5C and 5F). These results demonstrated the efficiency of the half-pulse UTE and IR-UTE sequences in measuring  $T_{2S}^*$  and detecting  $f_S$  changes in the presence of different proportions of LTCs.

Figure 6 shows  $T_{2S}^*$  and  $f_S$  measured (using half pulse excitation) from three WM long segments that were subjected to 1 h (no. 1), 2 h (no. 2), and 27 h (no. 3) D<sub>2</sub>O exchanges, respectively.  $T_{2S}^*$  was similar across all specimens measured with either UTE sequences or IR-UTE sequences with two different TIs (85 ms apart). The UTE sequence detected an increase of  $f_S$  from ~9.5% to ~47.1% in specimens no.1 through no. 3 (Fig. 6A), consistent with progressive D<sub>2</sub>O exchange. The IR-UTE signals of the specimen after 27 h D<sub>2</sub>O exchange (no. 3) could be characterized using single-component fitting at both TIs used, suggesting that  $f_S$  was less dependent on TI as the D<sub>2</sub>O exchange progressed.

Table 1 shows  $T_{2S}^*$  and  $f_S$  measured (using half pulse excitation) from three hemisphere slabs before and after 27 h exchange with  $D_2O$ . Before exchange with  $D_2O$ , the UTE signals showed single-component decay with a long  $T_2^*$  (~24 ms), and  $f_S$  and  $T_{2S}^*$  values could not be obtained. Meanwhile, the IR-UTE signals could be characterized through single-component  $T_2^*$  fitting at TI = 295 ms (i.e.,  $WM_L TI_{null}$ ) or through bicomponent  $T_2^*$  fitting at TI = 365 ms, with  $f_S$  values of less than 20% in all specimens. After exchange with  $D_2O$ , the UTE signals contained a  $f_S$  of ~40% and the IR-UTE signals showed single-component decay at both TIs used (i.e., TIs = 250ms and 380 ms).  $T_{2S}^*$  was in the range of 200 to 400  $\mu$ s either before or after exchange in all specimens. These results further support that protons that were unexchangeable with  $D_2O$  can be measured by UTE and IR-UTE sequences when used with half excitation pulses.

#### DISCUSSION

The present study used a bicomponent  $T_2^*$  fitting model as a simplified means to quantify  $f_S$  and  $T_{2S}^*$  in WM. In UTE images obtained with hard pulse excitation,  $f_S$  was found to be ~4% to 8% in native specimens and up to 86% in deuterated specimens. In IR-UTE images obtained with either hard or half pulse excitation,  $f_S$  was close to 100% at  $WM_L$  TI<sub>null</sub> and decreased significantly with increasing TI in native specimens but not in deuterated specimens.  $T_{2S}^*$  was found to be 150 to 400 µs in both native and deuterated specimens imaged using either hard pulse or half pulse excitations, and was consistent with previously reported values in myelin extract and nerve tissue (14,22).

#### Results Obtained With Hard or Half Excitation Pulses

This study first tested the UTE and IR-UTE sequences with hard excitation pulses in pure WM short blocks using a 5 mm-diameter solenoid coil that produces high SNR. The results showed that the  $f_S$  was strongly dependent on the choice of TI, whereas  $T_{2S}^*$  was not. When TI was increased by 5 ms and then by 25 ms from  $WM_L$  $\mathrm{TI}_{\mathrm{null}}$ ,  $f_S$  decreased from ~100% to ~62% and then to  $\sim$ 34%, respectively. Then, the sequences were used to image a native cerebral hemisphere slab using both hard and half excitation pulses. The UTE sequence failed to identify the WM<sub>S</sub> when used with half excitation pulses, and the IR-UTE sequence detected consistently lower  $f_S$ values using half excitation pulses than hard excitation pulses over a range of TIs that were offset from  $WM_L$ TI<sub>null</sub>. These results are consistent with the expectation that the half excitation pulses had lower efficiency than hard excitation pulses in detecting STCs. As compared with the half excitation pules, the hard excitation pulses have shorter durations  $(32 \,\mu s \, vs. \, 472 \,\mu s)$ , thus broader bandwidth (8.2 kHz vs. 2.7 kHz) and higher efficiency in exciting ultrashort T<sub>2</sub> components (20,23,25). However, with TI set at  $WM_L$  TI<sub>null</sub>,  $T^*_{2S}$  could be measured through single-component fitting of the multi-TE IR-UTE signals obtained using half excitation pulses, with the results similar to those obtained using hard excitation pulses. Furthermore, the IR-UTE sequences used with either hard or half excitation pulses were able to provide stable  $T_{2S}^*$  values and detect the differences in  $f_S$  when varying TIs were used. These results support the view that it is possible to use the IR-UTE sequences with half excitation pulses to detect STCs using clinical wholebody MR scanners.

