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## **Journal**

ACS Catalysis, 9(7)

#### **ISSN**

2155-5435

#### **Authors**

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# **Publication Date**

2019-07-05

#### DOI

10.1021/acscatal.9b01485

# **Supplemental Material**

https://escholarship.org/uc/item/56n0k4qz#supplemental

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The Surface Structure of Co<sub>3</sub>O<sub>4</sub>(111) under Reactive Gas Phase Environments

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**Abstract:** 

In this work, we thoroughly examined the structure of the Co<sub>3</sub>O<sub>4</sub>(111) surface in oxidative and

reductive conditions, i.e. in equilibrium with realistic pressures of O<sub>2</sub>/H<sub>2</sub>O and H<sub>2</sub>/H<sub>2</sub>O, using

density functional theory with self-interaction and dispersion corrections. We found that this

surface is, in fact, hydroxylated under most reaction conditions, and that subjecting the surface to

H<sub>2</sub> increases surface Co<sup>2+</sup> concentration. Large structural distortions facilitate the reduction and

stabilization of the Co-rich termination. At 423 K, hydroxylation readily occurs on both the O-

rich and Co-rich surfaces even at water pressure as low as 10<sup>-15</sup> bar, and non-dissociated water

molecules appear on the O-rich surface when water pressure is above ~10<sup>-11</sup> bar. Our approach

showed good agreement with hybrid functional calculations and vibrational spectroscopy

experiments. Under most catalytic conditions, where water is present as a reactant, product, or

impurity, we predict that the Co<sub>3</sub>O<sub>4</sub>(111) surface will be hydroxylated. Hydroxyls groups and

structural distortions undoubtedly play large roles in shaping the surface's catalytic properties

and interaction with supported structures. The results of the study show the necessity of the

inclusion of hydroxylation and surface Co concentration in computational studies of Co<sub>3</sub>O<sub>4</sub> and

provide surface structures under various conditions to aid in future studies on structure and

catalytic reactivity of Co<sub>3</sub>O<sub>4</sub>(111) used as a support or as an active phase.

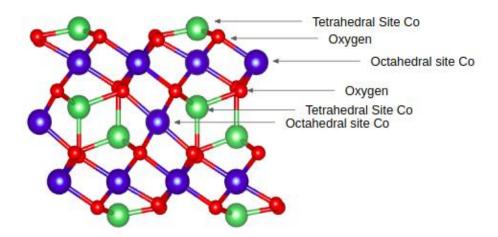
Keywords: Cobalt (II, III) oxide, DFT+U, hydroxylation, surface structure, surface stability

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#### 1. Introduction

Cobalt(II,III) oxide, Co<sub>3</sub>O<sub>4</sub>, is a versatile metal oxide capable of being both a catalyst support for single noble metal (NM) atoms, metal clusters, and a catalyst itself.<sup>1,2</sup> The catalytic properties of Co<sub>3</sub>O<sub>4</sub> have been studied extensively. In thermal catalysis, Co<sub>3</sub>O<sub>4</sub> is found to be active for low temperature CO/hydrocarbon oxidation, NO<sub>x</sub> reduction, N<sub>2</sub>O decomposition, and selective catalytic reduction of NO with NH<sub>3</sub>.<sup>1,3</sup> Most groups have attributed the material's extraordinary oxidation activity to its reducibility. On the other hand, NO reduction by H<sub>2</sub> can be performed over various NM<sub>1</sub>/Co<sub>3</sub>O<sub>4</sub> catalysts, where Co<sub>3</sub>O<sub>4</sub> acts as the support.<sup>4</sup> One promising application of the (111) surface is the oxidative dehydrogenation (ODH) of alkanes.<sup>5,6</sup> The (111) surface has recently been shown to have good selectivity for ethane ODH against combustion.<sup>7</sup>

Co<sub>3</sub>O<sub>4</sub> has the normal spinel structure and an indirect band gap. Co<sup>2+</sup> in tetrahedral (T<sub>d</sub>) sites are in the quadruplet spin state while Co<sup>3+</sup> in octahedral (O<sub>h</sub>) sites are in the singlet spin state.<sup>8</sup> Co<sup>2+</sup> ions aligned on the same (111) plane have the same spin direction, and these (111) planes are antiferromagnetically aligned in the bulk.<sup>9</sup> Experimentally, Co<sub>3</sub>O<sub>4</sub> particles calcined at 673 K expose (100), (110), and (111) surfaces in similar portions.<sup>10</sup> In the case of the (111) surface, it is polar and can be terminated by Co in tetrahedral sites (Co-T<sub>d</sub>), octahedral sites (Co-O<sub>h</sub>), O<sup>2-</sup>, or a mix of the three for a total of six unique bulk-truncated terminations (Figure 1). Thin films exposing O<sup>2-</sup> and Co-T<sub>d</sub> relax inwards to compensate for the surface dipole.<sup>11,12</sup> The relative stability of bulk-truncated low miller index Co<sub>3</sub>O<sub>4</sub> surfaces have been examined by various studies under purely oxygen-containing environments.<sup>10,13-15</sup>



**Figure 1. Side view of Co<sub>3</sub>O<sub>4</sub> through the [111] direction.** From truncating the Co<sub>3</sub>O<sub>4</sub> bulk, a total of six possible terminations of the (111) surface can be made. This surface is a type III surface by Tasker's definitions, where the group of the six labeled layers is stoichiometric when stacked but has a net dipole in the z-direction

Experimental characterization of the Co<sub>3</sub>O<sub>4</sub> (111) surface is a challenging task. Busca et al. obtained IR spectra of the surface after diluting the sample with KBr and activating under vacuum. <sup>16</sup> Vibrational modes in the fingerprint region (<1000 cm<sup>-1</sup>) were present, but the overall IR transmission window is severely limited. <sup>16</sup> High temperature evacuation or H<sub>2</sub> reduction were found to widen this window, but carry the risk of reducing the Co<sub>3</sub>O<sub>4</sub> sample to CoO. To mitigate this problem, alternative characterization techniques can be used, but they are generally limited to ultra-high vacuum (UHV). <sup>17</sup> Although x-ray photoelectron spectroscopy (XPS) can now be performed at ambient pressures, it is still difficult to identify types of surface O because OH, CO<sub>3</sub><sup>2-</sup>, and defective O all have similar O 2p binding energy. <sup>2,17–20</sup> Petitto et al. used XPS and high-resolution electron energy loss spectroscopy (HREELS) to study surface OH created on (111)-terminated single Co<sub>3</sub>O<sub>4</sub> crystals after annealing in H<sub>2</sub>O. <sup>21</sup> Schwarz et al. probed Co<sub>3</sub>O<sub>4</sub>(111) epitaxial films using XPS and infrared reflection absorption spectroscopy (IRAS) after exposing them first to D<sub>2</sub>O or CO. <sup>22,23</sup> To summarize, experimental characterization of the Co<sub>3</sub>O<sub>4</sub>(111) surface can only provide a limited view of the surface at low pressures. These

difficulties highlight the importance of theoretical studies for understanding surface structure, especially outside the limits of experimental characterization.

Although many computational studies have been performed on the Co<sub>3</sub>O<sub>4</sub>(111) surface using density functional theory (DFT), they generally lack three considerations. First, only a few studies regard the magnetic structure, even when it was shown that inconsistent magnetic coupling induces errors in calculated energy. <sup>24–26</sup> Second, mechanistic studies of reactions on the (111) surface generally do not consider the role of hydroxyls. <sup>5–7</sup> Low pressure H<sub>2</sub>O/D<sub>2</sub>O exposure experiments have shown that hydroxyls are stable on the surface up to 570 K. <sup>23</sup> Further, H<sub>2</sub> was also shown to dissociate easily over the (111) termination. <sup>27</sup> Therefore, any H<sub>2</sub>O or H<sub>2</sub> impurity in the gas environment will likely produce hydroxyls. Third, the electronic effects of surface reduction are not addressed in most studies. Early studies used generalized gradient approximation (GGA) exchange correlation functionals but did not include any self-interaction error correction for Co 3d electrons. <sup>13,28</sup> Conventional DFT predicts a bulk band gap that is too narrow. <sup>9,29</sup> Only a few studies have addressed electron localization upon reduction. <sup>24–26</sup>

The  $Co^{2+}$  and  $O^{2-}$  terminated surface is the most often used as the active site structure for the (111) termination.<sup>5-7</sup> In this study, we seek to provide insight into the structure of this surface under various reaction environments. We examine the structure of the surface in equilibrium with  $O_2$ ,  $H_2O$ , and  $H_2$  to address surface hydroxylation, vacancy formation, and possible change of the Co concentration, while taking a rigorous approach and minding both the bulk magnetic structure and electronic effects that occur through surface modifications.

# 2. Computational details

#### 2.1 Computational methods

Density functional theory (DFT) calculations were performed using the Vienna ab-initio simulation package (VASP), version 5.4.1.30-32 Exchange-correlation energy was calculated using the Perdew-Burke-Erzenhof (PBE) functional.<sup>33</sup> Spin polarization was applied in all calculations unless otherwise specified. Since PBE does not account for van der Waals (vdW) forces, the dDsC dispersion-correction method was used. 34,35 The projector-augmented-wave (PAW) method was used to describe the core electrons.<sup>36</sup> The one-electron wavefunctions are developed on a basis set of plane waves with a 450 eV energy cutoff. To correct for the large self-interaction error of the 3d electrons of Co, a Hubbard-like repulsion term was added using Dudarev's approach (DFT+U), with  $U_{\text{eff}} = U - J = 3.5 \text{ eV}$ . This value of  $U_{\text{eff}}$  is consistent with those found in literature as a value that properly describes the energy of formation of Co<sub>3</sub>O<sub>4</sub> from  $CoO.^{38}$  Due to the controversy on choosing the appropriate  $U_{eff}$ , calculations using  $U_{eff} = 2.0$  eV were also performed starting from  $U_{eff} = 3.5$  eV geometry. Single point calculations using the Heyd-Scuseria-Ernzerhof functional with dispersion (HSE06+dDsC) were also performed using geometry and charge density obtained from  $U_{\text{eff}} = 3.5 \text{ eV}$  as benchmarks for vacancy formation energies. 39,40 Previously, we benchmarked the energy of water adsorption on 11-layered Co-T<sub>d</sub>terminated slabs at the PBE+U+dDsC level (Ueff=3.5 eV), against HSE06+dDsC. We found that water is under bound by 0.15 eV/H<sub>2</sub>O when using PBE+U+dDsC. This same correction was applied to all water-adsorbed structures to construct the 1-D and 2-D surface stability diagrams in this work.

