

# UC Santa Cruz

## UC Santa Cruz Previously Published Works

### Title

Material/element-dependent fluorescence-yield modes on soft X-ray absorption spectroscopy of cathode materials for Li-ion batteries

### Permalink

<https://escholarship.org/uc/item/6519370g>

### Journal

AIP Advances, 6(3)

### ISSN

2158-3226

### Authors

Asakura, Daisuke  
Hosono, Eiji  
Nanba, Yusuke  
[et al.](#)

### Publication Date

2016-03-01

### DOI

10.1063/1.4943673

Peer reviewed

## Material/element-dependent fluorescence-yield modes on soft X-ray absorption spectroscopy of cathode materials for Li-ion batteries

Daisuke Asakura, Eiji Hosono, Yusuke Nanba, Haoshen Zhou, Jun Okabayashi, Chunmei Ban, Per-Anders Glans, Jinghua Guo, Takashi Mizokawa, Gang Chen, Andrew J. Achkar, David G. Hawthorn, Thomas Z. Regier, and Hiroki Wadati

Citation: *AIP Advances* **6**, 035105 (2016); doi: 10.1063/1.4943673

View online: <http://dx.doi.org/10.1063/1.4943673>

View Table of Contents: <http://aip.scitation.org/toc/adv/6/3>

Published by the [American Institute of Physics](#)

---

### Articles you may be interested in

[Modular soft x-ray spectrometer for applications in energy sciences and quantum materials](#)

*Review of Scientific Instruments* **88**, 013110 (2017); 10.1063/1.4974356

[Utility of the inverse partial fluorescence for electronic structure studies of battery materials](#)

*Applied Physics Letters* **100**, 193906 (2012); 10.1063/1.4711801

[Femtosecond X-ray magnetic circular dichroism absorption spectroscopy at an X-ray free electron laser](#)

*Review of Scientific Instruments* **87**, 033110 (2016); 10.1063/1.4944410

[High-efficiency in situ resonant inelastic x-ray scattering \(iRIXS\) endstation at the Advanced Light Source](#)

*Review of Scientific Instruments* **88**, 033106 (2017); 10.1063/1.4977592

[Temperature-dependent oxygen behavior of  \$\text{Li}\_x\text{Ni}\_{0.5}\text{Mn}\_{1.5}\text{O}\_4\$  cathode material for lithium battery](#)

*APL Materials* **4**, 116111 (2016); 10.1063/1.4968566

[SAXES, a high resolution spectrometer for resonant x-ray emission in the 400 – 1600eV energy range](#)

*Review of Scientific Instruments* **77**, 113108 (2006); 10.1063/1.2372731

---

# HAVE YOU HEARD?

Employers hiring scientists and  
engineers trust

**PHYSICS TODAY | JOBS**

[www.physicstoday.org/jobs](http://www.physicstoday.org/jobs)



## Material/element-dependent fluorescence-yield modes on soft X-ray absorption spectroscopy of cathode materials for Li-ion batteries

Daisuke Asakura,<sup>1,a</sup> Eiji Hosono,<sup>1</sup> Yusuke Nanba,<sup>1,b</sup> Haoshen Zhou,<sup>1</sup> Jun Okabayashi,<sup>2</sup> Chunmei Ban,<sup>3</sup> Per-Anders Glans,<sup>4</sup> Jinghua Guo,<sup>4,5</sup> Takashi Mizokawa,<sup>6,c</sup> Gang Chen,<sup>7</sup> Andrew J. Achkar,<sup>8</sup> David G. Hawthorn,<sup>8</sup> Thomas Z. Regier,<sup>9</sup> and Hiroki Wadati<sup>10</sup>

<sup>1</sup>Energy Technology Research Institute, National Institute of Advanced Industrial Science and Technology, Tsukuba, Ibaraki 305-8568, Japan

<sup>2</sup>Research Center for Spectrochemistry, The University of Tokyo, Bunkyo-ku, Tokyo, 113-0033, Japan

<sup>3</sup>The Chemical and Materials Science Center, National Renewable Energy Laboratory, Golden, Colorado, 80401, USA

<sup>4</sup>Advanced Light Source, Lawrence Berkeley National Laboratory, Berkeley, California, 94720, USA

<sup>5</sup>Department of Chemistry and Biochemistry, University of California, Santa Cruz, California, 95064, USA.

<sup>6</sup>Department of Complexity Science and Engineering, The University of Tokyo, Kashiwa, Chiba, 277-8561, Japan

<sup>7</sup>College of Physics/State Key Laboratory of Superhard Materials, Jilin University, Changchun, 130012, People's Republic of China

<sup>8</sup>Department of Physics and Astronomy, University of Waterloo, Waterloo, Ontario, N2L 3G1, Canada

<sup>9</sup>Canadian Light Source, University of Saskatchewan, Saskatoon, Saskatchewan, S7N 0X4, Canada

<sup>10</sup>Institute for Solid State Physics, The University of Tokyo, Kashiwa, Chiba, 277-8581, Japan

(Received 1 December 2015; accepted 25 February 2016; published online 7 March 2016)

We evaluate the utilities of fluorescence-yield (FY) modes in soft X-ray absorption spectroscopy (XAS) of several cathode materials for Li-ion batteries. In the case of total-FY (TFY) XAS for  $\text{LiNi}_{0.5}\text{Mn}_{1.5}\text{O}_4$ , the line shape of the Mn  $L_3$ -edge XAS was largely distorted by the self-absorption and saturation effects, while the distortions were less pronounced at the Ni  $L_3$  edge. The distortions were suppressed for the inverse-partial-FY (IPFY) spectra. We found that, in the cathode materials, the IPFY XAS is highly effective for the Cr, Mn, and Fe  $L$  edges and the TFY and PFY modes are useful enough for the Ni  $L$  edge which is far from the O  $K$  edge. © 2016 Author(s). All article content, except where otherwise noted, is licensed under a Creative Commons Attribution (CC BY) license (<http://creativecommons.org/licenses/by/4.0/>). [<http://dx.doi.org/10.1063/1.4943673>]