Table 1

# Comparison of $T_{2S}^*$ and $f_S$ in Native and Deuterated Specimens

This study included three  $D_2O$  exchange experiments. First, two WM mini-blocks were subjected to an identical 8 h five-pass D<sub>2</sub>O exchange procedure and imaged using hard excitation pulses. After exchange, the specimens showed a very narrow HDO peak in the FID spectrum at 11.7T, and an  $f_S$  as high as ~86% on UTE images at 3T versus  $\sim$ 7.9% in the native specimen. In the TE range of 10 µs to 4 ms, the UTE images of the deuterated specimen showed marked signal intensity decay with increasing TE, in contrast to the barely observable changes in the native specimen. This experiment confirmed significant water signal loss in the deuterated specimen, and proved that UTE sequences used with hard excitation pulses are sensitive to proportional water proton and myelin proton content changes in biological tissue. Second, three pure WM long segments were subjected to different durations of D<sub>2</sub>O exchange and imaged using half excitation pulses. When the specimens were imaged using UTE sequences,  $f_S$  was observed to increase from  $\sim 9.5\%$  (1 h, one-pass exchange) to  $\sim 47.1\%$  (27 h, fourpass exchange). When the specimens were imaged using IR-UTE sequences,  $f_S$  was shown to be less affected by the choice of TI as the  $D_2O$  exchange duration increased. This experiment demonstrated that UTE and IR-UTE sequences used with half excitation pulses are also sensitive to proportional water proton and myelin proton content changes in biological tissue. Lastly, three cerebral hemisphere slabs were imaged using half excitation pulses both before and after a 27 h four-pass  $D_2O$  exchange. Results obtained from these specimens further confirmed an increase in  $f_S$  after D<sub>2</sub>O exchange ( $f_S = \sim 39\% - 42\%$  vs. unmeasurable) when measured using UTE sequences. In all three experiments, when the UTE or IR-UTE sequences were used with either hard or half excitation pulses, the  $T_{2S}^*$  values were consistently in the range of 150 to 400 µs both before and after exchange with  $D_2O$ . Furthermore, these  $T_{2S}^*$  values were comparable with those obtained using the same protocols from myelin powder as well as a myelin paste in  $D_2O$  (22). These results jointly suggest that the ultrashort  $T_2^*$  signals survived  $D_2O$  exchange and provide evidence to support that 1) myelin, presumably the nonexchangeable methylene groups (14), is the major source of the ultrashort T<sub>2</sub><sup>\*</sup> signals seen in WM on IR-UTE images; and 2) half-pulse IR-UTE sequences are sensitive to myelin concentration changes and therefore could be used for direct quantification of myelin loss.

Because myelin has a tightly organized, closely packed, highly stable structure with a hydrophobic core (28), complete replacement of water protons in nervous tissue with deuterons may be difficult, and the efficiency may depend on the form and type of tissue the myelin is in. For example, Horch et al. detected a loss of ~95% of the LTCs in a frog sciatic nerve (peripheral nerve), but only ~68% in a rat optic nerve (cranial nerve) after both nerves were incubated in D<sub>2</sub>O buffer for 2 h (14). This might be associated with the structural and chemical differences of myelin tissue in the two nerves (29). Even after a rat optic nerve was immersed in D<sub>2</sub>O buffer for 10 days, only a loss of ~82% of LTCs was observed (14).

All the specimens used in the present study were much larger in size than the frog sciatic or rat optic nerves (~1 mm or smaller in diameter) used in the Horch et al.'s study (14). This might explain why our UTE experiments only detected an  $f_S$  of ~86% (instead of ~100%) in a WM short block (small specimen) after 8 h, five-pass D<sub>2</sub>O exchange, and ~39% to 47% in the larger specimens (one WM long segment and three cerebral hemisphere slabs) after 27 h, four-pass D<sub>2</sub>O exchange. The relatively lower  $f_S$  in the larger specimens relative to that in the small specimen might be due to difference in the efficiency of D<sub>2</sub>O exchange, the lower efficiency of half pulse excitation relative to hard pulse excitation, and the lower SNR obtained with the 76-mm surface coil compared to the 5 mm solenoid coil.