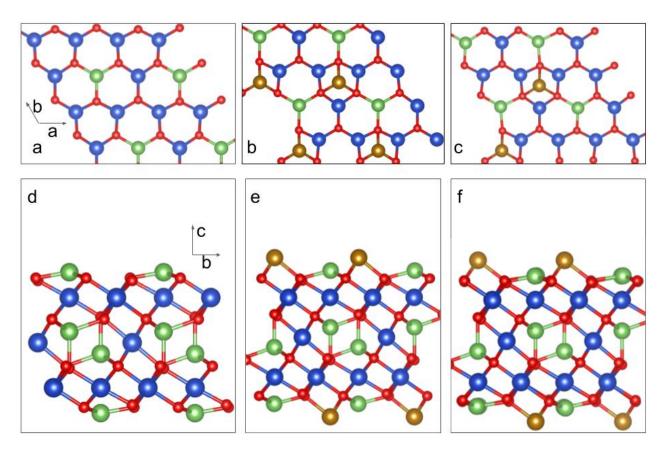
Since  $Co^{2+}$  ions in  $Co_3O_4$  are known to have antiferromagnetic ordering in the bulk, while  $Co^{3+}$  ions have no magnetic moment, careful considerations were taken to ensure that the bulk-

like section of the slab retains the correct Co<sup>2+</sup> magnetic ordering and spin state. Electron localization or delocalization were also properly ensured on the modified surface, so that no electrons were spuriously added to the other side of the slab. Specifically, the occupancy matrix control technique developed by Allen and Watson was used, where a spin-polarized VASP WAVECAR file containing desirable Co 3d projected electron occupations was generated and used as the starting guess for the system.<sup>41</sup>

The surface structures considered in this work were compared to each other by their surface Gibbs free energy<sup>42,43</sup>, under a given temperature and a set of gas pressures (Section S1). DFT energies used to develop the surface stability diagrams were all computed from structures in  $(2 \times 2)$  supercells (a, b, c axes shown on Figure 2 a and d) to ensure proper calculation of vdW contributions to electronic energy with dDsC. At PBE+U+dDsC level, assessing the adsorption energy of water in a  $(1 \times 1)$  cell induces an error within 0.025 eV/H<sub>2</sub>O. The outermost layer of O and Co atoms were displaced in the frequency calculations (Table S13). For a  $(2\times2)$  slab (a = b =11.4 Å,  $\gamma = 120^{\circ}$ ), the Brillouin zone was sampled using a 3 × 3 × 1 Monkhorst-Pack mesh.<sup>44</sup> Electronic energy difference between steps in each self-consistent field (SCF) cycle were converged to 10<sup>-6</sup> eV, and atomic positions were converged until the forces on unconstrained atoms are less than 0.02 eV/Å<sup>-1</sup>. A (1  $\times$  1) cell was used for HSE06+dDsC calculations, with a 3  $\times$  3  $\times$  1 Monkhorst-Pack mesh. At PBE+U+dDsC level, reducing the k-point mesh from 6  $\times$  6  $\times$ 1 to  $3 \times 3 \times 1$  in a  $(1 \times 1)$  cell only increases the total energy by <0.1 meV/atom. The lattice parameter and the atomic positions of the bulk cell were relaxed simultaneously. Structural and electron density visualizations were performed using the VESTA program.<sup>45</sup> Vibrational frequencies were calculated using finite differences with a step size of 0.015 Å. The calculated

v(OH) and v(OD) modes were rescaled by factors of 0.9789 and 0.9918 respectively; these factors were calculated using the same scheme as that specified by Freund, Sauer et al.<sup>46</sup>

#### 2.2 Bulk and surface models



**Figure 2.** Top (a, b, c, through the c axis) and side (d, e, f, through the a axis) views of (2x2) cells (area 1.13 nm²) of the three bare  $Co_3O_4(111)$  terminations considered in this work, containing (red) O, (green) Co-T<sub>d</sub> and (blue) Co-O<sub>h</sub> and additional surface Co (gold). a, d: O-rich termination,  $(Co^{Td})_x$ , has formula  $Co_{44}O_{64}$ , containing four exposed Co in  $C_{3v}$  sites on each surface, belonging to bulk  $T_d$  sites in the bulk. b, e: Co-rich termination,  $(Co^{Oh}Co^{Td})_x$ , has formula  $Co_{52}O_{64}$ , containing eight exposed Co in  $C_{3v}$  sites on each surface; four originally at bulk  $T_d$  sites and four originally at bulk  $O_h$  sites. c, f: stoichiometric termination,  $(Co^{Oh}_{0.5}Co^{Td})_x$ , has formula  $Co_{48}O_{64}$ , containing six exposed  $C_{3v}$  Co on each surface; two originally at bulk  $O_h$  sites and four at bulk  $T_d$  sites.

A cubic  $Co_{24}O_{32}$  cell was used to model the  $Co_3O_4$  bulk. The lattice parameter was found to be 8.06 Å, which matches the experimental lattice parameter, 8.08 Å.<sup>47</sup> A magnetic moment of 2.64  $\mu_B$  was found on the  $Co^{2+}$  ions in the bulk, consistent with values obtained by Selloni et al.,

but lower than the 3.26  $\mu_B$  measured by Roth.<sup>8,9</sup> As previously mentioned, Co<sub>3</sub>O<sub>4</sub>(111) is a Tasker type III polar surface with six different bulk-truncated terminations. To simulate the surface, three terminations, previously proposed by Zasada et al., were chosen (Figure 2).<sup>48</sup> (Co<sup>Td</sup>)<sub>x</sub> and (Co<sup>Oh</sup>Co<sup>Td</sup>)<sub>x</sub> were cut directly from the bulk and are oxygen-rich and cobalt-rich respectively, while (Co<sup>Oh</sup>0.5Co<sup>Td</sup>)<sub>x</sub> is stoichiometric.<sup>48</sup> A symmetric 11-layer slab was used to model (Co<sup>Td</sup>)<sub>x</sub>, and symmetric 13 layer slabs were used to model (Co<sup>Oh</sup>Co<sup>Td</sup>)<sub>x</sub> and (Co<sup>Oh</sup>0.5Co<sup>Td</sup>)<sub>x</sub>. In a (2 × 2) cell, they have formulas Co<sub>44</sub>O<sub>64</sub>, Co<sub>52</sub>O<sub>64</sub> and Co<sub>48</sub>O<sub>64</sub>, containing 3.56, 7.11, and 5.33 Co/nm<sup>2</sup> respectively. The (Co<sup>Td</sup>)<sub>x</sub> surface contains two unique types of surface O: those directly adjacent to surface Co<sup>Td</sup> (Figure 2a, red O bonded to 2 blue Co and 1 green Co) and those in isolation (Figure 2a, bonded to 3 blue Co). To simulate 5.33 Co/nm<sup>2</sup>, Co is added to three-fold adjacent O sites, generating another category of O bonded to both Co<sup>Td</sup> and Co<sup>Oh</sup> (Figure 2b, red O bonded to 2 blue Co, 1 green Co, and 1 gold Co). Finally, when all the three-fold O sites are occupied on the surface by Co at 7.11 Co/nm<sup>2</sup>, two types of O exist on the surface: isolated O and those between Co<sup>Td</sup> and Co<sup>Oh</sup>.

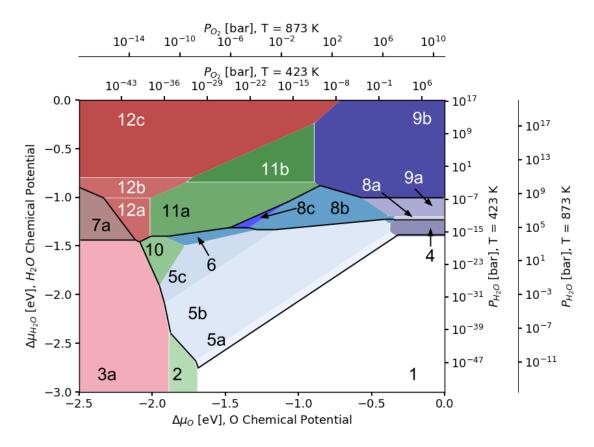
The antiferromagnetic alignment of surface Co ions on the  $(Co^{Oh}_{0.5}Co^{Td})_x$  termination was challenging to keep consistent; therefore, they were kept ferromagnetic and shifted by a constant correction per unit cell (see Figure S12). At least 15 Å of vacuum was used to prevent interactions between slabs in the z direction. For a given slab calculation, the five innermost layers were fixed to their bulk positions and all other atoms were relaxed. Adsorption of atoms and molecules were performed on only one side of the slab to prevent surface-surface interactions through the slab. Dipole corrections were applied in all slab calculations. H adsorption,  $H_2O$  adsorption, and oxygen vacancy formation were considered separately on all three surfaces, and the combination of up to  $1.78 H_2/nm^2$  (up to 4 H added per unit cell) and 2.78

 $H_2O/nm^2$  (up to 3  $H_2O$  added per unit cell) on the  $(Co^{Td})_x$  termination as a simplified case of  $H/H_2O$  co-adsorption. At given set of temperature and partial pressure, surfaces are compared to each other through approximate surface energy (see section S1).

