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EFFICIENCY OF EQUATIONS OF STATE FOR GASEOUS MIXTURES AT THE CRITICAL LOCUS. I. APPLICATION OF THE EQUATION OF BENEDICT, WEBB, AND RUBIN. II. FURTHER IMPROVEMENTS OF AN EQUATION OF STATE

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# University of California

# Ernest O. Lawrence Radiation Laboratory

# EFFICIENCY OF EQUATIONS OF STATE FOR GASEOUS MIXTURES AT THE CRITICAL LOCUS

 I. APPLICATION OF THE EQUATION OF BENEDICT, WEBB, AND RUBIN
 II. FURTHER IMPROVEMENTS OF AN EQUATION OF STATE

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Frank Jay Ackerman

(M.S. Thesis)

February 1, 1963

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## EFFICIENCY OF EQUATIONS OF STATE FOR GASEOUS MIXTURES AT THE CRITICAL LOCUS

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February 1, 1963

#### ABSTRACT

The critical state of gaseous mixtures is of practical interest in high-pressure reactions and petroleum engineering. A discussion of the critical state is appropriately based on the equation of state and rigorous thermodynamic relations.

The present investigation compares the eight-parameter equation of state proposed by Benedict, Webb, and Rubin with observed data and a two-parameter equation of Redlich and Kwong. Contrary to expectation, the more elaborate equation furnishes sometimes entirely unreasonable results. This failure is related to the fact that the values of the eight parameters are in general not sufficiently well defined by observations in a limited range of temperature and pressure.

An algebraic equation of state is desirable because it leads shortly to fugacity coefficients, which are the practically important objective in the study of the variables of state. The first part of this thesis illustrates the advantages and shortcomings of the equation of Redlich and Kwong. The improved equation obtained in Part II comes close to the accuracy of Pitzer's tables and at the same time permits an algebraic derivation of fugacity coefficients.

#### PART I. APPLICATION OF THE EQUATION OF BENEDICT, WEBB, AND RUBIN

#### A. Introduction

It would be useful to have a reliable method of predicting the critical temperature and pressure of any gas mixture, particularly as a reference point for the purpose of interpolation and correlation. Several methods of estimation using generalized correlations have been proposed. <sup>1, 2</sup> A broader basis for an estimation is furnished by a suitable equation of state.

The thermodynamic relations for the derivation of the critical state from an equation of state put the equation to a severe test because first and second derivatives are required, which are always more sensitive to deviations than the original function. Since many equations of state are not considered to be accurate in the critical region, their usefulness in a relation to find the critical state is questionable. Any equation considered must be applicable to mixtures as well as single components.

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#### B. The Critical State

Redlich and Kister used an equation of state<sup>3</sup>

$$P = RT/(V-b) - a/[T^{0.5}V(V+b)],$$

to predict the critical properties of several gas mixtures.<sup>4</sup> Parameters a and b can be expressed in terms of the critical temperature  $T_{C}$  and the critical pressure  $P_{C}$  as

$$a = 0.4278 R^2 T_C^{2.5} / P_C$$

and

$$b = 0.0867 RT/P_{C}$$
.

Their results were in reasonably fair agreement with experimental results; however, a more accurate prediction would be desirable. A better equation of state should furnish such results.

The eight-constant equation of Benedict, Webb, and Rubin<sup>5</sup> is used over considerable ranges of temperature and pressure to predict the properties of gases and gas mixtures. Benedict found that it described the critical state of pure components with reasonable accuracy for a number of gases within approximately 0.3°C and 0.4 atm of the experimental values.<sup>5</sup> Since Benedict's equation is generally well behaved for pure gases, it was expected that it would also furnish a good representation of the critical state for gas mixtures.

The critical state for a gas mixture is defined by:

$$(\partial \ln f/\partial y_1)_{P,T} = 0$$
 (1a)

$$(\partial^2 \ln f/\partial y_1^2)_{\rm PT} = 0 . \qquad (1b)$$

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From these two equations one can derive the thermodynamic equations for the critical lines. The two resulting equations are complicated expressions in terms of various partial differential quotients which, in the case considered here, must be solved simultaneously to yield the critical volume and temperature. The critical pressure is then found by substituting back into the equation of state. The equations are unwieldy and the rigorous solution, even with automatic computation methods, would be prohibitively time-consuming.

#### C. The Limiting Slopes of the Critical Lines

Actually a more elegant solution can be found. The limiting slopes of the critical-pressure line and the critical-temperature line with respect to the composition at the end points of these lines can be rigorously derived. According to Redlich and Kister, <sup>4</sup> we have

$$(dT/dy_1)_c = [(y_1-y_2) (dP/dy_1)_{V,T}^2/RE - (d^2P/dVdy_1)_T]/(d^2P/dTdV)_y$$
(2a)

and

$$(dP/dy_1)_c = (dP/dy_1)_{V,T} + (dP/dT)_{V,y}(dT/dy)_c$$
, (2b)

where  $y_1$  and  $y_2$  are mole fractions. The use of these relations requires only partial derivatives of the equation of state at the end points of the critical lines, i.e., at the critical points of the pure components.

We use the equation of Benedict, Webb, and Rubin in the form  

$$P = RTD + (B_0 RT - A_0 - C_0 / T^2) D^2 + (bRT - a) D^3 + aaD^6 + \frac{CD^3}{T^2} e^{-\gamma D^2} \times (1 + \gamma D^2), \quad (3)$$

with

$$D = 1/V.$$

For a mixture, the eight constants are assumed to depend on the constants of the constituent gases according to

$$A_{0}^{1/2} = \sum_{i} y_{i} A_{0i}^{1/2} ,$$
  

$$B_{0} = \sum_{i} y_{i} B_{0i} ,$$
  

$$C_{0}^{1/2} = \sum_{i} y_{i} C_{0i}^{1/2} ,$$
  

$$a^{1/3} = \sum_{i} y_{i} a_{i}^{1/3} ,$$
  

$$b^{1/3} = \sum_{i} y_{i} b_{i}^{1/3} ,$$
  

$$c^{1/3} = \sum_{i} y_{i} c_{i}^{1/3} ,$$

$$a^{1/3} = \sum_{i} y_{i} a_{i}^{1/3} ,$$
  

$$\gamma^{1/2} = \sum_{i} y_{i} \gamma_{i}^{1/2} .$$
(4)

and 

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The derivatives required by Eqs. (2a) and (2b) are:

$$\left(\frac{dP}{dy}\right)_{V,T} = (RT B_0^{\dagger} - A_0^{\dagger} - C_0^{\dagger}/T^2) D^2 + (b^{\dagger}RT - a^{\dagger}) D^3 + (aa^{\dagger} + aa^{\dagger}) D^6 + [c^{\dagger}(1+\gamma D^2) - C\gamma\gamma^{\dagger}D^4] D^3 e^{-\gamma D^2}/T^2,$$
(5)

where

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$$A_{0}' = \frac{dA_{0}}{dy_{1}} = 2 (y_{1}A_{0})^{0.5} + y_{2}A_{0}^{0.5}) (A_{0})^{0.5} - A_{0}^{0.5}), \text{ etc. ; (6)}$$

