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Simulation of axial tensile well deformation during reservoir compaction in offshore unconsolidated methane hydrate-bearing formation

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5

6 1. Introduction

7

8 Methane hydrate typically exists within the pores of unconsolidated formation under high pressure
9 and low-temperature conditions. As such, methane hydrate is a potential energy resource as it could
10 contain 500 gigatons of carbon [1] equivalent to 10 times the amount of world's undiscovered
11 conventional gas resources which are considered to be technically recoverable [2]. Field gas
12 production tests have been conducted in Canada [3], US [4], Japan [5], [6] and China [7], to assess the
13 feasibility of commercial gas production from methane hydrate reservoirs. One of the main challenges
14 toward sustainable gas production has been identified as well/formation integrity due to the
15 unconsolidated nature of the methane hydrate-bearing sediments [8]–[10]. Recent gas production tests
16 at the Nankai Trough in Japan show that sand production issue caused premature termination of the
17 gas production test [6], [11].

18

19 Earlier attempts to investigate well integrity in methane hydrate-bearing formation were documented

20 in Freij-Ayoub et al. [12], [13] where they assessed well integrity during heating-induced hydrate
21 dissociation. However, well integrity during reservoir compaction, which might be the cause of the
22 well failure/sand production at the Nankai Trough, was not assessed in their study. Subsequently, well
23 integrity in methane hydrate reservoirs during reservoir compaction was investigated by several
24 researchers ([14], [8], [15] and [9]). Rutqvist et al. [14] showed that the gap between the casing and
25 formation, which developed during well construction (e.g., poor cement job) would adversely affect
26 formation integrity around a horizontal well during gas production. Their work indicates the
27 importance of simulating the well construction processes for the assessment of wellbore integrity
28 during gas production. Qiu et al. [8] simulated 20-day gas production at the Nankai Trough and
29 showed that the casing, cement and screen could accumulate approximately 1% of plastic strain,
30 which they argue would be negligibly small to cause well failure. However, if gas production longer
31 than 20 days were simulated, greater reservoir compaction would occur, and there is a risk of well
32 failure. Yoneda et al. [9] also simulated gas production at the Nankai Trough and found that tensile
33 deformation of the well developed in the overburden layer due to reservoir compaction, which might
34 cause the tensile failure of the gravel pack to induce sand production. This highlights the importance
35 of analyzing the tensile deformation of the well during reservoir compaction. Finally, the work of Shin
36 and Satamarina [15] indicates that well integrity is affected by the change in formation permeability
37 during compaction. If the permeability of the soil is sensitive to porosity reduction during
38 depressurization, then reservoir compaction is inhibited as a low permeability zone develops around

39 the wellbore. This implies that different patterns of permeability change in response to porosity
40 change in the reservoir layer would develop different reservoir compaction profiles, which in turn
41 affects well integrity.

42

43 Table 1 shows the list of numerical work on well integrity during reservoir compaction [16], including
44 the ones introduced above. The main shortcomings of these studies are the omission of the well
45 construction process prior to simulating reservoir compaction ([9], [15], [17]–[27]) and omission of the
46 cement sheath ([14], [15], [20]–[24], [28]). Currently, the work by Xu
47 [29] is the only existing work which simulates detailed well construction processes, such as cement
48 shrinkage, and also models the integrity of the cement sheath. In Xu’s work, detailed well
49 construction processes including drilling, casing hanging, cementing, cement hardening/shrinkage and
50 casing landing are simulated. After the well construction, well integrity (i.e., casing and cement)
51 during different reservoir compaction profiles were assessed. In the simulation, however, the
52 depressurization (i.e., pore pressure) profile was specified by artificial step functions, and this may
53 have computed unrealistic reservoir compaction profiles. In addition, cement shrinkage values used in
54 the simulation were not representative of the Nankai Trough case, because such cement shrinkage
55 values were not investigated extensively at the time of his work. In addition, the casing-cement
56 interface friction behaviour was not calibrated against experimental data, and the simple Coulomb
57 friction model was employed. In this study, the prior work by Xu [29] is extended by employing (i)

58 depressurization profiles that are physically realistic for hydrate reservoirs; these consist of hydrate
59 dissociated (i.e., high permeability) and undissociated (i.e., low permeability) zones, and also
60 employing (ii) cement shrinkage volume estimated specifically for the Nankai Trough case [30]
61 together with (iii) calibrated casing-cement interface friction model. A parametric numerical study is
62 carried out with an axisymmetric finite element model considering different depressurization and
63 hydrate dissociation profiles in the reservoir.

64

65 Table 1 Existing numerical work on well integrity during reservoir compaction expanded from [16].

Authors	Casing	Cement	Formation	Well construction	Reservoir compaction
Bruno & Bovberg (1992)[20]	No	No	Yes	No	Yes
Hamilton et al. (1993)[21]	No	No	Yes	No	Yes
Fredrich et al. (2000)[22]	No	No	Yes	No	Yes
Sayers et al. (2006) [23]	No	No	Yes	No	Yes
Furui et al. (2011)[24]	No	No	Yes	No	Yes
Shin & Santamarina (2016)[15]	Yes	No	Yes	No	Yes
Chia & Bradley (1988)[25]	Yes	Yes	Yes	No	Yes
Yudovich et al. (1988)[27]	Yes	Yes	Yes	No	Yes
Chia & Bradley (1989)[26]	Yes	Yes	Yes	No	Yes
Li et al. (2003)[17]	Yes	Yes	Yes	No	Yes
Li et al. (2005)[18]	Yes	Yes	Yes	No	Yes
Jinnai & Morita (2009)[19]	Yes	Yes	Yes	No	Yes
Yoneda et al. (2018)[9]	Yes	Yes	Yes	No	Yes
Klar et al. (2010)[28]	Yes	No	Yes	Yes	Yes
Rutqvist et al. (2012)[14]	Yes	No	Yes	Yes	Yes
Qiu et al. (2015)[8]	Yes	Yes	Yes	Yes	Yes
Xu (2014)[29]	Yes	Yes	Yes	Yes	Yes

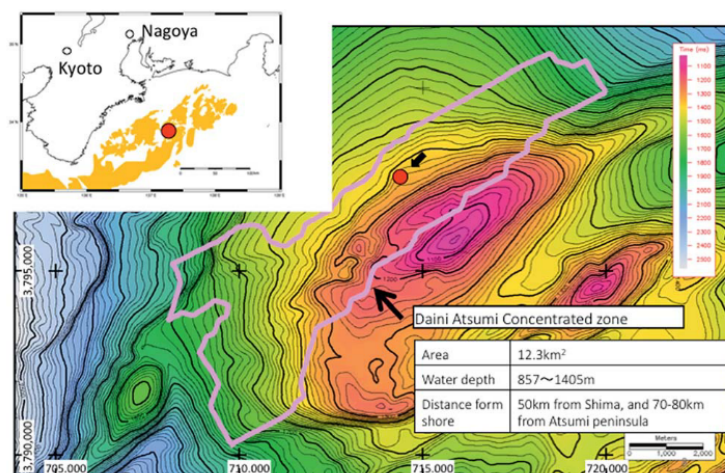
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67

68 This study focuses on the case of the Nankai Trough methane hydrate-bearing reservoir. Fig. 1 shows

69 the location of the Nankai Trough methane hydrate site for the 2013 gas production test [11]. The gas
70 production site was located on the north slope of the Daini Atsumi Knoll off the coast of Japan (Fig.
71 1a). Three wells were drilled by the drilling vessel D/V Chikyu in 2012: one production well (AT1-P)
72 and two monitoring wells (AT1-MC and AT1-MT1) (Fig. 1b). The methane hydrate reservoir layer is
73 located approximately 300 m below the seafloor, and it has roughly a 50 m thickness. The water depth
74 of the Nankai Trough site is approximately 1,000 m. As the depth of the reservoir layer from the
75 seafloor is relatively shallow, the formation consists of unconsolidated sand and clay. Such formation
76 is susceptible to large volumetric compaction upon depressurization. Therefore, caution must be taken
77 to assess the well integrity in response to reservoir compaction at the Nankai Trough site. Sand
78 production occurred during the six-day gas production test [10], [31], which could be attributed to
79 well failure during the gas production trial.

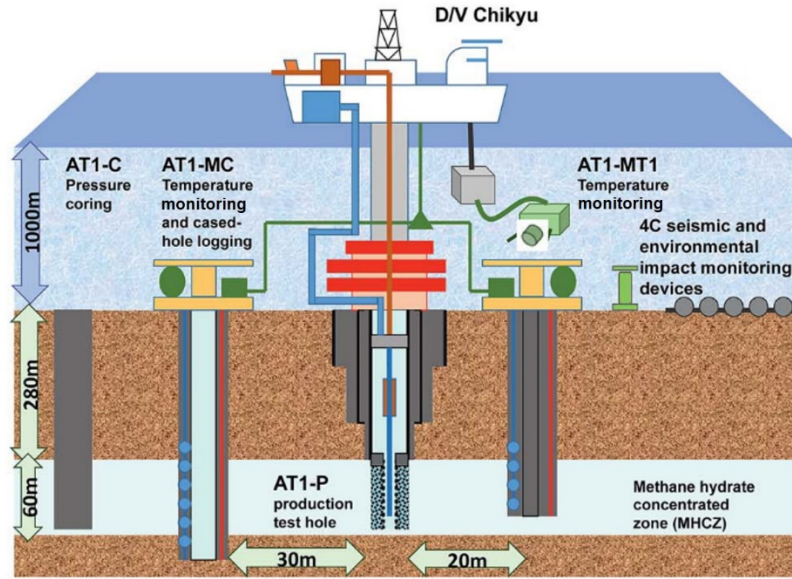
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81

82

(a)



83

84

(b)

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Fig. 1 Details of the Nankai Trough site for the 2013 gas production test [11]: (a) location of the test

86

site; (b) geometries of the wells and formation layers (courtesy of MH21).

87

88

Fig. 2 shows potential well failure mechanisms that may occur at the Nankai Trough site [16]. Uneven

89

hydrate dissociation, which occurs due to heterogeneous hydrate distribution around the well, could

90

cause bending or buckling failure due to the unbalanced lateral support from the reservoir layer. The

91

axial tension failure is caused by the elongation of the overburden formation in response to reservoir

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compaction, whereas axial compression occurs in the reservoir layer. A deviated well could suffer

93

bending failure at the reservoir layer boundaries. Finally, cement shrinkage could also cause well

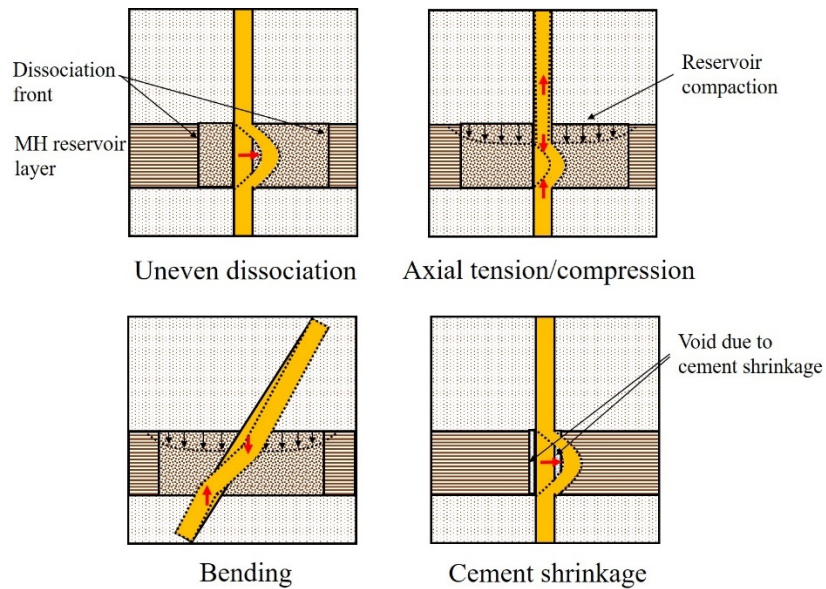
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failure by reducing lateral support for the well.

95

96 Among these well failure mechanisms, this paper examines well integrity due to the axial tension
 97 mechanism for the following reasons. First, the tensile strength of cement is approximately one-tenth
 98 of its compressive strength [32] and thus, the cement is much more likely to fail by tension than by
 99 compression. Second, the tension failure could propagate up to the seafloor with the progress of
 100 reservoir compaction, whereas the other failure types tend to be localized within the reservoir layer.
 101 Therefore, the tension mechanism of well failure is considered to be critical to the long-term
 102 sustainable gas production from hydrate reservoirs. As cement shrinkage can affect tensile
 103 deformation of the well by reducing the friction at the casing-cement interface, cement shrinkage is
 104 considered in this study in combination with reservoir compaction.

105



106

107

Fig. 2 Failure mechanisms of the well in hydrate reservoirs [16].

108

109 Focusing on the Nankai Trough case, the objectives of this study are as follows.

- 110 (i) to evaluate the effect of different reservoir compaction patterns on the tensile stress and strain
111 development of the casing/cement in the overburden layer via parametric numerical
112 simulations,
- 113 (ii) to evaluate the effect of cement shrinkage volume on tensile deformation of the well in the
114 overburden layer during reservoir compaction, and
- 115 (iii) to investigate the correlation between the tensile stress and strain development of the
116 casing/cement in the overburden layer and depressurization/hydrate dissociation patterns in
117 the reservoir.

118

119

120 2. Finite element modelling

121

122 2.1. Model geometry

123 Figure 3 shows a schematic diagram of the axisymmetric finite element model created for this study.

124 The total depth and radius of the model are 650 m and 600 m, respectively. The thickness of the

125 methane hydrate reservoir layer is 50 m, whereas the thicknesses of the overburden and underburden

126 layers are 300 m. A borehole with a radius of 0.312 m (12 1/4 inches) is drilled in the overburden

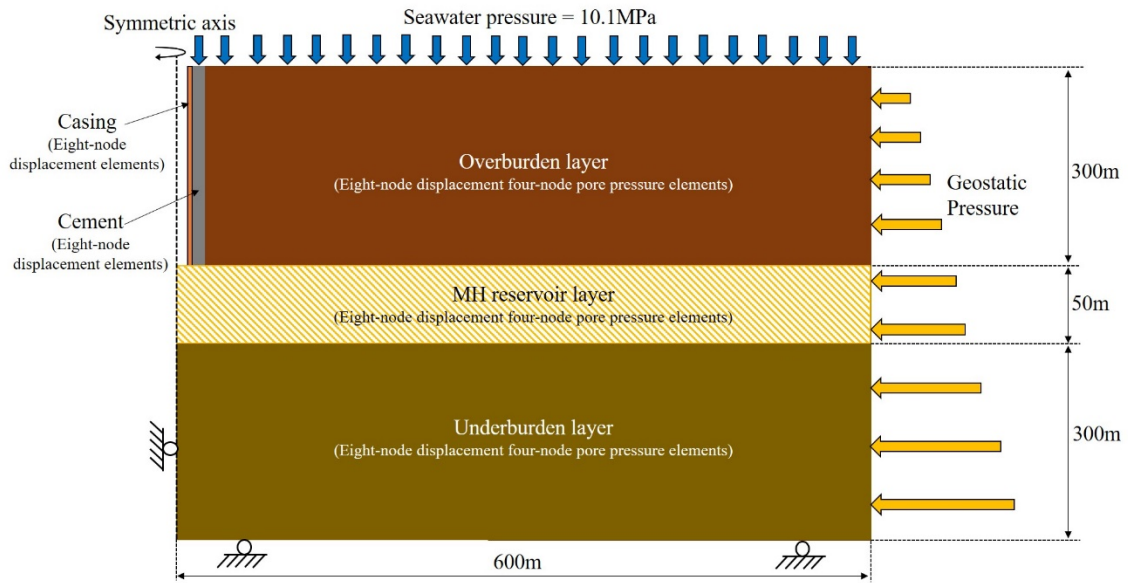
127 layer. The outer diameter and a wall thickness of the casing placed inside the borehole are 0.122 m (9

128 5/8 inches) and 0.01 m (0.4 inches), respectively. Cement is placed in the annulus between the casing
129 and formation. The roller boundary constraint is applied at the left and bottom edges of the model
130 whereas a constant distributed pressure load is applied at the top (i.e., hydrostatic pore pressure) and
131 right (i.e., geostatic stress) edges. The top of the model is assumed to be 1,000 m below the sea
132 surface.