Electronic-structure analyses for the electrode materials of Li-ion batteries (LIBs) are essential to understand the charge-discharge properties, which is capable of leading to improvements of the battery performance. The  $3d$  transition metals (TMs) are the major elements in the cathode materials for LIBs, such as TM oxides and phosphates. The TMs in the LIBs cathode materials experience the oxidation and reduction during charging/discharging process (i.e. Li-ion extraction/insertion). The redox chemistry determines the electrochemical properties of the electrode materials. Thus, element-selective electronic-structure analyses are indispensable for the application of LIBs.

<sup>a</sup>Author to whom correspondence should be addressed. Electronic mail: [daisuke-asakura@aist.go.jp](mailto:daisuke-asakura@aist.go.jp)

<sup>b</sup>Present affiliation: Inamori Frontier Research Center, Kyushu University, Fukuoka, 819-0395, Japan

<sup>c</sup>Present affiliation: Department of Applied Physics, Waseda University, Tokyo, 169-8555, Japan

In LIB research, TM  $K$ -edge X-ray absorption spectroscopy in hard X-ray region is widely used to investigate the electronic structures of cathode materials.<sup>1,2</sup> The redox reactions of TMs can be clarified by the edge/peak shift in the TM  $K$ -edge absorption spectra. However, the TM  $K$ -edge X-ray absorption spectroscopy is less sensitive to the  $3d$  electronic states of TMs than the TM  $L_{2,3}$ -edge (i.e. the  $2p$ - $3d$  transition) X-ray absorption spectroscopy, because the  $1s$ - $4p$  dipole transition is dominant at the TM  $K$  edges.

To obtain the detailed information of the  $3d$  orbitals, i.e. electronic-structure parameters such as charge-transfer energy, soft X-ray absorption spectroscopy (XAS) covering the TM  $L_{2,3}$  edges (i.e. the  $2p$ - $3d$  transition) has been widely used for various fields including solid state physics.<sup>3,4</sup> Total-electron yield (TEY) is the most conventional detection mode for the XAS experiments in the soft X-ray region, while the probing depth is approximately 5 nm at most. The surface sensitivity of the TEY mode is frequently inappropriate for investigating the redox reaction in the cathode materials. In contrast, total-fluorescence-yield (total FY: TFY) XAS can probe the bulk electronic structures as deep as  $\sim 100$  nm from the sample surface. However, the TM  $L_3$ -edge TFY XAS spectra tend to be distorted by self-absorption and saturation effects.<sup>5</sup> These effects are dependent on the density, particle size (for powder) and thickness (for thin film) of the sample and the measurement geometry such as normal/grazing incidence. In addition, the distortion in the line shapes of TM  $L_3$ -edge TFY XAS is enhanced in case where the target sample does not contain heavy elements such as rare earth.<sup>6</sup> Thus, powdered Li-TM oxides and related compounds used for cathode materials of LIB are frequently affected by these effects.<sup>7</sup> The distortion is significant when the O  $K_\alpha$  fluorescent emission is strongly dependent on the incoming photon energies at the TM  $L_{2,3}$  regions. To reduce the distortion in spectral line shapes, partial-fluorescence-yield (PFY) and inverse partial-fluorescence-yield (IPFY) modes have been utilized in previous works,<sup>6-14</sup> although the mechanisms of PFY and IPFY are still under debate.<sup>6,8,9,15-19</sup> The PFY detection mode using a silicon-drift detector (SDD) can select a specific emission, e.g. TM  $L_{\alpha,\beta}$ , suppressing the influence of the O  $K_\alpha$  emission. The distortions in the line shape of  $L_3$ -edge XAS can be further reduced by use of IPFY detection that is the inversed intensity of O  $K_\alpha$  emission.<sup>6-12</sup>

In this study, we select  $\text{LiFePO}_4$  and  $\text{LiNi}_{0.5}\text{Mn}_{1.5}\text{O}_4$  as model examples to investigate those effects in TFY mode and utilities of PFY and IPFY modes.  $\text{LiFePO}_4$  is a promising cathode material because of the low-cost constituent element and the highly stable charge-discharge properties.<sup>20</sup> In the olivine-type crystal structure, the Fe atoms are located in the  $\text{FeO}_6$  octahedrons and should be divalent. Spinel-type  $\text{LiNi}_{0.5}\text{Mn}_{1.5}\text{O}_4$  is an attractive cathode material in terms of the high voltage of  $\sim 4.8$  V vs.  $\text{Li/Li}^+$ .<sup>21</sup> The Ni and Mn atoms also form the (Ni,Mn) $\text{O}_6$  octahedrons with  $O_h$  symmetry. The formal valences of Ni and Mn for the initial state are considered as  $2+$  and  $4+$ , respectively. In addition, we performed the Cr  $L$ -edge XAS for  $\text{Li}_{0.35}\text{Mn}_{0.65}\text{Cr}_{0.35}\text{O}_2$  to judge the utility of FY modes at Cr  $L$  edges.

We demonstrate TFY, PFY and IPFY XAS for those as-synthesized typical cathode materials in order to discuss appropriate use of the FY modes for each material/element. In this study, we focus on trends in the TMs of the cathode materials as observed in the FY XAS spectra.