.

$$\left(\frac{d^{2}P}{dVdy_{1}}\right) = -D^{3} \left\{ 2 (RT B_{0}' - A_{0}' - C_{0}'/T^{2}) + 3(b'RT-a') \right\}$$

$$\begin{pmatrix} \frac{d^{2}P}{dVdy_{1}} \end{pmatrix} = -D^{3} \begin{cases} 2 (RT B_{0}' - A_{0}' - C_{0}'/T^{2}) + 3(b'RT-a') \\ \times D + 6(aa'+a'a)D^{4} \\ + De^{-\gamma D^{2}} \left[ C'(3+3\gamma D^{2}-2\gamma^{2}D^{4}) - C\gamma\gamma'D^{4}(7-2\gamma D^{2}) \right] / T^{2} \end{cases}; (7) \\ \left( \frac{dP}{dT} \right)_{V, Y} = RD + (B_{0}R + 2C_{0}/T^{3})D^{2} + bRD^{3} - 2CD^{3}e^{-\gamma D^{2}} \\ \times (1+\gamma D^{2})/T^{3}; (8) \end{cases}$$

$$\left(\frac{d^{2}P}{dTdV}\right)_{y} = -D^{2} \left\{ R + 2(B_{0}R + 2C_{0}/T^{3})D + 3bRD^{2} - 2D^{2}e^{-\gamma D^{2}} \left[ 3(1+\gamma D^{2}) + 2\gamma^{2}D^{4} \right] / T^{3} \right\}.$$
 (9)

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From these equations one can find the value of the slope at the end points  $(y_1 = 0 \text{ and } y_1 = 1)$  of the critical-temperature and critical-pressure lines.

#### D. Interpolation

The full critical-pressure and critical-temperature lines can now be approximately determined by means of an interpolation formula. As previously observed by Redlich and Kister, <sup>4</sup> it was found that a logarithmic-hyperbolic interpolation formula gives a reasonable approximation.

For the critical-pressure line, we abbreviate

$$p_1 = \lim_{y_1 \to 1} (dP/dy_1)_c$$
 (10)

and

$$p_2 = \lim_{y_1 \to 0} (dP/dy_1)_c .$$
(11)

The interpolation formula is then

$$ln P = y_1 ln P_1 + y_2 ln P_2$$
  
+ 
$$\frac{(ln P_1 - ln P_2 - p_2/P_2) (ln P_1 - ln P_2 - p_1/P_1)y_1 y_2}{(ln P_1 - ln P_2 - p_2/P_2)y_1 - (ln P_1 - ln P_2 - p_1/P_1)y_2}.$$
 (12)

If we replace P and p by T and t in Eqs. (10), (11), and (12), we have the formula used to calculate the critical-temperature line.

#### E. Results

With the available coefficients required in Eqs. (5) to (9), the critical temperatures and pressures as functions of the composition were obtained from Eq. (12) by automatic computation. For a comparison the critical lines were also calculated by means of Redlich and Kwong's equation of state. <sup>3</sup> The equations used are given by Redlich and Kister. <sup>4</sup> Results obtained from the calculations were compared with experimental data for 26 binary mixtures (methane with ethane, propane, n-butane, isobutane, pentane, heptane, decane, and nitrogen; ethane with propane, propene, butane, pentane, heptane, decane, benzene, and nitrogen; propene with butane, n-pentane, isopentane, benzene, and carbon dioxide; propene with 1-butene; butane with carbon dioxide; mith with with heptane and carbon dioxide; and carbon dioxide with sulfur dioxide). These results are presented in the figures.

Several sets of coefficients given by Opfell et al.<sup>6</sup> contain negative values for coefficients to be combined by the square-root formula (see Eq. 4). These sets have not been considered in the following.

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#### F. Discussion and Conclusions

An extensive comparison of the presented results does not reveal a clear advantage of either the equation of Benedict, Webb, and Rubin or that of Redlich and Kwong. This general result can be illustrated by examining some examples.

The critical-temperature line is always very well represented by either equation, Fig. 1 being a typical example. No further discussion of the critical temperature is necessary therefore.

The suitability of either equation, however, varies greatly for the critical-pressure lines of various systems. For many substances various authors have derived different sets of Benedict coefficients (Appendix C). At first we consider only those values of Benedict-Webb-Rubin coefficients that give the best fit. In certain cases the two equations are in good agreement with each other and with experimental data (Figs. 2, 5, 18, and 24). Sometimes the equation of Redlich and Kwong is better (Figs. 6 and 7); more often the better equation is that of Benedict et al. (Figs. 3, 4, 10, 12, 13, 16, 17, 19, 20, 22, and 25). In a few cases neither is very good (Figs. 8, 14, 15, and 26).

Considering now all proposed sets of Benedict coefficients, one finds the overall picture to be quite different. According to Figs. 2, 5, 8, 9, 10, 11, 18, 21, 24, and 26, one set is frequently in good agreement with Redlich and Kwong's equation. The deviations from the experimental data may be small (Figs. 2, 5, 18, and 24) or appreciable (Figs. 8, 9, 10, 11, 21, and 26). In many cases, some set of Benedict coefficients leads to larger deviations than does the equation of Redlich and Kwong (Figs. 2, 3, 4, 6, 7, 8, 12, 13, 17, 18, 21, 23, 24, 25, and 26). The slopes calculated from a few sets of Benedict constants are so far off that the interpolated curves are quite unreasonable (Figs. 13, 17, 18, 21, 25, and 26). These discrepancies are obviously due to entirely wrong values of one or both of the limiting slopes. In view of the thermodynamic connection of these slopes with the equation of state, the error lies with the coefficients, the assumed composition dependence of the coefficients, or both.

