

Lawrence Berkeley National Laboratory

LBL Publications

Title

Solar Neutrino Observations at the Sudbury Neutrino Observatory

Permalink

<https://escholarship.org/uc/item/59s3h8b7>

Author

Poon, AWP

Publication Date

2002-11-05

Peer reviewed

SOLAR NEUTRINO OBSERVATIONS AT THE SUDBURY NEUTRINO OBSERVATORY

A.W.P. Poon

Institute for Nuclear and Particle Astrophysics

Lawrence Berkeley National Laboratory, Berkeley, CA 94720

Representing the Sudbury Neutrino Observatory Collaboration*

*Q.R. Ahmad, R.C. Allen, T.C. Andersen, J.D. Anglin, J.C. Barton, E.W. Beier, M. Bercovitch, J. Bigu, S.D. Biller, R.A. Black, I. Blevis, R.J. Boardman, J. Boger, E. Bonvin, M.G. Boulay, M.G. Bowler, T.J. Bowles, S.J. Brice, M.C. Browne, T.V. Bullard, G. Bühler, J. Cameron, Y.D. Chan, H.H. Chen, M. Chen, X. Chen, B.T. Cleveland, E.T.H. Clifford, J.H.M. Cowan, D.F. Cowen, G.A. Cox, X. Dai, F. Dalnoki-Veress, W.F. Davidson, P.J. Doe, G. Doucas, M.R. Dragowsky, C.A. Duba, F.A. Duncan, M. Dunford, J.A. Dunmore, E.D. Earle, S.R. Elliott, H.C. Evans, G.T. Ewan, J. Farine, H. Fergani, A.P. Ferraris, R.J. Ford, J.A. Formaggio, M.M. Fowler, K. Frame, E.D. Frank, W. Frati, N. Gagnon, J.V. Germani, S. Gil, K. Graham, D.R. Grant, R.L. Hahn, A.L. Hallin, E.D. Hallman, A.S. Hamer, A.A. Hamian, W.B. Handler, R.U. Haq, C.K. Hargrove, P.J. Harvey, R. Hazama, K.M. Heeger, W.J. Heintzelman, J. Heise, R.L. Helmer, J.D. Hepburn, H. Heron, J. Hewett, A. Hime, J.G. Hykawy, M.C.P. Isaac, P. Jagam, N.A. Jelley, C. Jillings, G. Jonkmans, K. Kazkaz, P.T. Keener, J.R. Klein, A.B. Knox, R.J. Komar, R. Kouzes, T. Kutter, C.C.M. Kyba, J. Law, I.T. Lawson, M. Lay, H.W. Lee, K.T. Lesko, J.R. Leslie, I. Levine, W. Locke, S. Luoma, J. Lyon, S. Majerus, H.B. Mak, J. Maneira, J. Manor, A.D. Marino, N. McCauley, D.S. McDonald, A.B. McDonald, K. McFarlane, G. McGregor, R. Meijer, C. Mifflin, G.G. Miller, G. Milton, B.A. Moffat, M. Moorhead, C.W. Nally, M.S. Neubauer, F.M. Newcomer, H.S. Ng, A.J. Noble, E.B. Norman, V.M. Novikov, M. O'Neill, C.E. Okada, R.W. Ollerhead, M. Otori, J.L. Orrell, S.M. Oser, A.W.P. Poon, T.J. Radcliffe, A. Roberge, B.C. Robertson, R.G.H. Robertson, S.S.E. Rosendahl, J.K. Rowley, V.L. Rusu, E. Saettler, K.K. Schaffer, M.H. Schwendener, A. Schülke, H. Seifert, M. Shatkay, J.J. Simpson, C.J. Sims, D. Sinclair, P. Skensved, A.R. Smith, M.W.E. Smith, T. Spreitzer, N. Starinsky, T.D. Steiger, R.G. Stokstad, L.C. Stonehill, R.S. Storey, B. Sur, R. Tafirout, N. Tagg, N.W. Tanner, R.K. Taplin, M. Thorman, P.M. Thornewell, P.T. Trent, Y.I. Tserkovnyak, R. Van, R.G. Van, C.J. Virtue, C.E. Waltham, J.-X. Wang, D.L. Wark, N. West, J.B. Wilhelm, J.F. Wilkerson, J.R. Wilson, P. Wittich, J.M. Wouters, M. Yeh

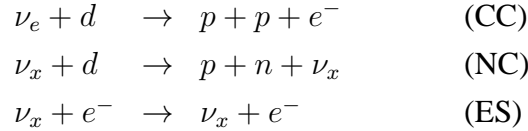
ABSTRACT

The Sudbury Neutrino Observatory (SNO) is a 1000-tonne heavy water Cherenkov detector. Its usage of D_2O as target allows the simultaneous measurements of the ν_e flux from 8B decay in the Sun and the total flux of all active neutrino species through the charged-current and the neutral-current interactions on the deuterons. Assuming the standard 8B shape, the ν_e component of the 8B solar neutrino flux is measured to be $\phi_e = 1.76_{-0.05}^{+0.05}(\text{stat.})_{-0.09}^{+0.09}(\text{syst.}) \times 10^6 \text{ cm}^{-2}\text{s}^{-1}$ for a kinetic energy threshold of 5 MeV. The non- ν_e component is found to be $\phi_{\mu\tau} = 3.41_{-0.45}^{+0.45}(\text{stat.})_{-0.45}^{+0.48}(\text{syst.}) \times 10^6 \text{ cm}^{-2}\text{s}^{-1}$. This 5.3σ difference provides strong evidence for ν_e flavor transformation in the solar neutrino sector. The total active neutrino flux is measured with the neutral-current reaction at a neutrino energy threshold of 2.2 MeV. This flux is determined to be $\phi_{\text{NC}} = 5.09_{-0.43}^{+0.44}(\text{stat.})_{-0.43}^{+0.46}(\text{syst.}) \times 10^6 \text{ cm}^{-2}\text{s}^{-1}$, and is consistent with solar model predictions. Assuming an undistorted 8B spectrum, the night minus day rate is $14.0 \pm 6.3(\text{stat.})_{-1.4}^{+1.5}(\text{syst.})\%$ of the average rate in the charged-current channel. If the total active neutrino flux is constrained to have no asymmetry, the night-day asymmetry in the ν_e flux is found to be $7.0 \pm 4.9(\text{stat.})_{-1.2}^{+1.3}(\text{syst.})\%$. A global analysis of all the available solar neutrino data in terms of matter-enhanced oscillations of two active flavors strongly favors the Large Mixing Angle (LMA) solution.

1 Introduction

For more than 30 years, solar neutrino experiments^{1,2,3,4,5,6} have been observing fewer neutrinos than what are predicted by the detailed models^{7,8} of the Sun. This deficit is known as the Solar Neutrino Problem. A comparison of the predicted and observed solar neutrino fluxes for these experiments are shown in Table 1. These experiments probe different parts of the solar neutrino energy spectrum, and show an energy dependence in the observed solar neutrino flux. These observations can be explained if the solar models are incomplete or neutrinos undergo flavor transformation while in transit to the Earth.

The Sudbury Neutrino Observatory⁹ was constructed to resolve this solar neutrino puzzle. It is capable of making simultaneous measurements of the electron-type neutrino (ν_e) flux from ${}^8\text{B}$ decay in the Sun and the flux of all active neutrino flavors through the following reactions:



The charged-current (CC) reaction on the deuteron is sensitive exclusively to ν_e , and the neutral-current (NC) reaction has equal sensitivity to all active neutrino flavors (ν_x ; $x = e, \mu, \tau$). Elastic scattering (ES) on electron is also sensitive to all active flavors, but with reduced sensitivity to ν_μ and ν_τ .

Because of its equal sensitivity to all active neutrinos, the neutral-current measurement can determine the total neutrino flux, hence resolving the Solar Neutrino Problem, *even if neutrinos oscillate*.¹⁰ SNO is currently the only experiment that can simultaneously observe the *disappearance* of ν_e and the *appearance* of another active species. This is illustrated in Figure 1. Another feature of the neutral-current interaction is its low kinematic threshold. By efficiently counting the free neutrons in the final state of the neutral-current reaction, the total active ${}^8\text{B}$ neutrino flux can be inferred for neutrinos with energy above the 2.2-MeV kinematic threshold.

Recent results^{11,12} from the measurements of the solar ${}^8\text{B}$ neutrino flux by the SNO detector using the CC, NC and ES reactions are presented in this paper. The results¹³ of a measurement of the day-night asymmetry of the neutrino event rates, which is predicted under certain neutrino oscillation scenarios, are also presented. Finally, the physics implications of these observations will be discussed.

Experiment	Measured Flux	SSM Flux ⁷
Homestake ¹	$2.56 \pm 0.16(\text{stat.}) \pm 0.16(\text{sys.}) \text{ SNU}$	$7.6^{+1.3}_{-1.1} \text{ SNU}$
SAGE ³	$70.8^{+5.3}_{-5.2}(\text{stat.})^{+3.7}_{-3.2} \text{ SNU}$	$128^{+9}_{-7} \text{ SNU}$
Galex ⁴	$77.5 \pm 6.2(\text{stat.})^{+4.3}_{-4.7}(\text{sys.}) \text{ SNU}$	$128^{+9}_{-7} \text{ SNU}$
GNO ⁵	$65.8^{+10.2}_{-9.6}(\text{stat.})^{+3.4}_{-3.6}(\text{sys.}) \text{ SNU}$	$128^{+9}_{-7} \text{ SNU}$
Kamiokande ²	$2.80 \pm 0.19(\text{stat.}) \pm 0.33(\text{sys.}) \times 10^6$	$5.05 \times 10^6 \left(\begin{smallmatrix} 1^{+0.20} \\ -0.16 \end{smallmatrix} \right)$
Super-Kamiokande ⁶	$2.35 \pm 0.03(\text{stat.})^{+0.07}_{-0.06}(\text{sys.}) \times 10^6$	$5.05 \times 10^6 \left(\begin{smallmatrix} 1^{+0.20} \\ -0.16 \end{smallmatrix} \right)$

Table 1: Summary of solar neutrino observations at different solar neutrino detectors. The measured fluxes at the radiochemical experiments are measured in Solar Neutrino Unit (SNU), which is defined as 1 capture per second per 10^{36} target atoms. The Kamiokande and Super-Kamiokande measurements are in units of $10^6 \text{ cm}^{-2} \text{ s}^{-1}$.



<u>Measure:</u>	<u>Transformation to active</u>
$\frac{\phi_{CC}}{\phi_{ES}} = \frac{\nu_e}{\nu_e + 0.15(\nu_\mu + \nu_\tau)}$	 $\phi_{CC}(\nu_e) < \phi_{ES}$
$\frac{\phi_{CC}}{\phi_{NC}} = \frac{\nu_e}{\nu_e + \nu_\mu + \nu_\tau}$	 $\phi_{CC}(\nu_e) < \phi_{NC}$

Figure 1: Using the measured solar neutrino fluxes from different reaction channels to provide “smoking gun” evidence of neutrino flavor transformation.

2 The Sudbury Neutrino Observatory

2.1 Physical Description of the SNO Detector

SNO⁹ is an imaging water Cherenkov detector located in the Creighton #9 mine, owned by the International Nickel Company (INCO) near Sudbury, ON, Canada. A barrel-shaped cavity with a height of 34 m and a diameter of 22 m was excavated at a depth of 2092 m (or 6000 meters of water equivalent) to house the detector. The muon flux traversing the detector is 67 day^{-1} .

Figure 2 shows a cross-sectional view of the SNO detector. The neutrino detection medium is 1000 metric tons of 99.92% isotopically pure D₂O contained in a 12-m diameter acrylic sphere. The acrylic vessel is constructed out of 122 ultraviolet transmitting acrylic panels. This sphere is surrounded by 7000 metric tons of ultra-pure H₂O contained in the cavity. This volume of H₂O shields the detector from high energy γ rays and neutrons originating from the cavity wall. A 17.8-m diameter stainless steel structure supports 9456 20-cm diameter inward-facing photomultiplier tubes (PMTs). A non-imaging light concentrator is mounted on each PMT, extending the total photocathode coverage to 55%. An additional 91 PMTs are mounted facing outward on the support structure to serve as a cosmic-ray veto. To cancel the vertical components of the terrestrial magnetic field, 14 horizontal magnetic compensation coils were built into the cavity wall. The maximum residual field at the PMT array is $<19 \mu\text{T}$, and the reduction in photo-detection efficiency is about 2.5% from the zero-field value.

A physics event trigger is generated in the detector when there are 18 or more PMTs exceeding a threshold of ~ 0.25 photo-electrons within a coincidence time window of 93 ns. All the PMT hits registered in the ~ 420 ns window after the start of the coincidence time window are recorded in the data stream. This widened time window allows scattered and reflected Cherenkov photons to be included in the event. The mean noise rate of the PMTs is ~ 500 Hz, which results in ~ 2 noise PMT hits in this 420 ns window. The instantaneous trigger rate is about 13-20 Hz, of which 5-8 Hz are physics triggers. The remaining triggers are diagnostic triggers for monitoring the well being of the detector. The trigger efficiency reaches 100% when the PMT multiplicity (N_{hits}) in the event window is ≥ 23 .