$\text{LiFePO}_4$  was synthesized with a solid-state reaction method.<sup>23</sup>  $\text{LiNi}_{0.5}\text{Mn}_{1.5}\text{O}_4$  was synthesized by an electrospinning method.<sup>24</sup> We confirmed the successful syntheses of the cathode materials using powder X-ray diffraction and electrochemical experiments. The fabrication process of  $\text{Li}_{0.35}\text{Mn}_{0.65}\text{Cr}_{0.35}\text{O}_2$  has been described elsewhere.<sup>22</sup> We also prepared a charged sample of  $\text{LiNi}_{0.5}\text{Mn}_{1.5}\text{O}_4$ . The electrochemical experiments were performed by the same procedures described in Ref. 24.

For  $\text{LiFePO}_4$  and  $\text{LiNi}_{0.5}\text{Mn}_{1.5}\text{O}_4$ , TEY, PFY, and IPFY XAS measurements were performed at BL-7A of the Photon Factory. The energy resolution of incident photon energies for the XAS measurements was  $E/\Delta E \sim 1500$ . An SDD was used for the PFY and IPFY detection. The energy width of SDD profile for fixed incident energies is approximately 100 eV. TFY XAS measurements were carried out at BL7.0.1 and BL8.0.1 of the Advanced Light Source with an energy resolution of  $E/\Delta E \sim 2000$  for the incident photon energies. A photodiode was used for the detection of TFY system.

For  $\text{Li}_{0.35}\text{Mn}_{0.65}\text{Cr}_{0.35}\text{O}_2$ , TEY, TFY PFY, and IPFY XAS experiments were performed at 11ID-1 Spherical Grating Monochromator (SGM) of the Canadian Light Source. The TFY spectra

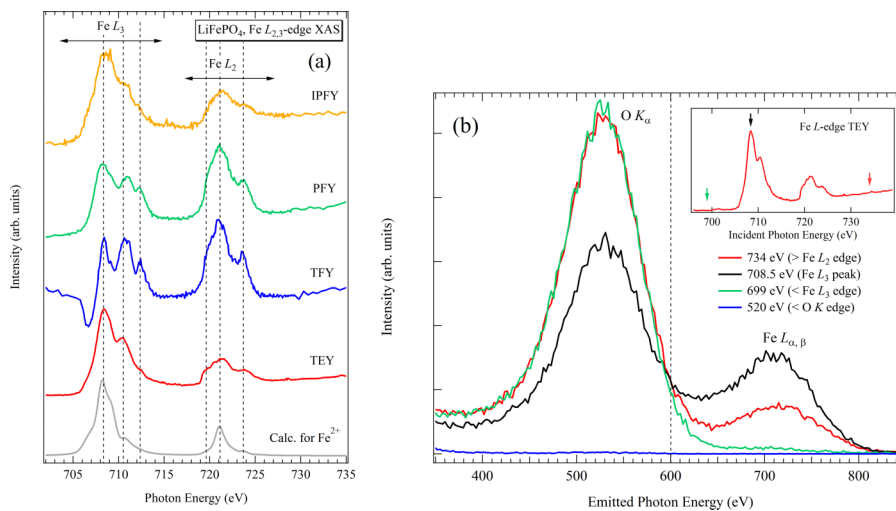


FIG. 1. (a) The Fe  $L_{2,3}$ -edge XAS spectra for LiFePO<sub>4</sub> taken by TEY, TFY, PFY and IPFY detection modes. Calculated spectrum for Fe<sup>2+</sup> HS state is also shown. (b) Emitted-photon-energy ( $h\nu_{\text{out}}$ ) profiles for several incident photon energies ( $h\nu_{\text{in}}$ ). The corresponding incident energies are indicated in the inset: TEY spectrum.

were measured with a channel plate detector. The IPFY spectra were measured using an SDD with an energy width of  $\sim 120$  eV. The energy resolution of incident photon energies was set to 5000. All the XAS measurements were performed at room temperature. Configuration-interaction full-multiplet (CIFM) calculations were adopted for the analysis of the  $L_{2,3}$ -edge XAS spectra.<sup>25–28</sup> Ligand-to-metal charge transfer (LMCT) was taken into account in the calculation.

Figure 1(a) shows the Fe  $L_{2,3}$ -edge TEY, TFY, PFY, and IPFY XAS spectra for LiFePO<sub>4</sub>. The main peak at 708.5 eV in the TEY spectrum is assigned to Fe<sup>2+</sup> high-spin (HS) state from the multiplet calculation while the peak at 710.5 eV is slightly larger than the calculated Fe<sup>2+</sup> HS spectrum. Most likely, a small fraction of Fe<sup>3+</sup> HS should coexist according to previous reports.<sup>29,30</sup> The TFY spectrum exhibits a dip-like structure at 706.5 eV and less pronounced  $L_3$  peaks at 708.4 and 710.5 eV. The  $L_3$  peak intensity is similar to that in the  $L_2$  edge. In the PFY spectrum corresponding to the Fe  $L_{\alpha,\beta}$  emission, the dip-like structure at 706.5 eV was not observed although the  $L_3/L_2$  intensity ratio is similar to that in the TFY spectrum and smaller than that in TEY. On the other hand, the IPFY spectrum which is the inverse of the O  $K_{\alpha}$  emission shows the TEY-like intensity ratio of  $L_3/L_2$ . Furthermore, the profile of the IPFY spectrum is almost the same as that of TEY, suggesting the bulk and surface electronic structures are identical. These results indicate that the saturation and self-absorption effects are emphasized at the Fe  $L$  edge of this material. IPFY can be the most appropriate method to detect the bulk Fe  $3d$  states. The results are consistent with the data collected by Yang *et al.*<sup>7</sup>