# 3. Results and discussion

# 3.1 C<sub>03</sub>O<sub>4</sub>(111) surface stability diagrams

For oxidative reactions, we consider the surface to be in equilibrium with  $O_2$  and  $H_2O$ . The surface stability diagram under realistic  $O_2$  and  $H_2O$  chemical potentials ( $\Delta\mu_{H2O} > -2.5$  eV,  $\Delta\mu_O > -2.5$  eV) is shown below:



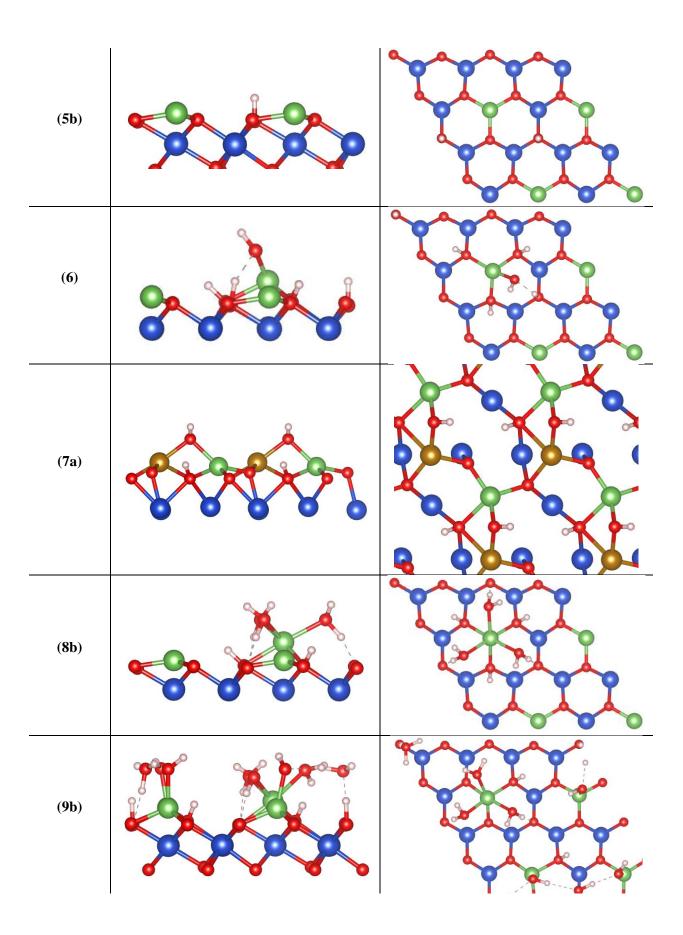
**Figure 3.** Surface stability diagram showing the most stable surface structure as a function of  $O_2$  and  $H_2O$  chemical potential and equivalent oxygen/water pressure at 423 K and 873 K. Regions are classified in 3 ways. I) By surface Co concentration: blue/grey/white: 3.56 Co/nm<sup>2</sup>; green: 5.33 Co/nm<sup>2</sup>, and red/brown: 7.11 Co/nm<sup>2</sup>. II) By general appearance of surface hydroxylation: bottom bold line separates non-hydroxylated regions and hydroxylated regions, middle bold line (only the segment running through the blue regions, or when  $\Delta\mu_O > -1.5$  eV) separates isolated hydroxyls and  $Co(H_2O)_3$  (hereby referred to as "triaqua") surface complexes (with eventually additional OH groups), and top bold line separates isolated triaqua surface complexes and large OH/H<sub>2</sub>O structures. III) By specific numbered structure (see Table 1 for descriptions, Figure 2 and Table 2 for representative geometries, and Table S10 for all geometries).

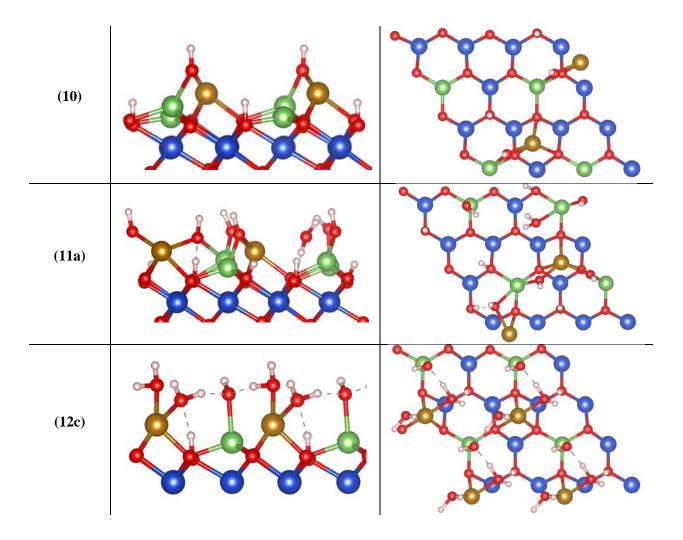
**Table 1.** Description of the specific types of surface structures observed in Figure 3 and 5

Label(s) on Figure 3 and/or Figure 5	Termination, coverage, and appearance
1	Bare (Co <sup>Td</sup> ) <sub>x</sub> surface
2	Bare (Co <sup>Oh</sup> <sub>0.5</sub> Co <sup>Td</sup> ) <sub>x</sub> surface
3 a, b	(Co <sup>Oh</sup> Co <sup>Td</sup> ) <sub>x</sub> surface with: a) 3.56 O <sub>vac</sub> /nm <sup>2</sup> and b) 7.11 O <sub>vac</sub> /nm <sup>2</sup>
4	(Co <sup>Td</sup> ) <sub>x</sub> with 1.78 OH/nm <sup>2</sup> , appearing as 1 dissociated H <sub>2</sub> O
5 a-c	$(\text{Co}^{\text{Td}})_x$ with: a) 0.89, b) 1.78, c) 2.67 OH/nm <sup>2</sup> , appearing as adsorbed H
6	$(\text{Co}^{\text{Td}})_x$ with 5.33 OH/nm <sup>2</sup> , appearing as 4 adsorbed H and 1 dissociated H <sub>2</sub> O
7 a-e	$(\text{Co}^{\text{Oh}}\text{Co}^{\text{Td}})_x$ with: a) 7.11 OH/nm², b) 8.00 OH/nm², c) 8.89 OH/nm², d) 9.78 OH/nm², e) 10.67 OH/nm², appearing as dissociated H <sub>2</sub>
8 a-c	(Co <sup>Td</sup> ) <sub>x</sub> with: a) 0 OH/nm <sup>2</sup> and 2.67 H <sub>2</sub> O/nm <sup>2</sup> b) 2.67 OH/nm <sup>2</sup> and 2.67 H <sub>2</sub> O/nm <sup>2</sup> , c) 2.67 H <sub>2</sub> O/nm <sup>2</sup> and 3.56 OH/nm <sup>2</sup> , appearing as one triaqua surface complex eventually surrounded by adsorbed H
9 a, b	(Co <sup>Td</sup> ) <sub>x</sub> with: a) 3.56 H <sub>2</sub> O/nm <sup>2</sup> , 3.56 OH/nm <sup>2</sup> and b) 5.33 OH/nm <sup>2</sup> , 4.44 H <sub>2</sub> O/nm <sup>2</sup> , each containing one triaqua surface complex and H-bonded OH/H <sub>2</sub> O chain. The triaqua complex is incorporated into the chain in structure 9b
10	(Co <sup>Oh</sup> <sub>0.5</sub> Co <sup>Td</sup> ) <sub>x</sub> with 3.56 OH/nm <sup>2</sup> , appearing as dissociated H <sub>2</sub> O
11 a, b	$(\text{Co}^{\text{Oh}}_{0.5}\text{Co}^{\text{Td}})_x$ with: a) 8.89 OH/nm <sup>2</sup> and 0.89 H <sub>2</sub> O/nm <sup>2</sup> , b) 10.89 OH/nm <sup>2</sup> and 1.78 H <sub>2</sub> O/nm <sup>2</sup> , appearing as dissociated and molecular H <sub>2</sub> O
12 a-c	(Co <sup>Oh</sup> Co <sup>Td</sup> ) <sub>x</sub> with: a) 8.89 OH/nm <sup>2</sup> , 0.89 H <sub>2</sub> O/nm <sup>2</sup> , b) 7.11 OH/nm <sup>2</sup> and 5.33 H <sub>2</sub> O/nm <sup>2</sup> , c) 7.11 OH/nm <sup>2</sup> and 7.11 H <sub>2</sub> O/nm <sup>2</sup> , appearing as dissociated and molecular H <sub>2</sub> O linked by H bonds.

**Table 2.** Side (through the a axis) and top (through the c axis) views of representative structures in regions 1-12 on the surface stability diagram (see Table 1, Figure 3, and Figure 5). 4 and 9b are adapted from ref <sup>49</sup>.