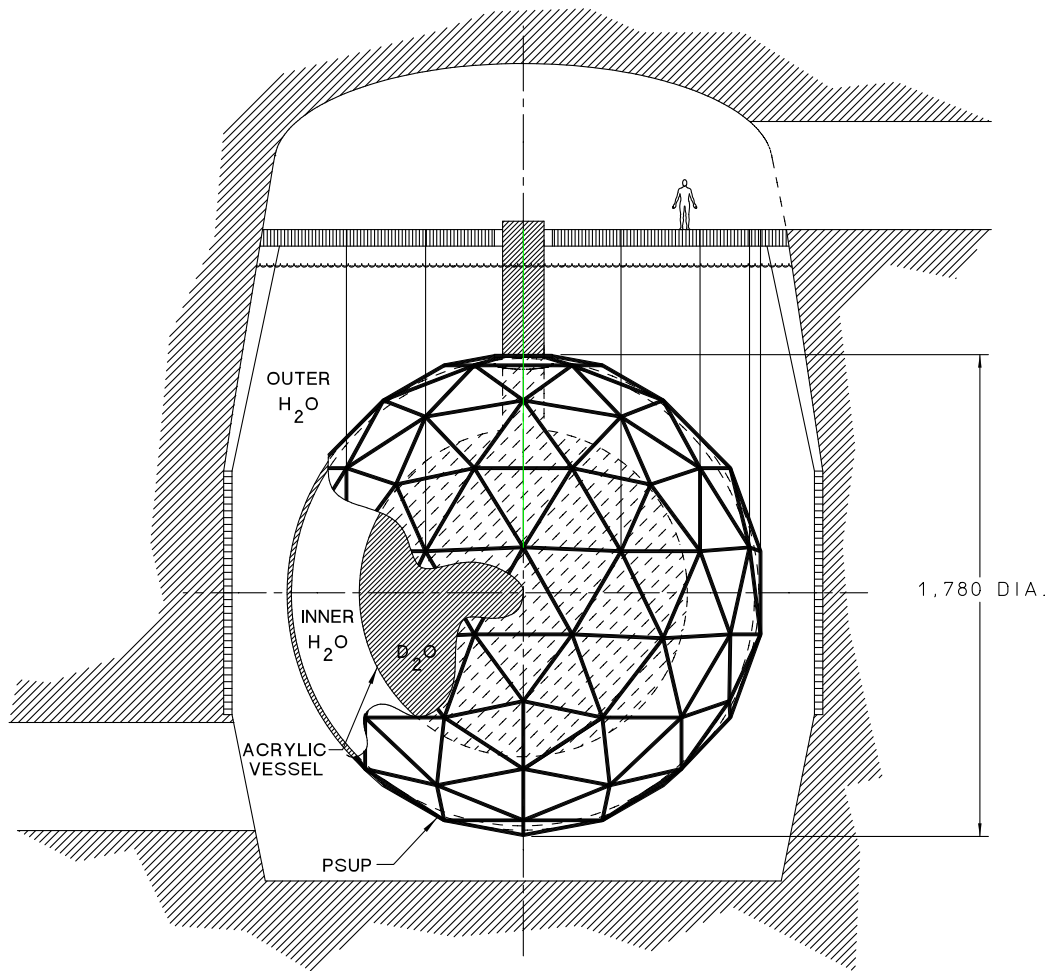


Figure 2: A cross-sectional view of the SNO detector. The outer geodesic structure is the PMT support (“PSUP”).

2.2 Solar Neutrino Physics Program at SNO

The solar neutrino physics program at SNO is designed to exploit its unique NC capability. Because the result of this NC measurement is a definitive statement on the Solar Neutrino Problem and flavor transformation of solar neutrinos, the SNO experiment has implemented a plan to make three independent NC measurements of the total ^8B active neutrino flux.

The first NC measurement was made with a pure D_2O target. The free neutron from the NC interaction is thermalized, and for $\sim 30\%$ of the time, a 6.25-MeV γ ray is emitted following the neutron capture by a deuteron. Only $\sim 50\%$ of the 6.25-MeV observed photopeak is above the neutrino analysis threshold, yielding a detection efficiency of $\sim 15\%$. The results from this NC measurement is presented in this paper. The second NC measurement is being made with 2 tonnes of NaCl added to the D_2O . The free neutron is readily captured by ^{35}Cl in this detector configuration, and a cascade of γ rays with a total energy of 8.6 MeV follow. The neutron detection efficiency is significantly enhanced, and $\sim 45\%$ of the NC events have a detectable signal above the analysis threshold. This phase of the experiment is scheduled to complete by the end of 2002. In the third NC measurement, discrete ^3He proportional counters will be installed inside the D_2O volume.¹⁴ The neutron detection efficiency of the proportional counter array is $\sim 40\%$. In this detector configuration, the detection of the CC and the NC signals are decoupled, and the covariance of the CC and NC signals that appear in the first two detector configurations is eliminated in this case.

3 Solar Neutrino Analysis at SNO

The data presented in this paper were recorded between November 2, 1999 and May 28, 2001. The corresponding livetime is 306.4 days. During this data acquisition period, the Sun was above the detector's horizon ("day" data set) for 128.5 days, and below the detector horizon ("night" data set) for 177.9 days. The target was pure D_2O during this period. Figure 3 summarizes the analysis procedure for extracting the CC, NC and the ES event rates in the SNO detector. In the following, each step in the analysis flow is discussed in detail.

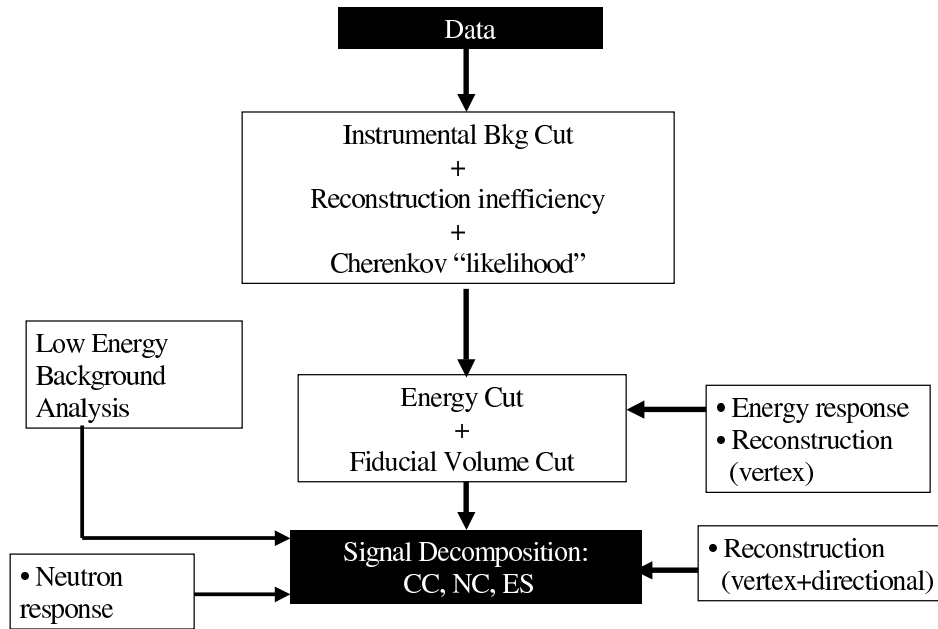


Figure 3: Simplified flow chart of solar neutrino analysis at SNO.

3.1 Data Reduction and Data Loss

After removing all the detector diagnostic triggers from the data stream, a significant portion of the remaining events are instrumental backgrounds. Examples of these backgrounds include electrical discharges in the PMTs (“flashers”) and light emission from discharges in insulating detector materials. Data reduction schemes were developed to remove these backgrounds.

3.1.1 Pass 0 Cut

The instrumental backgrounds have characteristic PMT time and charge distributions that are significantly different from Cherenkov light, and can be eliminated using cuts based on these distributions. For example, the discharge light emitted from a flasher PMT is detected across the detector ~ 70 ns after the initial discharge is registered. Some of these light-emitting instrumental backgrounds are localized near the water piping near the top of the detector. Veto PMTs were installed in this region in order to enhance the rejection efficiency of these non-Cherenkov events. Interference in the

electronics system can produce false events. Most of the registered electronic channel charges in these interference events are near the pedestal, and can be removed by a cut on the mean charge of the fired PMTs. Some of these electrical discharge or electronic interference background events also have different event-to-event time correlations from physics events, and time correlation cuts are used to remove these events. Two independent instrumental background rejection schemes are used. An event-by-event comparison of the data sets reduced by these two schemes shows a difference of $<0.2\%$.

The physics loss due to these instrumental background cuts is calibrated with a triggered ^{16}N 6.13-MeV γ -ray source¹⁵ and a triggered ^8Li 13-MeV endpoint β source¹⁶ deployed to the D_2O and H_2O volumes. Further tests of the N_{hits} dependence in the cuts are performed with an isotropic light source at various intensities. The physics acceptance of the instrumental background cuts, weighted over the fiducial volume, is measured to be $0.9966^{+0.0011}_{-0.0002}$. Instrumental background rejection is well over 99% at this stage.

3.1.2 High Level Cuts

After passing the instrumental background cuts, all events with $N_{hits} \geq 30$ (~ 3.5 MeV electron energy) are reconstructed. Once the event reconstruction information becomes available after the reconstruction, several high level physics cuts are applied to the Pass 0-reduced data set to further reduce the instrumental backgrounds. The efficiency of the reconstruction algorithm is calibrated with the ^{16}N and the ^8Li sources, and is found to be 0.9985 ± 0.0015 for neutrino events originated from the central 550 cm of the detector.

The high level cuts test the hypothesis that each event has the properties of electron Cherenkov light. The reconstruction figure-of-merit cuts test for the consistency between the time and angular expectations for an event fitted to the location of the reconstructed vertex and that based on the properties of Cherenkov light and the detector response.

Two parameters are used to further characterize Cherenkov light. The average opening angle between two hit PMTs ($\langle \theta_{ij} \rangle$), measured from the reconstructed vertex, is used to determine whether the topology of an event is consistent with Cherenkov light. The in-time ratio (ITR) is the ratio of the number of hit PMTs within an asymmetric time window around the prompt light peak to the number of calibrated PMTs in the event.

Channel	Pass-0	Reconstruction	High Level Cuts	Total
CC	$0.9966^{+0.0011}_{-0.0002}$	0.9985 ± 0.0015	0.9906 ± 0.0005	$0.9857^{0.0039}_{0.0021}$
ES	$0.9966^{+0.0011}_{-0.0002}$	0.9985 ± 0.0015	0.9903 ± 0.0007	$0.9854^{0.0040}_{0.0023}$
NC	$0.9966^{+0.0011}_{-0.0002}$	0.9985 ± 0.0015	0.9821 ± 0.0011	$0.9772^{0.0041}_{0.0023}$

Table 2: Physics acceptance at each of the instrumental background removal stages. A kinetic threshold of 5 MeV and a fiducial volume of the inner 550 cm of the detector are assumed.

Figure 4 shows the correlations between θ_{ij} and ITR for instrumental backgrounds and Cherenkov light events. As shown in the figure, this two dimensional cut has very high instrumental background rejection efficiency.

The total signal loss from the Pass 0 and the high level cuts are calibrated with the ^{16}N and the ^8Li sources. Because of the difference in the energy spectrum and spatial distribution for CC, ES and NC events, there is a slight difference in the data loss due to the Cherenkov likelihood cut for the different detection channels. Table 2 summarizes the physics acceptance after the Pass-0, reconstruction and high level cuts for the three neutrino detection channels.

The residual instrumental background contamination in the neutrino signal after the series of instrument background cuts is estimated by a bifurcated analysis, in which the signal contamination is obtained from cross calibrating the background leakage of two groups of orthogonal cuts. For the same fiducial volume and energy threshold, the instrumental background contamination is estimated to be <3 events (95% C.L.), or 0.1% of the final neutrino candidate data set. Table 3 summarizes the sequence of cuts that are used to reduce the raw data to 2928 neutrino candidate events.

3.2 Reconstruction and Energy Calibration

After all the instrumental background cuts have been applied to the data, energy and fiducial volume cuts are employed to reduce the physics backgrounds in the neutrino candidate data set. These physics backgrounds include low energy backgrounds from radioactive decays in the natural ^{238}U and ^{232}Th chains, and other backgrounds induced by cosmic rays. Because of progressively lower radioactivity towards the central region of the detector, one can place a fiducial volume cut and an energy threshold sufficiently high that most of the low energy backgrounds are removed from the neutrino candidate

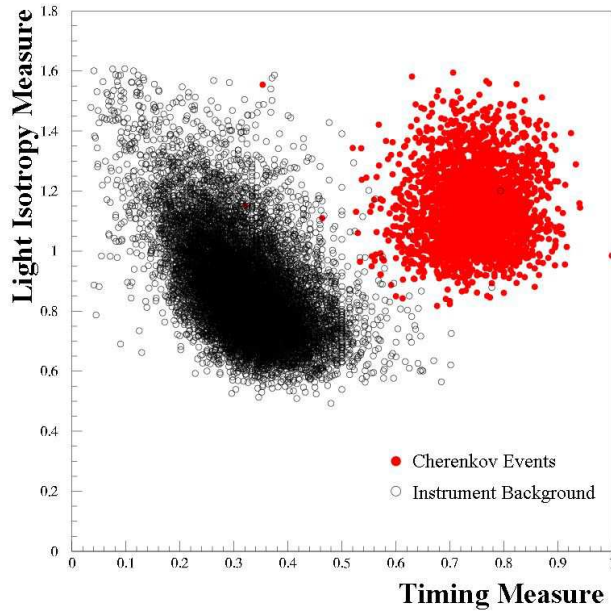


Figure 4: Separation of instrumental backgrounds and Cherenkov light events using the high level cuts.