133

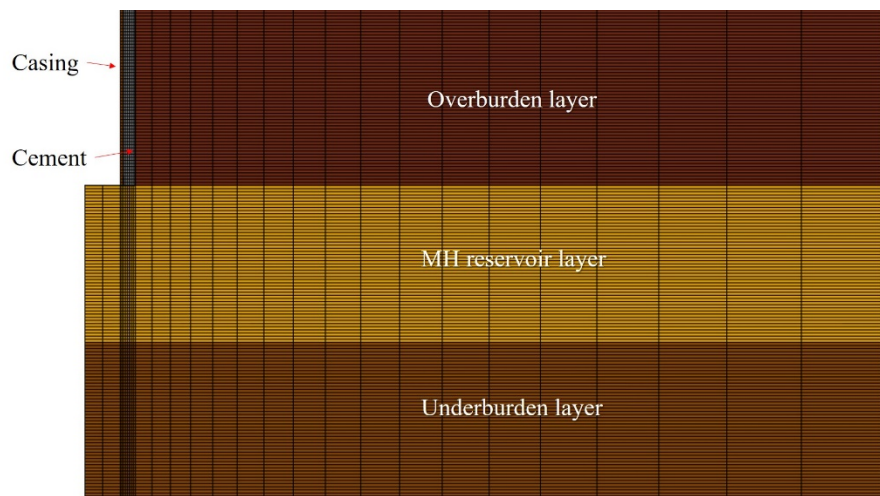
134 Fig. 4 shows the FE mesh of the model. The formation is discretized into 55,250 eight-node
135 displacement four-node pore pressure elements, whereas the casing and cement are discretized into
136 600 and 1,800 eight-node displacement elements, respectively. The horizontal length of the casing and
137 cement elements is uniformly set to be 5.0×10^{-3} m and 6.3×10^{-3} m, respectively, whereas the
138 horizontal length of the formation elements is increased exponentially with increasing radius from the
139 wellbore (5.3×10^{-2} m at the cement-formation interface and 53 m at the right edge of the model). The
140 vertical length of the mesh is uniformly set to 1 m regardless of the element type.

141



143 Fig. 3 The geometry of the axisymmetric finite element model.

144



146 Fig. 4 The mesh of the axisymmetric finite element model near the bottom of the wellbore (enlarged a
 147 hundred times in the horizontal direction).

148

149 2.2. Simulation steps

150 2.2.1. Initial conditions

151 The initial vertical stress distribution of the formation is derived from the *in situ* density measurement
152 data [33]. The initial void ratio distribution is also obtained from the same *in situ* density
153 measurement data. For the initial pore pressure distribution, the hydrostatic pore pressure distribution
154 with the seawater density of 1.027 g/cm³ is assumed. Two different initial horizontal effective stress
155 profiles (i.e., overconsolidated and normally consolidated overburden cases) are employed, as shown
156 in Fig. 5. The overconsolidated distribution is calculated via Equation 1:

$$\sigma'_h = (1 - \sin\phi')\text{OCR}^{\sin\phi'}\sigma'_v \quad (1)$$

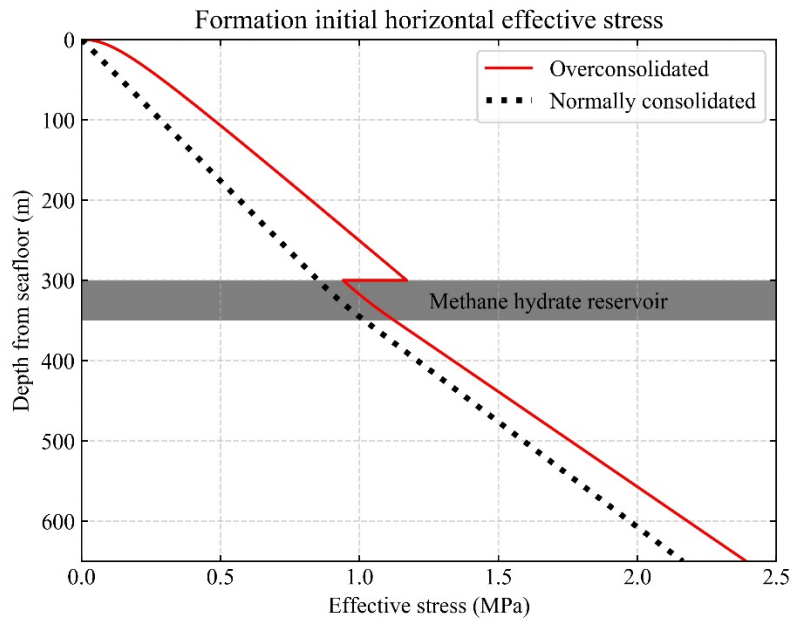
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158
159 where σ'_h is the horizontal effective stress, σ'_v is the vertical effective stress, ϕ' is the internal
160 friction angle, OCR is the overconsolidation ratio. The OCR values of the overburden layer are
161 derived from triaxial test data on formation core samples retrieved at the Nankai Trough [34], whereas
162 OCR = 1 is employed for the reservoir and underburden layers, in which case Equation 1 reduces to
163 Jaky's formula. This means that the reservoir and underburden layers are normally consolidated, and
164 this is consistent with the triaxial compression test results on reservoir and underburden sediment
165 cores recovered at the Nankai Trough [9]. Hence, the overconsolidated overburden case is more
166 representative of the actual Nankai Trough formation. For the normally consolidated overburden case,

167 the initial horizontal effective stress is calculated via $\sigma'_h = 0.4\sigma'_v$. The effect of different initial

168 horizontal effective stress distributions is investigated in Section 3.5.

169



170

171 Fig. 5 Initial horizontal effective stress distributions of the formation.

172

173 2.2.2. Well construction process

174 The construction process of the well is incorporated in the simulation. The modelling methodology of

175 the well construction process is identical to the one employed in [35]. The modelled construction

176 stages are listed in Table 2. The cement shrinkage volume of 0.75% is employed in the cement

177 shrinkage stage, which could be expected in the Nankai Trough scenario [30]. The volumetric strain

178 decrement (i.e., volumetric shrinkage) was generated via fictitious thermal contraction by reducing the

179 temperature of the cement elements.

180

181

Table 2 The well construction processes incorporated in the simulation.

Construction process	Duration (hour)
1. Drilling	14.4
2. Casing hanging	Immediate
3. Cementing	Immediate
4. Cement hardening/shrinkage	40.8
5. Casing landing	Immediate

182

183 2.2.3. Decoupled depressurization and hydrate dissociation process

184

The depressurization/hydrate dissociation stage is simulated in a decoupled manner by specifying the

185

pore pressure distribution in the reservoir layer, rather than simulating the actual depressurization and

186

dissociation processes in a thermo-hydromechanically coupled manner. This approach allows for

187

creating different reservoir compaction profiles. The analytical steady-state pore pressure distribution

188

shown below is employed to specify the pore pressure distribution in the reservoir layer:

189

$$u = C_1 \ln r + C_2 \quad (2)$$

190

191

where u is the pore pressure and r is the radius from the centre of the well. It is assumed that the

192

permeability of the hydrate dissociated zone ($0 \leq r \leq r_f$) is higher than that of the undissociated zone

193

($r > r_f$). Therefore, the above Equation 2 is applied to each zone separately while satisfying the

194

compatibility of the radial flow velocities at the boundary between the dissociated and undissociated

195

zones. By applying the remaining boundary conditions ($u = P_i$ at $r = r_o$, $u = P_o$ at $r = R_o$), the

196 values of the coefficients (C_1 and C_2) are obtained as follows:

197

$$C_1 = \begin{cases} (P_o - P_i) / \ln(r_f^{1-\alpha_p} R_o^{\alpha_p} / r_o) & (0 \leq r \leq r_f) \\ \alpha_p (P_o - P_i) / \ln(r_f^{1-\alpha_p} R_o^{\alpha_p} / r_o) & (r > r_f) \end{cases}$$
$$C_2 = \begin{cases} P_i - C_1 \ln r_o & (0 \leq r \leq r_f) \\ P_o - C_1 \ln R_o & (r > r_f) \end{cases}$$

198

199 where P_o is the hydrostatic pore pressure, P_i is the depressurized pore pressure in the wellbore, r_f
200 is the radius of the hydrate dissociation front, r_o is the radius of the wellbore, R_o is the radius where
201 hydrostatic pore pressure is recovered and α_p is the ratio of the permeability values between the
202 dissociated and undissociated zones. According to the literature, the value of α_p is dependent on the
203 hydrate saturation and it could be ~100 or higher [36]–[38]. In this study, it is set to a constant value
204 of 100. As to the value of r_f , coupled thermo-hydro(mechanical) simulations in the literature [5],
205 [28], [39] suggest that it is a fraction of R_o and increases with larger R_o . In this study, it is assumed
206 that $r_f = 0.5 R_o$.

207

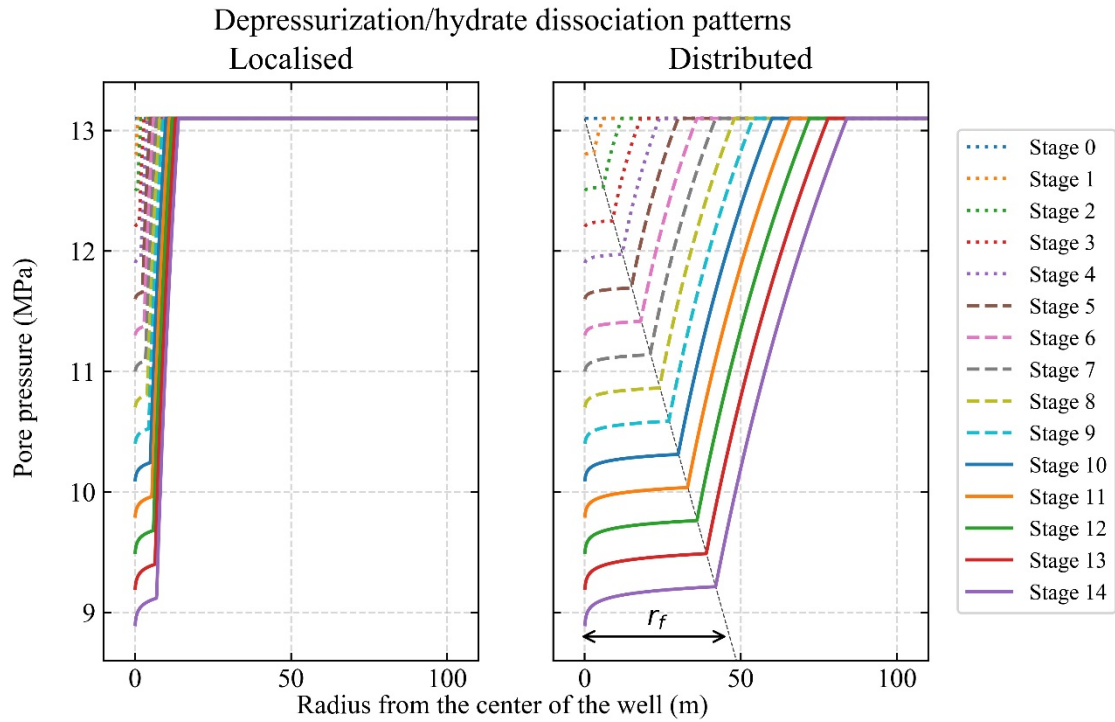
208 To model the progress of depressurization and hydrate dissociation, fourteen stages are considered in
209 the simulation. The values of P_i and r_f are linearly varied with time by ΔP_i and Δr_f in each stage from
210 the initial values of $P_i = P_o$ and $r_f = 0$. In the field, the rate of decrease of P_i depends on the speed of
211 depressurization specified by the operator, whereas the rate of increase of r_f depends on the speed of

212 hydrate dissociation, which is governed by the permeability of the reservoir and the heat supply from
213 the far-field. As changes in the formation permeability during hydrate dissociation are complex, the
214 rate of r_f increase may not be constant as assumed in this study. In order to estimate the increase rate
215 of r_f more precisely, it would be necessary to conduct thermo-hydromechanical coupled simulations,
216 similar to the ones presented in the literature [40]–[48].

217

218 Fig. 6 shows the simulated pore pressure profiles along the top of the reservoir layer in the case of
219 localized ($\Delta P_i = -0.3$ MPa and $\Delta r_f = 0.5$ m) and distributed ($\Delta P_i = -0.3$ MPa and $\Delta r_f = 3$ m) hydrate
220 dissociation cases. The former case represents low permeability hydrate-bearing reservoir scenarios,
221 which create a large difference in permeability between the dissociated zone and non-dissociated
222 zone. The latter case represents scenarios of hydrate-bearing formation with high permeability, which
223 results in less variation in permeability between the dissociated zone and non-dissociated zone. To
224 create various depressurization and hydrate dissociation profiles, different combinations of ΔP_i and
225 Δr_f values were employed (i.e., $\Delta P_i = -0.1, -0.2, -0.3, -0.4, -0.5, -0.6$ MPa and $\Delta r_f = 0.5, 1.0, 1.5, 2.0,$
226 $2.5, 3.0$ m). In total, 36 different depressurization and hydrate dissociation cases were simulated,
227 which are listed in Table 3.

228



229

230 Fig. 6 Simulated depressurization and hydrate dissociation (pore pressure) profiles in the reservoir

231 layer (localized depressurization and dissociation case ($\Delta P_i = -0.3$ MPa and $\Delta r_f = 0.5$ m) (left) and

232 distributed depressurization and dissociation case ($\Delta P_i = -0.3$ MPa and $\Delta r_f = 3$ m) (right)).

233

234

Table 3 The depressurization and hydrate dissociation cases simulated in this study.

Case#	ΔP_i (MPa)	Δr_f (m)	Case#	ΔP_i (MPa)	Δr_f (m)	Case#	ΔP_i (MPa)	Δr_f (m)
1		0.5	13		0.5	25		0.5
2		1.0	14		1.0	26		1.0
3	-0.1	1.5	15	-0.3	1.5	27	-0.5	1.5
4		2.0	16		2.0	28		2.0
5		2.5	17		2.5	29		2.5
6		3.0	18		3.0	30		3.0
7		0.5	19		0.5	31		0.5
8		1.0	20		1.0	32		1.0
9	-0.2	1.5	21	-0.4	1.5	33	-0.6	1.5
10		2.0	22		2.0	34		2.0
11		2.5	23		2.5	35		2.5
12		3.0	24		3.0	36		3.0

235

236

237 2.3. Constitutive models

238 Soils

239 The methane hydrate critical state model (MHCS model) [49] is employed to simulate the mechanical

240 behaviour of the soils at the site. The model parameters are calibrated against triaxial test data on

241 formation samples recovered at the Nankai Trough site [50], [51]. Selected calibration results are

242 presented in Fig. 7. Table 4 shows the calibrated values of the MHCS model parameters as well as

243 values of other formation parameters. It is noted that the MHCS parameters (i.e., m_1 , m_2 , A , B , C , D)

244 in the reservoir layer are set to zero in this study because of the pore pressure fixity in the simulation.