# Comparison of Magnetization Transfer and IR-UTE Measurements of Myelin Tissue

The restricted proton pool (STCs) in white matter include not only protons in myelin lipids but also those in proteins and other semisolid membrane components. According to the literature (13,14) and our own results, the ultrashort  $T_2^*$  signals (  $T_2^*,\ {\sim}300\,\mu s)$  primarily arise from protons in methyl  $(-CH_3)$  and methylene  $(-CH_2)$ groups of phospholipids in myelin tissue. If the observed short  $T_2^*$  signals were mainly from proteins, they would not survive D<sub>2</sub>O exchange because protons in amino  $(-NH_2)$  and carboxyl (-COOH) groups in the majority of proteins could be exchanged with deuterons in D<sub>2</sub>O. Protons in proteolipid proteins, which are highly hydrophobic, may survive  $D_2O$  exchange, but the proteolipid proteins-to-lipid ratio in brain myelin is about 1:9 (30). Therefore, the contribution from these protons, if any, would be very small. Magnetization transfer (MT) is another MR contrast mechanism that has been used to probe the restricted proton pool by measuring the magnetization exchange between the restricted proton pool and the free proton pool (10,31,32). Carboxyl (-COOH) groups on cholesterol and galactocerebrosides are thought to be the major interaction sites between the restricted pool and the free pool for magnetization exchange (33-35). Therefore, MT contrast and the measured contrast in this study provide different aspects of information of the restricted proton pool. Furthermore, MT ratio derived from MT imaging is an indirect measure, whereas the IR-UTE sequence provides a direct measure of STCs. Therefore, we believe the measured contrast is more specific to myelin lipid protons than MT ratio.

## Limitations

The WM<sub>L</sub> TI<sub>null</sub> was estimated in this study using the same IR-UTE sequence as that used for  $T_{2S}^*$  measurement, with the intention of eliminating sequence-dependent T<sub>1</sub> measurement inaccuracy (36). However, the T<sub>1</sub> signals of LTCs were modeled using a single-component fitting model, and the T<sub>2</sub><sup>\*</sup> signals were fitted using the magnitude images for simplicity. It has been reported that there are two or more T<sub>1</sub> relaxation pools in brain tissue that are associated with distinctive T<sub>2</sub> relaxation rates (37,38) and these might have different T<sub>1</sub>s. Therefore,