Analysis step	Number of events
Total event triggers	450 188 649
Neutrino data triggers	191 312 560
PMT hit multiplicity ($N_{hits} \geq 30$ cut)	10 088 842
Instrumental background (Pass 0) cuts	7 805 238
Cherenkov likelihood cuts	3 418 439
Fiducial volume cut	67 343
Energy threshold cut	3 440
Cosmic-induced background subtraction	2 928
Total events	2 928

Table 3: Number of candidate events remained after each data reduction step

data set. These cuts rely upon event reconstruction and detector energy calibration, which we will discuss in the following.

3.2.1 Reconstruction

The calibrated times and positions of the fired PMTs are used to reconstruct the vertex position and the direction of the particle. Two different reconstruction algorithms were developed. An event-by-event comparison shows excellent agreement between the data sets reconstructed by these two algorithms. The neutrino data presented in this paper are reconstructed using a maximum likelihood technique which uses both the time and angular characteristics of Cherenkov light. Vertex reconstruction accuracy and resolution for electrons are measured using Compton electrons from the ^{16}N γ -ray source, and their energy dependence is verified by the ^8Li β source. Compton scattered electrons from a 6.13-MeV γ ray are preferentially scattered in the forward direction relative to the incident γ -ray direction. Thus they provide information about the accuracy of the direction reconstruction. In order to minimize the effect of finite vertex resolution on this angular resolution measurement, only ^{16}N events that are reconstructed to more than 150 cm from the source are used in the measurement. At the ^{16}N energy (~ 5.5 MeV total electron energy), the vertex reconstruction resolution is 16 cm and the angular resolution is 26.7° . Reconstruction-related systematic uncertainties to the solar neutrino flux measurement are $_{-2.8}^{+2.9}\%$ for the CC channel and $_{-1.8}^{+1.8}\%$ for the NC channel.

3.2.2 Energy Estimator

Figure 5 shows the relationship between the detector calibration program and event-by-event energy estimation in the analysis. Optical calibration is obtained using a near-isotropic source of pulsed laser light^{17,18} at 337, 365, 386, 420, 500 and 620 nm. The light source is deployed to locations accessible by the source manipulator system on two orthogonal planes in the D_2O , and on a linear grid in the H_2O . Optical parameters of different optical media in the detector are obtained at these wavelengths.¹⁹ The attenuation lengths in D_2O and H_2O are found to be near the Rayleigh scattering limit. These optical parameters are inputs to the Monte Carlo/energy estimator engine.

The triggered ^{16}N source is used to provide the absolute energy calibration. The detector energy response to the photopeak of this source provides a normalization to the PMT photon collection efficiency used in the Monte Carlo, and establish the abso-

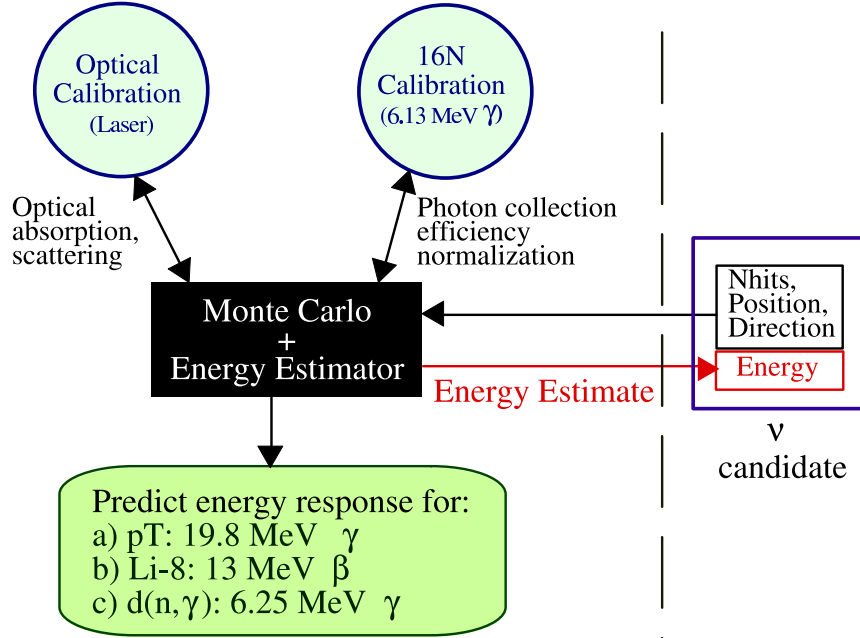


Figure 5: Calibration of the SNO detector and event-by-event energy estimator.

lute energy calibration. A long-term stability study of the detector response to the ^{16}N source shows a linear drift of $1.3\% \text{ year}^{-1}$. The cause of this effect is under investigation, and a drift correction is applied to the event-by-event energy estimator.

This tuned Monte Carlo is then used to make predictions for the energy response to different calibration sources. The pT source generates 19.8-MeV γ rays through the $^3\text{H}(p,\gamma)^4\text{He}$ reaction,²⁰ and is used to check the linearity of the energy response beyond the endpoint of the ^8B neutrino energy spectrum. To probe the regions that are inaccessible to the calibration source deployment system, the ^{252}Cf fission neutron source provides an extended distribution of 6.25-MeV γ rays from $d(n,\gamma)t$. Figure 6 shows a comparison of the Monte Carlo predictions and the detector responses to these sources.

The energy estimator uses the same input parameters (e.g. optical parameters) as the Monte Carlo. It assigns an effective kinetic energy T_{eff} to each event based upon its position, direction and the number of hit PMTs within the prompt (unscattered) photon peak. For an electron of total energy E_e , the derived detector energy response can be parameterized by a Gaussian:

$$R(E_{\text{eff}}, E_e) = \frac{1}{\sqrt{2\pi}\sigma_E(E_e)} \exp \left[-\frac{1}{2} \left(\frac{E_{\text{eff}} - E_e}{\sigma_E(E_e)} \right)^2 \right]$$

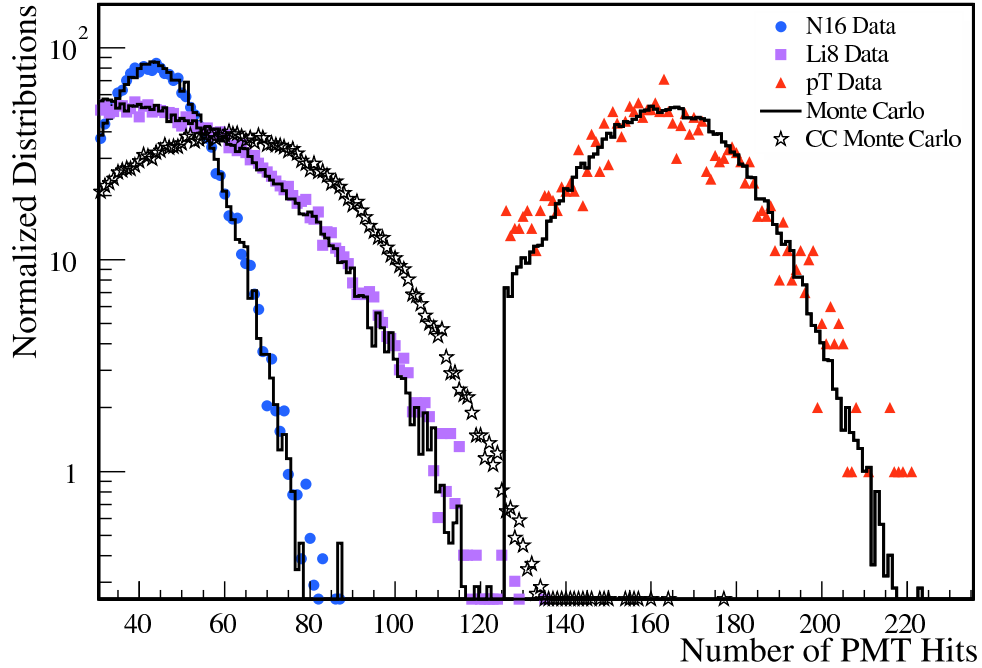


Figure 6: Comparison of the Monte Carlo predicted responses to different calibrated sources.

where $E_{\text{eff}} = T_{\text{eff}} + m_e$, and the energy resolution is given by

$$\sigma_E(E_e) = -0.0684 + 0.331\sqrt{E_e - m_e} + 0.0425(E_e - m_e) \text{ MeV}.$$

The systematic uncertainty on this absolute energy calibration is found to be $\pm 1.2\%$, which results in neutrino flux uncertainties of $^{+4.3}_{-4.2}\%$ for the CC channel and $^{+6.1}_{-6.2}\%$ for the NC channel. This is the most significant systematic uncertainty in the flux measurement. Other energy related systematic uncertainties to the flux include the energy resolution and the energy scale linearity. A summary of the systematic uncertainties can be found in Table 7.

A second energy estimator using N_{hits} is employed for validation purposes. These two energy estimators give consistent results in the neutrino flux measurement.

3.3 Neutron Efficiency

For the NC measurement of the solar neutrino flux, the probability that a neutron produced in the NC interaction will capture on deuterium and the detection efficiency of the 6.25-MeV γ ray emitted following the neutron capture have to be determined. The

capture probability depends on the isotopic abundances of various nuclei in the heavy water and the relevant capture cross sections for each. Near the boundary of the AV, the capture cross sections for neutrons in acrylic and light water become important as well. The detection efficiency of the 6.25-MeV γ ray can be readily deduced from the ^{16}N calibration data, since the latter produces predominantly 6.13-MeV γ rays.

Three different methods are used to calibrate the capture efficiency: a method which directly counts the number of neutrons measured with a ^{252}Cf source and compares to source expectations, a purely analytic calculation compared to Monte Carlo, and a multiplicity measurement which extracts the efficiency statistically from the distribution of the number of detected neutrons per fission.

The primary approach to measuring the capture efficiency is the “direct counting” method, in which the number of neutrons detected during a source run is compared to the total number expected to be generated by the source based on the known decay rate. This analysis was performed on ^{252}Cf source data taken at various locations in the detector. Because the ^{252}Cf source is not a triggered source, non-neutron backgrounds associated with the source and the detector were handled using reconstruction and instrumental background cuts, with a proper accounting of the associated acceptance of these cuts in the efficiency measurement. In Figure 7 the neutron capture efficiency on deuterium as a function of radial distance from the center of the detector is shown. The capture efficiency on deuterium for NC events is found to be $29.90 \pm 1.10\%$, and when the fiducial volume cut of 550 cm and energy threshold cut of $T_{\text{eff}} > 5.0$ MeV are included, the overall detection efficiency is $14.38 \pm 0.53\%$. The uncertainties on these numbers contribute a relative uncertainty of 3.68% in the extraction of the NC signal in the present analysis.

The analytic calculation uses known cross sections and isotopic abundances of the D_2O to calculate the expected efficiency, and propagates uncertainties on the microscopic parameters through the entire calculation. The largest uncertainties in the analytic model are on the fraction of H_2O in the heavy water, the capture cross sections on deuterium and ^{16}O , and the diffusion parameter which governs the total number of neutrons escaping the D_2O volume. For the NC signal, the total uncertainty determined in this way is $\sim 3\%$.

The third approach is provided primarily as a check because of its limited sensitivity. In this approach the multiplicity distribution of detected neutrons is fit with the

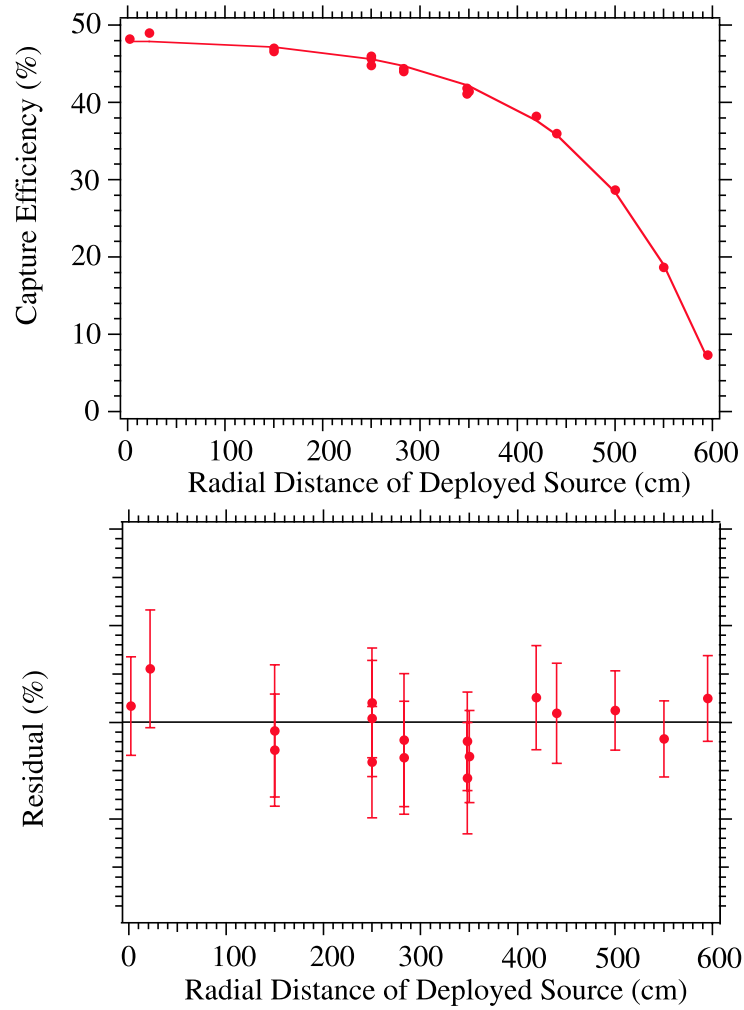


Figure 7: Final capture efficiencies measured with the direct counting method as a function of radial position, and the residuals from a fit to the radial dependence.