245 This would be a reasonable simplification considering that the hydrate is assumed to have no

246 enhancement effect on the bulk modulus in the MHCS model, i.e., the compaction behaviour of the

247 reservoir layer would not be affected by hydrate saturation. Although it is clear that this assumption of

248 the MHCS model has to be modified to reflect the dependence of bulk compressibility of hydrate-

249 bearing soil on hydrate saturation, this drawback would not generate significant errors in simulating

250 reservoir compaction for the following reasons. First, the majority of reservoir compaction occurs

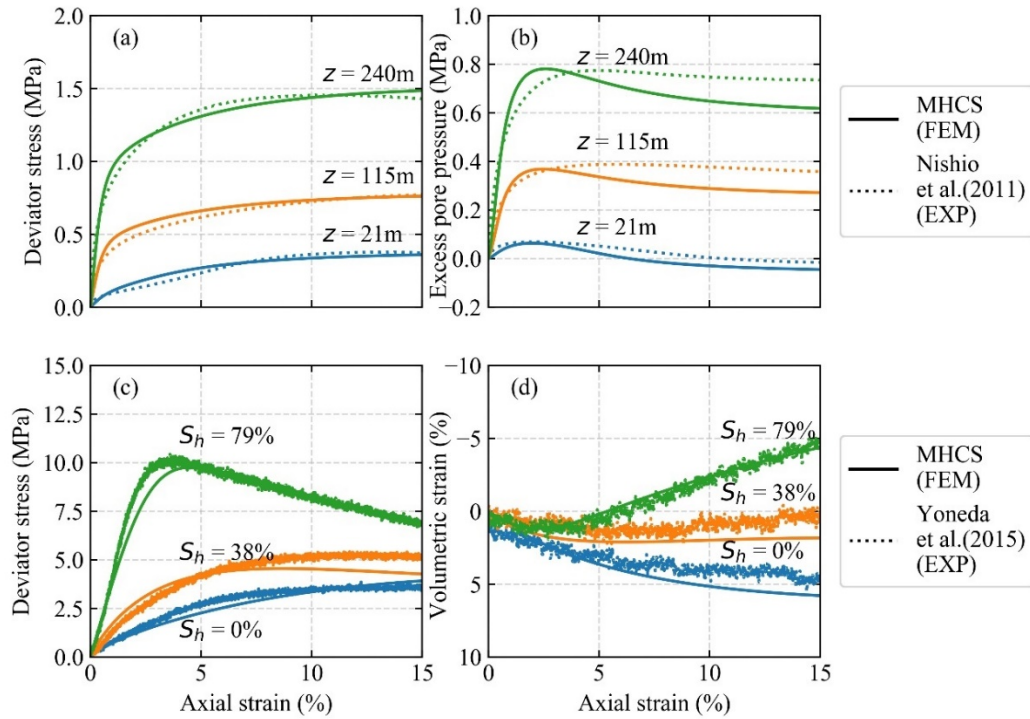
251 after hydrate-bearing soil undergoes yielding, where hydrate bond between soil particles breaks and

252 the enhancement effect of hydrate is lost. Second, the plastic compression is significantly larger than

253 the preceding elastic compression where the hydrate enhancement effect is still active. Therefore, it is

254 hypothesized in this study that hydrate saturation has negligible effects on the magnitude of reservoir
255 compaction and the simplified approach is taken where the MHCS parameters related to hydrate
256 saturation are all set to zero. This would be a reasonable assumption as compression tests on hydrate-
257 bearing soils revealed that the maximum change in elastic and plastic bulk compressibility were 68%
258 and 73%, respectively, for samples with hydrate saturation ranging between 18% and 85% [52], [53].
259 This is much smaller than an order of magnitude change, which can be considered negligible in
260 typical soil mechanics terms. The process of dissociation is not considered accordingly; only the
261 process of pore pressure propagation is considered. This simplified approach would simulate
262 scenarios closer to the worst case (the largest possible compaction) for well integrity than the fully-
263 coupled approach, which is convenient from the safety point of view. The values of the density and
264 void ratio of each layer of the formation are chosen based on the *in situ* measurement data at the
265 Nankai Trough [33]; the trend line for the raw density measurements is selected as the density
266 distribution and void ratios are back-calculated by assuming the constant grain density of clay and
267 sand particles (2.65 g/cm³).

268



269

270 Fig. 7 Calibration result of the MHCS model: (a) deviatoric stress vs. axial strain(clay); (b) excess
 271 pore pressure vs. axial strain (clay); (c) deviatoric stress vs. axial strain (sand); (d) volumetric strain
 272 vs. axial strain (sand).

273 Casing and cement

274 For the casing and cement elements, linear isotropic elasticity with the von Mises yield criteria
 275 (casing) and with the Mohr-Coulomb yield criteria (cement) is employed, respectively. The values of
 276 the casing and cement constitutive model parameters are listed in Table 5. These values are based on
 277 the actual casing and cement employed at the Nankai Trough [8].

278

279 Cement-casing and cement-formation interface

280 The interface behaviour in the contact tangential direction (i.e., interface friction) is modelled by an

281 interface friction constitutive model. The details of the interface friction constitutive model adopted in this
282 study are presented in Appendix A (supplementary material), and its verification is shown in Appendix B
283 (supplementary material). The parameters of the friction model are friction coefficient, cohesion and
284 ultimate elastic interface displacement. Table 6 lists the values of the parameters for the casing-cement and
285 cement-formation interfaces used in the model. The values of the friction coefficient and cohesion for the
286 casing-cement interface are obtained from the literature [54], whereas that of the ultimate elastic interface
287 displacement is calibrated to match the result of a laboratory experiment on a well specimen [55]. Details
288 of the calibration process are provided in Appendix C (supplementary material).

289

290 For the cement-formation interface, it is assumed that the interface friction coefficient is identical to
291 that of the underlying formation, and the mean friction coefficient value is calculated from the
292 calibrated values of the critical state frictional constant ($\mu = \tan(\sin^{-1}(3M/(6 + M)))$) to be 0.65
293 and it is used for the entire cement-formation interface. As to the interface cohesion, it is assumed to
294 be negligible as soil particles in the unconsolidated formation would not resist frictional force at zero
295 interface confining pressure (this is experimentally validated in the literature [54]). The value of the
296 ultimate elastic interface displacement is set to 0.25 mm. This is determined by a sensitivity analysis
297 where the reservoir compaction simulation, which is presented in the following sections, is performed
298 with varied values of the ultimate elastic interface displacement between 0.25 mm and 2.5 mm. It was
299 found that results (i.e., the development of stresses and strains in the casing and cement during

300 reservoir compaction) are identical regardless of the different values of the ultimate elastic interface
 301 displacement within the examined range. Therefore, the value is set to 0.25 mm in the study. This is
 302 supported by an experimental study [56], where the interface shearing between sand and a mortar
 303 plate is examined. The study shows that the value of the ultimate elastic interface displacement for the
 304 sand-mortar interface is approximately 0.3 mm.

305

306 The interface behaviour in the contact normal direction (i.e., interface pressure) is modelled by the
 307 ABAQUS inbuilt augmented Lagrange method, which is a combination of the linear penalty method
 308 and an augmentation iteration scheme. In the augmented Lagrange method, the contact pressure is
 309 calculated by multiplying the stiffness of the representative underlying elements with the interface
 310 penetration distance. The interface penetration is maintained below 0.1% of the characteristic element
 311 length of the model by iteratively augmenting the contact pressure.

312

313

Table 4 The parameter values of the MHCS model for the formation.

	Overburden clay	Methane hydrate reservoir	Underburden sand
Depth from the seafloor (m)	0~300	300~350	350~650
Saturated bulk density (kg/m ³)	1,750	1,750~2,000	2,000
Initial void ratio	1.31	1.31~0.717	0.717
Gradient of compression line, λ	0.18	0.10	0.10
Gradient of swelling line, κ	0.03	0.02	0.02
Critical state frictional constant, M	1.30	1.37	1.37
Poisson's ratio, ν	0.25	0.35	0.35
Subsurface constant, U	15	8	8
Stiffness enhancement constant, m_2	0	0	0
Hydrate degradation constant, m_1	0	0	0

Dilation enhancement constant, A	0	0	0
Dilation enhancement constant, B	0	0	0
Cohesion enhancement constant, C	0	0	0
Cohesion enhancement constant, D	0	0	0

314

315 Table 5 The parameter values of the constitutive models for the casing (von Mises) and cement

316

(Mohr-Coulomb).

	Casing	Cement
Density (kg/m^3)	7,897	1,198
Young's modulus (GPa)	200	0.131 (slurry) 3.81 (solid)
Poisson's ratio	0.28	0.49 (slurry) 0.20 (solid)
Yield stress (MPa)	379.5	N/A
Friction angle ($^\circ$)	N/A	30
Dilation angle ($^\circ$)	N/A	0
Cohesion (MPa)	N/A	2.72

317

318 Table 6 The parameter values of the interface friction constitutive model.

	Casing-cement interface	Cement-formation interface
Friction coefficient (-)	0.8	0.65
Cohesion (MPa)	3.0	0
Ultimate elastic interface displacement (mm)	0.5	0.25

319

320

321 3. Results

322

323 3.1. Effect of depressurization/hydrate dissociation patterns

324 In this section, the effect of depressurization/hydrate dissociation patterns on reservoir subsidence and

325 the stress and strain development of casing and cement is presented. The cement shrinkage volume of

326 0% and the overconsolidated overburden case are applied.

327

328 3.1.1. Formation deformation patterns

329 Fig. 8 shows the reservoir compaction profiles developed in the two different hydrate dissociation cases.

330 The one on the left-hand side shows the localized dissociation case ($\Delta r_f = 0.5$ m), whereas the right-

331 hand side one shows the distributed dissociation case ($\Delta r_f = 3$ m). It is noted that the depressurization

332 level is identical between these two cases ($\Delta P_i = -0.3$ MPa), but the pore pressure profiles within the

333 reservoir formation are different, causing different settlement profiles. The values of the maximum

334 subsidence (S_{max}) and the subsidence radius (R_s), which is defined as the radial distance where the

335 curvature of the subsidence distribution becomes maximum, are shown in the figures as circular and

336 square dots, respectively. It is found that the more the hydrate dissociation is localized, the smaller the

337 maximum subsidence and subsidence radius become. The distributed hydrate dissociation case would

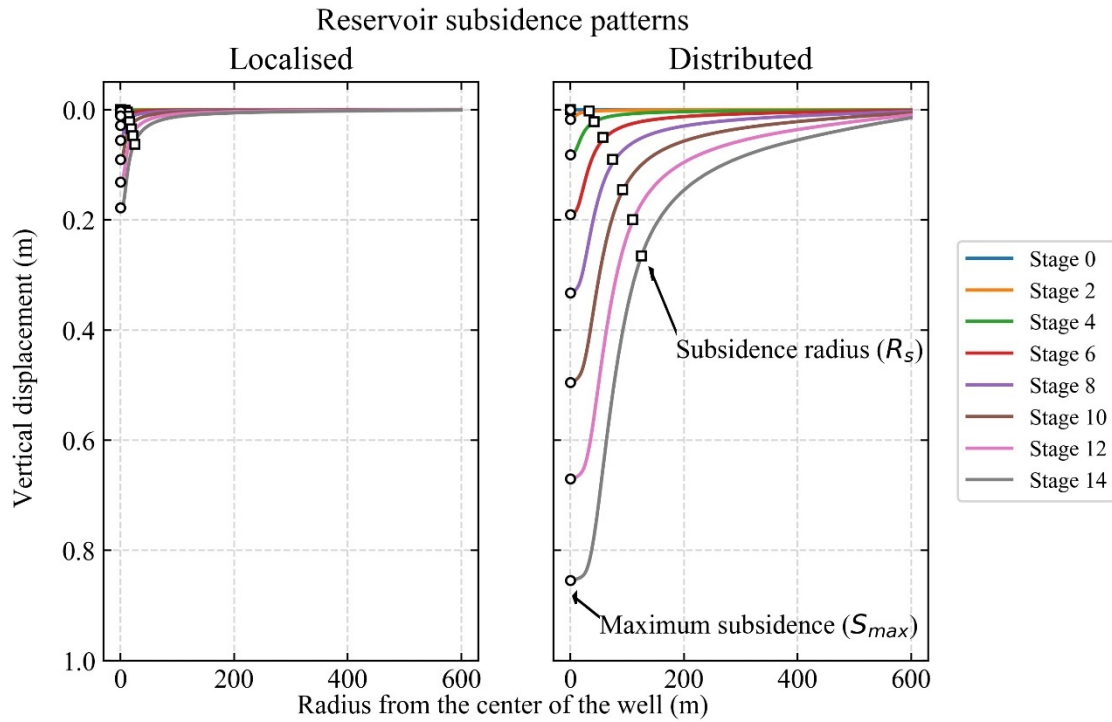
338 be analogous to reservoirs in which hydrate dissociation front advances quickly whereas the localized

339 case the slow progress of hydrate dissociation front. The former may be due to high absolute

340 permeability, low hydrate saturation, etc. and the latter the opposite. These two cases are simulated so

341 that a real pore pressure profile during actual depressurization would fall in between these two extreme

342 cases.



343

344 Fig. 8 Reservoir subsidence along the top of the reservoir layer (localized dissociation case ($\Delta P_i = -0.3$

345 MPa and $\Delta r_f = 0.5$ m) (left) and distributed dissociation case ($\Delta P_i = -0.3$ MPa and $\Delta r_f = 3$ m) (right)).