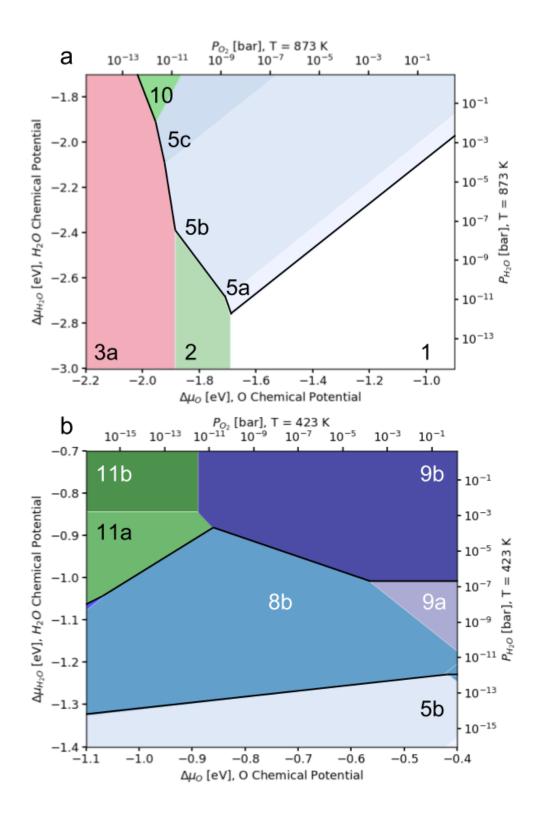
Number	Geometry, side view, through a axis	Geometry, top view, through c axis
(1)		
(2)		
(3)	C D D D D D D D D D D D D D D D D D D D	a D
(4)		





While in equilibrium with  $O_2$  and  $H_2O$ , the surface stability diagram is dominated by the  $(Co^{Td})_x$  termination. Only considering the  $(Co^{Td})_x$  regions, the diagram can be drawn into four sections based on the type of surface hydroxyls present. It should be noted that the structures listed in Table 1 are representatives of the various macroscopic regions on the stability diagram; it is possible that there are more stable isomers. These four regions are separated by bold black lines (Figure 3). The lowest line separates the non-hydroxylated regions below the line from strong surface hydroxyl groups. Under  $H_2O$  and  $O_2$  pressure, these regions can be interpreted as the equilibrated states of  $H_2O$  adsorption on surface vacancies. The coverage of isolated OH should be  $2\sim4$  OH/nm<sup>2</sup> (Table 1, Table 2 and Table S1). As mentioned in section 2.2, two types

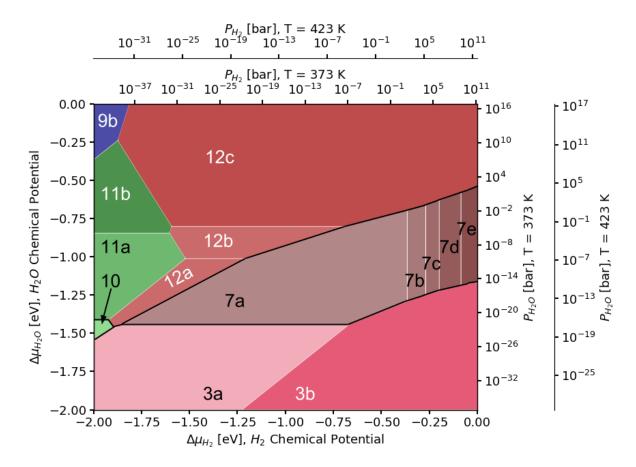
of O exist on the  $(Co^{Td})_x$  surface, ones adjacent to surface  $Co^{Td}$ , coordinated to the surface  $Co^{Td}$  and to subsurface  $Co^{Oh}$ , and ones in isolation, coordinated to 3 subsurface  $Co^{Oh}$ . It is more favorable to populate isolated O first. The line second-to-the-bottom, running across at roughly -1.3 eV, separates the isolated OH region to the one containing triaqua surface complexes, surrounded by OH groups coming from adsorbed H atoms. Structure 8b (Table 2, row 6) can serve as a representative for such OH-surrounded triaqua surface complexes, where one surface  $Co^{Td}$  is coordinated to three molecular  $H_2O$  and adjacent to three surface OH. Finally, the upper line, running across at roughly -1.1 eV, divides the triaqua surface complex-containing region below and the fully wetted region, containing extended H-bonded OH/ $H_2O$  structures. On the other hand, the diagram can also be divided by surface Co concentration. Regions in shades of blue/grey and white are Co poor  $(3.56 \ Co/nm^2)$ , regions in green are stoichiometric  $(5.33 \ Co/nm^2)$ , and regions in red and brown are Co-rich  $(7.11 \ Co/nm^2)$ . As expected, surface Co concentration increases as  $\Delta\mu_O$  decreases through all the major zones, i.e. when the gas environment shifts towards O poor conditions.



**Figure 4.** Portions of Figure 3 showing the  $H_2O/O_2$  surface stability diagram under realistic pressures at (top) 423 K and (bottom) 873 K. Regions are numbered and colored the same way as Figure 3. Only hydroxylated regions can be accessed at 873 K, while hydroxyls and triaqua complexes can form at 423 K. Higher surface Co concentration can be reached under low  $P_{O2}$ .

Although many possible terminations of this surface exist in theory, not all terminations can be accessed under realistic pressure and temperature. At 423 K (Figure 4, top), only a fraction of the hydroxylated regions (structures 5b, 8b, 9a, 9b, 11a, and 11b) can be visited at equilibrium; though, clear separations among isolated OH, triaqua complexes, and hydroxyl networks can still be observed when P<sub>H2O</sub> is increased. When O<sub>2</sub> pressure is between 1 bar and 10<sup>-4</sup> bar, and H<sub>2</sub>O pressure is above 10<sup>-7</sup> bar, the surface is predicted to be heavily hydroxylated, with a triaqua complex incorporated into the OH/H<sub>2</sub>O chain (9b). Triaqua complexes is predicted to be separated from the chain below 10<sup>-7</sup> bar (8b, 9a), and only isolated hydroxyls (5b) are stable below 10<sup>-12</sup> bar, a moisture level almost impossible to obtain. Overall, at 423 K, the surface stability diagram is dominated by triaqua complex regions with various amounts of additional H<sub>2</sub>O/OH groups. On the other hand, if the temperature is increased to 873 K (Figure 4, bottom), isolated hydroxyls will be exposed under 10<sup>-3</sup> bar. We note that, for the surface Co concentration to increase to above 3.56 Co/nm<sup>2</sup>, P<sub>O2</sub> needs to be reduced to below 10<sup>-11</sup> bar while P<sub>H2O</sub> is kept below 10<sup>-9</sup> bar, i.e. in UHV conditions.

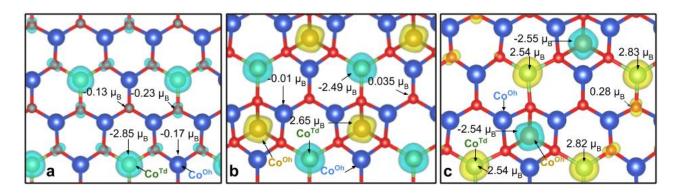
Finally, for reductive reactions involving H<sub>2</sub>, we consider the surface to be in equilibrium with H<sub>2</sub> and H<sub>2</sub>O. Under these conditions, the system temperature must be kept low, so that kinetics for bulk reduction is slow enough, and the near-surface region is kinetically trapped to retain structure like bulk Co<sub>3</sub>O<sub>4</sub>. By monitoring the surface oxidation state of NM<sub>1</sub>/Co<sub>3</sub>O<sub>4</sub> catalysts under NO/H<sub>2</sub> environments via ambient pressure (AP) XPS, Nguyen et al. observed that the near-surface regions generally do not fully reduce to CoO until above 500~600 K.<sup>4</sup> Under H<sub>2</sub>/H<sub>2</sub>O pressure, the surface stability diagram is as follows:



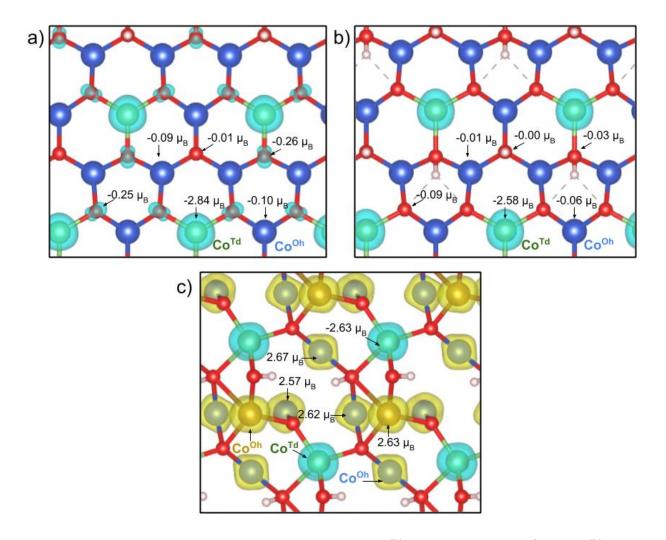
**Figure 5.** Surface stability diagram showing the most stable surface structure as a function of H<sub>2</sub> and H<sub>2</sub>O chemical potential and H<sub>2</sub>/H<sub>2</sub>O pressure at 423 K. The color scheme and numbering are the same as that in Figure 3. Regions are again classified in same 3 ways as those in Figure 3. I) By surface Co concentration: blue: 3.56 Co/nm<sup>2</sup>; green: 5.33 Co/nm<sup>2</sup>, and red/brown: 7.11 Co/nm<sup>2</sup>. II) By general appearance of surface hydroxylation: bottom bold line separates non-hydroxylated regions and hydroxylated regions and top bold line separates OH-only regions and molecular H<sub>2</sub>O-containing structures. III) By specific numbered structure (see Table 1 for descriptions, Table 2 for representative geometries, and Table S11 for all geometries).

Under  $H_2$  and  $H_2O$  pressure, separation of non-hydroxylated, hydroxylated, and hydrated regions under  $H_2/H_2O$  is also very apparent. When the  $H_2O$  chemical potential increases above - 1.5 eV, the surface transitions from a vacancy-covered state (representative structure: region 3a, see Table 2, row 1) to a hydroxylated state (representative structure: region 7a, see Table 2, row 5). When  $H_2O$  chemical potential increases to above roughly -0.75  $\sim$  -1.00 eV, the surface transitions to a fully hydrated state (representative structure: region 12c, see Table 2, row 10). We note that the bare  $(Co^{Oh}Co^{Td})_x$  termination does not appear on the stability diagram. Upon

reduction by H adsorption or  $O_{vac}$  formation, the surface undergoes a large relaxation (Figure S13). This relaxation makes vacancy formation much easier than on  $(Co^{Oh}Co^{Td})_x$  (Table S6), and H adsorption strong as well (Table S3), but not as strong as that on  $(Co^{Td})_x$ , which is expected, as the  $(Co^{Td})_x$  termination appears overly oxidized. If we again only consider realistic H<sub>2</sub> and H<sub>2</sub>O pressures, the surface is dominated by this heavily relaxed hydroxylated region (region 7a in Figure 5, see also Table 2 for structure).