our approach might induce insufficient nulling of LTC signals, as suggested by the observation that the value of  $f_S$  was not 100% in native WM in IR-UTE experiments at  $WM_L$   $TI_{null}$  (Fig. 2E). Fitting the signal in complex images would provide more accurate estimation of  $f_S$ . In addition, more advanced inversion recovery UTE sequences are currently under investigation by our team (39) and another group (40) to simultaneously suppress long  $T_2$  water signals with a broad range of  $T_1$ s. Second, our  $f_S$  values were relatively small compared with the known myelin tissue concentration in the wet mass of WM (41,42). This might be due to insufficient nulling of LTCs. In addition, one previous NMR spectroscopy experiment at 9.4T revealed multiple T<sub>2</sub><sup>\*</sup> components in purified bovine myelin extract (suspended in D<sub>2</sub>O), with 26.4% of the total signal having an effective lifetime of  $<\!25\,\mu s,~51.8\%$  of  $<\!0.1$  ms, and 91.6% of  $<\!1$  ms (at 20 °C) (21). The supershort  $T_2$  components ( $T_2$ , 0.01–0.1 ms) are probably not accessible by our UTE/IR-UTE imaging experiments at 3T. Third, a simplified bicomponent model was used in this study to fit the  $T_2^*$  signals in the magnitude IR-UTE images. The  $T_{2L}^*$  has limited biophysical meaning because it is associated with multiple water pools. Multicomponent models might be desirable in order to fully capture the tissue relaxation properties. However, such analysis can be very challenging and is prone to errors; thus, it needs further investigation. Fourth, plain D<sub>2</sub>O was used in this study to replace H<sub>2</sub>O in the specimens. Ideally D<sub>2</sub>O buffer should be used to mimic physiological osmotic pressure and pH effects. Nonetheless, the  $T_{2S}^*$  values before and after D<sub>2</sub>O exchange were relatively consistent in our study, suggesting that the  $T^*_{2S}$  was not significantly affected by incubation in plain  $D_2O$ , possibly because the methylene groups that are known to form hydrophobic nonpolar hydrocarbon chains (28) were not interchangeable with or inaccessible by  $D_2O$ (14). Fifth, cross-relaxation between macromolecules and water was not considered in this study. Such crossrelaxation can produce apparent shortening of  $T_1$  and  $T_2$ (31,43), but its effects on ultrashort  $T_2^*$  values are not fully understood. However, it is unlikely to have significantly affected the ultrashort  $T_2^*$  signals measured in this study. In this study, the  $WM_L\ TI_{null}$  was directly determined via fitting the IR-UTE images obtained with different TIs. Cross-relaxation-induced T<sub>1</sub>-shortening would only result in a shorter estimated  $WM_L$   $TI_{null}$ . Meanwhile, the ultrashort T<sub>2</sub><sup>\*</sup> values were largely consistent in specimens before and after different durations of D<sub>2</sub>O exchange, further suggesting our measurements were unlikely significantly contaminated by cross-relaxation between the restricted proton pool and the free water pool. This might be further confirmed through comparing the  $T_2^*$  fitting results obtained with UTE experiments with or without magnetization transfer preparation. Lastly, the 2D IR-UTE sequences typically employ slice-selective half excitation pulses for in vivo studies. Eddy currents associated with the slice-selective gradients may lead to out-of-slice signal contamination (20,44), thus affecting  $T_2^*$  analysis. This error can be reduced by measuring the slice-selective gradients, adding precompensation gradients, and/or using the measured readout gradients for re-gridding (45). In this study, we used nonselective gradients to minimize

eddy currents and speed up data acquisition. The use of nonselective pulses is valid as the focus of this paper is to further demonstrate that myelin protons are directly detectable with 2D IR-UTE sequences on a clinical MR scanner. Further work is needed to increase the robustness of this sequence by using advanced gradient calibration and image reconstruction techniques.

### CONCLUSION

The  $f_S$  values were significantly affected by the TIs used in IR-UTE sequences with either hard or half excitation pulses in native specimens but not in heavily deuterated specimens.  $T_{2S}^*$  was in the range of 150 to 400 µs in all measurements made with either hard or half-pulse excitation UTE and IR-UTE sequences. The major source of the ultrashort  $T_2^*$  signals seen in WM on IR-UTE images is likely to be myelin protons, presumably methylene groups. Half-pulse IR-UTE sequences have the potential for direct qualitative and quantitative myelin imaging.

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### SUPPORTING INFORMATION

Additional supporting information may be found in the online version of this article.

**Fig. S1.** Spectra from two white-matter (WM) short blocks examined either in the native condition or after 27 h immersion in deuterium oxide (D<sub>2</sub>O). (A) Spectrum from one WM short block examined in the native condition. The integrated area under the water peak: the total area = 1:1.13, suggesting a ~88.5% contribution from water to the total signal. (B) Spectrum from one WM short block examined after 27 h immersion in D<sub>2</sub>O (four changes of D<sub>2</sub>O with intervals of 1 h, 1 h, 2 h, and 23 h). The integrated area under the water peak: the total area = 1:4.07, suggesting a ~24.6% contribution from water to the total signal.

**Fig. S2.** Illustration of the procedure for defining the region of interest (ROI) in white matter for myelin proton  $T_2^*$  quantification in cerebral hemisphere slabs. The images were from the ~8 mm thick cerebral hemisphere slab. Yellow boxes show the final ROI. (A-H) Eight consecutive  $T_2$ -weighted fast spin echo (FSE) images (TR/TE = 3000 ms/40 ms, slice thickness = 1 mm, and no slice gap). (I) The maximum intensity projection (MIP) of (A) to (H). (J) The IR-UTE image (TR/TI/TE = 1000/300/2.2 ms, half pulse excitation, slice-selective gradient turned off) showing similar contrast between white matter (low signal) and grey matter (high signal) as (I). The ROI was first manually drawn on the IR-UTE image (J), copied to the multi-slice  $T_2$ -weighted FSE images (A-H) and then refined to exclude any gray matter contamination.