Figure 1(b) shows the emitted-photon-energy ( $h\nu_{\text{out}}$ ) profiles taken at three incident photon energies ( $h\nu_{\text{in}}$ ). Each profile was measured for 1 min. Although the energy width of the SDD profile is about 100 eV, the O  $K_{\alpha}$  and Fe  $L_{\alpha,\beta}$  emissions can be distinguished. The O  $K_{\alpha}$  and Fe  $L_{\alpha,\beta}$  emissions were not observed for  $h\nu_{\text{in}} = 520$  eV because 520 eV is lower than the O  $K$  edge absorption threshold. For  $h\nu_{\text{in}} = 699$  eV (below the Fe  $L_3$  edge), only the O  $K_{\alpha}$  emission is detected. The Fe  $L_{\alpha,\beta}$  emission (for  $600 < h\nu_{\text{out}} < 800$  eV) is enhanced for  $h\nu_{\text{in}} = 708.5$  eV corresponding to the Fe  $L_3$  edge, while O  $K_{\alpha}$  emission (for  $420 < h\nu_{\text{out}} < 600$  eV) is considerably reduced. The  $\sim 40\%$  reduction of O  $K_{\alpha}$  emission causes the apparent reduction of Fe  $L_3$  peak in the TFY spectrum. In addition, compared with the profile for  $h\nu_{\text{in}} = 708.5$  eV, the Fe  $L_{\alpha,\beta}$  and O  $K_{\alpha}$  emission intensities were reduced and unchanged for  $h\nu_{\text{in}} = 734$  eV, respectively. This is consistent with the Fe  $L_2$  TFY and PFY spectra which show high intensities.

As Fig. 2(a) shows, the Mn  $L_{2,3}$ -edge TEY spectrum is attributed to a Mn<sup>4+</sup> state. However, the TFY spectrum exhibits a considerable deformation at the Mn  $L_3$  edge, suggesting that the self-absorption and saturation effects strongly affect the Mn  $L_3$ -edge TFY spectrum. The IPFY

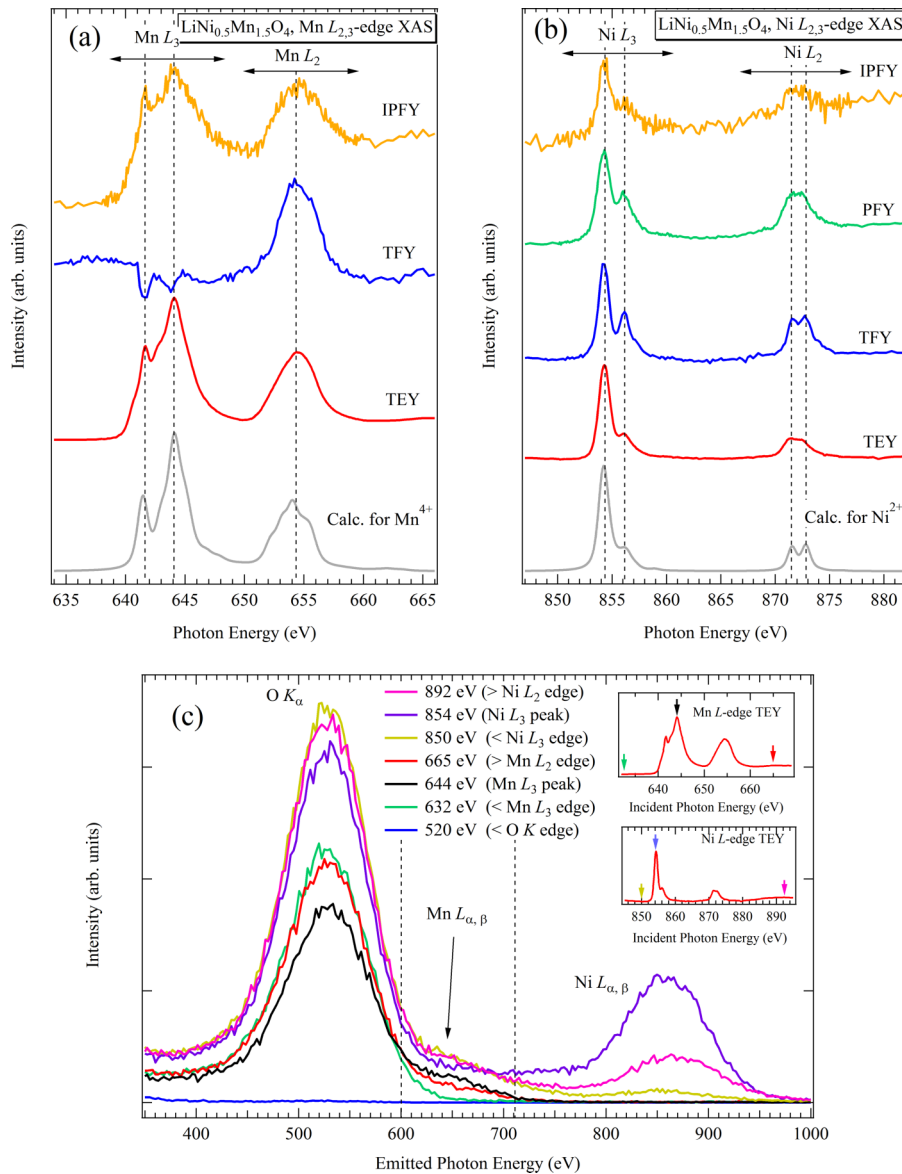


FIG. 2. (a) The Mn  $L_{2,3}$ -edge and (b) Ni  $L_{2,3}$ -edge XAS spectra for  $\text{LiNi}_{0.5}\text{Mn}_{1.5}\text{O}_4$  taken by TEY, TFY, PFY and IPFY detection modes. Calculated spectra are also shown. (c) Emitted-photon-energy ( $h\nu_{\text{out}}$ ) profiles for several incident photon energies ( $h\nu_{\text{in}}$ ). The corresponding incident energies are indicated in the insets: TEY spectra.

spectrum displays a clear Mn  $L_3$ -edge structure although the  $L_3/L_2$  intensity ratio is slightly smaller than that in TEY. The IPFY spectrum is also assigned to the  $\text{Mn}^{4+}$  state. The similarity between TEY and IPFY spectra indicates that the surface and bulk electronic structures are almost identical.