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Unquestionably the sixteen coefficients of the equation of Benedict et al. for binary mixtures can be chosen so that the critical lines are very well represented. However, the present study shows that the coefficients derived by previous authors from data in different ranges do not always give a good representation of the critical range, and sometimes, indeed, lead to entirely wrong results. A comparison of the two equations as applied to the heat content of gaseous mixtures leads to similar conclusions.<sup>7</sup>

The difficulty of adjusting a set of coefficients to all available data has been pointed out before.<sup>6</sup> The coefficients first derived by Benedict et al. are in general the best (curves A in Figs. 2, 3, 6, 12, and 24).<sup>5</sup> Opfell et al. recommend their coefficients only for the homogeneous region,<sup>8</sup> but their coefficients fit the data for ethane-propane better than the set of Benedict et al. (Fig. 10), though usually the opposite is true (Figs. 3, 6, 12, 13, 17, and 18). The wide variation of curves obtained from different sets of coefficients is illustrated in Figs. 23, 25, and 26.

The cause of the unexpected disadvantage of Benedict's equation in the present attempted application is really deep-seated. It results from the excessive number of coefficients in this equation. Of course, such an equation is supremely capable of being adapted to observations, but it is likely to lead to unreasonable results as soon as the region of observations is left behind.

The conclusion that a crude two-parameter equation is more reliable appears ridiculous at first. Yet there is much less danger of entirely unreasonable results, since the simpler function is necessarily restricted. The lesser adaptability and flexibility is actually an advantage in extending computations beyond the region of direct observation. As a matter of fact, the deviations resulting from the equation of Redlich and Kwong are seldom as high as in Fig. 8 and never unreasonable. They are usually small for two similar components (Figs. 2, 18, and 24).

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It may be concluded that the use of the equation of Benedict et al. for the critical properties can be recommended only if a method for the determination of the coefficients is found that is more definite than methods used to now.

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Fig. 1. Critical locus (temperature) of ethane-nitrogen<sup>9</sup> (full line: Redlich and Kwong; broken line: Benedict coeff.<sup>5,10</sup>).









Fig. 3. Critical pressure of methane-propane;  $\Delta^{12}$ ,  $O^{13}$  (Benedict coeff. A<sup>5</sup>, B<sup>8</sup>).



Fig. 4. Critical pressure of methane-butane  $^{14}$  (Benedict coeff.,  $A^5$ ,  $B^8$ ).



Fig. 5. Critical pressure of methane-isobutane<sup>15</sup> (Benedict coeff.<sup>5</sup>).



Fig. 6. Critical pressure of methane-pentane<sup>16</sup> (Benedict coeff. A<sup>5</sup>, B<sup>8</sup>).





Fig. 7. Critical pressure of methane-heptane<sup>17</sup> (Benedict coeff.<sup>5</sup>).



Fig. 8. Critical pressure of methane-decane<sup>18</sup> [Benedict coeff. for methane;<sup>5</sup> decane:  $A^6$  ( $\gamma = 0.4$ )  $B^6$ ( $\gamma = \infty$ )] and ethane-decane<sup>19</sup> [Benedict coeff. for methane;<sup>5</sup> decane<sup>6</sup> ( $\gamma = 0.4$ )].



Fig. 9. Critical pressure of methane-nitrogen<sup>20</sup> (Benedict coeff. for methane;<sup>5</sup> for ethane: A<sup>21</sup>, B<sup>10</sup>).

 $(x_{i},y_{i}) \in \mathbb{R}^{n \times n} (\mathbb{R}^{n \times n})$ 



















Fig. 14. Critical pressure of ethane-heptane<sup>27</sup> (Benedict coeff. A<sup>5</sup>, B: ethane<sup>11</sup>, propene<sup>5</sup>).



Fig. 15. Critical pressure of ethane-benzene<sup>28</sup> (Benedict coeff. for ethane:  $A^5$ ,  $B^{11}$ , for benzene<sup>29</sup>).



Fig. 16. Critical pressure of ethane-nitrogen<sup>9</sup> (Benedict coeff. for ethane<sup>11</sup>; for nitrogen: A<sup>10</sup>, B<sup>21</sup>).



Fig. 17. Critical pressure of propane-butane<sup>30</sup> (Benedict coeff. A<sup>5</sup>, B<sup>8</sup>).











Fig. 20. Critical pressure of propane-benzene<sup>33</sup> (Benedict coeff. for propane<sup>5</sup>; benzene<sup>29</sup>).



Fig. 21. Critical pressure of propane-carbon dioxide;  $O^{34} \Delta^{35}$  [Benedict coeff. for propane<sup>5</sup>; for carbon dioxide: A<sup>36</sup> (set A), B<sup>36</sup> (set B), C<sup>9</sup> (set II), D<sup>11</sup>, E<sup>9</sup> (set I)].


Fig. 22. Critical pressure of propene - 1-butene<sup>37</sup> (Benedict coeff. for propene<sup>5</sup>; 1-butene<sup>38</sup>).



Fig. 23. Critical pressure of butane-carbon dioxide; O<sup>39</sup>, A<sup>35</sup> [Benedict coeff. for butane<sup>5</sup>; carbon dioxide: A<sup>9</sup> (set II), E<sup>9</sup> (set I), B<sup>36</sup> (set A), C<sup>36</sup> (set B), D<sup>11</sup>].



Fig. 24. Critical pressure of pentane-heptane<sup>40</sup> [Benedict coeff. A<sup>5</sup>, B<sup>6</sup> ( $\gamma = \infty$ ), C<sup>6</sup> ( $\gamma = 0.4$ )].









## G. Appendices

## 1. Nomenclature

a,b	coefficients in the equation of Redlich and Kwong
D	= 1/V, density (mole/cm <sup>3</sup> )
f	fugacity
P	pressure (atm)
p ′	limiting slope of critical-pressure line
R	universal gas constant, 82.0567 (cm <sup>3</sup> -atm/mole-°K)
T	temperature (°K)
t	limiting slope of critical-temperature line
v	volume (cm <sup>3</sup> /mole)
y <sub>1</sub> ,y <sub>2</sub>	mole fractions
$A_0, B_0, C_0,$	coefficients for the equation of Benedict, Webb, and Rubin.
a, b, c, α, γ	Primed values indicate derivatives with respect to mole
÷	fraction, i.e., $B_0' = dB_0/dy$ .
	· · ·

Subscripts:

critical value:  $T_c, P_c$ , etc. с

component one in the mixture 1

component two in the mixture 2

. . .