346

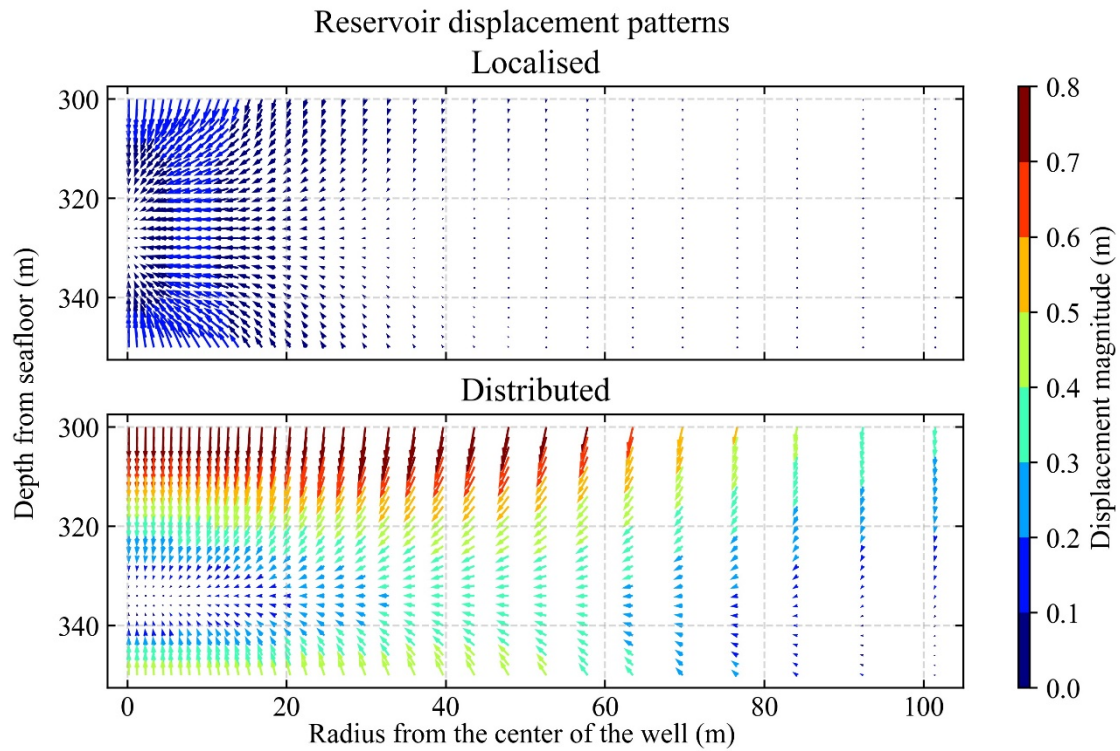
347 Fig. 9 shows the displacement patterns of the reservoir layer at the depressurization stage 14 in the two

348 different hydrate dissociation cases. The magnitudes of the displacement vectors in these figures are

349 normalised and scaled to increase their visibility. The reservoir layer deformation is concentrated near

350 the wellbore in the localised dissociation case, whereas it is more evenly spread radially in the

351 distributed dissociation case.



352

353 Fig. 9 Deformation patterns of the reservoir layer (localised dissociation case ($\Delta P_i = -0.3$ MPa and Δr_f

354 = 0.5 m) (top) and distributed dissociation case ($\Delta P_i = -0.3$ MPa and $\Delta r_f = 3$ m) (bottom)) at the

355 depressurization stage 14. The length of the arrows in each plot is normalized by the maximum

356 magnitude of displacement of respective subsidence cases.

357

358 Fig. 10 shows the overburden layer deformation patterns for the two cases. When the reservoir

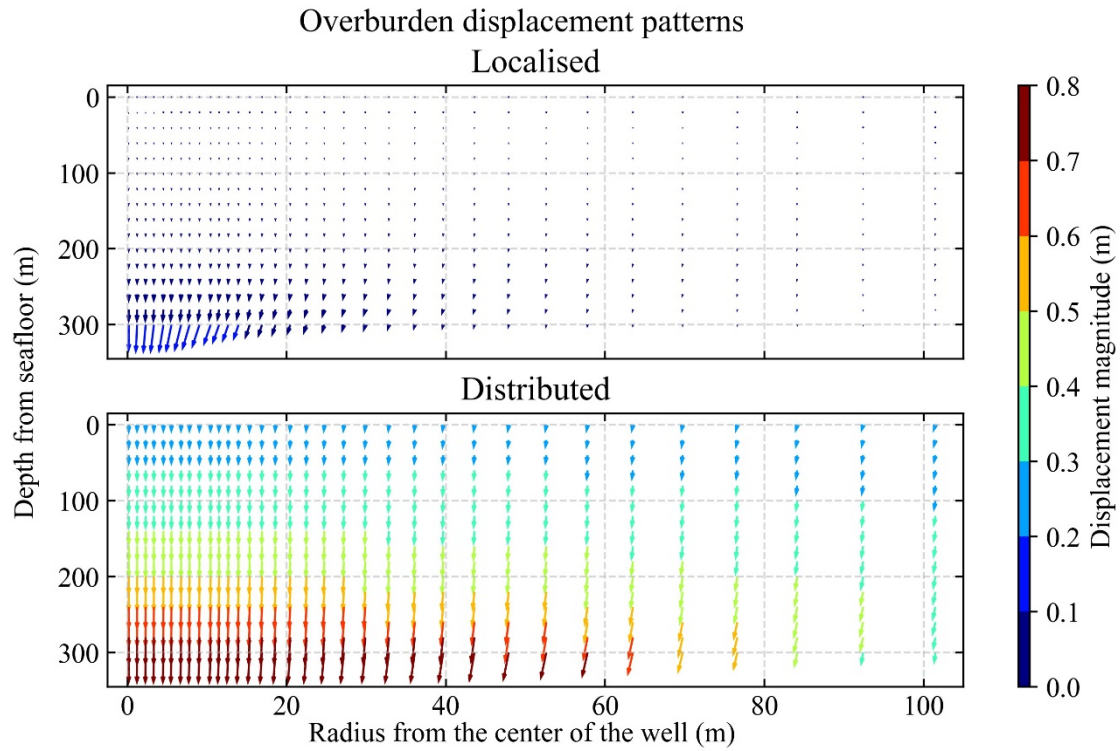
359 deformation is localised near the wellbore (i.e., localised dissociation case), the overburden layer

360 deformation is localised near the lower centre part. When the reservoir deformation is distributed (i.e.,

361 distributed dissociation case), it is more evenly spread in the vertical and horizontal directions. These

362 reservoir/overburden deformation patterns are found to have significant effects on well integrity, which

363 is described in the following sections.



364

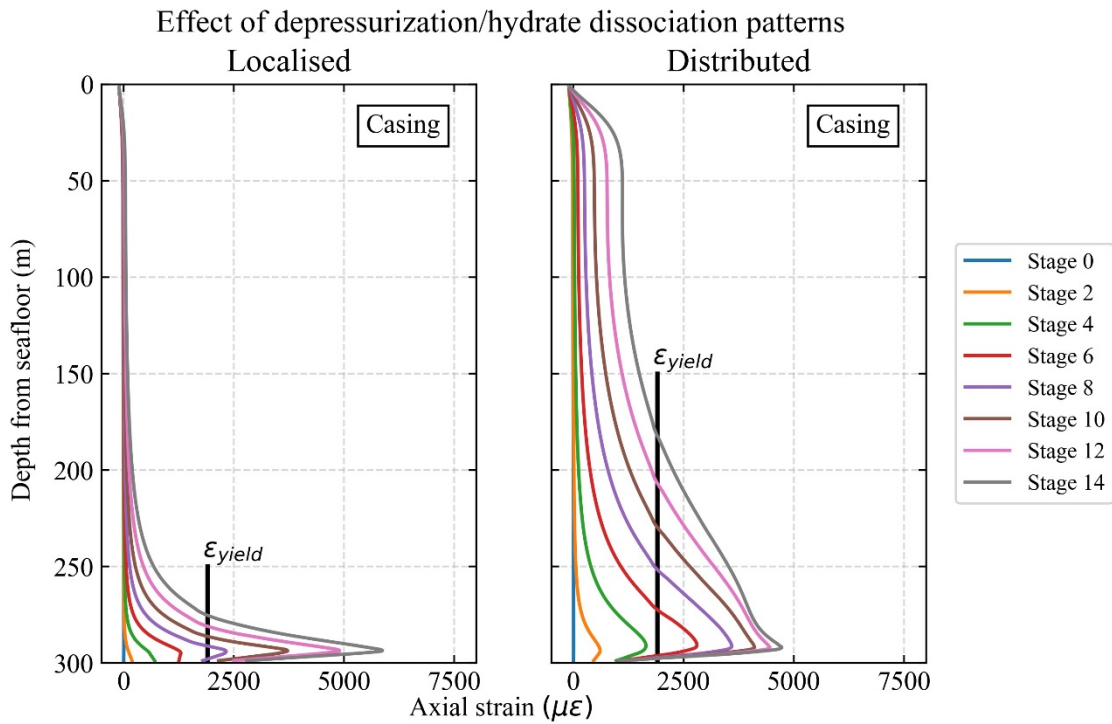
365 Fig. 10 Deformation patterns of the overburden layer (localized dissociation case ($\Delta P_i = -0.3$ MPa and
366 $\Delta r_f = 0.5$ m) (top) and distributed dissociation case ($\Delta P_i = -0.3$ MPa and $\Delta r_f = 3$ m) (bottom)) at the
367 depressurization stage 14.

368

369 3.1.2. Axial strain and stress development in the casing

370 Fig. 11 shows the axial strain development along the casing. In both localised and distributed
371 dissociation cases, the maximum axial strain level in the casing is developed near the bottom of the
372 overburden layer (approximately 290 m), and the value is roughly $5,000 \mu\epsilon$ at the dissociation stage 14.
373 However, the average axial strain level along the depth of the casing is greater in the distributed

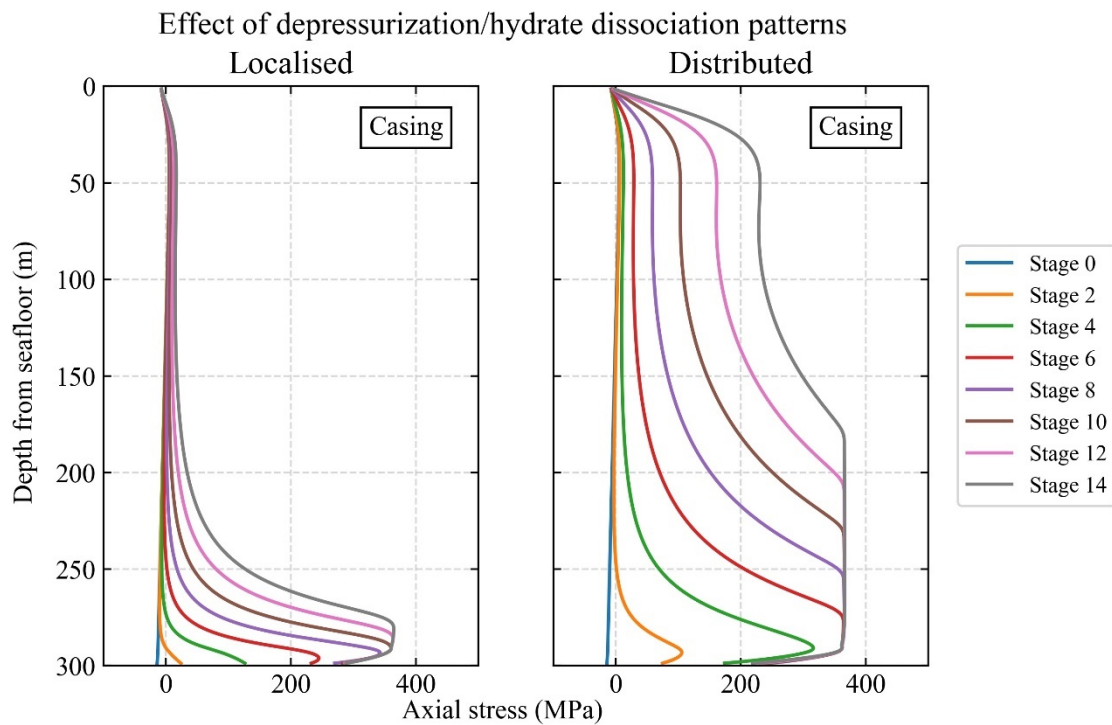
374 dissociation case than in the localised dissociation case. This corresponds to the overburden layer
 375 deformation pattern where the vertical displacement is more evenly spread over the depth of the
 376 overburden layer in the distributed dissociation case than in the localised dissociation case. The vertical
 377 lines in the figure refer to the yield strain level of the casing material, which indicate the zone of yielding
 378 in the casing. The detailed plastic strain profiles are presented in a later section.
 379



380
 381 Fig. 11 Axial strain profiles of the casing (localized dissociation case ($\Delta P_i = -0.3$ MPa and $\Delta r_f = 0.5$
 382 m) (left) and distributed dissociation case ($\Delta P_i = -0.3$ MPa and $\Delta r_f = 3$ m) (right)).

383
 384 Fig. 12 shows the axial stress development along the casing. The initial axial stress levels of the casing
 385 are slightly compressive due to the casing self-weight (i.e., gravitational body force is considered in the

386 simulation) and hydrostatic pressure from the seawater applied at the top of the casing during the well
 387 construction process. As is the case in the axial strain development, the maximum axial stress level is
 388 developed near the bottom of the overburden layer. The difference is that the axial stress level reaches
 389 a plateau once the deviator stress level exceeds the yield stress of the casing (379.5 MPa) and the area
 390 of the plateau extends upward with the progress of depressurization/hydrate dissociation stages. The
 391 area of axial stress plateau indicates the area of plastic strain development, and it covers the depths
 392 between 180 m and 290 m (i.e., 37% of the casing length) at the dissociation stage 14 in the distributed
 393 dissociation case. The average axial stress value of the casing is found to be greater in the distributed
 394 dissociation case than in the localised dissociation case.

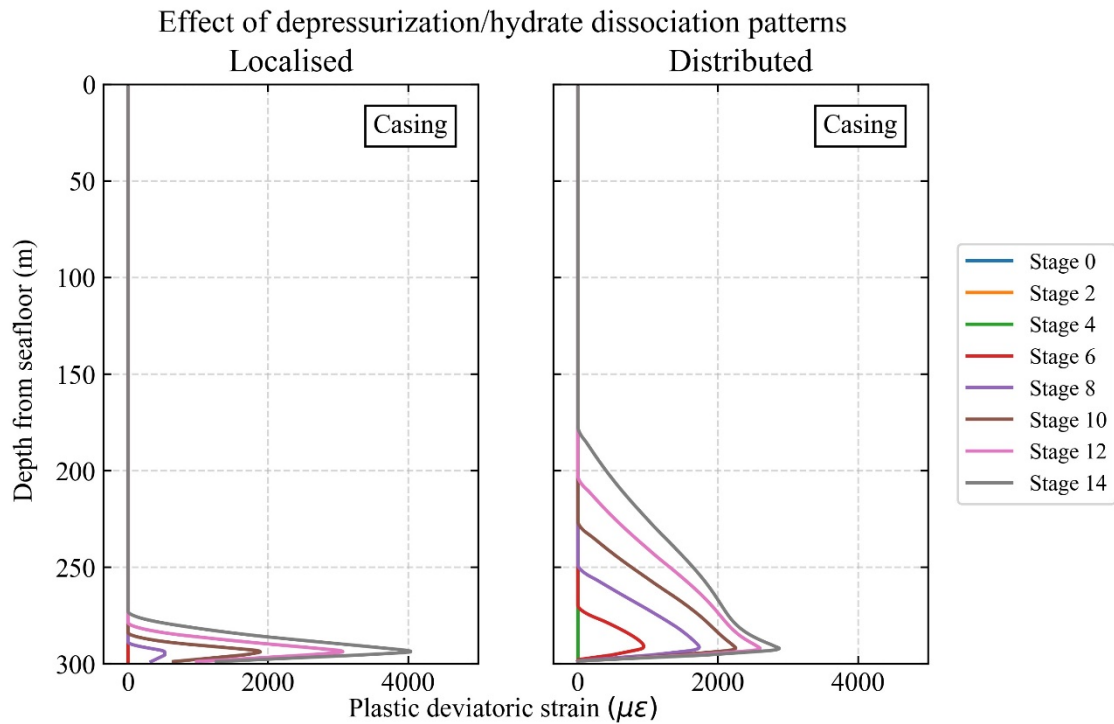


395
 396 Fig. 12 Axial stress profiles of the casing (localized dissociation case ($\Delta P_i = -0.3$ MPa and $\Delta r_f = 0.5$
 397 m) (left) and distributed dissociation case ($\Delta P_i = -0.3$ MPa and $\Delta r_f = 3$ m) (right)).