The Ni  $L_{2,3}$ -edge TEY spectrum exhibits the  $\text{Ni}^{2+}$  HS multiplet structure (Fig. 2(b)). In contrast to the Mn  $L$  edge, the distortions of the Ni  $L_3$  edge TFY and PFY spectra are suppressed. Although the  $L_3/L_2$  intensity ratios of TFY and PFY are slightly smaller than that for TEY, the TFY and PFY spectra are available for the analysis using the multiplet calculations. The IPFY spectrum is also similar to the TFY and PFY spectra. However, the signal-to-noise (S/N) ratio of IPFY spectrum is lower than those of TFY and PFY spectra. Thus, the IPFY detection mode is less advantageous for the analysis of Ni  $L_{2,3}$  edges of this material.

Figure 2(c) shows the  $h\nu_{\text{out}}$  profiles of emissions taken at several  $h\nu_{\text{in}}$ . The O  $K_{\alpha}$ , Mn  $L_{\alpha,\beta}$  and Ni  $L_{\alpha,\beta}$  emissions can be observed. The Mn  $L_{\alpha,\beta}$  emissions (for  $600 < h\nu_{\text{out}} < 720$  eV) are

enhanced at  $h\nu_{\text{in}} = 644$  eV. However, the degree of enhancement is lower than the degree of reduction of the O  $K_{\alpha}$  emission, leading to the deformation at the Mn  $L_3$  edge in the TFY spectrum. On the other hand, at the Ni  $L_3$  edge, the degree of reduction of the O  $K_{\alpha}$  emission becomes small and the Ni  $L_{\alpha,\beta}$  emission (for  $750 < h\nu_{\text{out}} < 970$  eV) increases considerably, such that the deformations at the Ni  $L_3$  edge in the TFY and PFY spectra are less pronounced.

The Mn  $L_3$ -edge TFY XAS spectrum for  $\text{LiNi}_{0.5}\text{Mn}_{1.5}\text{O}_4$  exhibits the large distortions in XAS line shapes, suggesting that the XAS line shapes of Mn  $L_3$  edge in Li-TM oxides (and phosphates) should be very sensitive to the self-absorption and saturation effects, while these effects depend on the density and composition of sample and experimental geometries. This is consistent with the Mn  $L_3$  edge/ $L_{\alpha,\beta}$  emission ( $\sim 640$  to  $\sim 660$  eV) being very near O  $K$  edge/ $K_{\alpha}$  emission ( $\sim 525$  to  $\sim 560$  eV) in photon energy. The Fe  $L_3$  edge XAS is also sensitive, although the influence of self-absorption and saturation effects should be smaller than that in the Mn  $L_3$  edge. Considering the present results, the IPFY detection mode can become powerful tools for the Mn and Fe  $L_{2,3}$ -edge measurements. In the case of Ni  $L_{2,3}$  edges, the TFY and PFY detection modes should be more appropriate than IPFY mode in terms of the S/N ratio. Most likely, Co  $L_{2,3}$  edges for Li-TM oxides (and phosphates) including Co would be positioned as the intermediate character between Fe and Ni  $L_{2,3}$  edges.

As Fig. 2(c) indicates, the  $h\nu_{\text{in}}$  dependence of O  $K_{\alpha}$  emission weakens as  $h\nu_{\text{in}}$  increases. To further verify this tendency, we also performed the Cr  $L$ -edge XAS for  $\text{Li}_{0.35}\text{Mn}_{0.65}\text{Cr}_{0.35}\text{O}_2$  (Fig. 3(a)). The TEY spectrum is attributed to the  $\text{Cr}^{3+}$  state. On the other hand, not only the  $L_3$  edge, but also the  $L_2$  edge is reversed with negative sign in the TFY spectrum, which is more problematic than the case of Mn  $L$ -edges (Fig. 3(b)). The Cr  $L$  edges ( $\sim 575$  to  $\sim 595$  eV) are so close to the O  $K$  edge ( $\sim 525$  to  $\sim 560$  eV) that the Cr  $L_{\alpha,\beta}$  emission cannot be separated from the O  $K_{\alpha}$  emission within the energy width of SDD ( $\sim 100$  eV) (Fig. 3(c)). The Cr  $L$ -edge IPFY spectrum in Fig. 3(a) is composed by extracting the structure below 570 eV in Fig. 3(c) as the O  $K_{\alpha}$  emission. However, the S/N ratio and  $L_3/L_2$  intensity ratio are not ideal. If the energy resolution of the detector is improved, for example by using a high-resolution soft X-ray emission spectrometer,<sup>31,32</sup> then the separation of Cr  $L_{\alpha,\beta}$  emissions from O  $K_{\alpha}$  emission should be facilitated.

