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THE ϕ/ω AND η/η' PRODUCTION RATIOS IN THE REACTION $\pi^+p\to\Delta^{++}$ + BOSON AT 3.7 GeV/c

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J. N. MacNaughton, and G. H. Trilling

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To be presented at the Lund International Conference on Elementary Particles, June 25-July 1, 1969.

THE
$$\phi/\omega$$
 AND η/η ! PRODUCTION RATIOS IN THE REACTION $\pi^+ p \to \Delta^{++}$ + BOSON AT 3.7 GeV/c*

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I. THE φ/ω RATIO

It is well known that the reaction $\pi p \to \Delta p$ is suppressed. In a new high-statistics experiment with $\pi^+ p$ at 3.7 GeV/c we are getting data which indicates just how enormously the reaction

$$\pi^+ p \rightarrow \Delta^{++} \phi$$
 (2⁺³₋₂ events) (1)

is suppressed relative to the reaction

$$\pi^+ p \rightarrow \Delta^{++} \omega$$
 (1324±100 events) . (2)

Here the center-of-mass momenta are $p_{\phi}=0.83$ GeV/c and $p_{\omega}=0.96$ GeV/c so that phase space corrections are small. We thus obtain a cross section ratio of:

$$\rho = \frac{\sigma(\pi^{+}p \to \Delta^{++}\phi)}{\sigma(\pi^{+}p \to \Delta^{++}\omega)} = \frac{1}{662} = 0.0015^{+0.0023}_{-0.0015}.$$

It is interesting to note that this ratio is in good agreement with the squared coupling constant ratio $g^2\phi\rho\pi/g^2\omega\rho\pi\approx 1/600$ determined collectively from the experimental ω and ϕ widths, the $\phi\to\rho\pi$ branching fraction, and the ratio between the phase spaces available for $\phi\to\rho\pi$ and $\omega\to\rho\pi$.

In our experiment we have analyzed about 60,000 four-prong events from an exposure of 180,000 pictures taken in the LRL 72-inch hydrogen bubble

chamber in a 3.7-GeV/c $\pi^{\dagger}p$ beam at the Bevatron. These yielded 15,066 events fitted and identified as

and 353 events fitted and identified unambiguously as

$$\pi^{\dagger} p \rightarrow K^{\dagger} K^{-} \pi^{\dagger} p \qquad (4)$$

In Fig. 1 we show a triangle plot of $M(K^+K^-)$ vs $M(p\pi^+)$. In Fig. 2 we show the projection on the $M(p\pi^+)$ axis. The K^+K^- spectrum, shown in Fig. 3, indicates clearly that we observe ϕ production (~ 18 events). However, there is no indication (see Fig. 4) that ϕ production occurs in association with Δ^{++} . Furthermore we note that ϕ production occurs at comparatively large t values (see Fig. 5), a result similar to that observed for π^-p interactions. We have examined the $\pi^+\phi$ and $p\phi$ mass distributions to find out whether the ϕ production here observed is associated with any higher mass meson or baryon system. No hint of any such effect is observed within our very limited ϕ data, which is consistent with uniform mass distributions. Thus the reaction $\pi^+p\to\pi^+\phi p$ appears to a nonperipheral process leading to three "particles" in the final state.

A. Mixing Angle Determination on a Quark Model

In the framework of the quark model the smallness of ρ can be interpreted in terms of the quark content of the ϕ and ω mesons; namely, they are almost entirely $\lambda\overline{\lambda}$ and $n\overline{n}$ + $p\overline{p}$ respectively. Alexander et al. have related R to the ωp mixing angle θ_{1} by:

$$\pm |R| = \frac{\langle \pi^{+} p/\phi \Delta^{++} \rangle}{\langle \pi^{+} p/\omega \Delta^{++} \rangle} = \frac{\cos \theta_{1} - \sqrt{2} \sin \theta_{1}}{\sin \theta_{1} + \sqrt{2} \cos \theta_{1}} = \tan (\theta_{0} - \theta_{1})$$

where θ_0 = arctan $(1/\sqrt{2})$ = 35.26° and $|R| = \sqrt{\rho}$. This relation gives a model-dependent determination of the mixing angle, but is independent of meson masses in the nonet. We find $|\theta_0 - \theta_1| = 2.2^{\circ} + 1.3^{\circ}$. In principle this method could yield information on whether the quadratic or linear form of the Gell-Mann/Okubo mass formula is appropriate for the bosons. 5,6 In practice however the values are sufficiently close; namely, $|\theta_1|$ (quadratic) = 39.9°±1.1° and $|\theta_1|$ (linear) = 37.1°±1.1°, that our data cannot distinguish between them. Our results are: $\theta_1(R = +\sqrt{\rho}) = 33.1^{\circ} + \frac{2.2^{\circ}}{-1.3^{\circ}}$ or $\theta_1(R = -\sqrt{\rho}) = 37.5^{\circ} + \frac{1.3^{\circ}}{-2.2^{\circ}}$. This result is displayed graphically in Fig. 6.