398

399 Fig. 13 shows the plastic deviatoric strain development of the casing. The area of the plastic strain
400 development is greater in the distributed dissociation case than in the localised dissociation case, and
401 the area of plastic deviatoric strain development corresponds to the area of the axial stress plateau (i.e.,
402 area of yielding) described earlier. The peak value of the plastic deviatoric strain profile is slightly
403 greater in the localised dissociation case at the depressurization stage 14 (4,000 $\mu\epsilon$ vs. 2,900 $\mu\epsilon$) because
404 of the localization of casing yielding in the bottom part of the well (270-300 m).

405



406

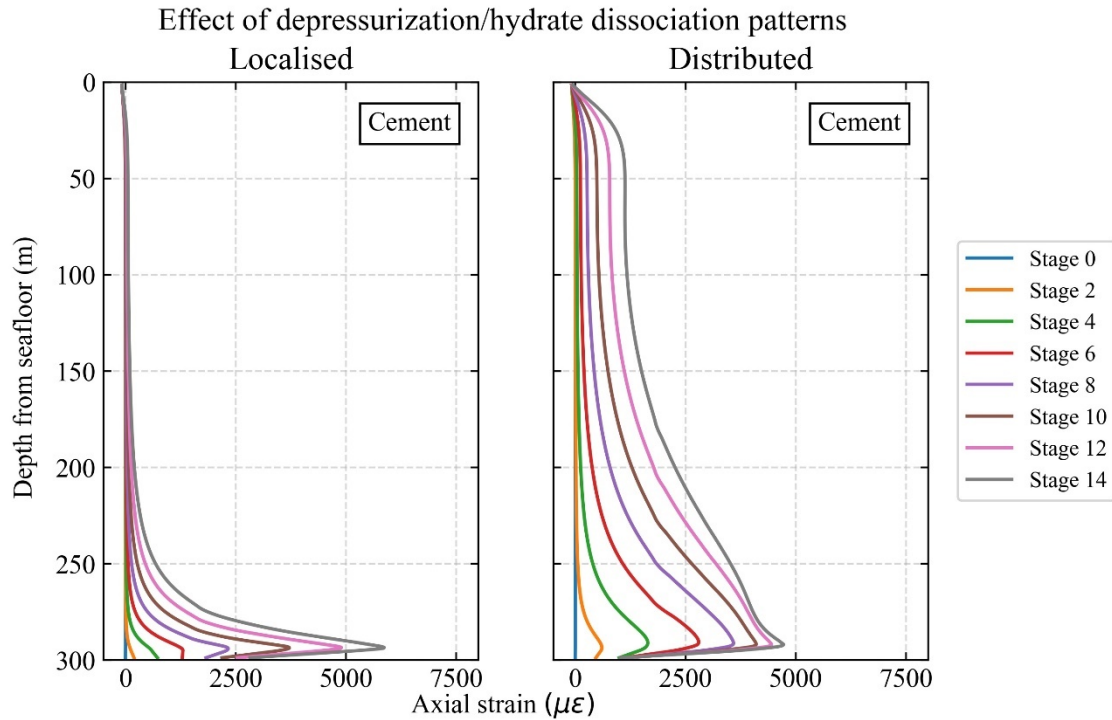
407 Fig. 13 Plastic deviatoric strain profiles of the casing (localized dissociation case ($\Delta P_i = -0.3$ MPa and
408 $\Delta r_f = 0.5$ m) (left) and distributed dissociation case ($\Delta P_i = -0.3$ MPa and $\Delta r_f = 3$ m) (right)).

409

410 3.1.3. Axial strain and stress development in cement

411 Fig. 14 shows the axial strain development of the well cement. They are identical to those of the casing,
412 which indicates that the interface slippage at the casing-cement interface is very limited in the simulated
413 reservoir subsidence cases. This also suggests that the axial strain distribution of the casing could be
414 estimated from that of the cement, which can be measured by strain sensors embedded in the cement.
415 Distributed measurement of the axial strain development of the well (with fibre optic sensing techniques,
416 for example) may be applicable for this purpose. An experimental study on the potential of distributed
417 fibre optic monitoring of well integrity is carried out by the authors in a separate study [55]. It is noted
418 that the small compressive strain (i.e., negative strain values) developed at the top of the well is caused
419 during well construction process (casing landing stage), where the casing is released from hanging and
420 compressed in the upper part of the well.

421



422

423 Fig. 14 Axial strain profiles of the cement (localized dissociation case ($\Delta P_i = -0.3$ MPa and $\Delta r_f = 0.5$

424 m) (left) and distributed dissociation case ($\Delta P_i = -0.3$ MPa and $\Delta r_f = 3$ m) (right)).

425

426 Fig. 15 shows the axial stress development of the cement. The initial axial stress levels of the cement

427 are compressive and change linearly with depths because of the self-weight of the cement and

428 hydrostatic seawater pressure applied at the top of the cement during the well construction process. Due

429 to the smaller stiffness of the cement relative to that of the casing, the axial stress increase (in tension)

430 in the cement is noted to be much smaller than that in the casing. In fact, the axial stress level in the

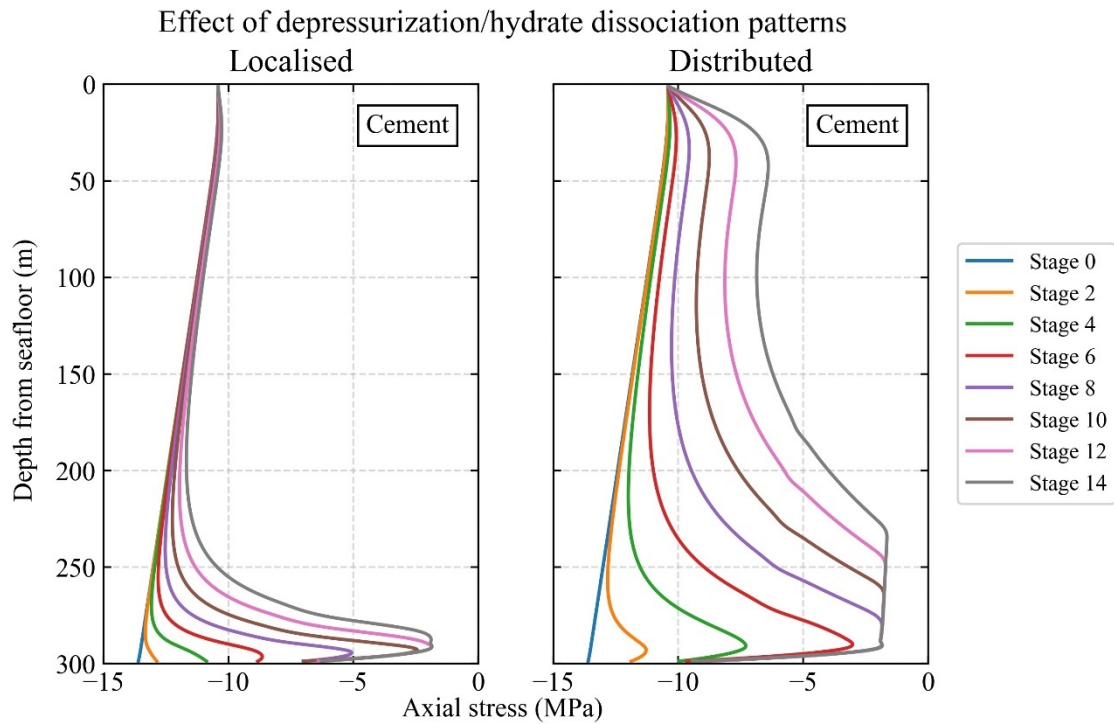
431 cement does not become tensile (i.e., positive values) throughout the simulated depressurization/hydrate

432 dissociation stages. The axial stress plateau is developed in the cement at the bottom part of the

433 overburden layer as well, while it remains in compression (approximately -2 MPa). This is because the

434 stress state in this area has reached the yield stress state governed by the Mohr-Coulomb criteria, which
 435 indicates that the cement fails in shear and not in tension in the simulated depressurization/hydrate
 436 dissociation stages. However, if the depth of the well from the sea surface (which is assumed to be 1,000
 437 m in this study) decreases, the initial compressive axial stress levels in the cement also decrease, which
 438 in turn could lead to the development of tensile failure prior to the development of shear failure. Hence,
 439 the initial stress state corresponding to the depth of the well from the sea surface would be an important
 440 factor in assessing the cement integrity.

441

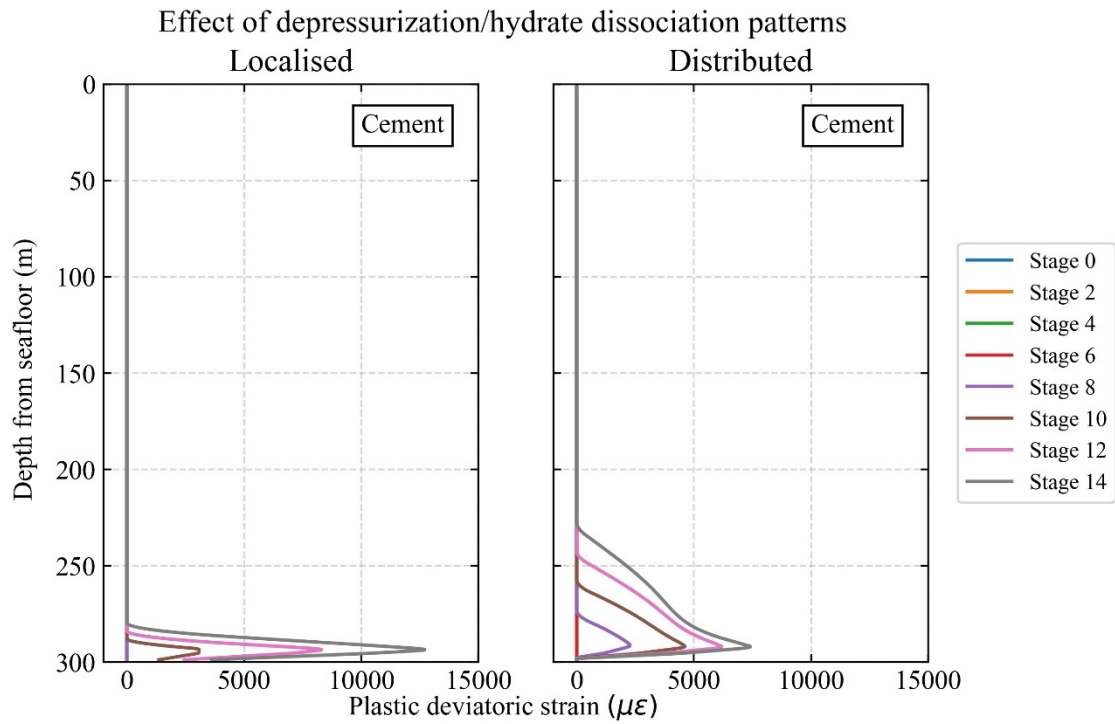


442

443 Fig. 15 Axial stress profiles of the cement (localized dissociation case ($\Delta P_i = -0.3$ MPa and $\Delta r_f = 0.5$
 444 m) (left) and distributed dissociation case ($\Delta P_i = -0.3$ MPa and $\Delta r_f = 3$ m) (right)).

445

446 Fig. 16 shows the plastic strain profiles in the cement, which are qualitatively similar to the ones for the
 447 casing shown earlier. The difference is that the peak value of the plastic deviatoric strain at the
 448 depressurization stage 14 is much greater in the cement than that in the casing (2,900 $\mu\epsilon$ (casing) vs.
 449 7,400 $\mu\epsilon$ (cement) in the distributed dissociation case and 4,000 $\mu\epsilon$ (casing) vs. 13,000 $\mu\epsilon$ (cement) in
 450 the localised dissociation case). This is because the area of yielding in the cement is localised within a
 451 smaller area than that in the casing, which reflects the brittleness of the shear failure of the cement.



452
 453 Fig. 16 Plastic deviatoric strain profiles of the cement (localized dissociation case ($\Delta P_i = -0.3$ MPa
 454 and $\Delta r_f = 0.5$ m) (left) and distributed dissociation case ($\Delta P_i = -0.3$ MPa and $\Delta r_f = 3$ m) (right)).

455
 456 3.2. Effect of cement shrinkage

457 Cement shrinkage occurs due to the capillary pressure development in the cement pores during the

458 cement hydration process. In the Nankai Trough formation case, the cement shrinkage volume could
459 potentially reach 0.75% [30]. Therefore, in this study, the volume of the cement elements is decreased
460 by 0.75% in the cement shrinkage stage to assess its effect on well integrity.

461

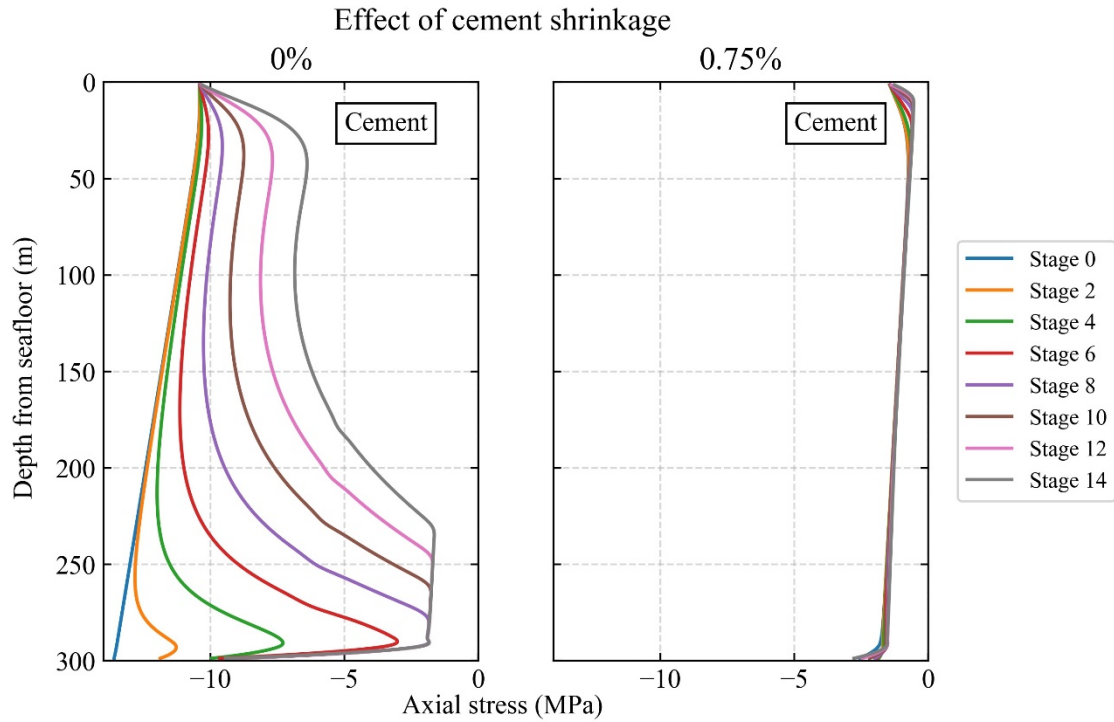
462 Fig. 17 shows the axial stress development of the cement with the cement shrinkage volume of 0% and
463 0.75%. The initial axial stress levels of the cement in the 0.75% shrinkage case are significantly larger
464 (less compressive) than in the 0% shrinkage case. This is because cement shrinkage during the well
465 construction process is simulated under the zero axial displacement condition (i.e., radial cement
466 shrinkage), which results in the decrease of the initial compressive axial stress generated by cement
467 self-weight and hydrostatic seawater pressure. It is noted that the axial stress levels do not become
468 tensile; instead, they reach limiting compressive stress values specified by the Mohr-Coulomb yield
469 surface.