Figure 4 summarizes the utilities of TFY, PFY and IPFY for each TM  $L$  edge of Li-TM oxides/phosphates. Note that the tendency is limited to powdered Li-TM oxides/phosphates with particle sizes between 10 nm and 10  $\mu\text{m}$  generally used as cathode materials of LIB. The IPFY mode can be useful for Cr, Mn, and Fe  $L$  edges for those materials. For  $L$  edges with higher energies, TFY and PFY are more suitable than IPFY. In addition, TFY and PFY should be available for Ti and V without any problems because these early TMs'  $L$  edges are located below the O  $K$  edge. Appropriate choice of the FY modes is inevitable to study the bulk electronic structures of Li-TM oxides and related materials. As Fig. 5 shows, the bulk electronic structure is often different from the surface electronic structure in particular for charged (or discharged) samples. In the case of  $\text{LiNi}_{0.5}\text{Mn}_{1.5}\text{O}_4$  (Fig. 5), the degree of oxidation of Ni due to the charge reaction is different between the surface and bulk. The Ni  $L$ -edge TFY XAS spectrum for the charged state shows enhanced peaks at 856 and 873 eV, indicating partial oxidation reaction of  $\text{Ni}^{2+} \rightarrow \text{Ni}^{3+/3.3}$  in the bulk region. On the other hand, in the surface-sensitive TEY mode, the line shape attributed to the  $\text{Ni}^{2+}$  state is almost maintained for the charged state, suggesting existence of redox-inactive Ni and/or some side reactions such as NiO formation near the surface. Thus, a combined use of TEY and FY is essential for the electronic-structure analyses of the electrode materials.

In conclusion, we demonstrated TEY, TFY, PFY and IPFY XAS studies for powdered several Li-TM oxides and a phosphate which are cathode materials for LIBs. The utilities of FY modes strongly depend on the element (edge) and material (composition). The FY modes should be carefully chosen for each material. The distortion of Ni  $L_3$ -edge TFY XAS line shapes is not large whereas Cr, Mn and Fe  $L_3$ -edge TFY XAS are greatly affected. We conclude that, in the FY XAS studies of the cathode materials, the IPFY XAS is most appropriate for the Cr, Mn, and Fe  $L$  edges and the TFY and PFY XAS modes are available enough for the Ni  $L$  edge which is far from the O  $K$  edge. This consideration will contribute to TM  $L$ -edge FY XAS studies of charged/discharged samples for novel electrode materials.

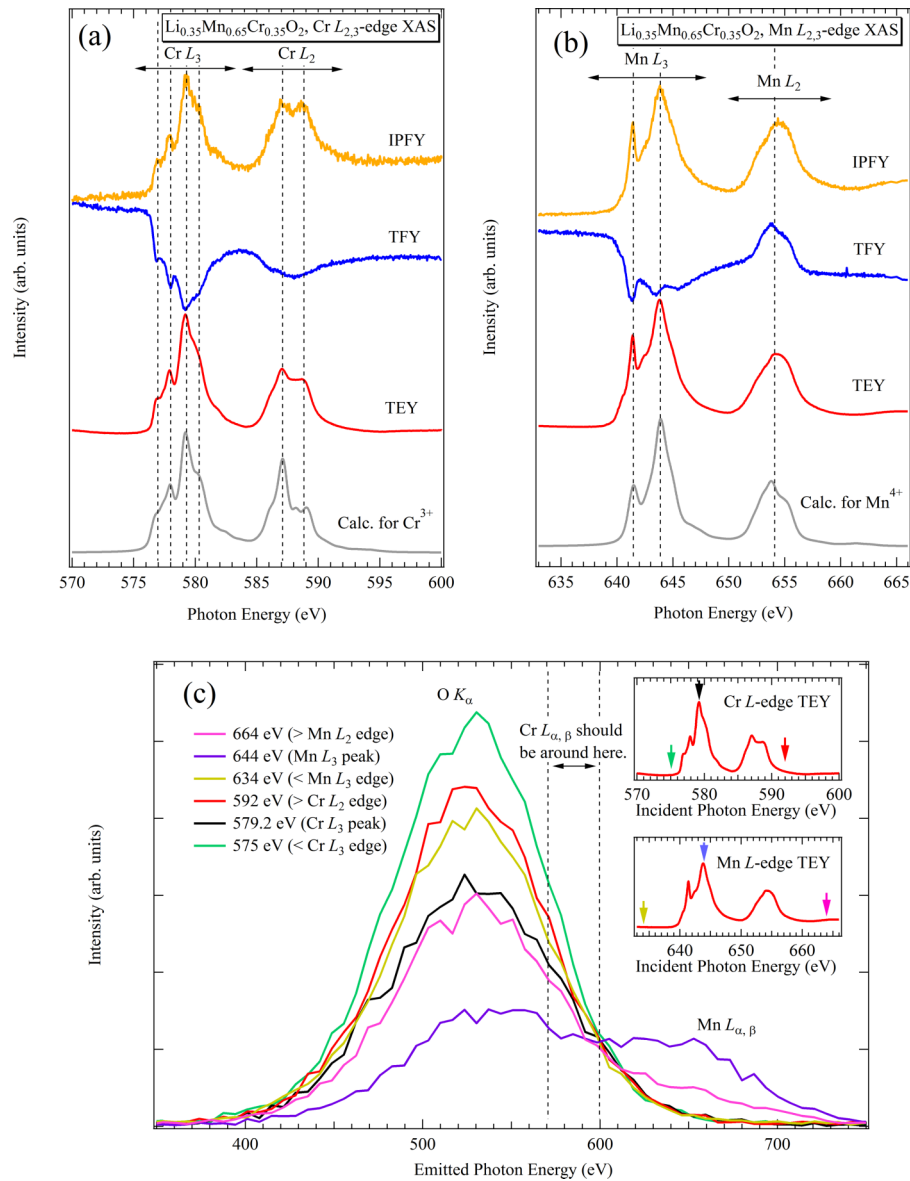


FIG. 3. (a) The Cr  $L_{2,3}$ -edge and (b) Mn  $L_{2,3}$ -edge XAS spectra for  $\text{Li}_{0.35}\text{Mn}_{0.65}\text{Cr}_{0.35}\text{O}_2$ . (c) Emitted-photon-energy ( $h\nu_{\text{out}}$ ) profiles for several incident photon energies ( $h\nu_{\text{in}}$ ). The Cr  $L_{\alpha,\beta}$  emission should be around 580 eV, which cannot be separated from the O  $K_{\alpha}$  emission in the present energy resolution.