Gas	T	P	D
	с (°К)	c (atm)	(g-mole/cm <sup>3</sup> )
Methane <sup>a</sup>	191.06	45.80	0.010098
Ethane <sup>a</sup>	305.43	48.18	0.006750
Ethylene <sup>b</sup>	282.40	50.50	0.008092
Propane <sup>a</sup>	369,97	42.00	0.004987
Propene <sup>b</sup>	365.00	45.60	0.005537
N-butane <sup>a</sup>	425.17	37.45	0.003922
Isobutane	408.06	36.00	0.003820
l-butene <sup>b</sup>	419.60	39.70	0.004171
N-pentane <sup>b</sup>	469.76	33.30	0.003216
Isopentane <sup>b</sup>	460.96	32.90	0.003243
Heptane <sup>a</sup>	540.17	27.00	0.002345
Decane <sup>a</sup>	619.46	21.19	0.001617
Benzene <sup>b</sup>	561.33	48.30	0.003802
Nitrogen <sup>b</sup>	126.20	33.50	0.011100
Carbon dioxide <sup>b</sup>	304.20	72.90	0.010634
Sulfur dioxide <sup>b</sup>	430.70	77.80	0.008180

2. Critical Data

<sup>a</sup>F.D. Rossini, <u>Selected Values of Physical and Thermodynamic</u> <u>Properties of Hydrocarbons and Related Compounds</u> (Carnegie Press, Pittsburgh, Pa., 1953).

<sup>b</sup>K. A. Kobe and R. E. Lynn Jr., Chem. Rev. <u>52</u>, 117 (1953).

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$\frac{Compound}{Methane} = \frac{Ref.}{8} \frac{10^{-6}A_0}{1.30224} = \frac{B_0}{2.25713} \frac{10^{-7}a}{4.94041} \frac{10^{-3}b}{3.8023} = \frac{10^{-12}C}{10^{-5}a} \frac{10^{-3}}{10^{-3}b} \frac{10^{-12}C}{10^{-5}a} \frac{10^{-3}}{10^{-3}b} \frac{10^{-12}C}{10^{-5}a} \frac{10^{-3}}{10^{-3}b} \frac{10^{-12}C}{10^{-5}a} \frac{10^{-3}}{10^{-3}b} \frac{10^{-3}}{10^{-3}b} \frac{10^{-12}C}{10^{-5}a} \frac{10^{-3}}{10^{-3}b} \frac{10^{-3}}{10^{-3}b} \frac{10^{-3}}{10^{-3}b} \frac{10^{-12}C}{10^{-5}a} \frac{10^{-3}}{10^{-3}b} \frac{10^{-3}}{10^{-3}b} \frac{10^{-12}C}{10^{-5}a} \frac{10^{-3}}{10^{-3}b} \frac{10^{-3}}{10^{-3}b} \frac{10^{-12}C}{10^{-5}a} \frac{10^{-3}}{10^{-3}b} \frac{10^{-3}}{10^{-3}b}$	. /
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Isobutane 5 10.2332 137.548 84.9991 193.779 42.4376 286.034 10.7417 34.0019	
1-butene 38 9.06051 116.028 92.7332 168.214 34.8175 274.943 9.10966 29.5962	
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Isopentane 5 12,7967 160.057 174.642 375.652 66.8157 695.058 17.0014 46.3026	•
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$ Decane  6 \begin{cases} \gamma=0.4 \\ \gamma=\infty \end{cases} \begin{array}{cccc} 25.2325 \\ 19.7449 \end{array} \begin{array}{cccccc} -64.5222 \\ 42.1866 \end{array} \begin{array}{ccccccccccccccccccccccccccccccccccc$	
Benzene 29 6.51013 50.3020 343.016 557.047 76.6343 1176.52 7.00159 29.3016	
$ CO_{2} \begin{array}{c} 11 \\ 9 \\ \begin{cases} A \\ I \\$	3 6 8 6
N2         10         1.05370         40.7437         0.805943         2.51040         2.32782         0.728376         1.27211         5.3002           N2         21         1.19257         45.8013         0.588940         1.49013         1.98165         0.548110         2.91569         7.5004	8
SO <sub>2</sub> 42 2.12054 26.1827 79.3879 84.4395 14.6542 113.362 0.719604 5.9239	0

<sup>a</sup>Constants are used with pressure expressed in atmospheres, temperature in K, and volume in cm<sup>3</sup>/gram-mole

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## PART II. FURTHER IMPROVEMENTS OF AN EQUATION OF STATE

#### A. Introduction

Numerous attempts have been made to formulate an equation of state for gases. This fact indicates both the need for such a relation and the incomplete success of the attempts described in the literature.

The discussion in Part I illustrates a shortcoming of an otherwise excellent relation: The equation of Benedict, Webb, and Rubin<sup>5</sup> contains so many individual coefficients that it becomes a flexible interpolation aid, but it is not very suitable for predictions beyond the range of the data used for the computation of the coefficients.

On the other hand, a two-parameter equation such as that of Redlich and Kwong does not provide sufficient flexibility.<sup>3</sup> It is necessarily limited by the well-known deviations from the theorem of corresponding states.

A successful compromise is the tables of Pitzer and coworkers. 43,44,45 They are based essentially on three individual parameters, and represent the data (with the exception of highly polar substances) with considerable accuracy.

Generalized charts and tables have the important advantage of a potential accuracy higher than that of reasonable algebraic equations;<sup>46</sup> however, they are not very suitable for the most important purpose of an equation of state. Indeed, the practical demand for an equation of state does not stem from a need for representing P-V-T relations themselves. What is really needed is information on fugacity coefficients, particularly those of mixtures. Their derivation from tables is not satisfactory, since numerical integration and still much more numerical differentiation of empirical data severely depresses the accuracy. For these reasons Redlich and Dunlop<sup>47</sup> developed an improvement of the equation of Redlich and Kwong, introducing Pitzer's acentric factor as a third individual coefficient.

In the present attempt an algebraic expression approaching the accuracy of Pitzer's tables is sought. The results of Part I corroborate

the advantage of the choice of the equation of Redlich and Kwong as a starting point. A reasonably good fit in the critical region was thought to be especially desirable, since the equation of Redlich and Kwong leads to large discrepancies in this region, although it is fairly good at much higher pressures.

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B. The Deviation Functions

In this work, we represent  $z^{(0)}$  of Pitzer's tables by the function  $Z_{RK}$  of Redlich and Kwong and a deviation function  $Z_0$  according to

$$z^{(0)} = Z_{\rm RK} + Z_0 , \qquad (1)$$

and  $z^{(1)}$  of Pitzer's tables by an algebraic function

$$\mathbf{z}^{(1)} = \mathbf{Z}_{1} \quad . \tag{2}$$

Thus the compressibility factor for any gas represented in Pitzer's tables by

$$Z = z^{(0)} + \omega \cdot z^{(1)}$$
 (3)

is now represented by the algebraic equations

$$Z = Z_{RK} + Z_0 + \omega \cdot Z_1 .$$
 (4)

Pitzer's acentric factor<sup>44</sup>

$$\omega = \log (P_C/p_S) - 1.00$$
 (5)

is computed from the vapor pressure  $P_S$  at the reduced temperature  $T_r = 0.7$  and the critical pressure  $P_C$ .