470

471 In the subsequent depressurization stages, the cement is stretched in the axial direction due to reservoir
472 compaction, causing the reduction in the initial compressive axial stress levels. In the 0% shrinkage
473 case, the stress plateau (i.e., Mohr-Coulomb yield surface) is reached in the depressurization stage 8,
474 whereas the cement has already yielded in the depressurization stage 0 in the 0.75% shrinkage case as
475 mentioned earlier, and the axial stress level in the cement remains constant at approximately -2 MPa
476 throughout the subsequent depressurization stages. The negative residual axial stress values show that

477 the plastic deformation of the cement in the 0.75% shrinkage case occurs in shear but not in tension, as
 478 is the case for the 0% shrinkage scenario.

479

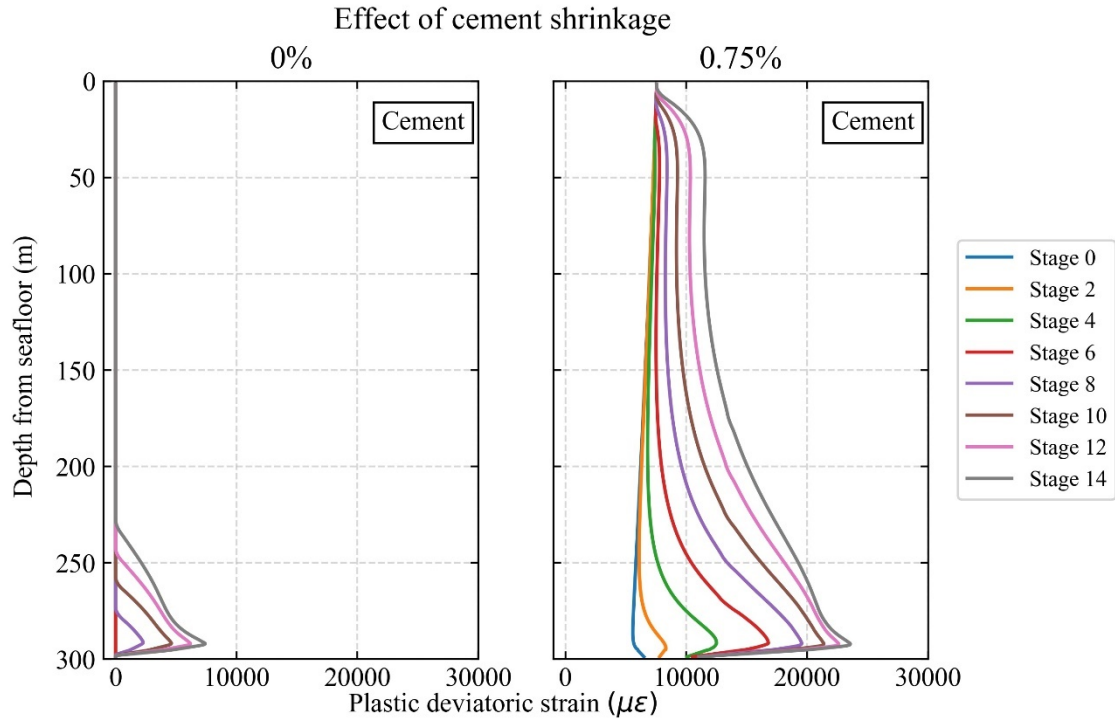


480

481 Fig. 17 Axial stress profiles of the cement with the cement shrinkage volume of 0% and 0.75% ($\Delta P_i =$
 482 -0.3 MPa and $\Delta r_f = 3$ m).

483

484 Fig. 18 shows the plastic deviatoric strain development of the cement. The average plastic strain level
 485 of approximately $6,600 \mu\epsilon$ is already developed in the cement due to the cement shrinkage volume of
 486 0.75%, and it increases with the progress of reservoir subsidence to the maximum value of $24,000 \mu\epsilon$ at
 487 the subsidence stage 14 at the bottom part of the well. This maximum plastic deviatoric strain value is
 488 three times greater than that in the 0% cement shrinkage case ($7,400 \mu\epsilon$ vs. $24,000 \mu\epsilon$).



490

491 Fig. 18 Plastic deviatoric strain profiles of the cement with the cement shrinkage volume of 0% and

492 0.75% ($\Delta P_i = -0.3$ MPa and $\Delta r_f = 3$ m).

493

494 Fig. 19 shows the axial strain development of the casing and cement with the cement shrinkage volume

495 of 0% and 0.75%. It is found that, unlike the axial stress development, the axial strain development of

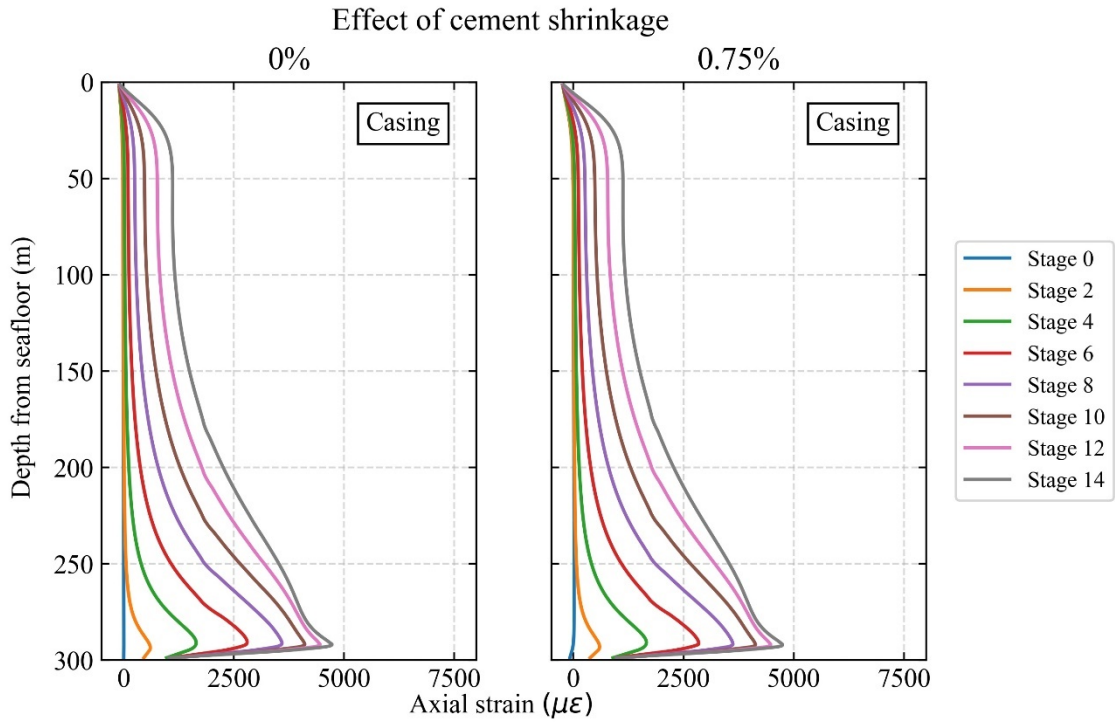
496 the cement is not affected by the cement shrinkage. Also, the axial strain profiles of the casing and

497 cement are found to be identical to each other. These results suggest that the interface slippage is not

498 induced at either the formation-cement or cement-casing interface by the cement shrinkage volume of

499 0.75%.

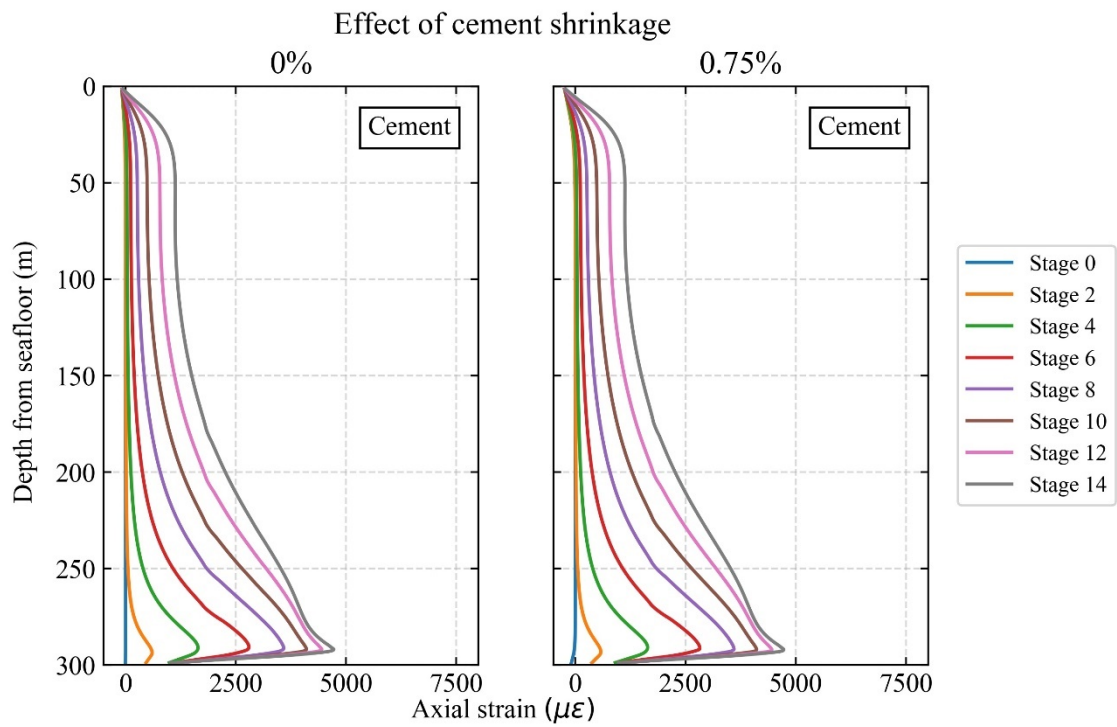
500



501

502

(a)



503

504

(b)

505 Fig. 19 Axial strain profiles of (a) the casing and (b) the cement with the cement shrinkage volume of

506 0% and 0.75% ($\Delta P_i = -0.3$ MPa and $\Delta r_f = 3$ m).

507

508 A back-of-the-envelope calculation is performed with the analytical solution for the cavity
509 expansion/contraction of an elastic cylinder. The decrease in the radial effective stress at the cement-
510 formation interface due to cement shrinkage can be calculated by Equation 3:

511

$$\Delta\sigma'_r = G \left(\left(\frac{r_c}{r_o} \right)^2 - 1 \right) \frac{\Delta V_{cement}}{100} \quad (3)$$

512

513 where $\Delta\sigma'_r$ is the change in the radial effective stress, G is the shear modulus of the formation, r_c is
514 the outer radius of the casing, r_o is the radius of the wellbore and ΔV_{cement} is the volume shrinkage
515 of the cement in percent. Equation 3 is valid for small shrinkage volume ($\Delta V_{cement} \ll 100\%$). The
516 value of shear modulus of the overburden layer at 200 m below the seafloor is approximately 40 MPa
517 and the value of r_c/r_o is 0.7857. By setting the value of ΔV_{cement} to 0.75%, the decrease in the radial
518 effective stress is calculated to be $\Delta\sigma'_r = -0.115$ MPa. The corresponding decrease in the ultimate
519 interface shear stress at the cement-formation interface is $\Delta\tau_{ult} = \mu\Delta\sigma'_r = -0.092$ MPa ($\mu = 0.8$). This
520 decrease in the interface shear resistance is too small to initiate interface slippage. The above discussion
521 is relevant to the cement-formation interface. As to the casing-cement interface, the interface pressure
522 increases (rather than decreases) due to cement shrinkage, which reduces the potential of interface
523 slippage. This is because the cement shrinkage induces inward radial displacement, where the cement

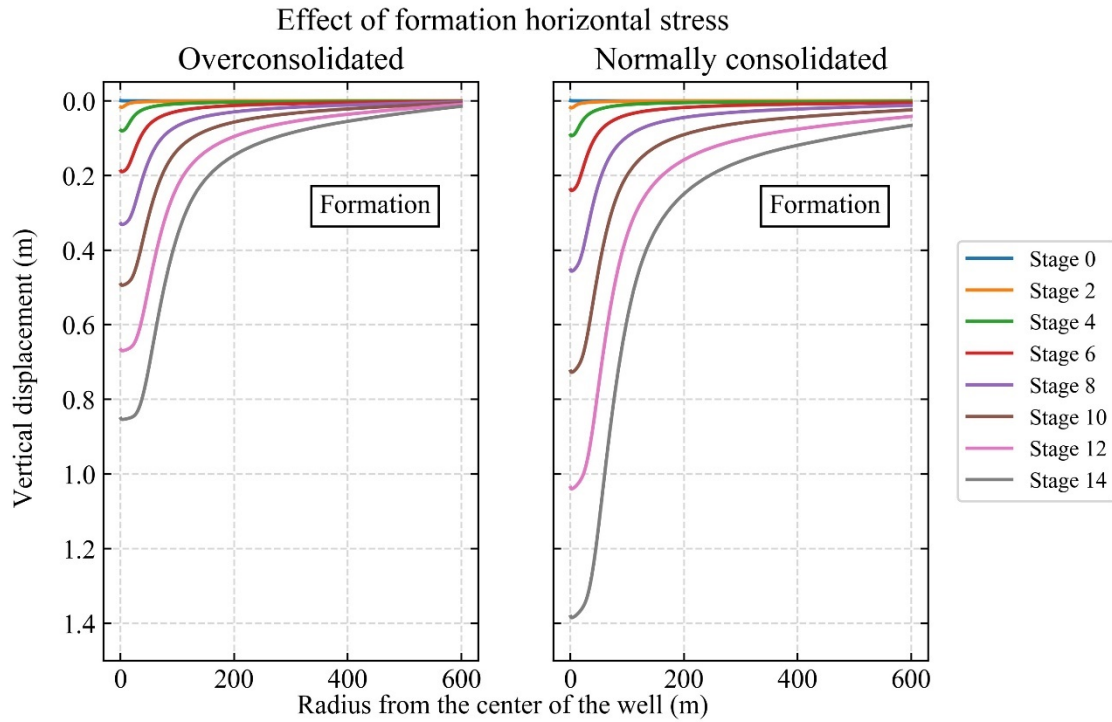
524 tries to separate from the formation but at the same time press against the casing wall. Therefore, the
525 cement shrinkage volume of 0.75% does not affect the shaft friction development at the cement-
526 formation or casing-cement interface, and hence the axial strain development of the casing and cement
527 is not altered by the 0.75% cement shrinkage either.