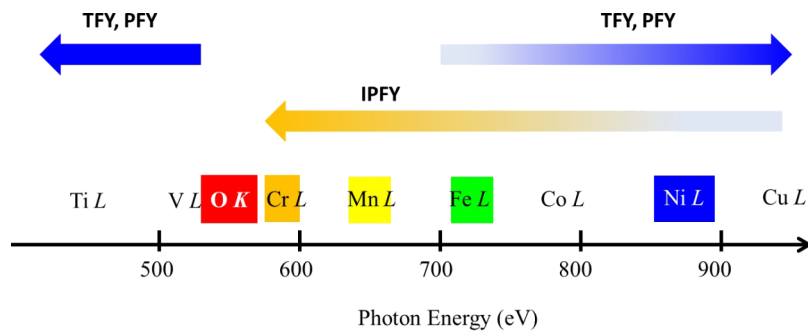


FIG. 4. A scheme of utilities of TFY, PFY and IPFY for each TM  $L$  edge of powdered Li-TM oxides/phosphates (Ti, V, Co and Cu  $L$  edges are not investigated here). The color scales of the arrows indicate the utilities of each XAS mode.



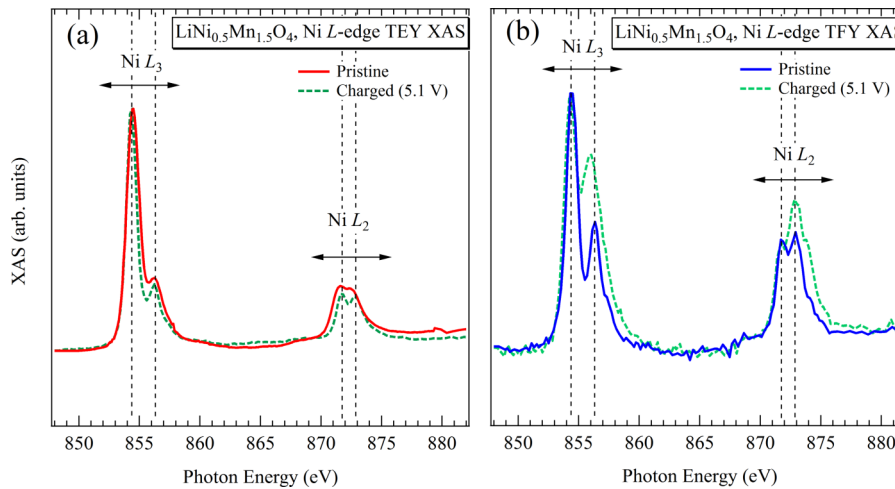


FIG. 5. The Ni *L*-edge XAS spectra for the pristine and charged  $\text{LiNi}_{0.5}\text{Mn}_{1.5}\text{O}_4$  recorded by (a) the TEY and (b) TFY modes. The spectra were normalized at 854.2 eV to emphasize the differences at 856 and 873 eV in (b).

## ACKNOWLEDGEMENTS

This work was performed under the approvals of the Photon Factory Program Advisory Committee (Proposal Nos. 2012G107, 2013G169 and 2014G100) and the Advanced Light Source (Proposal No. 04907R). The Advanced Light Source is supported by the Director, Office of Science, Office of Basic Energy Sciences, of the U.S. Department of Energy under Contract No. DE-AC02-05CH11231. Research described in this paper was partly performed at the Canadian Light Source, which is supported by CFI, NSERC, the University of Saskatchewan, the Government of Saskatchewan, WD, NRC, and CIHR. The synthesis work at NREL was supported by the NREL LDRD Program. This work was also conducted based on the Japan-U.S. cooperation project for research and standardization of Clean Energy Technologies.