II. THE η/η RATIO

In this same experiment we have observed 139±12 events ascribed to 7

and 39±10 events for

where we used a "narrow" Δ^{++} definition $M(p\pi^{+}) = 1160-1280 \text{ MeV}$.

Using the Particle Data Tables, ⁸ we find the 139 observed η events correspond to 594 ± 58 η events where all decay modes are taken into account. Using the data of Rittenberg, ⁹ we find our observed η ' events correspond to 127 ± 33 η ' events where again all decay modes are taken into account. These two numbers give $\rho=4.7\pm1.3$, where ρ is the ratio of $\Delta^{++}\eta$ events to $\Delta^{++}\eta$ ' events.

Again, we have used the formula of Alexander et al. to evaluate the η/η^{τ} mixing angle. To make this comparison we multiply the above number of $\Delta^{++}\eta$

events by a factor such that the result corresponds to the same outgoing center-of-mass momentum in both cases, namely $p_{cm}=0.853~\text{GeV/c}$. The magnitude of this factor is obtained from the shape of the $\Delta^{++}\eta$ cross section as determined by Brown et al. ¹⁰ Inclusion of the appropriate phase space and flux factors with this factor gives an overall correction factor of 0.93 to be applied to ρ . Thus we find from $R=\pm\sqrt{0.93}~\rho=\tan(\theta_0-\theta)$ that $\theta(R=+\sqrt{\rho})=-29.0^{\circ}\pm3.3^{\circ}$ and $\theta(R=-\sqrt{\rho})=-80.4^{\circ}\pm3.3^{\circ}$. In this case the linear Gell-Mann/Okubo formula gives $|\theta|=23.7^{\circ}\pm0.3^{\circ}$ and the quadratic formula gives $|\theta|=10.4^{\circ}\pm0.2^{\circ}$. Thus our result is in agreement with the linear GMO formula, which is the conclusion obtained by Benson et al. ⁶

Of course, this conclusion depends on the validity of the quark model⁴ and on the identification of the $\eta^*(958)$ as the ninth member of the pseudoscalar nonet rather than some other candidate such as the E(1420).