528

529 3.3. Effect of the initial horizontal stress of the formation

530 The simulation results presented in the earlier sections are computed with an assumption that the
531 overburden clay layer is overconsolidated, which would be reasonable according to the triaxial test
532 results on formation samples recovered at the Nankai Trough [34]. However, the actual stress state of
533 the Nankai Trough formation contains uncertainty due to the fact that the site is located in the subduction
534 zone where the geologic conditions are complex. Also, the formation samples examined in the triaxial
535 tests were found to be significantly disturbed during sampling, which decreases the reliability of the
536 estimation. Therefore, additional simulations for the normally consolidated overburden case are
537 conducted. The reservoir and underburden layers are assumed to be normally consolidated regardless
538 of the simulation cases. The difference in the horizontal stress profiles between the consolidated and
539 normally consolidated overburden cases are shown earlier in Fig. 5. It is noted that the cement shrinkage
540 volume was set to 0% for both cases.

541



542

543 Fig. 20 Reservoir subsidence along the top of the reservoir layer in the overconsolidated and normally

544 consolidated cases ($\Delta P_i = -0.3$ MPa and $\Delta r_f = 3$ m).

545

546 Fig. 20 shows the reservoir subsidence profiles for the overconsolidated and normally consolidated

547 overburden cases. The maximum reservoir subsidence at the depressurization/hydrate dissociation stage

548 14 increases from 0.85 m (overconsolidated case) to 1.4 m (normally consolidated case). This is because

549 the ratio of vertical effective stress to horizontal effective stress (i.e., K_0 value) becomes smaller in the

550 normally consolidated overburden case (0.40) relative to the overconsolidated case (0.44), which results

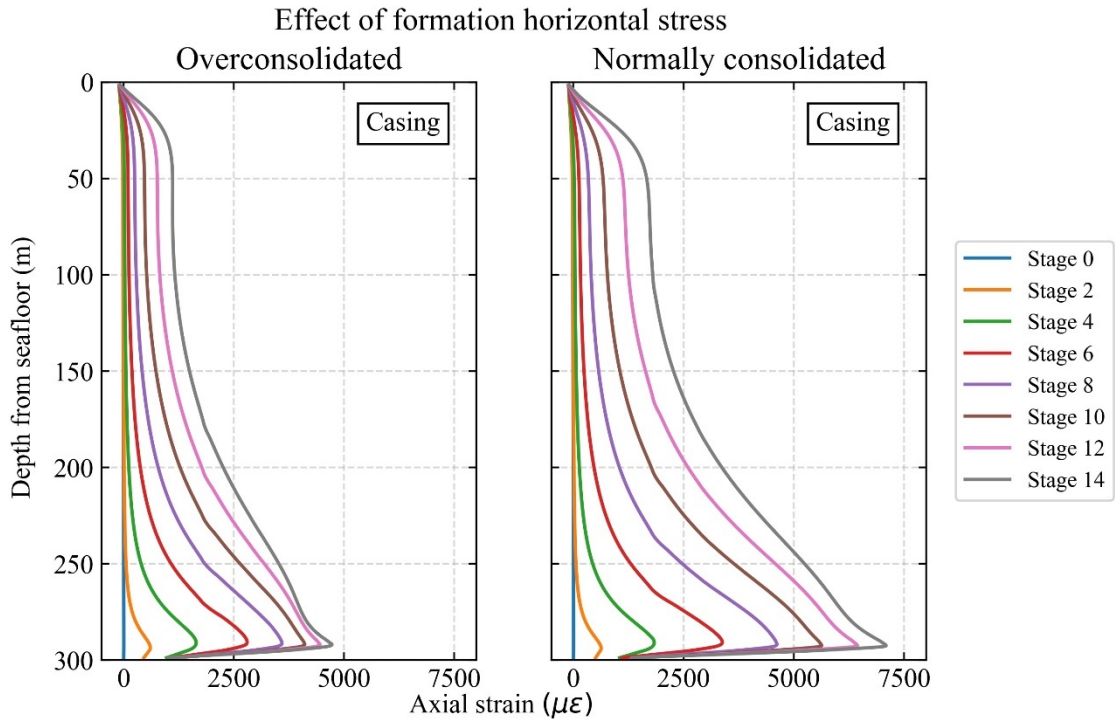
551 in greater magnitudes of volumetric plastic strain (relative to plastic deviatoric strain) computed in the

552 MHCS model.

553

554 Because the reservoir subsidence increases in the normally consolidated overburden case compared to
555 the overconsolidated overburden case, the axial and plastic deviatoric strain development of the casing
556 and cement also become greater. Fig. 21 shows the axial strain development of the casing and cement.
557 The maximum axial strain level in the casing increases from approximately 4,700 $\mu\epsilon$ (overconsolidated
558 overburden case) to 7,100 $\mu\epsilon$ (normally consolidated overburden case), and so does the maximum axial
559 strain level in the cement. However, the axial strain profiles of the casing and cement are still identical
560 in both overconsolidated and normally consolidated overburden cases, indicating that the change in the
561 radial effective stress between the overconsolidated and normally consolidated overburden cases (K_0
562 value change from 0.44 to 0.40) does not induce interface slippage at the casing-cement interface. Fig.
563 22 shows the plastic deviatoric strain development of the casing and cement. It is found that the
564 maximum plastic deviatoric strain levels in the casing and cement (at stage 14) increase from 2,900 $\mu\epsilon$
565 (casing) and 7,400 $\mu\epsilon$ (cement) (overconsolidated overburden case) to 5,300 $\mu\epsilon$ (casing) and 18,000 $\mu\epsilon$
566 (cement) (normally consolidated overburden case), respectively. The area where the casing and cement
567 develop plastic strain increase significantly for the normal consolidated overburden case as well.
568 Therefore, the initial horizontal stress levels of the formation have significant effects on well integrity
569 during reservoir compaction.

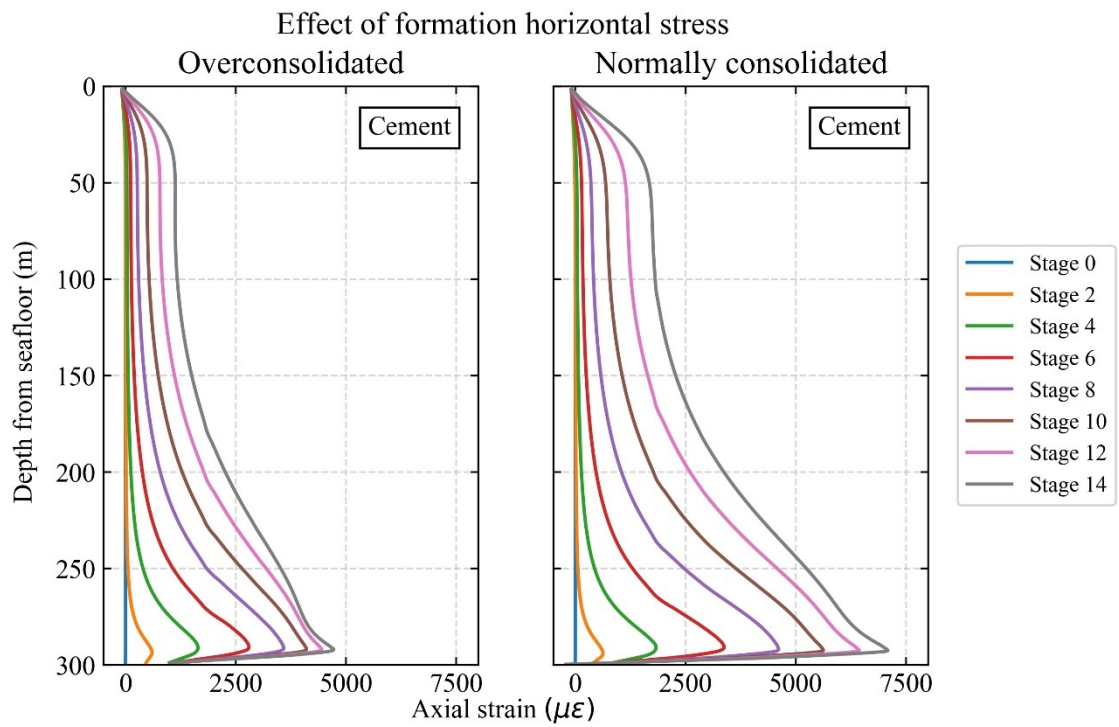
570



571

572

(a)



573

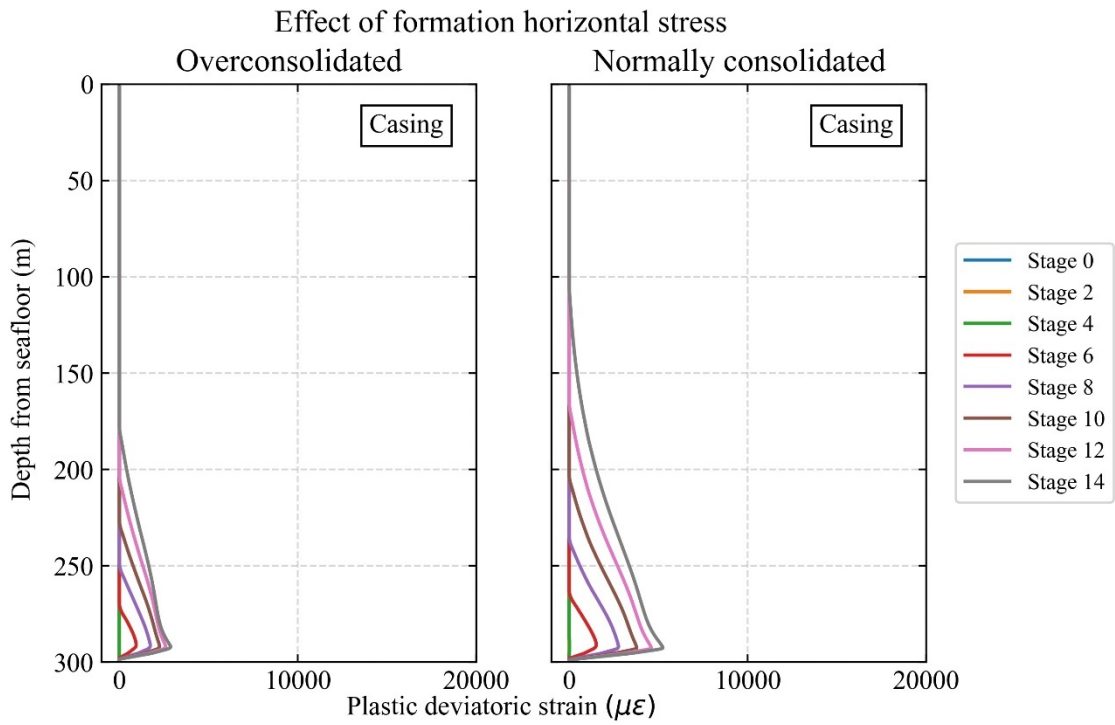
574

(b)

575 Fig. 21 Axial strain profiles of (a) the casing and (b) the cement in the overconsolidated and normally

576

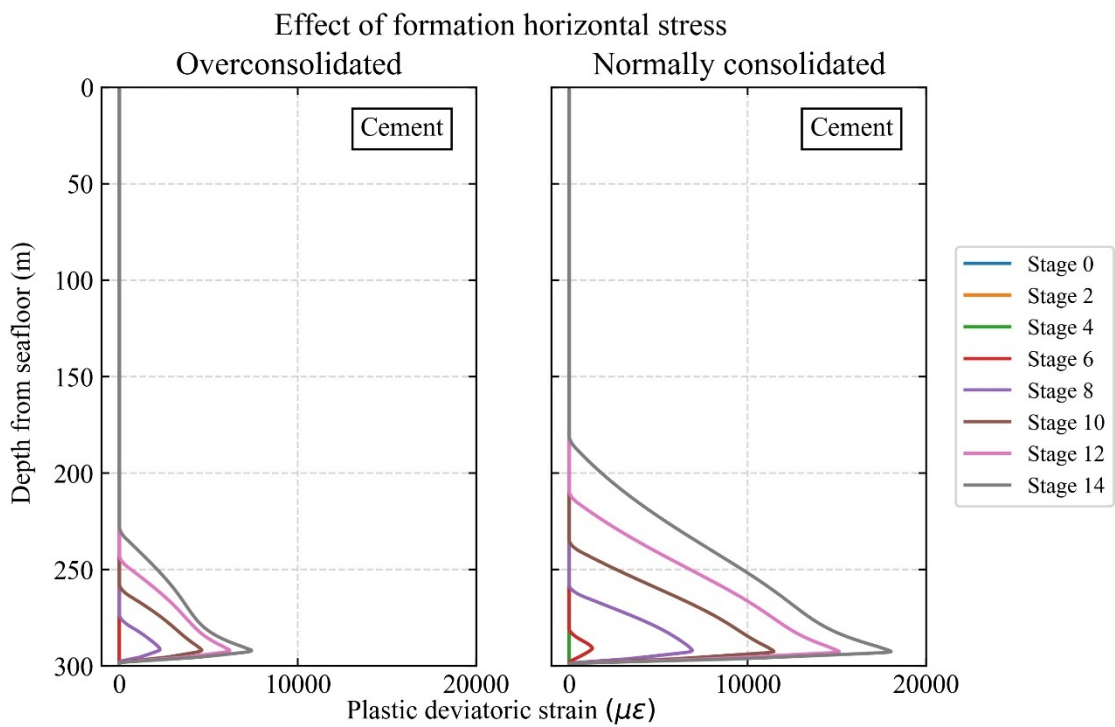
consolidated cases ($\Delta P_i = -0.3$ MPa and $\Delta r_f = 3$ m).



577

578

(a)



579

580

(b)

581 Fig. 22 Plastic deviatoric strain profiles of (a) the casing and (b) the cement in the overconsolidated
582 and normally consolidated cases ($\Delta P_i = -0.3$ MPa and $\Delta r_f = 3$ m).

583

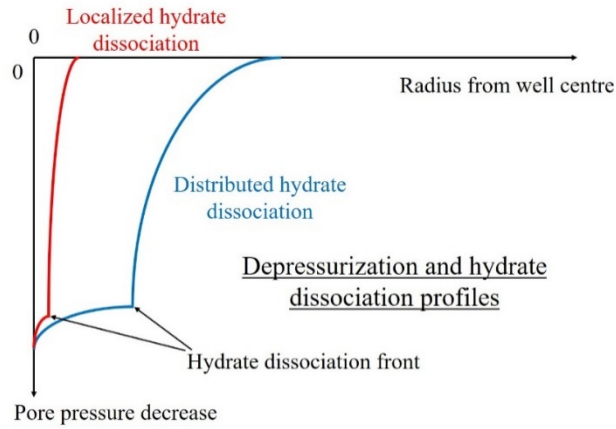
584

585 4. Discussion

586

587 The results shown in the previous section indicate that the pattern of depressurization/hydrate
588 dissociation scenarios (i.e., localized and distributed dissociation cases) influences the distributions of
589 stresses and strains in the casing and cement. Fig. 23 shows a schematic diagram summarizing how
590 different hydrate dissociation patterns (at the same pressure drawdown) result in different reservoir
591 subsidence profiles. In general, the localized dissociation case induces smaller values of maximum
592 reservoir subsidence and subsidence radius than the distributed dissociation case. This is because
593 when the radius of hydrate dissociation front (r_f) is small, pressure drawdown does not propagate afar
594 in the reservoir in the radial direction, resulting in smaller reservoir volume subjected to compaction.