Component	^{232}Th (g/g)	^{238}U (g/g)
D ₂ O	3.7×10^{-15}	4.5×10^{-14}
H ₂ O	3.7×10^{-14}	4.5×10^{-13}
AV	1.9×10^{-12}	3.6×10^{-12}

Table 4: The target radio-purity levels for different components in the SNO detector.

expression

$$P(d) = \sum_{N=1}^{\infty} \sum_{r=d}^{\infty} \frac{r! \epsilon^d \epsilon^{r-d}}{d!(r-d)!} \frac{\exp\left(-\frac{(r-N\mu)^2}{2N\sigma^2}\right)}{2N\pi\sigma^2} \frac{\lambda^N \exp^{-\lambda}}{N!}$$

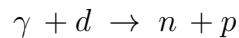
where $P(d)$ is the probability to detect d neutrons per unit time given that r neutrons were generated according to Poisson statistics for the ^{252}Cf fission rate λ and the known ^{252}Cf multiplicity per fission (mean $\mu = 3.79 \pm 0.006$ and width $\sigma^2 = 1.57 \pm 0.02$). In this case, both the source fission rate λ and the neutron detection efficiency ϵ are derived from a maximum likelihood fit. The results from this approach and the direct counting approach are consistent.

3.4 Low Energy Backgrounds

Low levels of ^{232}Th and ^{238}U can be found naturally in all materials. The SNO detector has been designed with very stringent radiopurity targets for different components in the detector. Table 4 lists the ^{232}Th and ^{238}U target levels for the D₂O, acrylic vessel (AV), and H₂O. At these radiopurity levels, the background to the NC signal is approximately 10% of the NC signal predicted by the Standard Solar Model.

Radioactive decays of the daughters in the natural ^{232}Th and ^{238}U chains are the dominant backgrounds in the neutrino signal window. These decays can contribute to the backgrounds in two different ways.

A γ ray with energy in excess of the binding energy of the deuteron ($=2.2$ MeV) can potentially photodisintegrate a deuteron:



The free neutron in the final state is indistinguishable from that in the NC signature. In the natural ^{232}Th and the ^{238}U chains, γ rays that can photodisintegrate the deuterons are

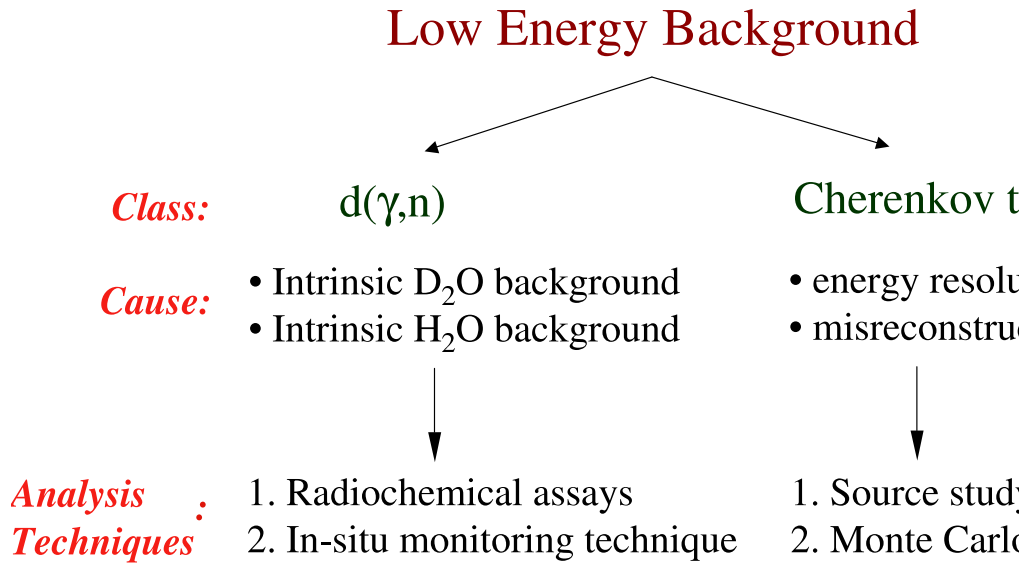


Figure 8: Low Energy Background Analysis. The techniques that are used in understanding these backgrounds are shown.

emitted in the decays of ^{208}Tl and ^{214}Bi . To measure the contribution of this photodisintegration background to the NC signal, it is necessary to determine the concentration of the different isotopes in the detector. “Ex-situ” radiochemical assays and “in-situ” Cherenkov light pattern recognition techniques are used to measure the contribution from this photodisintegration background.

A small fraction of decays with large Q values (e.g. ^{208}Tl from the ^{232}Th chain and ^{214}Bi from the ^{238}U chain) inside the fiducial volume might generate enough light in the detector and register a calibrated energy above the analysis threshold. Decays that are originated from outside the fiducial volume might get reconstructed to within the fiducial volume due to finite reconstruction resolution. Low energy radioactive sources and simulations are used to understand this class of *Cherenkov tail* backgrounds.

Figure 8 shows the general strategy of the low energy background analysis. In the following, we shall discuss the analysis techniques mentioned above in more details.

3.4.1 Photodisintegration Background

D₂O and H₂O Radioactivity

Several techniques were developed to measure the ^{232}Th and ^{238}U concentration in the D_2O and the H_2O . These techniques sample different daughters in the ^{232}Th and ^{238}U chains, and can be broadly classified into two categories:

- **Ex-situ:** These are techniques that involve chemical removal of radioisotopes from the D_2O and the H_2O , and their decays are counted in a system external to the SNO detector. These techniques include the extraction of Ra isotopes using MnO_x beads (the “ MnO_x ” technique²²), and the extraction of Ra, Th and Pb isotopes using HTiO membranes (the “HTiO” technique²³). Typical assays circulated approximately 400 tonnes of water through the extraction media. These techniques provide isotopic identification of the decay daughters and contamination levels in the assayed water volumes, presented in Fig. 10 (a). Secular equilibrium in the ^{238}U decay chain is broken by the ingress of long-lived (3.8 day half-life) ^{222}Rn in the experiment. Measurements of this background are made by periodically extracting and cryogenically concentrating ^{222}Rn from the ^{238}U chain (the “Rn assay”²⁴). The Rn results are presented (as mass fractions in $\text{g}(\text{U})/\text{g}(\text{D}_2\text{O})$) in Fig. 10(b).

The MnO_x Technique

In the MnO_x technique, D_2O or H_2O is passed through polypropylene columns that contain beads coated with a manganese oxide compound (MnO_x), which extracts Ra from the flowing water. After a large volume of water has passed through the columns, they are removed and dried. The dried column is then attached to a gas flow loop on an electrostatic counter (ESC). The Rn produced from Ra decay is swept from the columns into the ESC where it decays. The charged Po ions from the decay of Rn are carried by the electric field onto an α counter where the decays of the Po are detected, and their α energy spectra are collected. For the ^{232}Th chain, the relevant Po α decays are ^{216}Po (6.8 MeV α) and ^{212}Po (8.8 MeV α), whereas the relevant ones for the U chain are ^{218}Po (6.0 MeV α) and ^{214}Po (7.7 MeV α).

The HTiO Technique

In this technique, D_2O or H_2O is passed through hydrous titanium oxide (HTiO) trapped on filtration fibers. The HTiO ion-exchanger is first deposited onto a

microfiltration membrane. Then the columns containing the loaded filters are used to extract ^{224}Ra from the Th chain and ^{226}Ra from the U chain from a large volume of D_2O or H_2O . After extraction, Ra is eluted with nitric acid, and subsequently concentrated down to ~ 10 ml of eluate. This is then mixed with liquid scintillator and counted using $\beta - \alpha$ delayed coincidence counters.²⁷ For the ^{232}Th chain, the coincidences of the β -decay of ^{212}Bi and the α -decay of ^{212}Po are counted, whereas the coincidences of the β -decay of ^{214}Bi and the α -decay of ^{214}Po are counted for the ^{238}U chain.

Rn Assays

Measurements of ^{226}Ra concentration in the D_2O and the H_2O in the MnO_x and the HTiO techniques are not sufficient to determine the total radioactive background from the ^{238}U chain. This is because even a small leak of the underground laboratory air (~ 3 pCi/l of ^{222}Rn) can lead to significant disequilibrium between ^{226}Ra and ^{214}Bi . The Rn assay technique was developed to tackle this problem. In this technique, water drawn from discrete sample points in the detector is flowed through a degasser to liberate Rn. The Rn is purified and collected in a cryogenic collector. The subsequent α decays are counted in a Lucas cell scintillator (ZnS) chamber on a 2.54-cm diameter photomultiplier tube. Since there is a delay of many ^{220}Rn lifetimes between the preparation of the Lucas cells and their subsequent counting, this method is sensitive only to ^{222}Rn decays.

- **In-situ:** This is a technique that uses pattern recognition on the Cherenkov light distribution to determine the concentration of ^{232}Th and ^{238}U in the D_2O and the H_2O .^{25,26} The ^{208}Tl decay (from the ^{232}Th chain) has a Q value of ~ 4.9 MeV, and the ^{214}Bi decay (from the ^{238}U chain) has a Q value of 3.27 MeV. Almost every ^{208}Tl decay emits a 2.614 MeV γ , several low energy γ 's and a β with an endpoint of $\sim 1-1.8$ MeV, whereas there is a unique branch in the ^{214}Bi decay that produces a single β with an endpoint energy of 3.27 MeV. A single particle produces more Cherenkov photons than multiple particles with the same total energy because of the Cherenkov kinetic threshold. Also the light isotropy is different for these two cases because multiple particles generate a more isotropic light pattern (multiple Cherenkov cones), whereas a single particle gives a more directional light distribution due to light emission within a single Cherenkov cone. Therefore, by selecting a set of energy and fiducial volume cuts, it is possible to

separate ^{208}Tl and ^{214}Bi decays statistically by their differences in light isotropy. The light isotropy parameter that has been developed ($\langle\theta_{ij}\rangle$) is the average angle between all hit PMT pairs within the prompt light time window in an event. Figure 9 shows the difference in θ_{ij} between ^{208}Tl and ^{214}Bi decays. In this in-situ analysis, Cherenkov events fitted within 450 cm from the detector center and extracted from the neutrino data set provide a time-integrated measure of these backgrounds over the same time period of the neutrino analysis. Statistical separation of the ^{208}Tl and ^{214}Bi events is obtained by a maximum likelihood fit of the Cherenkov signal isotropy.

Results from the *ex situ* and *in situ* methods are consistent with each other as shown on the right hand side of Figs. 10(a) and 10(b). For the ^{232}Th chain, the weighted mean (including additional sampling systematic uncertainty) of the two determinations is used for the background analysis. The ^{238}U chain activity is dominated by Rn ingress which is highly time dependent. Therefore the *in situ* determination was used for this activity as it provides the appropriate time weighting. For the present data set, the time-averaged equivalent equilibrium ^{238}U and ^{232}Th concentration in the D_2O are found to be

$$\begin{aligned} ^{232}\text{Th} & : 1.63 \pm 0.58 \times 10^{-15} \text{ g Th / g D}_2\text{O} \\ ^{238}\text{U} & : 17.8_{-4.3}^{+3.5} \times 10^{-15} \text{ g U / g D}_2\text{O}. \end{aligned}$$

Once the ^{232}Th and ^{238}U concentrations are measured, the photodisintegration background can be determined through Monte Carlo calculations that use the photodisintegration cross section, properties of γ -ray and neutron transport, and the detector response function as inputs. The expected number of observed photodisintegration neutrons are 18.4 ± 6.5 counts for the ^{232}Th background and $25.9_{-6.3}^{+5.0}$ counts for the ^{238}U background in the pure D_2O phase of the experiment.