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TRIANGLE PLOT W/ERRORS

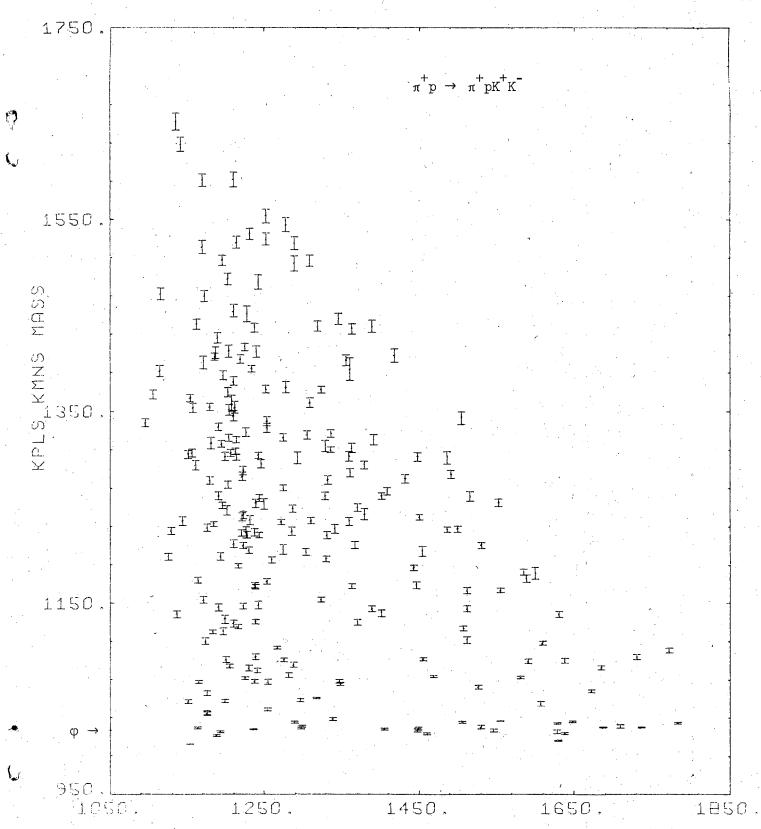


Fig. 1

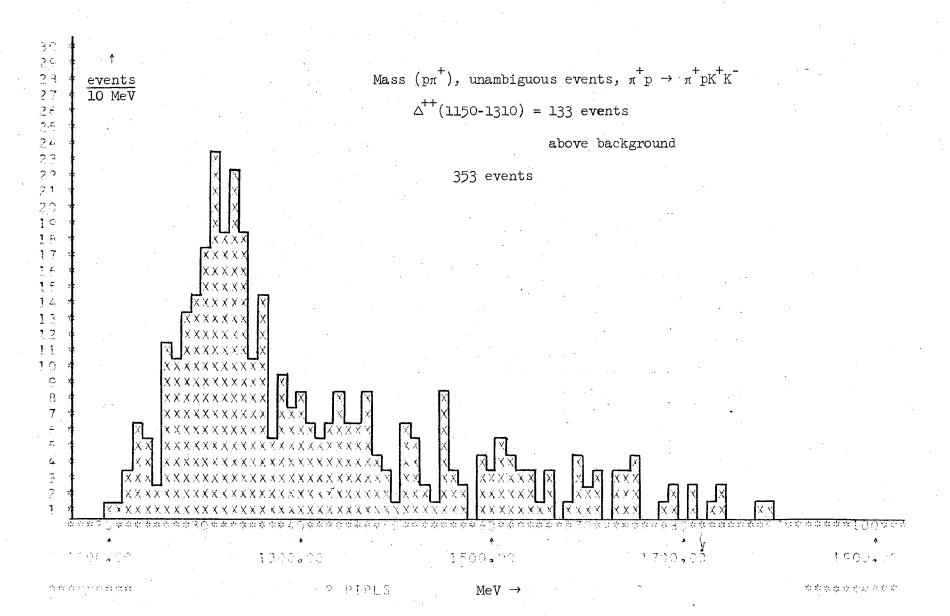


Fig. 2

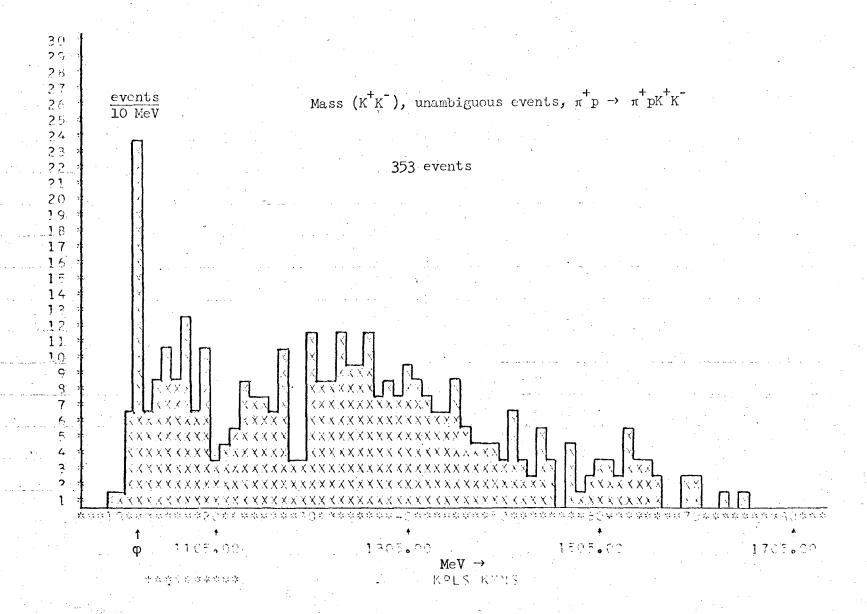