**Figure 6.** Top view of spin density difference on a) bare  $(Co^{Td})_x$ , b) bare  $(Co^{Oh}Co^{Td})_x$ , and c)  $(Co^{Oh}_{0.5}Co^{Td})_x$ , at isosurface level 0.04, where pale blue represents negative and yellow represents positive spin density difference. Key:  $Co^{Td}$ : green, surface  $Co^{Oh}$ : gold, subsurface  $Co^{Oh}$ : blue. Surface Co (green) on the  $(Co^{Td})_x$  termination appear to be slightly more oxidized than bulk  $Co^{2+}$ , with slightly higher projected magnetic moments. Charge for this offstoichiometry is compensated by small density differences on surface O. No projected magnetic moments are found on surface O in the  $(Co^{Oh}Co^{Td})_x$  termination. All the added Co appear to be  $C^{o2+}$  based on projected magnetic moments, while Co originally on the surface appear slightly reduced. Less Co can be added to the surface to form  $(Co^{Oh}_{0.5}Co^{Td})_x$ ; these Co also appear to be slightly reduced  $Co^{2+}$  and quenches spin differences on some surface O.



**Figure 7.** Top view of spin density difference on a)  $(Co^{Td})_x$  with 1.78 OH/nm², b)  $(Co^{Td})_x$  with 7.11 OH/nm², and c)  $(Co^{Oh}Co^{Td})_x$  with 7.11 OH/nm². Key:  $Co^{Td}$ : green, surface  $Co^{Oh}$ : gold, subsurface  $Co^{Oh}$ : blue, H: white. Electrons from adsorbing H on  $(Co^{Td})_x$  are transferred to surface O (red), subsurface  $Co^{Oh}$ , and surface  $Co^{Td}$ . On the other hand, electrons from H adsorption on  $(Co^{Oh}Co^{Td})_x$  are transferred mostly to subsurface  $Co^{Oh}$ .

Besides surface relaxation, adsorbate-induced charge transfer is very different on the bare Co-rich  $(Co^{Oh}Co^{Td})_x$  as well, relative to Co-poor  $(Co^{Td})_x$ . On  $(Co^{Td})_x$ , initial  $H_2$  adsorption results in delocalization of charge onto surface O and subsurface  $Co^{Oh}$ ; the small projected magnetic moments on surface O disappear (Figure 6a, 7a), while subsurface Co's projected magnetic moments decrease in absolute value (from -0.177  $\mu_B$  to -0.088  $\mu_B$ , Figure 6a, 7a). At a higher OH concentration on  $(Co^{Td})_x$ , charge is added to surface  $Co^T$ , surface O, and subsurface  $Co^{Oh}$  as well

(Figure 6a, 7b). On the other hand, on the  $(Co^{Oh}Co^{Td})_x$  surface, the charge is almost entirely transferred to subsurface  $Co^{Oh}$  (Figure 6b, 7c). This result can be rationalized by again considering the overall surface stoichiometry; due to the overoxidized state of the  $(Co^{Td})_x$  termination,  $H_2$  adsorbs strongly to balance the stoichiometry. Further, the  $(Co^{Oh}Co^{Td})_x$  slab is closer to that of CoO, whereas the  $(Co^{Td})_x$  is closer to  $Co_2O_3$ , making near-surface Co atoms in the  $(Co^{Td})_x$  slab closer to  $Co^{3+}$ , and near-surface Co in  $(Co^{Oh}Co^{Td})_x$  closer to  $Co^{2+}$ . By the spectrochemical series,  $Co^{2+}$  induces a much smaller crystal field splitting than  $Co^{3+}$ , and as a result, subsurface Co on reduced  $(Co^{Td})_x$  have magnetic moment smaller than 1  $\mu_B$ , while those on  $(Co^{Oh}Co^{Td})_x$  have magnetic moment closer to 2.6  $\mu_B$  (Figure 7b, 7c).

## 3.2 Comparison to experimental results reached through vibrational spectroscopy

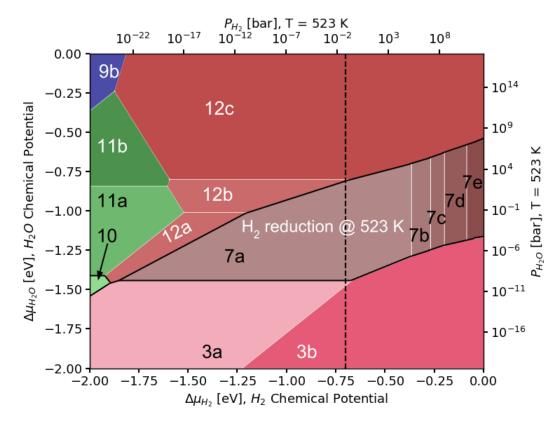
Experimental characterization of surface hydroxyls on Co<sub>3</sub>O<sub>4</sub> by vibrational spectroscopy is difficult, but not impossible. In Busca et al.'s *ex situ* IR investigation of air-calcined, then reduced Co<sub>3</sub>O<sub>4</sub>, the  $\nu$ (OH) IR window was open when the surface is reduced in 200 Torr of H<sub>2</sub> at 523 K.<sup>16</sup> Upon evacuation at 373 K,  $\nu$ (OH) on the reduced Co<sub>3</sub>O<sub>4</sub> were observed at 3680, 3658, 3630, and ~3580 cm<sup>-1</sup>, with an additional broad band centered at 3330 cm<sup>-1</sup>.<sup>16</sup> Assuming the particle did not reshape during the H<sub>2</sub> treatment, some OH modes remaining after evacuation should correspond to those observed on the equilibrated reduced (111) surface.

Vibrational spectroscopy after treatment only by water has also been performed by various groups on this surface. Petitto et al. performed ex situ HREELS on a single (111) terminated Co<sub>3</sub>O<sub>4</sub> crystal after annealing at 623 K under 6.6 x 10<sup>-5</sup> Pa H<sub>2</sub>O, cooling down to 110 K, and flashing to 400 K to remove molecular water, all in the absence of H<sub>2</sub> reduction.<sup>21</sup> They observed a single v(OH) peak located at 3650 cm<sup>-1</sup>.<sup>21</sup> Through *in situ* IRAS at 500 K under 3.0 ×  $10^{-7}$  mbar D<sub>2</sub>O, Schwarz et al. reached a similar result on Co<sup>2+</sup> and O<sup>2-</sup> terminated Co<sub>3</sub>O<sub>4</sub>(111)

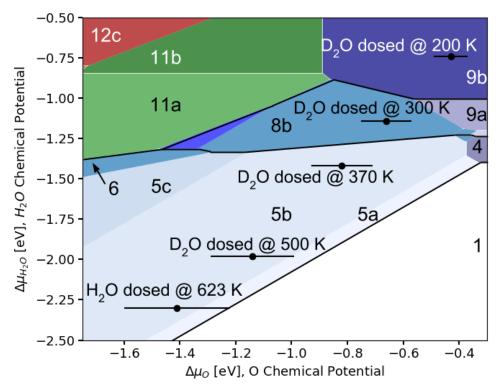
thin films.<sup>22</sup> They found a single v(OD) peak at 2650 cm<sup>-1</sup>. They also observed the surface's hydroxylation behavior changing drastically in the interval of 200 K – 500 K under the same pressure. Two broad hydroxyl peaks appear at 2715 cm<sup>-1</sup> and 2550 cm<sup>-1</sup> at 200 K, while an array of sharp, discrete peaks of variable intensity appear in the range 2730 cm<sup>-1</sup> to 2500 cm<sup>-1</sup> between 300 K and 400 K.

Table 3. Gas phase conditions reached in vibrational spectroscopy above 500 K

Reference	Conditions: $T[K]$ and $P_i[bar]$	Conditions: $\Delta \mu_i$ [eV]
Busca et al. 16	$T = 523 \text{ K}, P_{H2} = 0.27 \text{ bar}$	$\Delta\mu_{H2}=\text{-}0.70~\text{eV}$
Petitto et al. <sup>21</sup>	$T = 623 \text{ K}, P_{H2O} = 6.6 \text{ x } 10^{-10} \text{ bar}$	$\Delta\mu_{\rm H2O} = -2.30~eV$
	$P_{O2} = 4 \times 10^{-10} \sim 4 \times 10^{-16} \text{ bar}$	$\Delta \mu_{\rm O} = -1.22 \sim -1.60 \text{ eV}$
Schwarz et al. <sup>22,23</sup>	$T = 200 \text{ K}, P_{D2O} = 3 \text{ x } 10^{-11} \text{ bar}$	$\Delta\mu_{D2O} = -0.74 \text{ eV}$
	$P_{O2} = 10^{-10} \sim 10^{-16} \text{ bar}$	$\Delta \mu_{\rm O} = -0.37 \sim -0.49 \; {\rm eV}$
	$T = 300 \text{ K}, P_{D2O} = 3 \text{ x } 10^{-11} \text{ bar}$	$\Delta\mu_{\rm D2O} = -1.14~{\rm eV}$
	$P_{O2} = 10^{-10} \sim 10^{-16} \text{ bar}$	$\Delta \mu_{\rm O} = -0.57 \sim -0.75 \; {\rm eV}$
	$T = 370 \text{ K}, P_{D2O} = 3 \text{ x } 10^{-11} \text{ bar}$	$\Delta\mu_{\rm D2O} = -1.43~{\rm eV}$
	$P_{O2} = 10^{-10} \sim 10^{-16} \text{ bar}$	$\Delta \mu_{\rm O} = -0.71 \sim -0.93 \; eV$
	$T = 500 \text{ K}, P_{D2O} = 3 \text{ x } 10^{-11} \text{ bar}$	$\Delta\mu_{\rm D2O} = -1.98~{\rm eV}$
	$P_{O2} = 10^{-10} \sim 10^{-16} \text{ bar}$	$\Delta \mu_0 = -0.99 \sim -1.29 \text{ eV}$



**Figure 8.** Gas phase environmental conditions during vibrational characterization by Busca et al. represented as a dashed line on the H<sub>2</sub>O/H<sub>2</sub> surface stability diagram (Figure 5), as H<sub>2</sub>O chemical potential is not well-defined. The corresponding termination should contain 7.11 Co/nm<sup>2</sup> and 7.11 OH/nm<sup>2</sup> (structure 7a, brown) because an unrealistic high pressure of H<sub>2</sub>O must be applied to generate a mixed OH/H<sub>2</sub>O layer (structure 12c, red).