Daughters in the ^{232}Th and ^{238}U chain in the H_2O and the acrylic vessel (AV) can still produce a photodisintegration background. This happens when the γ 's with an energy >2.2 MeV enter the D_2O volume and photodisintegrate a deuteron. *Ex situ* assays and *in situ* pattern recognition techniques described above were used to measure the ^{232}Th and ^{238}U concentrations in the H_2O . The photodisintegration background inside the fiducial volume due to radioactivity in the H_2O was then determined. The time-averaged equivalent equilibrium ^{238}U and ^{232}Th concentration in the H_2O are found to

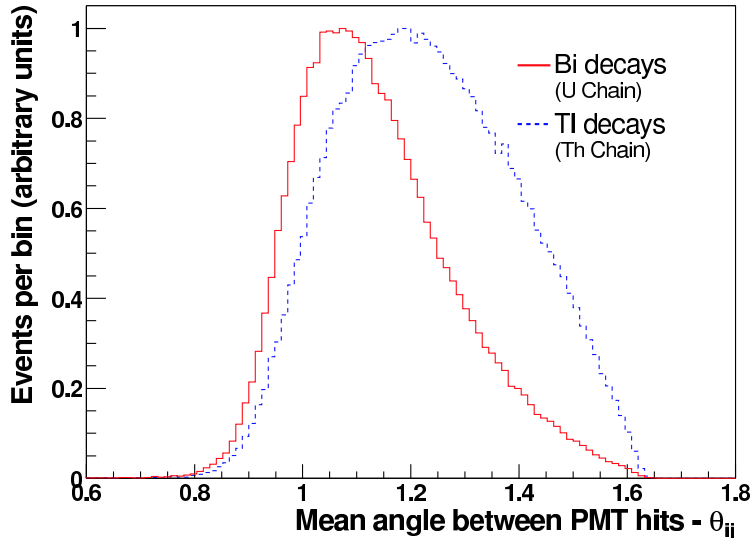


Figure 9: Difference in θ_{ij} between ^{208}Tl and ^{214}Bi decays.

be

$$^{232}\text{Th} : 14.2 \pm 6.6 \times 10^{-14} \text{ g Th / g H}_2\text{O}$$

$$^{238}\text{U} : 75.5 \pm 33.0 \times 10^{-14} \text{ g U / g H}_2\text{O}.$$

At these concentrations, the number of observed photodisintegration neutrons in the fiducial volume are determined to be $5.6^{+3.6}_{-2.2}$ counts for the ^{232}Th chain and $5.6^{+4.2}_{-2.9}$ counts for the ^{238}U chain.

Acrylic Vessel Radioactivity

Prior to the construction of the acrylic vessel, neutron activation analyses were done on samples of the acrylic panels used in the vessel. These results, along with that of a dust density measurement on the vessel subsequent to its construction, indicate that the radioactive load of the acrylic vessel to be $7.5^{+1.7}_{-1.3} \mu\text{g}$ of ^{232}Th and $15 \pm 15 \mu\text{g}$ of ^{238}U . Studies of the Cherenkov data revealed a radioactive hotspot on the acrylic vessel. The radioactivity of this hotspot is equivalent to $10^{+8.6}_{-3.6} \mu\text{g}$ of ^{232}Th . Combining the results from radioassays and the hotspot analysis, the expected number of observed photodisintegration background events due to radioactivity in the acrylic vessel is $15.8^{+6.0}_{-6.8}$ counts.

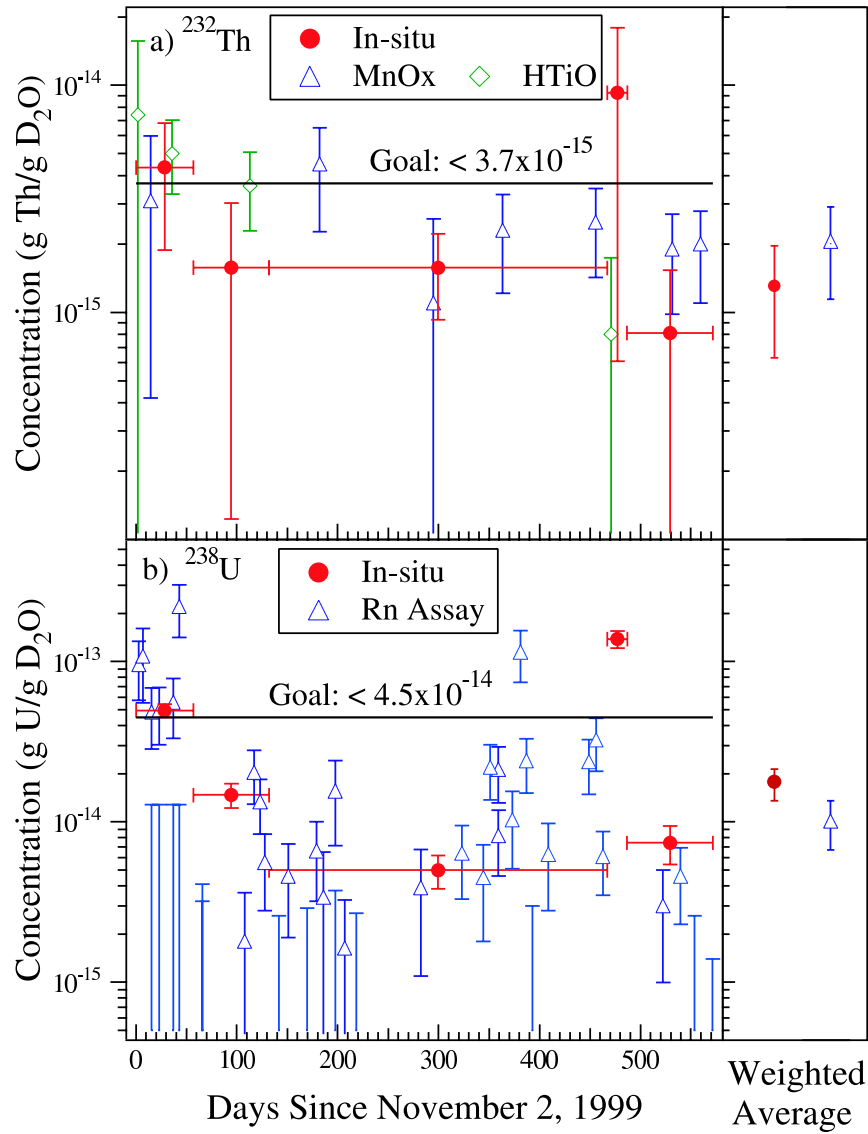


Figure 10: Thorium (a) and uranium (b) backgrounds (equivalent equilibrium concentrations) in the D_2O deduced by *in situ* and *ex situ* techniques. The MnO_x and $HTiO$ radiochemical assay results, the Rn assay results, and the *in situ* Cherenkov signal determination of the backgrounds are presented for the period of this analysis on the left-hand side of frames (a) and (b). The right-hand side shows time-integrated averages including an additional sampling systematic uncertainty for the *ex situ* measurement.

	^{232}Th (counts)	^{238}U (counts)	Total
D ₂ O photodisintegration	18.4±6.5	25.9 ^{+5.0} _{-6.3}	44.3 ^{+8.2} _{-9.1}
AV photodisintegration	14.2 ^{+5.8} _{-6.6}	1.6±1.6	15.8 ^{+6.0} _{-6.8}
H ₂ O photodisintegration	5.6 ^{+3.6} _{-2.2}	5.6 ^{+4.2} _{-2.9}	11.2 ^{+5.5} _{-3.6}
Atmospheric ν 's and sub-Cherenkov threshold μ 's	Not applicable		4 ± 1
Spontaneous Fission	Not applicable		≪ 1
$^2\text{H}(\alpha, \alpha)\text{pn}$	Not applicable		2 ± 0.4
$^{17}\text{O}(\alpha, \text{n})$	Not applicable		≪ 1
Terrestrial and reactor $\bar{\nu}$'s	Not applicable		1 ⁺³ ₋₁
External neutrons	Not applicable		≪ 1
Total			78 ± 12

Table 5: Summary of neutron background ($E > 5.5$ MeV) in the fiducial volume ($R < 550$ cm.) for the pure D₂O phase of the experiment

Table 5 is a summary of the photodisintegration background for the data set used in this analysis. In this table, we have also listed the contributions from other sources of neutron backgrounds. Some of these are cosmic ray events and atmospheric neutrinos. To reduce these backgrounds, an additional neutron background cut imposed a 250-ms veto (in software) following every event in which the total number of PMTs which registered a hit was greater than 60. The contribution from these additional sources are very small compared to the photodisintegration background.

3.4.2 Cherenkov Tail Background

Low energy backgrounds from Cherenkov events in the neutrino signal region were evaluated by using acrylic encapsulated sources of U and Th deployed throughout the detector volume and by Monte Carlo calculations. In the following, we discuss how this class of background is determined for different origins of radioactivity.

Cherenkov Tail Background from D₂O Radioactivity

The Cherenkov tail background arising from internal radioactivity in the D₂O were

determined by a Monte Carlo study. A Monte Carlo study is justified because the detector response in the D₂O is well calibrated, and the internal radioactivity is well measured by the *ex situ* and *in situ* techniques. In this Monte Carlo calculation, the number of Cherenkov background events is normalized to the number of photodisintegration (pd) neutrons in the neutrino signal window and are found to be:

$$\begin{aligned}
 {}^{208}\text{Tl}: & \quad 0.162 \begin{matrix} +0.092 \\ -0.030 \end{matrix} \text{ Cherenkov tail event/pd neutron event} \\
 {}^{214}\text{Bi}: & \quad 0.670 \begin{matrix} +0.460 \\ -0.125 \end{matrix} \text{ Cherenkov tail event/pd neutron event.}
 \end{aligned}$$

Given the number of detected photodisintegration neutrons (Th: 18.4 ± 6.5 , U: $25.9^{+5.0}_{-6.3}$), the expected Cherenkov tail backgrounds are:

$$\begin{aligned}
 {}^{232}\text{Th}: & \quad 3.0 \begin{matrix} +2.0 \\ -1.3 \end{matrix} \text{ counts} \\
 {}^{238}\text{U}: & \quad 17.4 \begin{matrix} +12.4 \\ -5.3 \end{matrix} \text{ counts} \\
 \text{Total}: & \quad 20.4 \begin{matrix} +12.6 \\ -5.5 \end{matrix} \text{ counts.}
 \end{aligned}$$

As a consistency check, this Monte Carlo technique was used to make predictions on the energy spectrum of a variety of background sources, including a stainless-steel encapsulated Th source and a Rn spike in the detector during a water circulation pump failure. Figure 11 compares the predicted energy spectrum for the stainless-steel encapsulated Th source run to the actual calibration source data.

Cherenkov Tail Background from External Radioactivity

In principle, one can perform Monte Carlo calculations to predict the Cherenkov tail background in the fiducial volume for radioactivity from the acrylic vessel, the H₂O and the PMT support geodesic. However, most of the events in the neutrino signal window with origins in the regions external to the fiducial volume come from the tail of the energy response and the reconstruction response functions. These response functions are not as well understood outside the fiducial volume.

A technique was developed to determine this external (acrylic vessel, H₂O and PMT support geodesic) Cherenkov tail backgrounds in the neutrino signal window using the low energy background radial distribution (R) calibrated by low energy acrylic encapsulated sources. The principle of this technique is to use the acrylic source data to generate the R^3 probability distribution functions (PDFs) for the low energy backgrounds in the AV, H₂O and the PMT regions. The R^3 distribution in these regions in the neutrino data was then fitted to a linear combination of these PDFs by the extended maximum likelihood method. One difficulty of this technique is the lack of sufficient

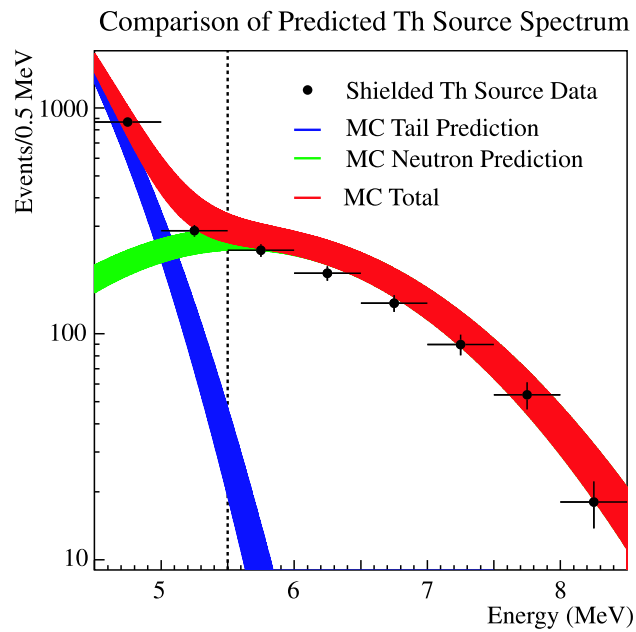


Figure 11: Comparison of Monte Carlo prediction with systematic uncertainties to a stainless-steel encapsulated Th source at the center of the D₂O. The β 's in the decays are blocked by the stainless steel shroud, while the γ 's emitted in the decays escaped the stainless steel container. The predictions, shown here as a $\pm 1\sigma$ band, have been normalized by the livetime and the source strength.