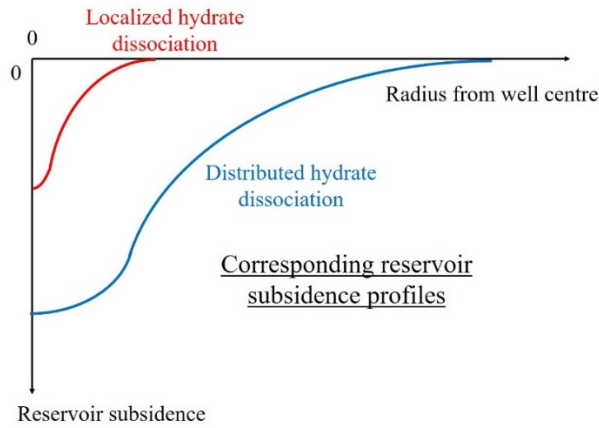
595



596

597

(a)



598

599

(b)

600 Fig. 23 Effect of different depressurization/hydrate dissociation patterns on reservoir subsidence

601 characteristics: (a) input pore pressure profiles; (b) output reservoir subsidence profiles.

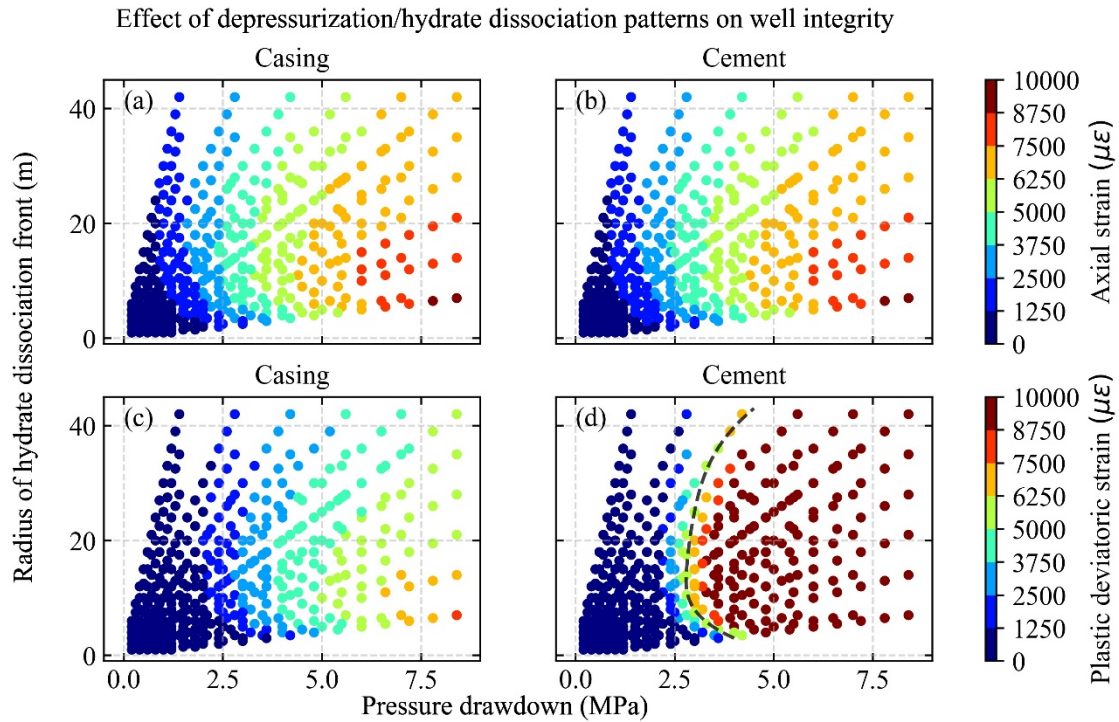
602

603 In this section, contour plots are created to evaluate the maximum axial and plastic deviatoric strains

604 as well as the plasticity length developed in the casing and cement at different values of pressure

605 drawdown ($\sum \Delta P_i$) and radius of hydrate dissociation front ($\sum \Delta r_f$).

606



607

608 Fig. 24 The effect of depressurization/hydrate dissociation patterns on well integrity: (a) casing axial

609 strain; (b) cement axial strain; (c) casing plastic deviatoric strain; (d) cement plastic deviatoric strain.

610

611 Fig. 24 shows the change in the maximum axial and plastic deviatoric strain levels in the casing and

612 cement subjected to different depressurization/hydrate dissociation patterns. A data point on the

613 contour plot is extracted from each of the fourteen depressurization/hydrate dissociation stages in each

614 of the thirty-six simulation cases (i.e., 504 data points in total). Results show that the larger the

615 pressure drawdown and the smaller the radius of hydrate dissociation front are, the greater the

616 maximum axial strain levels in the casing and cement become (Fig. 24a and b). For example, when

617 the radius of hydrate dissociation front is only 5 m as the pressure drawdown of 8 MPa is maintained,

618 the maximum axial strain levels in the casing and cement could both reach 10,000 $\mu\epsilon$ (i.e., 1%). This

619 level of strain does not cause failure in the casing, which is ductile enough to withstand up to several
620 tens of percent strain, but potentially not in the cement which is a much more brittle material than the
621 casing. Fig. 24c and d show the maximum plastic deviatoric strain levels in the casing and cement.
622 Plastic deviatoric strain levels change gradually in the casing with the change of pressure drawdown
623 and radius of hydrate dissociation. On the other hand, large values of plastic deviatoric strain develop
624 rapidly in the cement when the pressure drawdown and radius of hydrate dissociation front exceed the
625 values along the dashed line shown in Fig. 24d. Hence, the plastic strain level increases significantly.
626 This indicates that the cement failure could be highly localized in the form of a shear band.

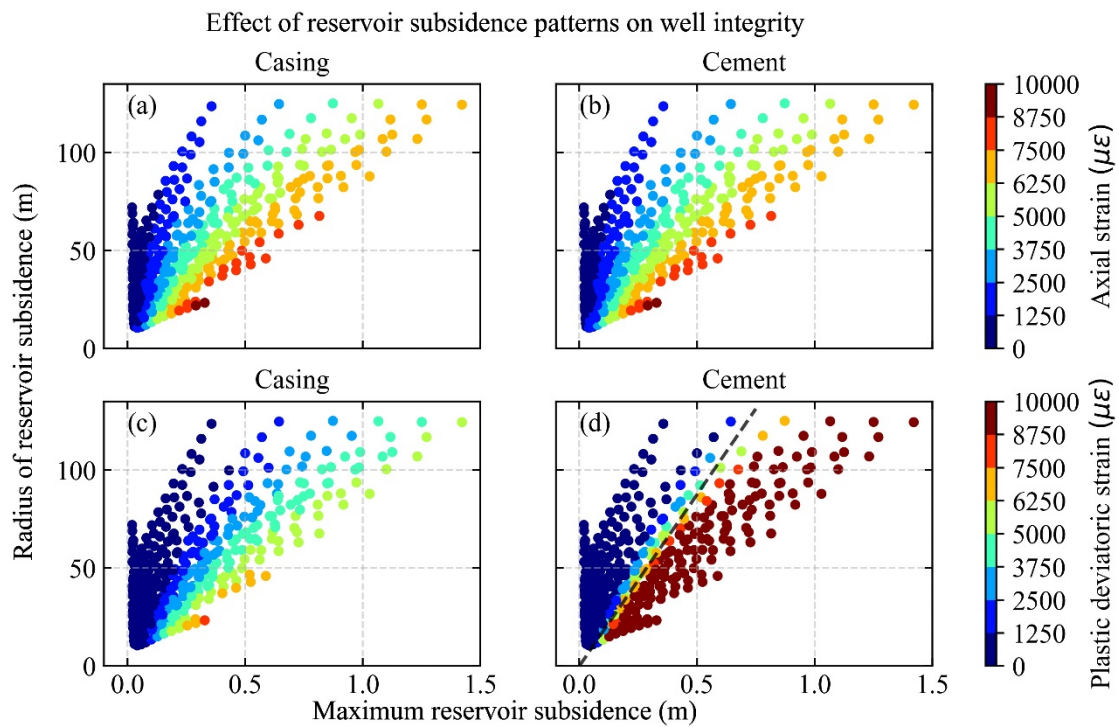
627

628 These results suggest that, in order to avoid the development of large axial strain levels in the casing
629 and cement, the pressure drawdown may have to be kept at a low level until the radius of hydrate
630 dissociation front increases above a certain value. For instance, when the axial strain level of 10,000
631 $\mu\epsilon$ needs to be avoided, the pressure drawdown would have to be temporarily held at 6 MPa until the
632 radius of hydrate dissociation front reaches 25 m.

633

634 It is noted that the simulated pressure drawdown and the radius of hydrate dissociation front are
635 assumed to increase simultaneously and linearly with time in this study. This may not be realistic
636 considering that it is usual to perform a rapid pressure drawdown in the field practice, which would
637 not cause a noticeable increase in the radius of hydrate dissociation front. The pressure drawdown will

638 be maintained while the radius of hydrate dissociation front increases to produce gas from the hydrate
 639 reservoir. Hence, the effect of the path of pressure drawdown and changes in the radius of hydrate
 640 dissociation front has to be investigated extensively considering the stress/strain development of the
 641 well.
 642



643
 644 Fig. 25 The effect of reservoir subsidence characteristics on well integrity: (a) casing axial strain; (b)
 645 cement axial strain; (c) casing plastic deviatoric strain; (d) cement plastic deviatoric strain.

646

647

648 Fig. 25 shows the change in the maximum axial and plastic deviatoric strain levels in the casing and

649 cement with changing reservoir subsidence characteristics. The results show that the larger the

650 maximum reservoir subsidence and the smaller the subsidence radius are (i.e., the more the reservoir
651 compaction is localised), the greater the maximum axial strain levels in the casing and cement become
652 (Fig. 25a and b). As to the plastic strain development, the maximum plastic deviatoric strain level in
653 the casing changes gradually with changes in the reservoir subsidence characteristics. On the other
654 hand, large plastic deviatoric strain levels develop abruptly in the cement. In this study, the two
655 distinct areas of the damaged (below the line) and undamaged (above the line) cement are identified
656 in Fig. 25d. The line of separation can be approximated by Equation 4:

$$R_s = 175 S_{max} \quad (4)$$

658
659 where R_s is the radius of formation subsidence and S_{max} is the maximum formation subsidence.

660
661 Although the proposed line separates the damaged and undamaged cement areas clearly, the position
662 and shape of the line could be affected by the initial hydrate distribution in the reservoir, which is
663 highly heterogeneous in the field. In this study, the effect of hydrate saturation on reservoir
664 compaction is not considered (i.e., hydrate saturation values in the reservoir are uniformly set to zero).
665 This assumption may be acceptable because it is assumed in the soil model that hydrate saturation has
666 negligible effects on the compressibility of the hydrate-bearing soil. However, the shear resistance
667 will be enhanced by the presence of hydrate, which helps the reservoir resist inward displacement

668 during depressurization/hydrate dissociation through the cavity contraction mechanism. Therefore, the
669 effect of hydrate saturation on the characteristics of reservoir subsidence can be examined in a future
670 study by conducting a fully coupled thermo-hydromechanical simulation that computes deformations
671 in the reservoir with complex hydrate saturation profiles. Also, a sensitivity analysis of the MHCS
672 model parameters may be conducted to investigate their effect on the risk plots. This will allow
673 incorporating the uncertainty of parameter values as they were determined through trial-and-error
674 curve fitting against laboratory experiment data on formation samples without proper optimization.

675

676

677 5. Conclusions

678

679 In this study, a parametric study of well integrity under different reservoir subsidence patterns for the
680 case of the Nankai Trough methane hydrate reservoir is carried out by conducting a series of finite
681 element simulations. The well construction processes are incorporated prior to the reservoir
682 subsidence stages to investigate the effect of cement shrinkage on well integrity during reservoir
683 compaction. Also, the effect of the initial horizontal stress profile of the formation (i.e.,
684 overconsolidated and normally consolidated overburden cases) on well integrity is assessed. The
685 model parameters for the simulations (soil, cement, casing and the interfaces) are calibrated against
686 relevant laboratory test data. The primary findings of this study are presented below:

687

688 (i) Various reservoir subsidence profiles are simulated for the Nankai Trough case to examine the
689 scenarios when the maximum reservoir subsidence and the radius of formation subsidence
690 vary between 0.01 m and 1.42 m, and 10.5 m and 125 m, respectively. These subsidence
691 profiles correspond to the pressure drawdown of between 0.1 MPa and 8 MPa, and hydrate
692 dissociation front radius of between 0.5 m and 42 m. The largest maximum axial strain levels
693 developed in the casing and cement are both 9,500 $\mu\epsilon$, and the largest plastic deviatoric strain
694 levels are 7,700 $\mu\epsilon$ (casing) and 29,000 $\mu\epsilon$ (cement). With these levels of strains, the casing
695 would still be far from failure (which would require $\sim 30\%$ strain), and the plasticity gradually
696 spread in a region with depths of approximately 100 m to 300m. On the other hand, localised
697 failures such as shear band may develop in the cement.

698

699 (ii) A large pressure drawdown combined with a small radius of hydrate dissociation front, which
700 corresponds to higher ratios of reservoir subsidence to the lateral extent of subsidence, is
701 found to induce the largest levels of axial and plastic deviatoric strain in the casing and
702 cement. Therefore, such a ratio could be used to predict cement damage (e.g., cement could
703 be damaged if the ratio exceeds 175). These results indicate that the well integrity would be
704 the most vulnerable in the initial stages of hydrate dissociation after rapid depressurization. In
705 order to prevent well failure, the pressure drawdown may need to be kept at a low level (e.g.,

706 several MPa) until hydrate dissociation front advances to a certain radius (e.g., a couple of
707 tens of metres).

708

709 (iii) The effect of cement shrinkage during wellbore construction on wellbore stability was
710 examined for the two scenarios. Cement shrinkage volume of 0.75% is found to develop
711 approximately 6,600 $\mu\epsilon$ plastic deviatoric strain in the cement prior to reservoir subsidence,
712 and it increases to the maximum value of 24,000 $\mu\epsilon$ by the time reservoir subsidence reaches
713 0.85 m. Compared to the 0% shrinkage case, the maximum plastic deviatoric strain increases
714 by more than 200% (7,400 $\mu\epsilon$ vs. 24,000 $\mu\epsilon$) due to the cement shrinkage of 0.75%. The
715 effect of initial in-situ horizontal stress on wellbore stability was also examined. The slight
716 decrease in the initial horizontal stress levels of the formation ($K_\theta = 0.44$ vs. 0.40) was found
717 to increase the maximum plastic deviatoric strain level in the cement by more than 100%.
718 Results suggest that the underestimation of cement shrinkage and overestimation of formation
719 horizontal stress could have contributed to well failure at the Nankai Trough site.

720

721 The developed well integrity contour plots, as shown in Fig. 24 and Fig. 25, will be useful for evaluating
722 the risk of casing and cement damage during gas production from methane hydrate reservoirs, provided
723 that either hydrate dissociation front or maximum reservoir subsidence and subsidence radius data are
724 available. Coupled thermo-hydromechanical simulations of hydrate dissociation-induced reservoir

725 compaction will not only provide such data (which are difficult to obtain through field measurements)
726 but could also be used to update the risk plots. It can incorporate the effect of highly heterogeneous
727 distributions of hydrate saturation on the development of pore pressure (depressurization) and
728 subsidence profiles in the reservoir layer during gas production. Therefore, coupled thermo-
729 hydromechanical simulations for the well integrity analysis needs to be conducted in future studies.

730

731

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739

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