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Electron spin resonance study of the local environment for the Gd^{3+} and Eu^{2+} ions in $\text{Ca}_{1-x}\text{R}_x\text{B}_6$ ($\text{R} = \text{Gd}, \text{Eu}$) ($0.0001 \leq x \leq 0.30$)

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Abstract

The environment of $\text{Gd}^{3+}/\text{Eu}^{2+}$ ($4f^7, S = \frac{7}{2}$) in $\text{Ca}_{1-x}\text{R}_x\text{B}_6$ ($\text{R} = \text{Gd}, \text{Eu}; 0.0001 \leq x \leq 0.30$) is studied by electron spin resonance (ESR). For $x \leq 0.001$ the spectra show Lorentzian shape (insulating phase). As x increases, the spectra present a superposition of Lorentzian and Dysonian resonances (coexistence of insulating and metallic phases). For $x \geq 0.01$, the line shape becomes pure Dysonian (metallic phase). Thus, the intermediate concentration regime of $\text{Ca}_{1-x}\text{R}_x\text{B}_6$ is intrinsically inhomogeneous. These compounds show no weak ferromagnetism.

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1. Introduction

The cubic system $\text{Ca}_{1-x}\text{R}_x\text{B}_6$ ($\text{R} = \text{rare earths}$; space group 221, $\text{Pm}\bar{3}\text{m}$), specially La, was extensively investigated since the reported weak-ferromagnetism (WF) at $T_C \sim 600\text{--}800\text{ K}$ by Young et al. [1]. Later reports on strong sample dependent WF [2] and doubts about its intrinsic nature were raised [3]. It was argued that the intrinsic WF could be hidden by the FM of Fe and Ni impurities [4] and that CaB_6 is a $\sim 1\text{ eV}$ -gap semiconductor. Hence, an ESR study, probing the local $\text{R} = \text{Gd}^{3+}/\text{Eu}^{2+}$ environment is of great interest to understand the magnetic/non-magnetic and metallic/non-metallic properties in $\text{Ca}_{1-x}\text{R}_x\text{B}_6$.

2. Experiments

Single crystals of $\sim 1 \times 0.5 \times 0.3\text{ mm}^3$ of $\text{Ca}_{1-x}\text{R}_x\text{B}_6$ ($\text{R} = \text{Gd}, \text{Eu}$) ($0.0001 \leq x \leq 0.30$) were grown as described

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in Ref. [1]. The structure, phase purity and orientation were checked by powder and Laue X-ray diffraction. The ESR experiments were done in a Bruker X-band (9.479 GHz) spectrometer with a TE_{102} room-T cavity coupled to a T-controller of helium gas flux system for $4.2 \leq T \leq 300\text{ K}$. The $\text{Gd}^{3+}/\text{Eu}^{2+}$ concentration was obtained from Curie–Weiss fits of $\chi(T)$ ($2 \leq T \leq 300\text{ K}$) measured in a Quantum Design SQUID DC-magnetometer.

3. Experimental results

Fig. 1 shows the room-T Eu^{2+} ESR for $\text{H} \parallel [001]$, $x = 0.023, 0.07$ and 0.30 crystals. The spectra for $x = 0.023$ and 0.07 were simulated by the superposition of two Eu^{2+} spectra: a resolved fine structure of Lorentzian resonances (fsL) corresponding to Eu^{2+} ions in an insulating phase [5] and a single Dysonian (D) resonance associated to Eu^{2+} ions in a metallic phase [6]. For $x = 0.30$ the spectrum was simulated with a D resonance. The insets show the spectra for (a) three crystals of the same $x = 0.03$ batch, and (b) an as-grown and 950°C five days annealed/

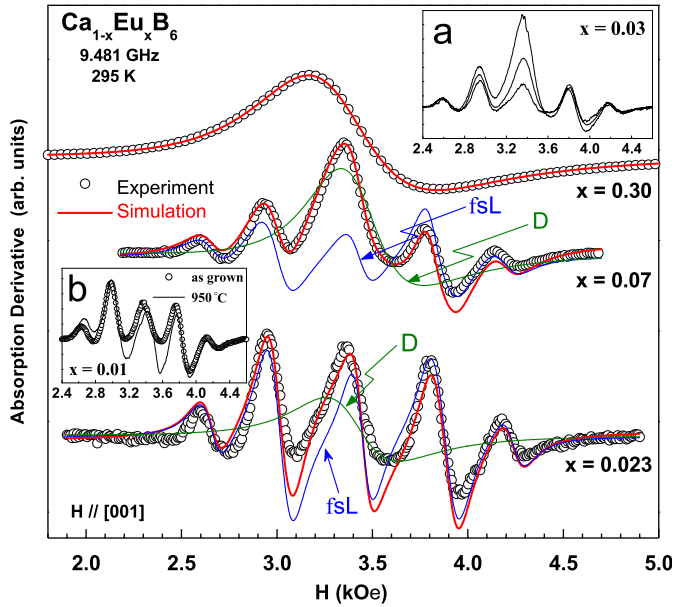


Fig. 1. ESR of Eu^{2+} . The fsL (blue) spectrum corresponds to the seven resonances in a cubic insulating phase. The D resonance corresponds to the metallic phase. For $x = 0.023$ and 0.07 , S (red) are the superposition of fsL and D. For $x = 0.30$ S is given by a D resonance. Insets: spectra for (a) three crystals of the same $x = 0.03$ batch, and (b) an as-grown and 950°C five days annealed/liquid N_2 quenched $x = 0.01$ crystal.

liquid N_2 quenched $x = 0.01$ crystal. These results confirm the highly metastable inhomogeneous coexistence of two different local environments for the Eu^{2+} ions in the range of $0.023 \lesssim x \lesssim 0.15$.