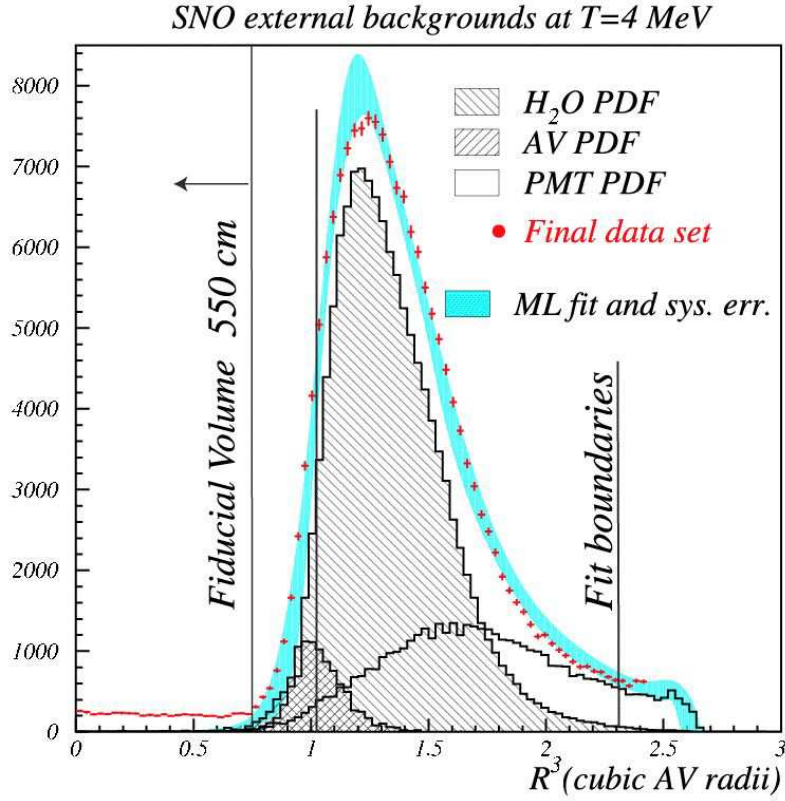


Figure 12: External Cherenkov tail R^3 fit at $T_{\text{eff}} > 4$ MeV. The extended maximum likelihood method was used in the fit, and the band represents the systematic uncertainties.

statistics at the intended neutrino analysis threshold ($T_{\text{eff}} = 5$ MeV). In order to overcome this, the analysis was proceeded at a lower energy ($T_{\text{eff}} = 4$ MeV) and then extrapolated to 5 MeV. The basic assumptions here are that there is no correlation between R^3 and energy, and that the reconstruction does not get worse with higher energy. The R^3 fit was performed for the range $1.02 < (R/600)^3 < 2.31$. Figure 12 shows the results of this R^3 fit. The band shown in this plot is the range of the systematic uncertainties. The data points are not centered in this band because the systematic uncertainties are not normally distributed. Table 5 is a summary of the Cherenkov tail background ($T_{\text{eff}} > 5$ MeV) in the fiducial volume ($R < 550$ cm.) for the pure D_2O phase of the experiment.

Several consistency checks of this external Cherenkov tail analysis were performed. The radial PDF for H_2O background was compared to the radial distribution of a radon spike during a water circulation pump failure. This is shown in Figure 13. The expected

	^{232}Th (counts)	^{238}U (counts)	Total
D ₂ O	$3.0^{+2.0}_{-1.3}$	$17.4^{+12.4}_{-5.3}$	$20.4^{+12.6}_{-5.5}$
AV	Not applicable		$6.3^{+2.9}_{-6.3}$
H ₂ O	Not applicable		$2.8^{+3.9}_{-2.8}$
PMT	Not applicable		$16.0^{+11.1}_{-8.0}$
Total	Not applicable		$45.5^{+17.5}_{-11.9}$

Table 6: Summary of Cherenkov tail background ($E > 5.5$ MeV) in the fiducial volume ($R < 550$ cm.) for the pure D₂O phase of the experiment. The high energy γ background contribution in the fiducial volume is included in the PMT $\beta\gamma$ entry.

background rate in Table 6 was also checked against the Monte Carlo calculations and the results of a simultaneous fit of ν signals and backgrounds in N_{hits} , R^3 and angular correlation to the Sun. All these results are consistent.

3.5 Signal Extraction

In order to test the null hypothesis, the assumption that there are only electron neutrinos in the solar neutrino flux, the extended maximum likelihood method is used in extracting the CC, ES and neutron (i.e. NC+background) contributions in the candidate data set. Background contributions are constrained to the measured values discussed above. The undistorted ^8B spectrum from Ortiz *et al.*²⁸ is assumed in the signal decomposition. Data distributions in T_{eff} , the volume-weighted radial variable $(R/R_{\text{AV}})^3$ where $R_{\text{AV}} = 600$ cm is the radius of the acrylic vessel, and $\cos\theta_{\odot}$ are simultaneously fitted to the probability density functions (PDFs) generated from Monte Carlo simulations. $\cos\theta_{\odot}$ is the angle between the reconstructed direction of the event and the instantaneous direction from the Sun to the Earth. This distribution is shown in Figure 14(a) for the analysis threshold of $T_{\text{eff}} \geq 5$ MeV and fiducial volume selection of $R \leq 550$ cm, where R is the reconstructed event radius. The forward peak ($\cos\theta_{\odot} \sim 1$) arises from the strong directionality in the ES reaction. The $\cos\theta_{\odot}$ distribution for the CC reaction, before accounting for the detector response, is expected to be $(1-0.340\cos\theta_{\odot})$.²⁹ Fig. 14(b) shows the distribution of events in the volume-weighted radial variable $(R/R_{\text{AV}})^3$, and Figure 14(c) shows the kinetic energy spectrum of the selected events.

H₂O PDF from the acrylic sources and the radon data

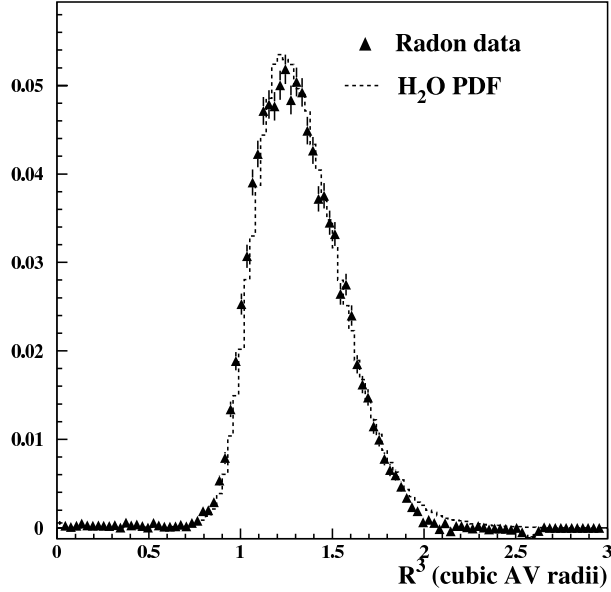


Figure 13: Comparison between the PDF obtained from the radon spike in the detector and the H₂O PDF derived from acrylic source data.

The extraction yields $1967.7^{+61.9}_{-60.9}$ CC events, $263.6^{+26.4}_{-25.6}$ ES events, and $576.5^{+49.5}_{-48.9}$ NC events *. The uncertainties given above are statistical only, and the systematic uncertainties are summarized in Table 7. The extracted counts for each neutrino detection channel above can be converted into integrated fluxes above the kinetic energy threshold of 5 MeV. Assuming an undistorted ${}^8\text{B}$ neutrino spectrum²⁸ and using only cross sections for ν_e , the measured neutrino fluxes from each of the channels are (in units of $10^6 \text{ cm}^{-2}\text{s}^{-1}$):

$$\phi_{\text{CC}} = 1.76^{+0.06}_{-0.05}(\text{stat.})^{+0.09}_{-0.09}(\text{syst.})$$

$$\phi_{\text{ES}} = 2.39^{+0.24}_{-0.23}(\text{stat.})^{+0.12}_{-0.12}(\text{syst.})$$

$$\phi_{\text{NC}} = 5.09^{+0.44}_{-0.43}(\text{stat.})^{+0.46}_{-0.43}(\text{syst.}).$$

The CC and ES results reported here are consistent with our earlier results¹¹ for $T_{\text{eff}} \geq 6.75$ MeV. The excess of the NC flux over the CC and ES fluxes implies neutrino flavor transformations.

*We note that this rate of neutron events also leads to a lower bound on the proton lifetime for “invisible” modes (based on the free neutron that would be left in deuterium³⁰) in excess of 10^{28} years, approximately 3 orders of magnitude more restrictive than previous limits.³¹ The possible contribution of this mechanism to the solar neutrino NC background is ignored.

A direct test of the null hypothesis, i.e. there are only ν_e in the solar neutrino flux, can be readily performed with a simple change of variables:

$$\begin{aligned}\phi_{\text{CC}} &= \phi_e \\ \phi_{\text{NC}} &= \phi_e + \phi_{\mu\tau} \\ \phi_{\text{ES}} &= \phi_e + \epsilon\phi_{\mu\tau},\end{aligned}$$

where ϕ_e is the flux of the electron component, $\phi_{\mu\tau}$ is the flux of non-electron component, and $\epsilon=0.1559$ is the ratio of the elastic scattering cross sections for $\nu_{\mu\tau}$ and ν_e above the kinetic threshold of 5 MeV. Assuming an undistorted ${}^8\text{B}$ energy spectrum, a maximum likelihood extraction using these transformed variables gives the fluxes of the electron and non-electron components as:

$$\begin{aligned}\phi_e &= 1.76_{-0.05}^{+0.05}(\text{stat.})_{-0.09}^{+0.09}(\text{syst.}) \\ \phi_{\mu\tau} &= 3.41_{-0.45}^{+0.45}(\text{stat.})_{-0.45}^{+0.48}(\text{syst.})\end{aligned}$$

Combining the statistical and systematic uncertainties in quadrature, $\phi_{\mu\tau}$ is $3.41_{-0.64}^{+0.66}$, which is 5.3σ above zero, and provides strong evidence for flavor transformation consistent with neutrino oscillations.^{32,33} Adding the Super-Kamiokande ES measurement of the ${}^8\text{B}$ flux⁶ $\phi_{\text{ES}}^{\text{SK}} = 2.32 \pm 0.03(\text{stat.})_{-0.07}^{+0.08}(\text{syst.})$ as an additional constraint, we find $\phi_{\mu\tau} = 3.45_{-0.62}^{+0.65}$, which is 5.5σ above zero. Figure 15 shows the flux of non-electron flavor active neutrinos *vs.* the flux of electron neutrinos deduced from the SNO data. The three bands represent the one standard deviation measurements of the CC, ES, and NC rates. The error ellipses represent the 68%, 95%, and 99% joint probability contours for ϕ_e and $\phi_{\mu\tau}$.

Removing the constraint that the solar neutrino energy spectrum is undistorted, the signal decomposition is repeated using only the $\cos\theta_\odot$ and $(R/R_{\text{AV}})^3$ information. The total flux of active ${}^8\text{B}$ neutrinos measured with the NC reaction is

$$\phi_{\text{NC}} = 6.42_{-1.57}^{+1.57}(\text{stat.})_{-0.58}^{+0.55}(\text{syst.})$$

which is in agreement with the shape constrained value above and with the standard solar model prediction⁷ for ${}^8\text{B}$ ν 's, $\phi_{\text{SSM}} = 5.05_{-0.81}^{+1.01}$.

3.6 Day-Night Analysis

The solar neutrino flux results presented above provide strong evidence for neutrino flavor transformation independent of solar model predictions. One of the possible mechanism for this flavor transformation is mass-induced neutrino oscillations. For some

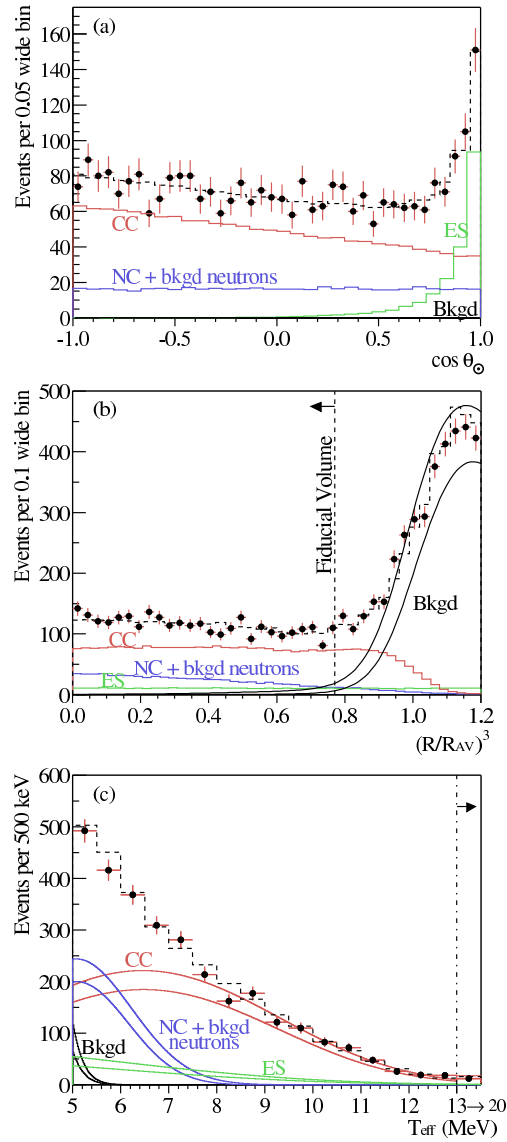


Figure 14: (a) Distribution of $\cos \theta_{\odot}$ for $R \leq 550$ cm. (b) Distribution of the volume weighted radial variable $(R/R_{AV})^3$. (c) Kinetic energy for $R \leq 550$ cm. Also shown are the Monte Carlo predictions for CC, ES and NC + background neutron events scaled to the fit results, and the calculated spectrum of Cherenkov background (Background) events. The dashed lines represent the summed components, and the bands show $\pm 1\sigma$ uncertainties. All distributions are for events with $T_{\text{eff}} \geq 5$ MeV.