The form of the functions  $Z_0$  and  $Z_1$  should be such as to preserve the inherent good features of the original equation. These include the boundary conditions (1)  $Z \rightarrow 1$  as  $T_r \rightarrow \infty$  and (2) the limiting volume is 0.26  $V_C$  for  $P_R \rightarrow \infty$ . Therefore, the functions  $Z_0$ and  $Z_1$  must disappear for  $T \rightarrow \infty$  and be finite for  $P \rightarrow \infty$ . For simplicity and ease of handling, it is desired that  $Z_0$  and  $Z_1$  be functions only of the reduced pressure and temperature. Such additional functions do not interfere with the critical conditions

$$\frac{\partial V}{\partial P} = 0; \quad \frac{\partial^2 V}{\partial P^2} = 0.$$
 (6)

In addition, the functions should be of a form that can be integrated without serious difficulty. A series expansion for the deviation functions is not suitable because of poor convergence and because the boundary conditions cannot be conveniently satisfied. With these considerations in mind, functions  $Z_0$  and  $Z_1$  were evolved, first by inspection, and then by trials of the various forms with the aid of a least-squares computer program to determine the coefficients.

### C. Results

The functions finally adopted are

$$Z_{0} = -A_{1} P_{r}^{3} / \left\{ 1.0 + A_{2} (T_{r} - 1.0)^{2} + A_{3} \left[ P_{r} - A_{4} - A_{5} (T_{r} - 1.0) \right]^{4} \right\}$$
$$+ \left\{ \frac{B_{1} P_{r} \cdot (T_{r} - B_{2} - B_{3} P_{r} + B_{4} P_{r} T_{r}^{2}) (1.0 - B_{5} P_{r} + B_{6} T_{r} P_{r})}{\left[ 1.0 + B_{7} (T_{r} - B_{8} - B_{9} P_{r} - B_{10} P_{r} T_{r})^{4} \right]} \right\}$$
$$+ B_{11} T_{r}^{3} P_{r}^{3} / (T_{r}^{4} + B_{12} P_{r}^{4})$$

$$Z_{1} = \frac{T_{r}P_{r}(T_{r}-1.0-0.049P_{r})(C_{1}+C_{2}P_{r}-C_{3}T_{r}P_{r}+C_{4}T_{r})}{\left[T_{r}^{4}+C_{5}(T_{r}-C_{6}-C_{7}P_{r}+C_{8}T_{r}P_{r})^{4}\right]}$$
(7)

The numerical values for the coefficients are given in Table I. Set I was determined by a least-squares computation of the separate functions based on Pitzer's tables. Set II is a least-squares fit to the data of actual gases (232 points for nitrogen,  $^{49,50}$  methane,  $^{51,52}$  hydrogen sulfide,  $^{53}$  propane,  $^{54}$  carbon dioxide,  $^{55,56}$  sulfur dioxide,  $^{57}$  and water  $^{58,59}$ ). However, Pitzer's tables are based on more extensive experimental material. Therefore we may assume that the coefficients of Set I are more reliable in general application. In the following, the coefficients of Set I are used in all calculations. Table II shows the considerable improvement obtained in the critical region. The improvement in other regions can be seen from Table III which lists deviations from observed compressibility factors for Eq. (7), for Pitzer's tables, and for Dunlop's function. For clarity's sake, deviations smaller than 0.003 are suppressed.

Figures 27-33 compare graphically the deviations of Eq. (7) with those of Pitzer's tables. These values were obtained by means of a computer program of Mr. L. D. Sortland, which interpolates between Pitzer's tabulated data.

		$\begin{array}{ccc} A_{1} & 0.035 \\ A_{2} & 14137.6 \\ A_{3} & 1397.124 \\ A_{4} & 1.030 \\ A_{5} & 13.440 \end{array}$		
		Set I		Set II
в,		0.00260913		0.00476544
B <sub>2</sub>		3.19325		5,15397
B <sub>3</sub>		1.77486		0.344342
B <sub>4</sub>	· · · · · · · · · · · · · · · · · · ·	0.434418		0.422500
B <sub>5</sub>		0.144392		0.155967
B <sub>6</sub>		0.00704658		0.00648025
B <sub>7</sub>		616.830	9	37.761
в <sub>8</sub>		1.00122	· .	0.788177
B <sub>q</sub>		0.0112141		0.0423118
B <sub>10</sub>		0.0495574		0.0436974
B <sub>11</sub>		0.000442593		0.00050277
B <sub>12</sub>		0.0602768		0.0559964
C		0.825714		0.926890
C <sub>2</sub>		0.00736587		0.00691554
$\bar{c_3}$		0.00255204		-0.00339345
C <sub>4</sub>		0.00115729		0.00115571

 $C_4$ 

C<sub>5</sub>

C<sub>6</sub>

C<sub>7</sub>

с<sub>8</sub>

Table I. Coefficients for Eq. (7)

-44-

0.101212

2.46596

0.220411

0.0161963

0.109996

1.28600

1.03462

.

0.0161573

					<u></u>	Z <sub>c</sub>	· · · · · ·	
Substance	T	Р	ω	Observed	- <u></u>	Calc	ulated	
	<u>(°K)</u>	(atm)			Pitzer	Redlich-Kwong	Redlich-Dunlop	Present work
N <sub>2</sub>	126.3	33.54	0.040	0.292	0.288	0.333	0.339	0.290
CH <sub>4</sub>	191.0	45.79	0.013	0.290	0.290	0.333	0.334	0.290
н <sub>2</sub> s	373.6	88.87	0.100	0.2833	0.283	0.333	0.315	0.288
C <sub>3</sub> H <sub>8</sub>	370.0	42.01	0.152	0.2766	0.279	0.333	0.327	0.287
co <sub>2</sub>	304.2	72.80	0.225	0.2746	0.273	0.333	0.323	0.285
so <sub>2</sub>	430.66	77.808	0.2325	0.2697	0.272	0.333	0.314	0.285
н <sub>2</sub> о	647.3	218.4	0.348	0.2276	0.263	0.333	0.236	0,283
								. ,