**Figure 9.** Gas phase environmental conditions during vibrational characterization under  $H_2O/D_2O$  pressure, overlaid with the  $H_2O/O$  surface stability diagram of Figure 3. These conditions fall under three main regions:  $(Co^T)_x$  with triaqua-complex-incorporated  $OH/H_2O$  chains (structure 9b), triaqua surface complex with additional OH (8b) for  $D_2O$  dosed at 200 K and 300 K, and  $(Co^T)_x$  with isolated OH (structure 5b) for  $D_2O$  exposure at 370 K and beyond. Gas phase conditions in characterization experiments are marked as dots with error bars for O chemical potential.

In our previous work, we identified three main structures observed on  $(Co^{Td})_x$  surface upon exposure to 3 x  $10^{-8}$  mbar  $D_2O$  at 300 K, without considering any surface reduction. The three motifs are: i) isolated hydroxyls, ii) surface  $Co(H_2O)_3$  complexes, and iii) extended hydroxyl networks.<sup>49</sup> Here, accounting for the effect of  $O_2$  in the atmosphere, we provide an  $O_2$ -equilibrated view of the hydroxylation behavior observed on  $Co_3O_4(111)$  films during and after the experiments described. Through interpolation of experimental thermochemistry data, the conditions produced by the  $H_2$ ,  $H_2O$ , and  $D_2O$  treatments can be converted to a set of water and hydrogen or oxygen chemical potentials (Table 3) and overlaid on the stability diagrams (Figure 8, 9). Since no *in situ* spectra are available under the conditions of Busca et al.'s and Petitto et

al.'s experiments, we assume that cooling and evacuation does not disrupt the surface structures generated during H<sub>2</sub> and H<sub>2</sub>O treatments. To approximate the oxygen content in the atmosphere observed in the water exposure experiments, we consider three magnitudes above and below the vacuum level in the experiments to calculate a possible range of chemical potentials for the conditions. For the H<sub>2</sub>-exposure experiment, we consider the range of water pressures that can be realistically reached to set a possible range of water chemical potentials.

In *ex situ* IR after  $H_2$  reduction and evacuation, multiple types of hydroxyl groups remain on the surface. <sup>16</sup> This result can be attributed to the large variation of exposed surfaces on a  $Co_3O_4$  particle. <sup>10</sup> Our calculations show that for the (111) surface, the  $(Co^{Oh}Co^{Td})_x$  termination (7.11  $Co/nm^2$ ) with 7.11  $OH/nm^2$  (structure 7a) is most likely exposed (Figure 8). <sup>16</sup> Two types of OH groups can be found on this surface: a bidentate OH group and a tetradentate OH group (see Table 2), which we found to exhibit  $3667\sim3668$  cm<sup>-1</sup> and  $3527\sim3528$  cm<sup>-1</sup>  $\nu(OH)$  modes respectively. These modes are close to two reported values (3680 cm<sup>-1</sup> and 3580 cm<sup>-1</sup>) out of many modes found by Busca et al. It is suggested that other modes are associated with other terminations of  $Co_3O_4$ . We did not consider the hydroxyl overlayer structure (structure 12c) in our assessment because  $P_{H2O}$  must be unrealistically high at 523 K for the fully hydroxylated structure to exist.

Since no  $H_2$  is present in the gas phase in the preparation of  $Co_3O_4(111)$  films and  $D_2O/H_2O$  exposure experiments of Schwarz et al., we considered the surfaces to be in equilibrium with an oxygen pressure 3 magnitudes about the vacuum level. With  $P_{O2}$  in consideration, we predict the following structures.  $D_2O$  exposure at 200 K, compared to our previous investigation, still corresponds to a hydrogen bonded hydroxyl network containing a triaqua complex (structure 9b). Next,  $D_2O$  exposure at 300 K corresponds to triaqua complexes

immediately surrounded by OD (structure 8b). The  $Co(D_2O)_3+OD$  structure shows modes in the range 2715~2723 cm<sup>-1</sup>, belonging to the v(OD) modes of the triaqua surface complex, and modes in the range 2547~2551 cm<sup>-1</sup> belonging to OD surrounding the complex. This result shows that the triaqua surface complex is robust and insensitive to surrounding OH. Finally, the conditions obtained by  $H_2O$  or  $D_2O$  adsorption experiments 370 K and above by Petitto et al. and Schwarz et al. fall under the sparsely hydroxylated region (structure 5b). Surface tridentate modes are located at 2684 cm<sup>-1</sup>/2681 cm<sup>-1</sup> for OD and 3637 cm<sup>-1</sup>/3633 cm<sup>-1</sup> for OH. These results match those obtained by Petitto et al. after purging physisorbed water (3645 cm<sup>-1</sup>) and the OD bands are 25-35 cm<sup>-1</sup> away from  $D_2O$  adsorption modes found by Schwarz et al. (2657 cm<sup>-1</sup>).

Previously, through a comparison between  $D_2^{18}O$  and  $D_2^{16}O$  exposure IRAS, it was shown that O of isolated surface OD generated through  $D_2O$  exposure above 500 K originate from the surface, since the  $\nu(OD)$  modes are not modified upon change of O isotope in water. Here, since we assume that no  $H_2$  is present in the atmosphere, the only way to generate only these hydroxyls is through vacancy re-filling by  $D_2O$ , and diffusion of the proton to surface O atoms. Upon  $D_2^{18}O$  dissociation at vacancy, one  $^{18}OD$  and one  $^{16}OD$  are generated, but since the surface is still rich in  $^{16}O$ ,  $\nu(^{16}OD)$  will be the predominant peaks in the IR spectrum. Another hypothesis is that eventual OD groups from water have low intensity in the spectrum.  $^{49}$ <sup> $^{49}$ <sup>

#### 3.3 Comparison among $U_{eff} = 3.5 \text{ eV}$ , 2.0 eV, and HSE06

The DFT+U approach is a very popular method used to describe strongly correlated systems as it makes a good compromise between accuracy and feasibility. The most popular adaptation is the rotationally-invariant method of Dudarev et al., where only one quadratic on-

site self-interaction penalty is added to the DFT energy to enforce integer occupancies in each spin orbital of the orbitals of interest.<sup>37</sup> Choosing an appropriate  $U_{eff}$  value to use is a very controversial topic in DFT modeling of strongly correlated systems.<sup>50</sup>

Approaches taken in literature can be divided to two categories: directly computing  $U_{eff}$  through linear response to correct DFT's unphysical energy response to addition and removal of fractional numbers of electrons and 2) fitting some empirical properties of the system. Studies using both approaches have been performed for the  $Co_3O_4$  bulk and surfaces. Chen et al. calculated  $U_{eff}$  of  $Co^{2+}$  and  $Co^{3+}$  in bulk  $Co_3O_4$  and used the weighted average, 5.9 eV, to study bulk properties; however, the approach predicted the ferromagnetic ground state being more favorable. The approach also does not predict well surface and bulk redox properties. The  $U_{eff}$  determined through this approach is also unique to the calculation method; a later study performed using PAW pseudopotentials in VASP, instead of Troullier-Martins pseudopotentials in Quantum Espresso found  $U_{eff}$  of  $Co^{2+}$  and  $Co^{3+}$  at bulk sites to be 4.3 eV and 5.02 eV.

Since the Co linear response  $U_{eff}$  does not predict well redox properties, authors interested in reactive properties tend to choose smaller values. Bajdich et al. and Chen et al. have noted that  $U_{eff} \sim 3$  eV is optimal to properly describe reactivity. Jiang et al. used  $U_{eff} = 2.0$  eV to match bulk band gap with experiment. Recently, Hu et al., using results obtained from RPA as benchmarks, proposed the usage of two different values of  $U_{eff}$ , applied to the Co and surface Co atoms respectively, to examine the reactivity of oxygen-rich  $Co_3O_4(110)$ . To ensure that our choice of  $U_{eff} = 3.5$  eV is appropriate for  $Co_3O_4(111)$ , we compared energies of water adsorption and O vacancy formation generated at this  $U_{eff}$  to those obtained using  $U_{eff} = 2.0$  eV and HSE06+dDsC. We chose the value of  $U_{eff} = 2.0$  eV in addition to our initial choice  $U_{eff} = 3.5$  eV, as it is a very popular value used in many DFT studies on the reactivity of  $Co_3O_4$  surfaces. Table

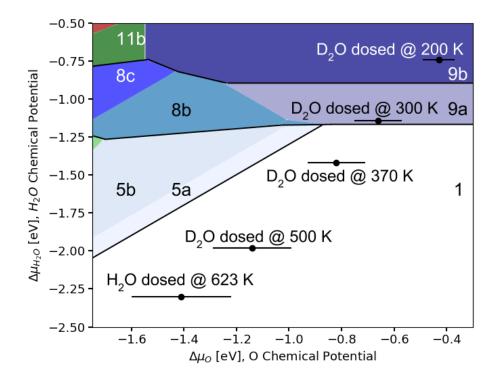
4 compares energy of vacancy formation, for  $O_{vac}=3.56~\text{nm}^{-2}$  and energy of adsorption of water at  $3.56~\text{H}_2\text{O}/\text{nm}^2$  and  $7.11~\text{OH}/\text{nm}^2$ , reached using the three methods on the  $(\text{Co}^{\text{Td}})_x$  termination of  $\text{Co}_3\text{O}_4$ .