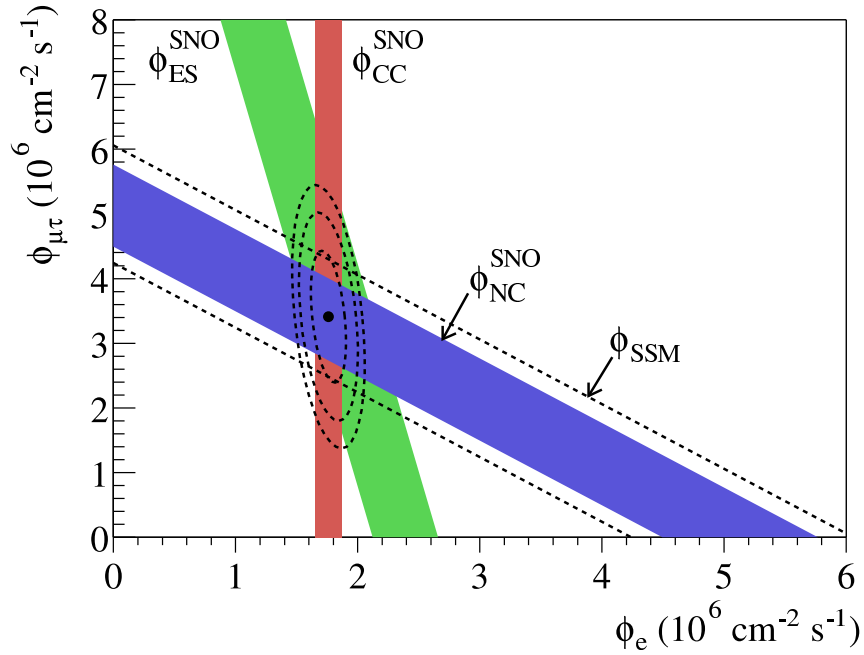


Figure 15: Flux of ${}^8\text{B}$ solar neutrinos which are μ or τ flavor vs. flux of electron neutrinos deduced from the three neutrino reactions in SNO. The diagonal bands show the total ${}^8\text{B}$ flux as predicted by the SSM⁷ (dashed lines) and that measured with the NC reaction in SNO (solid band). The intercepts of these bands with the axes represent the $\pm 1\sigma$ errors. The bands intersect at the fit values for ϕ_e and $\phi_{\mu\tau}$, indicating that the combined flux results are consistent with neutrino flavor transformation assuming no distortion in the ${}^8\text{B}$ neutrino energy spectrum.

Source	CC Uncert. (percent)	NC Uncert. (percent)	$\phi_{\mu\tau}$ Uncert. (percent)
Energy scale †	-4.2,+4.3	-6.2,+6.1	-10.4,+10.3
Energy resolution †	-0.9,+0.0	-0.0,+4.4	-0.0,+6.8
Energy non-linearity †	± 0.1	± 0.4	± 0.6
Vertex resolution †	± 0.0	± 0.1	± 0.2
Vertex accuracy	-2.8,+2.9	± 1.8	± 1.4
Angular resolution	-0.2,+0.2	-0.3,+0.3	-0.3,+0.3
Internal source pd †	± 0.0	-1.5,+1.6	-2.0,+2.2
External source pd	± 0.1	-1.0,+1.0	± 1.4
D ₂ O Cherenkov †	-0.1,+0.2	-2.6,+1.2	-3.7,+1.7
H ₂ O Cherenkov	± 0.0	-0.2,+0.4	-0.2,+0.6
AV Cherenkov	± 0.0	-0.2,+0.2	-0.3,+0.3
PMT Cherenkov †	± 0.1	-2.1,+1.6	-3.0,+2.2
Neutron capture	± 0.0	-4.0,+3.6	-5.8,+5.2
Cut acceptance	-0.2,+0.4	-0.2,+0.4	-0.2,+0.4
Experimental uncertainty	-5.2,+5.2	-8.5,+9.1	-13.2,+14.1
Cross section ²¹	± 1.8	± 1.3	± 1.4

Table 7: Systematic uncertainties on fluxes. The experimental uncertainty for ES (not shown) is -4.8,+5.0 percent. † denotes CC vs. NC anti-correlation.

values of the mixing parameters, spectral distortions and a measurable dependence on solar zenith angle are expected.^{34,35,36} This solar zenith angle dependence might be caused by interaction with the matter of the Earth (the MSW effect) and would depend not only on oscillation parameters and neutrino energy, but also on the neutrino path length and e^- density through the Earth.

To look for this distinctive signature of neutrino oscillation, a solar neutrino flux analysis similar to that described above was performed for two solar zenith angle θ_z bins: $\theta_z > 0$ (“day”) and $\theta_z < 0$ (“night”). During the development of this analysis, the data were partitioned into two sets of approximately equal livetime (split at July 1, 2000), each having substantial day and night components. Analysis procedures were refined during the analysis of Set 1 and fixed before Set 2 was analyzed with the same procedures. The latter thus served as an unbiased test. Unless otherwise stated, the analysis presented in the following is for the combined data set.

For each neutrino interaction channel, the asymmetry ratio (\mathcal{A}) of the measured day flux (ϕ_D) and night flux (ϕ_N):

$$\mathcal{A}_i = 2 \frac{\phi_N - \phi_D}{\phi_N + \phi_D} \quad i = \text{CC,NC,ES}$$

was determined. Figure 16 shows the day and night energy spectrum for all the accepted events above the kinetic energy threshold of 5 MeV and inside the fiducial volume of the inner 550 cm. In the extraction of the neutrino fluxes, backgrounds were subtracted separately for the two zenith angle bins. The results were then corrected for the orbital eccentricity by normalizing to an Earth-Sun distance of 1 AU. Table 8 is a summary of these extractions, where the day and night fluxes are given for the combined data. A χ^2 consistency test of the six measured fluxes between Sets 1 and 2 yielded a chance probability of 8%. A similar test done directly on the three asymmetry parameters gave a chance probability of 2%. No systematic has been identified, in either signal or background regions, that would suggest that the differences between Set 1 and Set 2 are other than a statistical fluctuation. For the combined analysis, \mathcal{A}_{CC} is $+2.2\sigma$ from zero, while \mathcal{A}_{ES} and \mathcal{A}_{NC} are -0.9σ and -1.2σ from zero, respectively. Note that \mathcal{A}_{CC} and \mathcal{A}_{NC} are strongly statistically anti-correlated ($\rho = -0.518$), while \mathcal{A}_{CC} and \mathcal{A}_{ES} ($\rho = -0.161$) and \mathcal{A}_{ES} and \mathcal{A}_{NC} ($\rho = -0.106$) are moderately anti-correlated. Table 9 gives the systematic uncertainties on the asymmetry parameters.

Systematic checks were made on a set of signals that are continuously present in the detector in order to look for any diurnal variation in the detector response. These studies include livetime verification using the detector diagnostic triggers (pulsed at

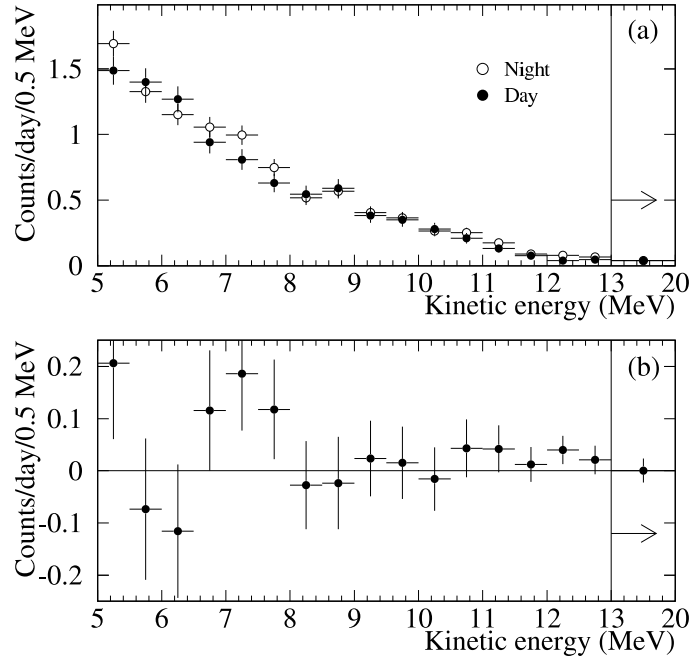


Figure 16: (a) Energy spectra (signals + background) for day and night bins. The final energy bin extends from 13.0 to 20.0 MeV. (b) Difference between the spectra (night - day). The day rate was 9.23 ± 0.27 events/day, and the night rate was 9.79 ± 0.24 events/day.

signal	Combined		$\mathcal{A}(\%)$
	ϕ_D	ϕ_N	
CC	$1.62 \pm 0.08 \pm 0.08$	$1.87 \pm 0.07 \pm 0.10$	$+14.0 \pm 6.3^{+1.5}_{-1.4}$
ES	$2.64 \pm 0.37 \pm 0.12$	$2.22 \pm 0.30 \pm 0.12$	$-17.4 \pm 19.5^{+2.4}_{-2.2}$
NC	$5.69 \pm 0.66 \pm 0.44$	$4.63 \pm 0.57 \pm 0.44$	$-20.4 \pm 16.9^{+2.4}_{-2.5}$

Table 8: The results of signal extraction for the two zenith angle bins. The fluxes shown here have been normalized to an Earth-Sun distance of 1 AU. An undistorted ^8B spectrum was assumed in the signal decomposition. The systematic uncertainties (combined set) include a component that cancels in the formation of the \mathcal{A} . Except for the dimensionless \mathcal{A} , the units are $10^6 \text{ cm}^{-2} \text{ s}^{-1}$. Flux values have been rounded, but the asymmetries were calculated with full precision.

Systematic	$\delta\mathcal{A}_{CC}$	$\delta\mathcal{A}_{ES}$	$\delta\mathcal{A}_{NC}$
Long-term energy scale drift	0.4	0.5	0.2
Diurnal energy scale variation	1.2	0.7	1.6
Directional energy scale var.	0.2	1.4	0.3
Diurnal energy resolution var.	0.1	0.1	0.3
Directional energy resolution var.	0.0	0.1	0.0
Diurnal vertex shift var.	0.5	0.6	0.7
Directional vertex shift var.	0.0	1.1	0.1
Diurnal vertex resolution var.	0.2	0.7	0.5
Directional angular recon. var.	0.0	0.1	0.1
PMT β - γ background	0.0	0.2	0.5
AV+H ₂ O β - γ background.	0.0	0.6	0.2
D ₂ O β - γ , neutrons background.	0.1	0.4	1.2
External neutrons background.	0.0	0.2	0.4
Cut acceptance	0.5	0.5	0.5
Total	1.5	2.4	2.4

Table 9: Effect of systematic uncertainties on \mathcal{A} (%). For presentation, uncertainties have been symmetrized and rounded.

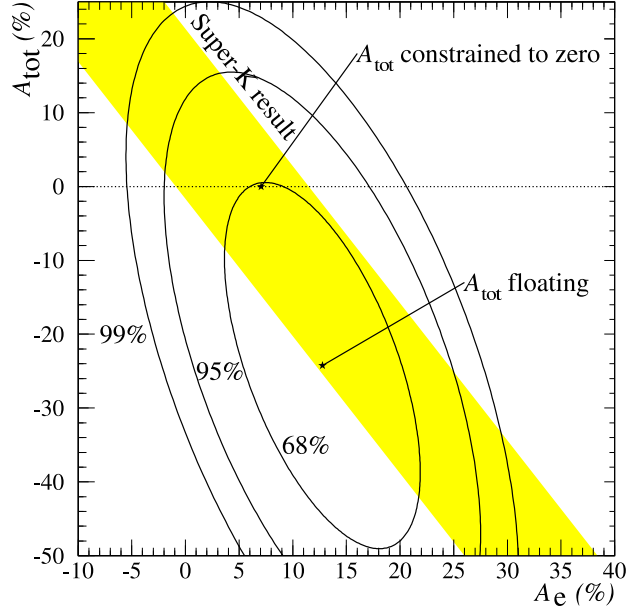


Figure 17: Joint probability contours for \mathcal{A}_{tot} and \mathcal{A}_e . The points indicate the results when \mathcal{A}_{tot} is allowed to float and when it is constrained to zero. The diagonal band indicates the 68% joint contour for the Super-K \mathcal{A}_{ES} measurement ($\mathcal{A}_{ES} = 3.3\% \pm 2.2\%$ (stat.) $^{+1.3\%}_{-1.2\%}$ (sys.).⁶)

5 Hz), variation in the observed muon rate, variation in the detector response to the muon-induced secondary neutrons, variation in the detector response to the radioactive hotspot on the acrylic vessel, and a neutrino signal extraction based on an east-west division instead of a zenith angle division. These studies do not show any significant diurnal variation in the detector response. In the study of the diurnal variation of the radioactive hotspot, a limit of 3.5% on the rate asymmetry was determined. Because of its steeply falling energy spectrum, a 0.3% limit was set for the diurnal variation of the detector's energy scale.

The asymmetry ratio in the ν_e flux \mathcal{A}_e and the total active neutrino flux \mathcal{A}_{total} can be readily determined by a change of variables in the signal extraction (see Sec. 3.5). Figure 17 shows the \mathcal{A}_e vs. \mathcal{A}_{tot} joint probability contours. Forcing $\mathcal{A}_{tot} = 0$, as predicted by active-only models, yielded the result of $\mathcal{A}_e = 7.0\% \pm 4.9\%$ (stat.) $^{+1.3\%}_{-1.2\%}$ (sys.).