Table II. Results for the critical point

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			NI	TROGEN		
TR	PR		Z(OBS)	<b>∆Z(EQ 7)</b>	$\Delta Z(PITZER)$	$\Delta Z (DUNLOP)$
1.0000	1.0000		0.2923	·····	0.0043	-0.0491
2.0227	0.3043		0.9933			
2.0227	0.6086		0.9874 :		2	
2.0227	1.5216		0.9731	1		0.0042
2.0227	3.0430		0.9660			0.
2.0227	6.0860		1.0183	-0.0051		-0.0128
2.9027	0.3043		1.0017		· ·	
2.9027	0.6086		1.0036			
2.9027	1.5216		1.0105			
2.9027	3.0430	-	1.0272		-0.0070	· .
2.9027	6.0860	. :	1.0845	-0.0054	-0.0064	•
3.7827	0.3043		1.0044			
3.7827	0.6086	• •	1.0090	0.0030		,
3.7827	1.5216	:	1.0229	0.0067	· · ·	0.0040
3.7827	3.0430		1.0469	0.0087		0.0055
3.7827	6.0860	1	1.1022			0.0067
4.6627	0.3043		1.0040			
4.6627	0.6086	1.	1.0094		0.	· . ·
4.6627	1.5216		1.0240	0.0065	0.0076	0.0032
4.6627	3.0430		1.0474	0.0088	0.0239	0.0041
4.6627	6.0860	:	1.0979		0.0292	0.0048
5.5427	0.3043		1.0044		0.0046	
5.5427	0.6086		1.0090			
5.5427	1.5216		1.0231	0.0062	0.0257	
5.5427	3.0430		1.0441	0.0076	0.0784	,
5.5427	6.0860		1.0891	•	0.0916	
0.8346	0.1014		0.9380	·	0.0068	
0.8346	0.2029		0.8708		0.0118	
0.8346	0.3043		0.7878	-0.0034	0.0038	-0.0064
1.0106	0.1014		0.9670			· ·
1.0106	0.2029		0.9319	0.0034		
1.0106	0.3043		0.8949	0.0049		
1.0106	0.4058		0.8518	· · · · · ·	· ·	
1.0106	0.8118		0.6515	0.0088	0.0054	-0.0035
1.0106	1.2173		0.2425		0.0037	-0.0441
1.0106	1.6230		0.2845	-0.0085		-0.0458
1.0106	2.0288		0.3328	-0.0067		-0.0569
1.2306	0.4058		0.9299	0.0063		0.0055
1.2306	0.8115	1 T	0.8526	0.0104		0.0063
1.2306	1.2173		0.7656	0.0082		
1.2306	1.6230		0.6826	0.0067		-0.0108
1.2306	2.0288		0.6155			-0.0241
1.4507	0.4058	:	0.9600			
1.4507	0.8115	:	0.9204	0.0034		0.0043
1.4507	1.2173	1 A A A A A A A A A A A A A A A A A A A	0.8799		-0.0033	
1.4507	1.6230		0.8459		-0.0038	
1.4507	2.0288		0.8148			-0.0037

Table III. Deviations from the observed compressibility factor. (Deviations below 0.003 are omitted)

$\begin{array}{c ccccccccccccccccccccccccccccccccccc$	2(DUNLOP) -0.0477 -0.0128 -0.0499
TRPR $Z(OBS)$ $\Delta Z(EQ 7)$ $\Delta Z(PITZER)$ $\Delta Z$ 1.00001.00000.29030.97951.62800.29700.97951.62801.48600.90721.62802.97000.8527-0.0096	2(DUNLOP) -0.0477 -0.0128 -0.0499
1.0000         1.0000         0.2903           1.6280         0.2970         0.9795           1.6280         1.4860         0.9072           1.6280         2.9700         0.8527         -0.0096	-0.0477 -0.0128 -0.0499
1.6280       0.2970       0.9795         1.6280       1.4860       0.9072         1.6280       2.9700       0.8527       -0.0096	-0.0128 -0.0499
1.6280 1.4860 0.9072 1.6280 2.9700 0.8527 -0.0096	-0.0128
1.6280 2.9700 0.8527 -0.0096	-0.0128 -0.0499
	-0.0499
1.6280 7.4300 0.9812 -0.0128	0 0070
1.8020 0.2970 0.9862	0.0072
1.8020 1.4860 0.9410 -0.0032	0 0073
1.8020 2.9700 0.9108 -0.0062	-0.0072
1.8020 7.4300 1.0120 -0.0099	-0.0353
1.9760 0.2970 0.9913	
1.9760 1.4860 0.9639	
1.9760 2.9700 0.9496	-0.0036
1.9760 7.4300 1.0395 -0.0071 -0.0037	-0.0248
2.1500 0.2970 0.9947 -	
2.1500 1.4860 0.9788 -0.0039	
2.1500 2.9700 0.9758 -0.0047	
2.1500 7.4300 1.0610 -0.0054 -0.0049	-0.0182
2.3250 0.2970 0.9972	
2.3250 1.4860 0.9908 0.0031 -0.0033	
2.3250 2.9700 0.9949 -0.0051	
2.3250 7.4300 1.0786 -0.0042	-0.0126
2.5000 0.2970 0.9989	
2.5000 1.4860 0.9986 0.0035	
2.5000 2.9700 1.0079	
2.5000 7.4300 1.0916 -0.0031	-0.0089
2.6740 0.2970 1.0002	
2.6740 1.4860 1.0043 0.0037	
2.6740 2.9700 1.0177 0.0040	
2.6740 7.4300 1.1008 -0.0040	-0.0065

				(continueu)		
			HYDROGE	N SULFIDE		
2. · · ·	TR	PR	Z(OBS)	△Z(EQ 7)	$\Delta Z(PITZER)$	$\Delta z$ (DUNLOP)
	1.0000	1.0000	0.2836	-0.0050		-0.0426
	0.7431	0.0766	0.9219	-0.0138	-0.0035	-0.0114
	0.8323	0.0766	0.9503	-0.0032	0.0041	
· .	0.8323	0.1531	0.8984	-0.0049	0.0081	-0.0039
	0.8323	0.7657	0.1199			0.0129
	0.9215	0.0766	0.9650		-0.0036	,
	0.9215	0.1531	0.9290		-0.0046	
	0.9215	0.3829	0.8038			
· · · ·	0.9215	0.7657	0.1265	-0.0075		-0.0069
	0.9215	1.5314	0.2378	-0.0036		0.0103
	1.0107	0.0766	0.9742			
	1.0107	0.1531	0.9500	0.0032		
	1.0107	0.3829	0.8633	0.0048		
	1.0107	0.7657	0.6775	0.0059	0.0034	
	1.0107	1.5314	0.2648	-0.0143		-0.0198
	1.0107	3.8290	0.5444			0.0685
	1.1000	0.0766	0.9803			
	1.1000	0.1531	0.9612			
	1.1000	0.3829	0.9004	0.0059		0.0044
	1.1000	0.7657	0.7828	0.0103		0.0074
	1.1000	1.5314	0.4638		-0.0076	
	1.1000	3.8290	0.5606			0.0259
	1.1891	0.0766	0.9851		•	
	1.1891	0.1531	0.9705			
	1.1891	0.3829	0.9251	0.0049		0.0047
	1.1891	0.7657	0.8441	0.0108	·	0.0098
	1.1891	1.5314	0.6561	0.0110	-0.0037	0.0080
	1.1891	3.8290	0.5965			