Fig. 2 displays the Gd^{3+} ESR at 4.2 K for $\text{H} \parallel [001]$, $x = 0.0013$ and 0.01 crystals. The spectra were simulated by the superposition of two Gd^{3+} spectra: a resolved fsL resonances corresponding to Gd^{3+} ions in an insulating phase and a single D resonance associated to Gd^{3+} ions in a metallic phase. For $x = 0.01$ the spectrum was simulated with a D resonance. Similarly to the Eu doped crystal, these results confirm the coexistence of two different local environments for the Gd^{3+} ions in the range of $0.001 \lesssim x \lesssim 0.003$.

4. Analysis and discussion

The ESR of $\text{R} = \text{Gd}^{3+}/\text{Eu}^{2+}$ in $\text{Ca}_{1-x}\text{R}_x\text{B}_6$ show three concentration regimes. For low x the line shapes are Lorentzian, thus, the environment for $\text{Gd}^{3+}/\text{Eu}^{2+}$ ions is insulating and the fine structure can be resolved. The isotropic g -value and anisotropy of the fine structure [7,8] indicate that the $\text{Gd}^{3+}/\text{Eu}^{2+}$ local symmetry is cubic. For intermediate x the ESR present a superposition of fsL and D resonances. The line shape starts to show Dysonian shape (metallic phase), i.e., the microwave skin-depth is comparable to the size of our crystals. For higher x the resonance is Dysonian (metallic phase) and the fine structure can no longer be resolved. The g -value and line width are T-independent down to ~ 10 K [8], indicating that there are no direct/indirect magnetic interactions between the R ions. For small x each Eu^{2+} in $\text{Ca}_{1-x}\text{R}_x\text{B}_6$ gives rise

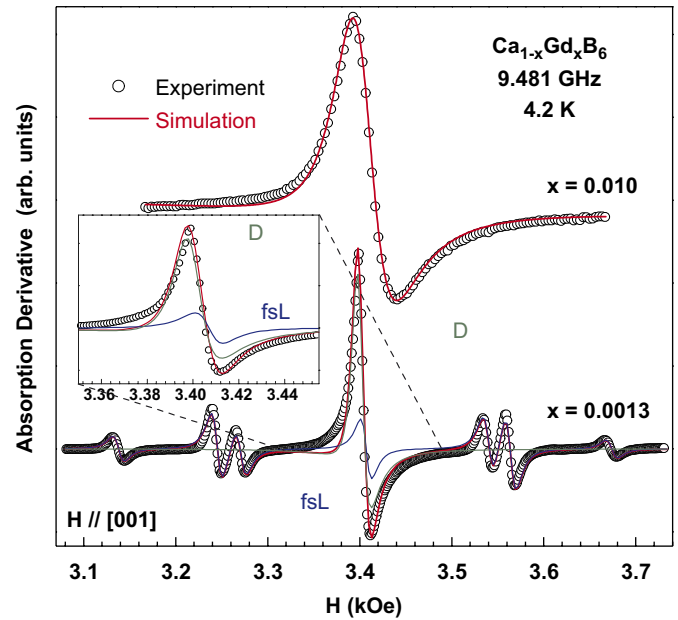


Fig. 2. ESR of Gd^{3+} . The fsL (blue) spectrum corresponds to the seven resonances in a cubic insulating phase. The D (green) resonance corresponds to the metallic phase. For $x = 0.0013$, S (red) are the superposition of fsL and D. For $x = 0.01$, S is given by a D resonance. The inset is a blow up of the central part of the spectrum.

to an impurity bound state in the semiconductor gap which is localized within about one unit cell. As x increases the number of impurity states increases, starts to overlap, and eventually form a percolative network. Our data indicate coexistence of insulating and metallic phases for Gd and Eu concentrations of $x \approx 0.001$ and 0.02 , respectively. These values are well below the percolation threshold for nearest neighbor (nn), next-to-nearest neighbors (nnn) and next to nnn in a simple cubic lattice ($x = 0.307, 0.137$ and 0.099). The Dysonian line shape indicates that the size of the Eu/Gd rich regions should be of the order of the skin-depth which is about $1 \mu\text{m}$ for the resistivity of pure EuB_6 and GdB_6 . According to Ref. [9] we may associate the fsL and D spectra with regions rich in Ca^{2+} and Eu^{2+} , respectively.

5. Conclusions

In $\text{Ca}_{1-x}\text{R}_x\text{B}_6$ ($\text{R} = \text{Eu}/\text{Gd}$), as a function of x , an evolution from insulating to a metallic character of the compound is verified from the change of the $\text{Eu}^{2+}/\text{Gd}^{3+}$ ESR line shapes. The percolative transition between these two regimes is estimated at $x \sim 0.14$ for Eu doped samples, indicating that nnn bonds contribute to the percolative network. For lower x , a highly metastable inhomogeneous coexistence of insulating and metallic phases is observed. This coexistence was also found in $\text{Ca}_{1-x}\text{Gd}_x\text{B}_6$, however, due to the Gd^{3+} doping extra electron, the percolative interval is found at a much lower level of Gd concentration ($x \sim 0.0015$). All measured $\text{Ca}_{1-x}\text{R}_x\text{B}_6$ ($\text{R} = \text{Eu}, \text{Gd}$) crystals presented a WF ≤ 0.5 emu/mol, i.e., much smaller than that of La doped CaB_6 [1].

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