Region	χ_{min}^2/dof	ϕ_B	$\mathcal{A}_e(\%)$	Δm^2	$\tan^2 \theta$	c.l.(%)
LMA	57.0/72	5.86	6.4	5.0×10^{-5}	0.34	—
LOW	67.7/72	4.95	5.9	1.3×10^{-7}	0.55	99.5

Table 10: Best fit points in the MSW plane for global MSW analysis using all solar neutrino data. ϕ_B is the best-fit ^8B flux for each point, and has units of $10^6 \text{ cm}^{-2} \text{ s}^{-1}$. Δm^2 has units of eV^2 . \mathcal{A}_e is the predicted asymmetry for each point.

4 Analysis of Neutrino Mixing Parameters

Using the day and night energy spectrum in Figure 16, an analysis to constrain the allowed MSW mixing parameters under the 2ν flavor mixing framework was performed. The radial distribution R and the solar angular correlation ($\cos \theta_\odot$) information were omitted. The expected number of counts were calculated for each of the energy bins by convoluting the ^8B neutrino spectrum,²⁸ the MSW survival probability, the neutrino interaction cross sections and the energy response of the SNO detector. With the high energy *hep* neutrino flux fixed at $9.3 \times 10^3 \text{ cm}^{-2} \text{ s}^{-1}$,⁷ the total ^8B flux ϕ_B , the difference Δm^2 between the squared masses of the two neutrino mass eigenstates, and the mixing angle θ are the only free parameters in a χ^2 fit of the MSW model to the observed data. Figure 18 shows the allowed regions in the Δm^2 - θ space at the 90%, 95%, 99% and 99.73% confidence levels.

Additional information from other solar neutrino experiments can be used in constructing the allowed Δm^2 - θ regions. In Figure 19, the solar neutrino fluxes measured by the Cl experiment¹ and the Ga experiments,^{3,4,5} along with the day and night neutrino energy spectra from the Super-Kamiokande experiment⁶ are combined with the day and night energy spectra from SNO in a fit of the three free parameters (ϕ_B , Δm^2 and θ). The *pp*, *pep*, ^7Be and *hep* neutrino fluxes were fixed at their predicted values in the Standard Solar Model.⁷ Table 10 summarizes the best fit points in the Δm^2 - θ plane. It is clear from this table and Figure 19 that this global analysis strongly favors the Large Mixing Angle (LMA) region, and excludes the region of $\tan \theta > 1$.

5 Conclusions

The results that are presented in this paper are truly groundbreaking. The neutral-current measurement is the first direct measurement of the total active ^8B neutrino flux.

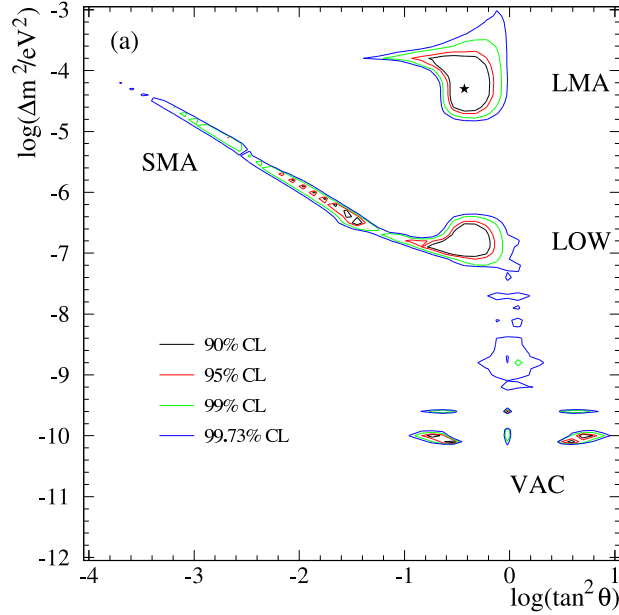


Figure 18: Allowed regions of the MSW plane determined by a χ^2 fit to SNO day and night energy spectra. The star indicates the best fit.

The neutrino fluxes measured by this channel provide strong evidence for neutrino flavor transformation, thereby solving the long-standing Solar Neutrino Problem. The significance of ν_e transforming into ν_μ or ν_τ is at 5.3σ level. This neutral-current measurement also verified the Standard Solar Model prediction of the ^8B neutrino flux.

The analysis presented in this paper is also the first direct measurement of the day-night asymmetries in the ν_e flux and the total neutrino flux. When combining the day and night energy spectra from SNO with results from other solar neutrino experiments in a global 2- ν flavor analysis of the MSW oscillation parameters, the LMA solution is strongly favored and the “dark side” ($\tan \theta > 1$) is excluded.

In the coming years, SNO will continue to make significant contributions to solar neutrino physics by refining the neutral-current measurement and the day-night asymmetry with different experimental techniques. It will also attempt to look for other possible signatures of neutrino oscillation (e.g. spectral distortion) and other rare processes (e.g. nucleon decays, anti- ν) in its future physics program.

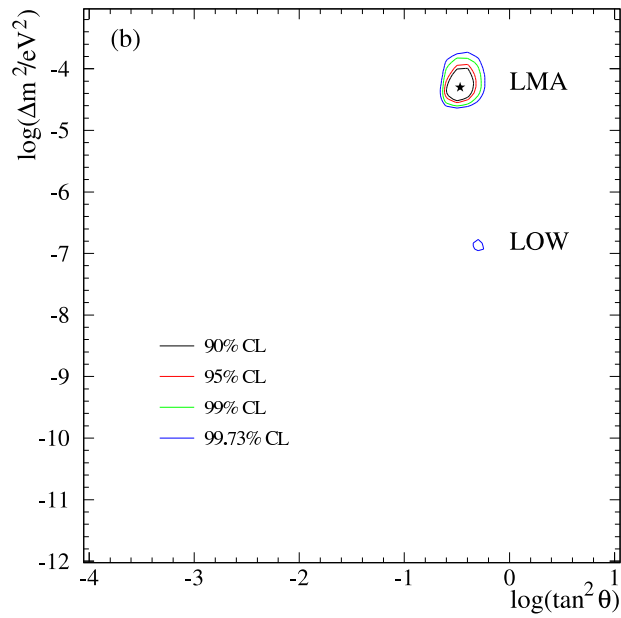


Figure 19: Allowed regions of the MSW plane determined by a χ^2 fit to SNO day and night energy spectra, the solar neutrino fluxes from the Cl and Ga experiments and the day and night energy spectra from the Super-Kamiokande experiment. The star indicates the best fit.

6 Acknowledgements

This research was supported by the Natural Sciences and Engineering Research Council of Canada, Industry Canada, National Research Council of Canada, Northern Ontario Heritage Fund Corporation and the Province of Ontario, the United States Department of Energy and in the United Kingdom by the Science and Engineering Research Council and the Particle Physics and Astronomy Research Council. Further support was provided by INCO, Ltd., Atomic Energy of Canada Limited (AECL), Agra-Monenco, Canatom, Canadian Microelectronics Corporation, AT&T Microelectronics, Northern Telecom and British Nuclear Fuels, Ltd. The heavy water was loaned by AECL with the cooperation of Ontario Power Generation. We thank the SNO technical staff for their strong contributions.

References

- [1] B. T. Cleveland *et al.*, *Astrophys. J.* **496**, 505 (1998).
- [2] K.S. Hirata *et al.*, *Phys. Rev. Lett.* **65**, 1297 (1990); K.S. Hirata *et al.*, *Phys. Rev.* **D44**, 2241 (1991); K.S. Hirata *et al.*, *Phys. Rev.* **D45**, 2170E (1992); Y. Fukuda *et al.*, *Phys. Rev. Lett.* **77**, 1683 (1996). This is the combined results from Kamiokande II and III.
- [3] J.N. Abdurashitov *et al.*, *J.Exp.Theor.Phys.* **95**, 181 (2002); J.N. Abdurashitov *et al.*, *Zh.Eksp.Teor.Fiz.* **122** 211 (2002).
- [4] W. Hampel *et al.*, *Phys. Lett.* **B447**, 127 (1999). This is the combined results from GALLEX I, II, III, and IV.
- [5] M. Altmann *et al.*, *Phys. Lett.* **B490**, 16 (2000). This is the most recent published result. The GNO Collaboration provided update at the Neutrino 2002 conference: $65.4 \pm 6.5(\text{stat.}) \pm 3.0(\text{sys.})$ SNU.
- [6] S. Fukuda *et al.*, *Phys. Rev. Lett.* **86**, 5651 (2001). Since the publication of our results, the Super-Kamiokande collaboration has published their latest results in S. Fukuda *et al.*, *Phys. Lett.* **B539**, 179 (2002). The global neutrino analysis presented in this paper used the results in the Super-Kamiokande results from 2001.
- [7] J.N. Bahcall, M.H. Pinsonneault, and S. Basu, astro-ph/0010346 v2 (2001).

- [8] A.S. Brun, S. Turck-Chièze, and J.P. Zahn, *Astrophys. J.* **525**, 1032 (1999); S. Turck-Chièze *et al.*, *Ap. J. Lett.*, **555** (2001).
- [9] The SNO Collaboration, *Nucl. Instr. and Meth.* **A449**, 172 (2000).
- [10] H.H. Chen, *Phys. Rev. Lett.* **55**, 1534 (1985).
- [11] Q.R. Ahmad *et al.* (the SNO Collaboration), *Phys. Rev. Lett.* **87**, 071301 (2001).
- [12] Q.R. Ahmad *et al.* (the SNO Collaboration), *Phys. Rev. Lett.* **89**, 011301 (2002).
- [13] Q.R. Ahmad *et al.* (the SNO Collaboration), *Phys. Rev. Lett.* **89**, 011302 (2002).
- [14] M.C. Browne *et al.*, *IEEE Trans. on Nucl. Sci.* **46**, 873 (1999).
- [15] M.R. Dragowsky *et al.*, *Nucl. Instr. Meth.* **A481**, 284 (2002).
- [16] N.J. Tagg *et al.*, *Nucl. Instr. Meth.* **A489**, 178 (2002).
- [17] R.T. Ford, *Nitrogen/Dye Laser System for the Optical Calibration of SNO*, M.Sc. thesis, Queen's University (1993).
- [18] R.T. Ford, *Calibration of SNO for Detection of ^8B Neutrinos*, Ph.D. thesis, Queen's University (1998).
- [19] B.A. Moffat, *The Optical Calibration of the Sudbury Neutrino Observatory*, Ph.D. thesis, Queen's University (2001).
- [20] A.W.P. Poon *et al.*, *Nucl. Instr. and Meth.* **A449**, 172 (2000).
- [21] Cross section uncertainty includes g_A uncertainty (0.6%), difference between NSGK (S. Nakamura, T. Sato, V. Gudkov and K. Kubodera, *Phys. Rev.* **C63** 034617 (2001)) and BCK (M. Butler, J.-W. Chen and X. Kong, *Phys. Rev.* **C63** 035501 (2001)) in cross section calculations (0.6%), radiative correction uncertainties (0.3% for CC, 0.1% for NC, A. Kurylov, M. J. Ramsey-Musolf and P. Vogel, *Phys. Rev.* **C65** 055501(2002)), uncertainty associated with neglect of real photons in SNO (0.7% for CC), and theoretical cross section uncertainty (1%, S. Nakamura *et al.*, [arXiv:nucl-th/0201062](https://arxiv.org/abs/nucl-th/0201062), (to be published)).
- [22] T.C. Andersen *et al.*, *Measurement of Radium Concentration in Water with Mn-Coated Beads at the Sudbury Neutrino Observatory*, nucl-ex/0208010, submitted to NIM for publication.
- [23] T.C. Andersen *et al.*, *A Radium Assay Technique Using Hydrous Titanium Oxide Adsorbent for the Sudbury Neutrino Observatory*, nucl-ex/0208015, submitted to NIM for publication.

- [24] I. Levine *et al.*, *Measurement of ^{222}Rn Dissolved in Water at the Sudbury Neutrino Observatory*, to be submitted to NIM.
- [25] X. Chen, D.Phil. thesis, Oxford University (1997).
- [26] G. McGregor, D.Phil. thesis, Oxford University (2002).
- [27] R.K. Taplin, D.Phil. thesis, Oxford University (1998).
- [28] C.E. Ortiz *et al.*, *Phys. Rev. Lett.* **85**, 2909 (2000).
- [29] J.F. Beacom and P. Vogel, *Phys. Rev. Lett.* **83**, 5222 (1999).
- [30] V.I. Tretyak and Yu.G. Zdesenko, *Phys. Lett.* **B505**, 59 (2001).
- [31] J. Evans and R. Steinberg, *Science*, **197**, 989 (1977).
- [32] Z. Maki, N. Nakagawa, and S. Sakata, *Prog. Theor. Phys.*, **28**, 870 (1962).
- [33] V. Gribov and B. Pontecorvo, *Phys. Lett.* **B28**, 493 (1969).
- [34] S. P. Mikheyev and A. Y. Smirnov in '86 *Massive Neutrinos in Astrophysics and in Particle Physics*, Proceedings of the Moriond Workshop, edited by O. Fackler and J. Tran Thanh Van, Editions Frontières, Gif-sur-Yvette, 1986, p. 335.
- [35] A. J. Baltz and J. Weneser, *Phys. Rev.* **D37**, 3364 (1988).
- [36] M. C. Gonzalez-Garcia, C. Peña-Garay and A. Y. Smirnov, *Phys. Rev.* **D63**, 113004 (2001).