Table III. (continued)

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		Table III.	(continued)					
PROPANE								
TR	PR	Z(OBS)	ΔZ(EQ 7)	$\Delta Z(PITZER)$	AZ (DUNLOP)			
1.0000	1.0000	0.2769	-0.0105		-0.0622			
0.8400	0.1620	0.8964		0.0117				
0.8400	0.8100	0.1226			0.0083			
0.8850	0.1620	0.9143		0.0036				
0.8850	0.3240	0.8068	-0.0071		-0.0064			
0.8850	0.8100	0.1240	-0.0046		-0.0065			
0.8850/	1.6200	0.2503	0.0085	0.0079	-0.0050			
0.9300	0.1620	0.9272						
0.9300	0.3240	0.8408			-0.0041			
0.9300	0.8100	0.1291	0.0124		-0.0207			
0.9300	1.6200	0.2444	-0.0059		-0.0320			
0.9740	0.1620	0.9376						
0.9740	0.3240	0.8665			<u>х</u>			
0.9740	0.8100	0.5400	-0.0142	-0.0169	-0.0272			
0.9740	1.6200	0.2527	-0.0136		-0.0470			
0.9740	3.2400	0.4604			-0.0925			
1.0200	0.1620	0.9460		-0.0032				
1.0200	0.3240	0.8870						
1.0200	0.8100	0.6606		-0.0039	-0.0114			
1.0200	1.6200	0.2751	-0.0164	-0.0034	-0.0592			
1.0200	3.2400	0.4678			-0.1003			
1.2000	0.1620	0.9683						
1.2000	0.3240	0.9360						
1.2000	0.8100	0.8365	0.0045	-0.0038				
1.2000	1.6200	0.6612	0.0162		-0.0094			
1.2000	3.2400	0.5647	-0.0060	-0.0057	-0.0911			
1.3800	0.1620	0.9808						
1.3800	0.3240	0.9620		1				
1.3800	0.8100	0.9065		-0.0059				
1.3800	1.6200	0.8219		-0.0111	-0.0052			
1.3800	3.2400	0.7347		0.0037	-0.0454			
1.3800	8.1000	1.0051		0.0063	-0.1934			

Table III. (continued)							
<u></u>	•	CARBON	DIOXIDE				
TR	PR	Z(OBS)	∆Z(EQ 7)	△Z(PITZER)	<b>∆Z</b> ( DUNLOP		
1.0000	1.0000	0.2749	-0.0108		-0.0551		
1.0400	0.1863	0.9443	0.0035				
1.0400	0.9310	0.6193	-0.0035	-0.0138	-0.0166		
1.1310	0.1863	0.9577					
1.1310	0.9310	0.7581	0.0087	-0.0103	-0.0071		
1.1310	1.8630	0.4785			-0.0358		
1.2410	0.1863	0.9702			•		
1.2410	0.9310	0.8377		-0.0042	-0.0099		
1.2410	1.8630	0.6837	0.0174	-0.0037	-0.0163		
1.2410	4.6600	0.6928	0.0043		-0.1125		
1.3510	0.1863	0.9793					
1.3510	0.9310	0.8890		-0.0057	-0.0073		
1.3510	1.8630	0.7873		-0.0134	-0.0189		
1.3510	4.6600	0.7587	-0.0044	0.0056	-0.1047		
1.4420	0.1863	0.9843					
1.4420	0.9310	0.9194		-0.0032	-0.0040		
1.4420	1.8630	0.8449	-0.0069	-0.0134	-0.0166		
1.4420	4.6600	0.8166	-0.0080	0.0063	-0.0881		
1.5700	0.1863	0.9886					
1.5700	0.9310	0.9459	-0.0041	•	-0.0035		
1.5700	1.8630	0.8987	-0.0125	-0.0091	-0.0141		
1.5700	4.6600	0.8819	-0.0236	-0.0037	-0.0708		
1.6790	0.1863	0.9930					
1.6790	0.9310	0.9627	-0.0042				
1.6790	1.8630	0.9307	-0.0138	-0.0079	-0.0111		
1.6790	4.6600	0.9249	-0.0386	-0.0087	-0.0592		

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		Table III.	(continued)		
		SULFUR	DIOXIDE		
TR	PR	Z(OBS)	$\Delta Z(EQ 7)$	$\Delta Z(PITZER)$	△Z(DUNLOP)
1.0000	1.0000	0.2700	-0.0156	-0.0035	-0.0592
0.6580	0.0257	0.9532	-0.0173	-0.0039	-0.0053
0.6810	0.0257	0.9608	-0.0123		
0.7040	0.0257	0.9667	-0.0087		
0.7040	0.0514	0.9313	-0.0185	-0.0048	-0.0033
0.7280	0.0257	0.9710	-0.0066		
0.7280	0.0514	0.9407	-0.0135		
0.7280	0.0771	0.9082	-0.0218	-0.0054	-0.0039
0.7500	0.0257	0.9742	-0.0051		
0.7500	0.0514	0.9478	-0.0101		
0.7500	0.0771	0.9206	-0.0151		
0.7500	0.1027	0.8905	-0.0222		-0.0032
0.8670	0.0771	0.9541	-0.0035		
0.8670	0.1543	0.9043	-0.0080		
0.8670	0.2310	0.8495	-0.0140		-0,0044
0.8670	0.3090	0.7843	-0.0248	-0.0054	-0.0117
0.9820	0.0771	0.9702			
0.9820	0.1543	0.9396			
0.9820	0 2310	0.9080			
0.0020	0.2000	0 8743			-0.0031
0.9820	0.3860	0 8387			-0.0041
0.9020	0.4430	0.0001	-0 0037		-0.0059
0.9020	0.5400	0 7588	-0.0045	-0.0034	-0.0075
0.9020	0.6170	0 7133	-0.0054	-0.0047	-0.0096
0.9020	0.6040	0.6608	-0.0078	-0.0041	-0.0133
0.9820	0.7710	0.5093	-0.0116	-0.0073	-0.0184
0.9020	0.9490	0.5110	-0 0223	-0.0224	-0.0304
1 0000	0 2570	0.9036	-0.0225	0.0224	0.0304
1.0000	0 5140	0.7914		-0.0032	-0.0057
1.0000	0.7710	0.6431	-0.0038	-0.0051	-0.0125
1.0000	2 0600	0.3107	-0.0131		-0.0284
1.0000	2.2100	0.3335	-0.0175	-0-0056	-0.0323
1.0000	2 5700	0.3630	-0.0166	-0.0078	-0.0298
1.0000	2.9700	0.3049	-0.0137	+0.0077	-0.0243
1.0000	2.0000	0 4221	-0.0147	-0.0110	-0.0214
1.0000	3.3400	0.4519	-0 0143	-0.0127	-0.0164
1.0000	3.5400	0.4910		-0.0159	-0 0119
1.0000	3 9600	0.5080	-0.0163	-0.0181	-0.0064
1.1000	9.0000	0.9216	0.0048	0.0101	-0.0055
1 1000	1 2850	0.5912	0.0247		0.0000
1.1000	1 9270	0.4000	-0.0143		-0.0545
1.1000	2 5700	0 4188	-0.0209	-0.0067	-0.0793
1.1000	3.2100	0.4713	-0.0213	-0,0169	-0.0890
1 1000	3 8600	0.5240	-0.0185	-0,0170	-0-0901
1 2150	0 6420	0.0047	0.0044	040110	-0-0038
1 2150	1 2950	0.0071	0.0185	0,0044	-0.0035
1 2150	1 0270	0 6461	0.0261		-0.0128
1 2160	2.5700	0.5817	0.0062	-0,0056	-0.0514
1 3160	2 2100	0.5910	-0.0055	-0.0057	-0.0834
1 2150	3 9600	0.2017	-0.0042	0.0001	-0.1025
1.2120	0000 ·	0.0101	0.0072		0 • 1 V 2 J