Fig. 3

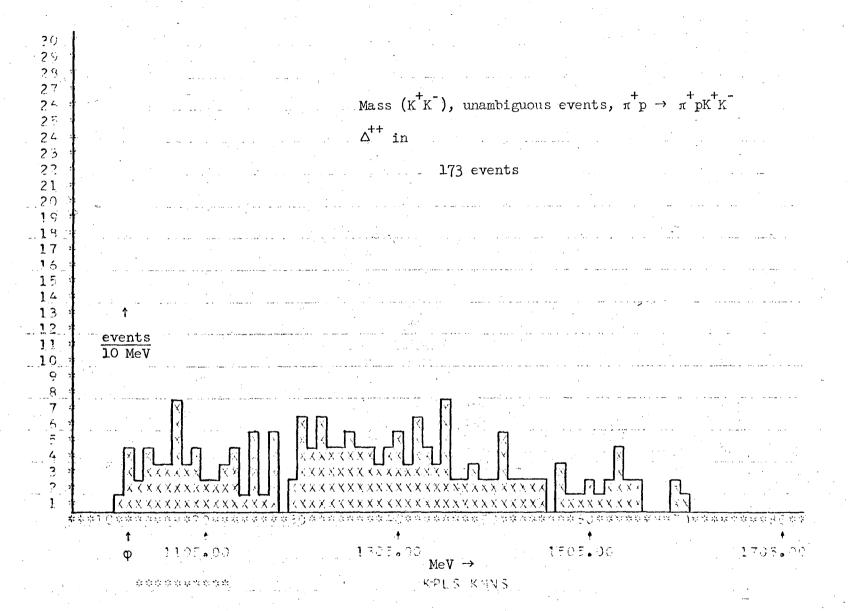
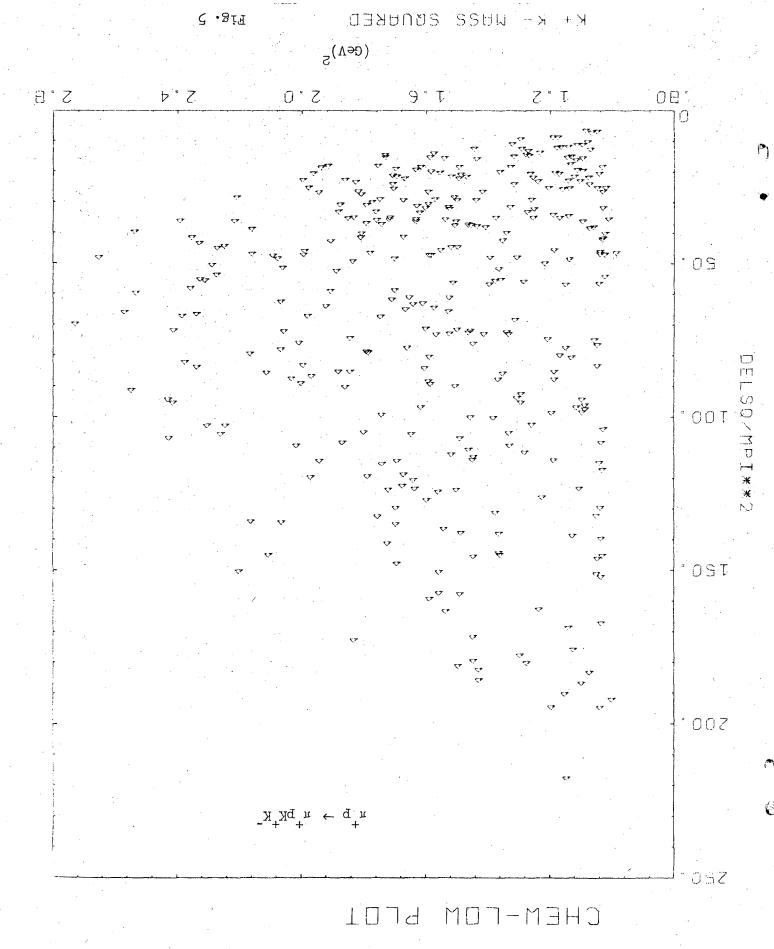
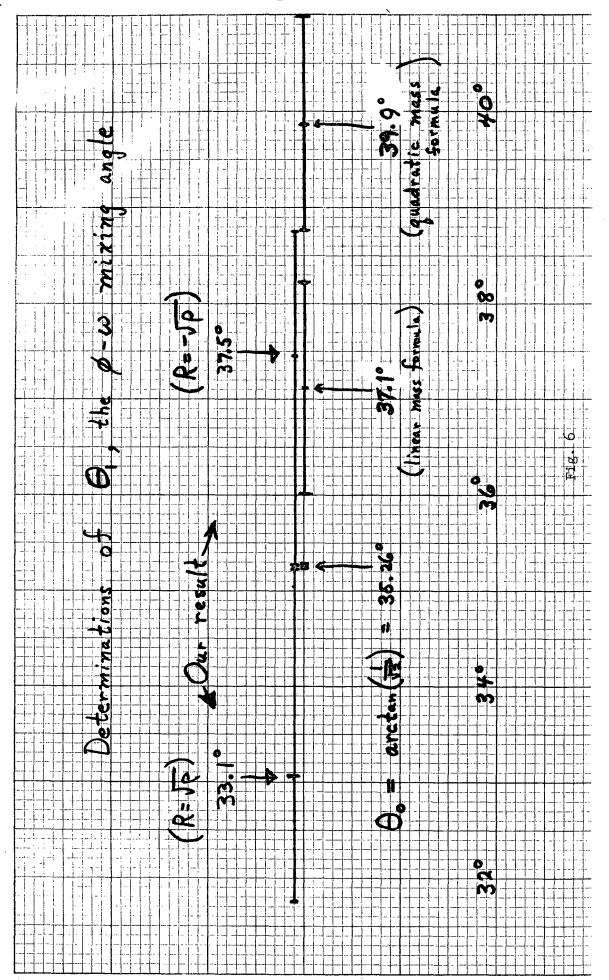


Fig. 4





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