- <sup>1</sup> Y. Shiraishi, I. Nakai, T. Tsubata, T. Himeda, and F. Nishikawa, *J. Solid State Chem.* **133**, 587 (1997).
- <sup>2</sup> A. Ito, Y. Sato, T. Sanada, M. Hatano, H. Horie, and Y. Osawa, *J. Power Sources* **196**, 6828 (2011).
- <sup>3</sup> F. de Groot and A. Kotani, *Core Level Spectroscopy of Solids (Advances in Condensed Matter Science)* (CRC Press Taylor & Francis Group, Boca Raton, FL, 2008).
- <sup>4</sup> F. M. F. de Groot, *J. Electron Spectrosc. Relat. Phenom.* **67**, 529 (1994).
- <sup>5</sup> S. Eisebitt, T. Boske, J.-E. Rubensson, and W. Eberhardt, *Phys. Rev. B* **47**, 14103 (1993).
- <sup>6</sup> H. Wadati, A. J. Achkar, D. G. Hawthorn, T. Z. Regier, M. P. Singh, K. D. Truong, P. Fournier, G. Chen, T. Mizokawa, and G. A. Sawatzky, *Appl. Phys. Lett.* **100**, 193906 (2012).
- <sup>7</sup> S. Yang, D. Wang, G. Liang, Y. M. Yiu, J. Wang, L. Liu, X. Sun, and T.-K. Sham, *Energy Environ. Sci.* **5**, 7007 (2012).
- <sup>8</sup> A. J. Achkar, T. Z. Regier, H. Wadati, Y.-J. Kim, H. Zhang, and D. G. Hawthorn, *Phys. Rev. B* **83**, 081106 (2011).
- <sup>9</sup> A. J. Achkar, T. Z. Regier, E. J. Monkman, K. M. Shen, and D. G. Hawthorn, *Sci. Rep.* **1**, 182 (2011).
- <sup>10</sup> S. I. Bokarev, M. Dantz, E. Suljoti, O. Kühn, and E. F. Aziz, *Phys. Rev. Lett.* **111**, 083002 (2013).
- <sup>11</sup> P. Wernet, K. Kunnus, S. Schreck, W. Quevedo, R. Kurian, S. Techert, F. M. F. de Groot, M. Odelius, and A. Fohlich, *J. Phys. Chem. Lett.* **3**, 3448 (2012).
- <sup>12</sup> D. Perk and T. Regier, *Environ. Sci. Technol.* **46**, 3163 (2012).
- <sup>13</sup> C. Yogi, D. Takamatsu, K. Yamanaka, H. Arai, Y. Uchimoto, K. Kojima, I. Watanabe, T. Ohta, and Z. Ogumi, *J. Power Sources* **248**, 994 (2014).
- <sup>14</sup> H. Wadati, D. G. Hawthorn, T. Z. Regier, G. Chen, T. Hitosugi, T. Mizokawa, A. Tanaka, and G. A. Sawatzky, *Appl. Phys. Lett.* **97**, 022106 (2010).
- <sup>15</sup> E. F. Aziz, M. H. Rittmann-Frank, K. M. Lange, S. Bonhommeau, and M. Chergui, *Nature Chem.* **2**, 853 (2010).
- <sup>16</sup> T. Z. Regier, A. J. Achkar, D. Peak, J. S. Tse, and D. G. Hawthorn, *Nature Chem.* **4**, 765 (2012).
- <sup>17</sup> F. M. F. de Groot, *Nature Chem.* **4**, 766 (2012).
- <sup>18</sup> E. F. Aziz, K. M. Lange, S. Bonhommeau, and M. Chergui, *Nature Chem.* **4**, 767 (2012).
- <sup>19</sup> R. Kurian, K. Kunnus, P. Wernet, S. M. Buterin, P. Glatzel, and F. M. F. de Groot, *J. Phys. Condens. Matter* **24**, 452201 (2012).
- <sup>20</sup> A. Padhi, K. Nanjundaswamy, and J. B. Goodenough, *J. Electrochem. Soc.* **144**, 1188 (1997).
- <sup>21</sup> Q. M. Zhong, A. Bonakdarpour, M. J. Zhang, Y. Gao, and J. R. Dahn, *J. Electrochem. Soc.* **144**, 205 (1997).
- <sup>22</sup> F. Du, Z.-F. Huang, C.-Z. Wang, X. Meng, and G. Chen, *J. Appl. Phys.* **102**, 113906 (2007).
- <sup>23</sup> C. Ban, W.-J. Yin, H. Tang, S.-H. Wei, Y. Yan, and A. C. Dillon, *Adv. Energy Mater.* **2**, 1028 (2012).

- <sup>24</sup> E. Hosono, T. Saito, J. Hoshino, Y. Mizuno, M. Okubo, D. Asakura, K. Kagesawa, D. Nishio-Hamane, T. Kudo, and H. S. Zhou, *CrystEngComm* **15**, 2592 (2013).
- <sup>25</sup> Y. Nanba and K. Okada, *J. Phys. Soc. Jpn.* **79**, 114722 (2010).
- <sup>26</sup> Y. Nanba and K. Okada, *J. Electron Spectrosc. Relat. Phenom.* **185**, 167 (2012).
- <sup>27</sup> Y. Nanba, D. Asakura, M. Okubo, Y. Mizuno, T. Kudo, H. S. Zhou, K. Amemiya, J.-H. Guo, and K. Okada, *J. Phys. Chem. C* **116**, 24896 (2012).
- <sup>28</sup> Y. Nanba, D. Asakura, M. Okubo, H. S. Zhou, K. Amemiya, K. Okada, P.-A. Glans, C. A. Jenkins, E. Arenholz, and J.-H. Guo, *Phys. Chem. Chem. Phys.* **16**, 7031 (2014).
- <sup>29</sup> X. Liu, J. Liu, R. Qiao, Y. Yu, H. Li, L. Suo, Y.-s. Hu, Y.-D. Chuang, G. Shu, F. Chou, T.-C. Weng, D. Nordlund, D. Sokaras, Y. J. Wang, H. Lin, B. Barbiellini, A. Bansil, X. Song, Z. Liu, S. Yan, G. Liu, S. Qiao, T. J. Richardson, D. Prendergast, Z. Hussain, F. M. F. de Groot, and W. Yang, *J. Am. Chem. Soc.* **134**, 13708 (2012).
- <sup>30</sup> S. Kurosumi, N. Nagamura, S. Toyoda, K. Horiba, H. Kumigashira, M. Oshima, S. Furutsuki, S.-i. Nishimura, A. Yamada, and N. Mizuno, *J. Phys. Chem. C* **115**, 25519 (2011).
- <sup>31</sup> G. Ghiringhelli, A. Piazzalunga, C. Dallera, G. Trezzi, L. Braicovich, T. Schmitt, V. N. Strocov, R. Betemps, L. Patthey, X. Wang, and M. Grioni, *Rev. Sci. Instrum.* **77**, 113108 (2006).
- <sup>32</sup> Y. Harada, M. Kobayashi, H. Niwa, Y. Senba, H. Ohashi, T. Tokushima, Y. Horikawa, S. Shin, and M. Oshima, *Rev. Sci. Instrum.* **83**, 013116 (2012).
- <sup>33</sup> C. Piamonteze, F. M. F. de Groot, H. C. N. Tolentino, A. Y. Ramos, N. E. Massa, J. A. Alonso, and M. J. Martínez-Lope, *Phys. Rev. B* **71**, 020406(R) (2005).