Table III. (continued)

		Table III.	(continued)		
		W	ATER		
TR	PR	Z(OBS)	$\Delta Z(EQ 7)$	$\Delta Z(PITZER)$	
1.0000	1.0000	0.2279	-0.0550	-0.0370	-0.0628
0.6000	0.0010	0.9973		0.0331	
0.6000	0.0020	0.9946		0.0317	
0.6000	0.0030	0.9919	-0.0037	0.0302	
0.6000	0.0040	0.9893	-0.0049	0.0289	
0.6000	0.0050	0.9866	-0.0061	0.0274	0.0030
0.6000	0.0060	0.9860	-0.0053	0.0281	0.0057
0.7000	0.0100	0.9857	-0.0045	0.0084	
0.7000	0.0200	0.9715	-0.0088	0.0060	0.0055
0.7000	0.0300	0.9565	-0.0137		0.0077
0.7000	0.0400	0.9402	-0.0198		0.0089
0.8000	0.0400	0.9671	-0.0051	0.0039	0.0052
0.8000	0.0800	0.9312	-0.0122	0.0052	0.0088
0.8000	0.1200	0.8893	-0.0237		0.0083
0.8000	0.1600	0.8393	-0.0419	-0.0115	
0.9000	0.1000	0.9472			0.0048
0.9000	0.2000	0.8858	-0.0095	-0.0042	0.0061
0.9000	0.3000	0.8121	-0.0237	-0.0111	0.
0.9000	0.4000	0.7207	-0.0479	-0.0275	-0.0119
1.0000	0.2000	0.9268			
1.0000	0.4000	0.8407	-0.0041	-0.0048	
1.0000	0.6000	0.7397	-0.0091	-0.0104	
1.0000	0.8000	0.6045	-0.0189	-0.0191	-0.0064
1.1000	0.2000	0.9480			-0.0060
1.1000	0.4000	0.8931			-0.0092
1.1000	0.6000	0.8349		-0.0037	-0.0092
1.1000	0.8000	0.7730	0.0071	-0.0053	-0.0052
1.1000	1.0000	0.7062	0.0148	-0.0040	0.0033
1.1000	1.2000	0.6287	0.0214	-0.0064	0.0123
1.2000	0.2000	0.9599	-0.0046	-0.0042	-0.0106
1.2000	0.4000	0.9224	-0.0043	-0.0043	-0.0156
1.2000	0.6000	0.8853		-0.0040	-0.0173
1.2000	0.8000	0.8478	0.0035	-0.0038	-0.0163
1.2000	1.0000	0.8098	0.0100		-0.0130
1.2000	1.2000	0.7710	0.0171		-0.0081
1.3000	0.2000	0.9665	-0.0091	-0.0084	-0.0144
1.3000	0.4000	0.9399	-0.0102	-0.0090	-0.0203
1.3000	0.6000	0.9158	-0.0080	-0.0075	-0.0224
1.3000	0.8000	0.8929	-0.0037	-0.0058	-0.0220
1.3000	1.0000	0.8704		-0.0035	-0.0201
1.3000	1,2000	0.8482	0.0076		-0.0171

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Fig. 27. Deviations from the observed compressibility factor for nitrogen.



Fig. 28. Diviations from the observed compressibility factor for methane.







Fig. 30. Deviations from the observed compressibility factor for propane.











Fig. 33. Deviations from the observed compressibility factor for water.

Considerable improvement over Dunlop, especially in the critical region, is seen in Tables II and III. For water (Table III) Dunlop's values are better. This is to be expected from the manner of choosing the respective functions and the acknowledged poor representation of water by the Pitzer tables. However, even here, Dunlop's values are not significantly better.

The diagrams show that Pitzer's tables give a better representation than Eq. (7). But the algebraic representation comes very close to the results of the tables. As Fig. 27 shows, use of Eq. (7) beyond the range of the tables gives very good results. The data given for nitrogen in Table III at the reduced temperatures 4.6627 and 5.5427 under Z (Pitzer) are extrapolated, since Pitzer's tables do not cover these high temperatures. The results of this direct extrapolation are poor.

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# D. Appendix: Nomenclature

A, B, C	Constants in the deviation functions
P	Pressure
P	Critical pressure
Pr	Reduced pressure
Ps	Vapor pressure at a reduced temperature $T_r = 0.7$
Т	Temperature
T <sub>r</sub>	Reduced temperature
V	Volume
V <sub>c</sub>	Critical volume
Z	PV/RT, the compressibility factor
$z^{(0)}, z^{(1)}$	Values from Pitzer's tables
Z <sub>RK</sub>	Compressibility factor as calculated by using the equation
	of Redlich and Kwong
z <sub>0</sub> , z <sub>1</sub>	Deviation functions
ω	Pitzer's acentric factor

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