**Table 4.** Comparison of water adsorption energy and energy of vacancy formation calculated using  $U_{eff} = 3.5 \text{ eV}$ , 2.0 eV, and HSE06+dDsC (water adsorption data adapted from  $^{49}$ )

Description	Energy [eV], U <sub>eff</sub> =3.5 eV	Energy [eV], $U_{\it eff}$ =2.0 eV	Energy [eV], HSE06+dDsC
Vacancy formation, $(Co^T)_x$ with 3.56 $O_{vac}/nm^2$ , energy of $O_{vac}$ formation, $E_{form}$ per $O_{vac}$	2.07	3.01	2.31
Water adsorption, $(Co^T)_x$ with 7.11 OH/nm <sup>2</sup> and 3.56 H <sub>2</sub> O/nm <sup>2</sup> in the "pair" configuration (see SI, Figure S13, top), $E_{ads}$ per H <sub>2</sub> O molecule	-1.02	-0.91	-1.08
Water adsorption, $(Co^T)_x$ with 7.11 OH/nm <sup>2</sup> and 3.56 H <sub>2</sub> O/nm <sup>2</sup> in the "hexagonal" configuration (see SI, Figure S13, bottom), E <sub>ads</sub> per H <sub>2</sub> O molecule	-1.00	-0.96	-1.10

As seen in the table above, the energy of vacancy formation decreases as  $U_{\rm eff}$  increases. Upon generation of a single  $O_{\rm vac}$  on this termination, electrons are transferred to the Co atoms immediately surrounding the vacancy. Increasing  $U_{\rm eff}$  increases the penalty placed on delocalization while facilitating electron localization. Selcuk and Selloni also observed that O vacancy formation energy on the (110) surface, in the case where electrons are localized, heavily depends on  $U_{\rm eff}$ . On the other hand, the adsorption of water is an electron-neutral process, which explains why the change in  $E_{\rm ads}$  is not very significant. These adsorption results obtained from  $U_{\rm eff} = 3.5$  eV do not strongly differ from those at  $U_{\rm eff} = 2$  eV but agree better with results from HSE06.

For further confirmation, we recalculated using  $U_{\text{eff}} = 2.0$  eV the surface energy of structures shown in Figure 9 and Table 4 and constructed an analogous surface stability diagram combined shown on Figure 10.



**Figure 10.** Surface structure as a function of oxygen and water chemical potential calculated using  $U_{eff} = 2.0$  eV. Gas phase conditions observed in vibrational characterization are marked as dots with error bars for O chemical potential. No hydroxylation is predicted to exist in D<sub>2</sub>O/H<sub>2</sub>O exposure conditions at 370 K and above when using  $U_{eff} = 2.0$  eV.

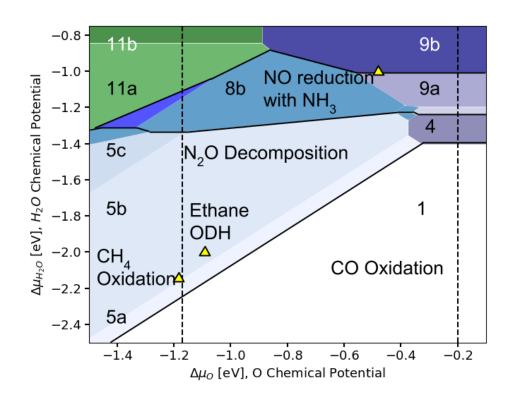
The lower  $U_{eff}$  value heavily decreases the energy of adsorption of hydrogen and weakens slightly water adsorption. Interestingly, the stability diagram reached from  $U_{eff} = 2.0 \text{ eV}$  does not predict hydroxylation through  $D_2O$  exposure or annealing in  $H_2O$ , in contrast with experiments. The value  $U_{eff} = 3.5 \text{ eV}$  hence produces results that match much better the spectroscopic experiments. Like vacancy formation, H adsorption adds electrons to the surface and becomes expectedly more favorable with increasing  $U_{eff}$ . Altogether, these test results validate the theoretical method used in the present paper.

# 3.4 Predicted surface structure under catalytic reaction conditions

Co<sub>3</sub>O<sub>3</sub>'s applications in thermal catalysis can be divided into three categories. First, Co<sub>3</sub>O<sub>4</sub> has been known as an effective CO and hydrocarbon oxidation catalyst. <sup>1,54–56</sup> Xie et al. found that CO oxidation can occur over Co<sub>3</sub>O<sub>4</sub> nanorods at a temperature as low as -77 °C, but is easily poisoned by water. 54,55 The activity appears to depend on the exposed surfaces and catalyst shape.<sup>57</sup> Second, Co<sub>3</sub>O<sub>4</sub> can also be used as a catalyst or support for NO<sub>x</sub> reduction.<sup>1,3,4,58,59</sup> N<sub>2</sub>O decomposition to N2 and O2 over Co3O4 was found to occur at 873 K, and selective catalytic reduction (SCR) reaction of NO with NH<sub>3</sub> and O<sub>2</sub> was found to occur over Co<sub>3</sub>O<sub>4</sub> at 423 K.<sup>3,58</sup> NO reduction with H<sub>2</sub> can also be performed over single noble metal (NM) atoms anchored on  $\text{Co}_3\text{O}_4$  to produce  $\text{N}_2$  or  $\text{N}_2\text{O}.^{4,60}$  Third, one promising application of the (111) surface is the oxidative dehydrogenation (ODH) of alkanes.<sup>5–7</sup> This reaction was explored by Tyo et al., Liu et al, and Fung et al.<sup>5-7</sup> It was found that the alkane activation occurs over one isolated surface O and one O adjacent to surface Co (see Section 2.2), and that substitution of surface Co-T<sub>d</sub> with other metals lowers the activation barrier of alkanes by increasing surface reducibility.<sup>61</sup> Liu et al. showed that Co<sub>3</sub>O<sub>4</sub> particles primarily exposing the (111) facet have higher conversion and selectivity for ethane ODH than particles that expose higher miller index surfaces such as (311).<sup>62</sup> Overall, reactions using Co<sub>3</sub>O<sub>4</sub> as a catalyst or support and containing O<sub>2</sub> in the atmosphere are performed below 900 K, and reactions using H<sub>2</sub>/H<sub>2</sub>O as reactant/product are performed below 500-600 K.

**Table 5.** Gas phase conditions observed in catalytic reactions

Reaction	Conditions: $T[K]$ and $P_i[bar]$	Conditions: $\Delta \mu_i$ [eV]
$2CO + O_2 \rightarrow 2CO_2^{54}$	$T = 196 \text{ K}, P_{O2} = 0.02 \text{ bar}$	$\Delta\mu_{O}$ = -0.20 eV
$2N_2O \rightarrow 2N_2 + O_2^{58}$	$T = 873 \text{ K}, P_{O2} = 2 \text{ x } 10^{-3} \text{ bar}$	$\Delta\mu_O = -1.17 \text{ eV}$
$2C_2H_6 + O_2 \rightarrow 2C_2H_4 + 2H_2O$	$T = 873 \text{ K}, P_{O2} = 0.018 \text{ bar}$	$\Delta\mu_{\rm O}$ = -1.09 eV
$2C_2H_6 + 7O_2 \rightarrow 4CO_2 + 6H_2O^7$	$P_{H2O} = 0.033 \text{ bar}$	$\Delta\mu_{\rm H2O} = -2.00~\rm eV$
$CH_4 + 2O_2 \rightarrow CO_2 + 2H_2O$	$T = 873 \text{ K}, P_{O2} = 0.0017 \text{ bar}$	$\Delta\mu_{\rm O}$ = -1.18 eV
$2CH_4 + 3O_2 \rightarrow 2CO + 4H_2O^{63}$	$P_{H2O} = 0.0050 \text{ bar}$	$\Delta\mu_{\rm H2O} = -2.14~eV$
$2NO + 2H_2 \rightarrow N_2 + 2H_2O$	$T = 423 \text{ K}, P_{H2} = 0.0095 \text{ bar}$	$\Delta \mu_{\rm H2} = -0.66 \; \rm eV$
$2NO + H_2 \rightarrow N_2O + H_2O^4$	$P_{H2O} = 0.040 \text{ bar}$	$\Delta\mu_{\rm H2O} = -0.85~eV$
	$T = 523 \text{ K}, P_{H2} = 0.00049 \text{ bar}$	$\Delta\mu_{H2} = -0.98 \text{ eV}$
	$P_{H2O} = 0.0483 \text{ bar}$	$\Delta\mu_{\rm H2O} = -1.08~eV$
$4NO + 4NH_3 + O_2 \rightarrow 4N_2 + 6H_2O^3$	$T = 423 \text{ K}, P_{O2} = 0.020 \text{ bar}$	$\Delta\mu_{\rm O}$ = -0.48 eV
	$P_{H2O} = 6 \times 10^{-4} \text{ bar}$	$\Delta\mu_{\rm H2O} = -1.00~\rm eV$

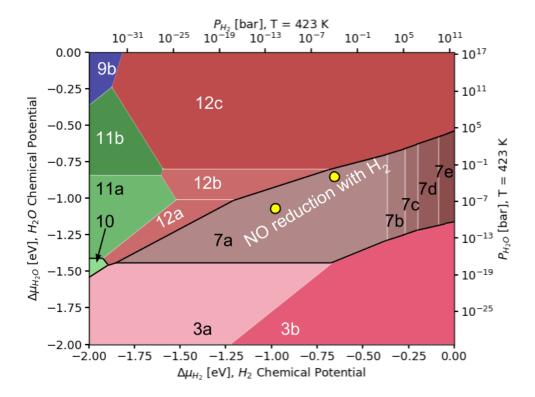


**Figure 11.** Surface stability diagram as a function of  $O_2$  and  $H_2O$  chemical potential. Conditions in low temperature CO oxidation,  $N_2O$  decomposition, ethane ODH,  $CH_4$  oxidation, and NO reduction with  $NH_3$  are marked on the diagram as yellow triangles if  $P_{H2O}$  is well-defined and dashed lines otherwise.

Reactive conditions considered in this work are listed on Table 5. H<sub>2</sub>, O<sub>2</sub>, and H<sub>2</sub>O are assumed to be ideal gases, and their partial pressures are estimated using the reported reaction conversion and selectivity, assuming the reactions take place isobarically at 1 bar. Co<sub>3</sub>O<sub>4</sub> is most commonly used for oxidative reactions without H<sub>2</sub>; therefore, in most gas phase reactions, the Co<sub>3</sub>O<sub>4</sub> surfaces are in equilibrium with H<sub>2</sub>O and O<sub>2</sub>. Let us start this discussion with low temperature CO oxidation. Since the water chemical potential is not clearly defined for that reaction, conditions are indicated as a dashed vertical line on the right side of Figure 11. However, the nature of the Co<sub>3</sub>O<sub>4</sub>(111) surface can change quite drastically when water is present. In these conditions, from our calculations (Figure 11), hydroxyls will populate the surface when water chemical potential is greater than -1.45 eV, which, at 196 K, corresponds roughly to a pressure higher than 10<sup>-30</sup> bar; in other words, any water present will contaminate the surface. As the reaction temperature increases at the same O<sub>2</sub> pressure, the surface becomes less susceptible to hydroxylation. CO TPD studies have also shown that lattice oxygen does not appear to be reactive; therefore, it is crucial for surface Co<sup>2+</sup> to be exposed during CO oxidation to adsorb reactants.<sup>64</sup> At 196 K, under the O<sub>2</sub> pressure applied in the CO oxidation reaction, the surface already becomes covered with an triaqua-complex-containing H<sub>2</sub>O/OH overlayer (structure 9b) when P<sub>H2O</sub> is in the order of 10<sup>-15</sup> bar. Both experimentally determined and calculated energy of CO adsorption falls far above -1.1 eV, indicating that any water impurities will completely block the reaction, in agreement with experiments.<sup>64,65</sup> Conditions for N<sub>2</sub>O decomposition also correspond to a line on Figure 11, but at a lower O chemical potential, mainly from the much higher temperature. The surface structure during N<sub>2</sub>O decomposition is also dependent on the amount of H<sub>2</sub>O impurities in the gas phase; though, the only accessible regions, i.e.  $P_{H2O} > 10^{-12}$  bar and above, are the bare  $(Co^{Td})_x$  and  $(Co^{Td})_x$  terminations with  $\leq 1.78$ 

OH/nm<sup>2</sup> (5a, 5b). The transition from bare to hydroxylated can occur very easily, but no hydroxyl cluster or overlayer can form.

Reactions that generate H<sub>2</sub>O, as a product or side product, also carry the threat that the Co<sub>3</sub>O<sub>4</sub> surface will hydroxylate. In the conditions associated to the ODH of ethane, the surface is found to be slightly hydroxylated, with 1.78 OH/nm<sup>2</sup> (structure 5b). CH<sub>4</sub> oxidation would also fall under the same region on the phase diagram. Both reactions occur at high temperatures, so it is reasonable to expect the surfaces to have low degrees of hydroxylation. Finally, for NO reduction with NH<sub>3</sub>, the conditions for this reaction fall at the boundary between the zone of triaqua-complex-containing OH/H<sub>2</sub>O chains (9b) and triaqua complex/OH/H<sub>2</sub>O chain existing separately (8b). In other words, surface Co is unlikely to be exposed during this reaction.



**Figure 12.** Surface stability diagram as a function of H<sub>2</sub> and H<sub>2</sub>O chemical potential and hydrogen/water pressure at 423 K. Conditions in NO reduction with H<sub>2</sub> found in the literature are marked on the diagram as yellow circles.

On the other hand, the environment during NO reduction with H<sub>2</sub> equilibrates the surface with H<sub>2</sub>O and H<sub>2</sub> (Figure 12). In this reaction, the Co<sub>3</sub>O<sub>4</sub> plays the role of support for noble metal atoms (Pd, Pt, and Rh; the conditions here refer to NO reduction over Pd/Co<sub>3</sub>O<sub>4</sub>), and it was observed through AP-XPS that when the temperature is moderate enough, surface transition to CoO is inhibited.<sup>4,60</sup> In these conditions, the (Co<sup>Oh</sup>Co<sup>Td</sup>)<sub>x</sub> surface with 7.11 OH/nm<sup>2</sup> (structure 7a) would be exposed. The high coverage of hydroxyl groups present might, however, play a major role in modifying the catalysts support interaction and the nature of the charge transfer between the support and the metallic nanoparticle.<sup>66</sup>

#### 4. Conclusion

In this study, we investigated the surface structure of three terminations of  $\text{Co}_3\text{O}_4(111)$  in equilibrium with a wide range of gas environments and mapped out detailed surface stability diagrams with respect to vacancy formation and hydroxylation of the surface either from hydrogen or water adsorption. We found that all three terminations are accessible, but the Copoor ( $(\text{Co}^{\text{Td}})_x$ , 3.56  $\text{Co/nm}^2$ ) and Co-rich ( $(\text{Co}^{\text{Oh}}\text{Co}^{\text{Td}})_x$ , 7.11  $\text{Co/nm}^2$ ) terminations are the most prevalent. When equilibrated with oxygen and water, we found that the Co-poor termination exhibits three stages of hydroxylation – formation of isolated OH at low  $\Delta\mu_{\text{H2O}}$ , transition to isolated OH/H<sub>2</sub>O clusters when  $\Delta\mu_{\text{H2O}}$  > -1.3 eV, and finally to a fully hydrated state when  $\Delta\mu_{\text{H2O}}$  >-1.0 eV.

We found that if the surface is in equilibrium with oxygen, the Co-poor/O-rich surface  $(Co^{Td})_x$  is predominantly exposed, and it is very difficult to produce the stoichiometric  $(Co^{Oh}_{0.5}Co^{Td})_x$  or the Co-enriched  $(Co^{Oh}Co^{Td})_x$  termination. Since we found it to be very easy to populate the  $(Co^{Td})_x$  surface with hydroxyls through  $H_2$  adsorption or  $H_2O$  dissociation over a

vacancy, we predict that hydroxyls will appear during oxidative reactions that produce water or even when water is present as an impurity at low partial pressure. We compared the calculations results concerning hydroxylation of  $(Co^{Td})_x$  to those obtained from vibrational spectroscopy on  $Co_3O_4$  single crystals and  $Co_3O_4(111)$  films and conclude that at low and medium temperature triaqua surface complexes surrounded by OH/H<sub>2</sub>O groups are observed, while only isolated OH are present at higher temperature. In a similar way, unless the environment contains very little water, the O-poor termination  $(Co^{Oh}Co^{Td})_x$  is predicted also be hydroxylated.

We found that stabilization of the stoichiometric termination  $(Co^{Oh}_{0.5}Co^{Td})_x$  and Co-rich termination  $(Co^{Oh}Co^{Td})_x$  is only possible under vacuum at high temperature and in hydrogen and water at low temperatures. The  $H_2/H_2O$  conditions are the most important for catalytic applications. We found that reduction of this Co-rich termination through H adsorption and  $O_{vac}$  formation is easy due to the large relaxation observed at 7.11  $OH/nm^2$  or 3.56  $O_{vac}/nm^2$ . At 423 K and below and under realistic  $H_2/H_2O$  pressure, the  $(Co^{Oh}Co^{Td})_x$  termination with 7.11  $OH/nm^2$  is highly stable. The OH groups on that termination exhibit vibrational frequencies that fall under the range observed in *ex situ* IR experiments. We believe that this termination should be considered when modeling reactions over  $Co_3O_4$  in low temperature reducing environments.

We validated our results through changing the  $U_{eff}$  value and comparing with results reached through hybrid functionals. We found good agreement with both macroscopic experimental observations and HSE06+dDsC calculations using  $U_{eff}=3.5$  eV. Hence, the various terminations of  $Co_3O_4$  should be hydroxylated under a large range of conditions. It is therefore of high importance to include a correct coverage in hydroxyl groups for a realistic modelling of the  $Co_3O_4$  surface in calculations; the bare surface model usually used in reaction pathway calculations might not be adequate. The effect of hydroxyls in the modification of

reaction mechanisms and their interactions with supported metal atoms/clusters must also be considered. Further studies will focus on the interaction among various surface features discussed in this work, and the implications of these structures on the surface's reactivity and properties of the supported active sites.

## 5. Acknowledgements

The calculations in this work used computational and storage services associated with the Hoffman2 Shared Cluster provided by UCLA's Institute for Digital Research and Education Group (IDRE). This work used the Extreme Science and Engineering Discovery Environment (XSEDE), which is supported by National Science Foundation grant number ACI-1548562. Specifically, it used the Bridges system, which is supported by NSF award number ACI-1445606, at the Pittsburgh Supercomputing Center (PSC).<sup>67,68</sup> Startup funding from the Department of Chemical and Biomolecular Engineering of UCLA is acknowledged.

## **6.** Supplementary materials:

- 1) Surface energy calculation details, 1-dimensional surface stability diagrams, energies of adsorption, energies of vacancy formation, top/side views of geometries, and other details available in supporting information (1)
- 2) Geometries of all structures referenced in this work are available in supporting information(2)

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# TOC graphic:

