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### **Title**

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By Shape Memory Alloys

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DETERMINATION OF THE RECOVERY STRESSES  
DEVELOPED BY SHAPE MEMORY  
ALLOYS

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The mechanical shape memory effect is exhibited by several alloys all of which undergo thermoelastic martensitic transformation [1]. The memory effect is demonstrated by inducing a macroscopic strain within a certain range in a specimen while in martensitic - state and then heating to revert martensite to the high temperature phase (htp). During reversion, the specimen tends to recover its original unstrained shape. Strain recovery starts at the  $A_s$  temperature and finishes at the  $A_f$  temperature\*.

One of the most interesting applications of shape memory alloys is in energy conversion devices that utilize the memory effect in converting low-grade heat sources into useful mechanical work [2 - 4]. In evaluating the performance of a particular alloy in such devices it becomes essential to determine the levels of recovery stresses developed during the martensite  $\rightarrow$  htp transformation upon heating. There have been a few investigations dealing with experimental measurements of the recovery stresses developed by Ni - Ti alloys near the equi-atomic composition [5, 6]. However, no attempt has been made as yet to predict these stresses. In the present paper, a simple relation-

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\*  $A_s$  = temperature at which the martensite  $\rightarrow$  htp transformation starts during heating.

$A_f$  = temperature at which this transformation finishes during heating.

\*\* A recovery stress is defined as the external stress required to balance the internal stresses that tend to revert martensite to the htp.



ship has been developed between the recovery stresses, temperature and strain

For the discussions which follow, it is instructive to identify a reference point from which the macroscopic strain is measured i.e.; a zero - point strain. This can conveniently be taken to be the point at which the material consists entirely of the htp.

Consider a polycrystalline shape memory alloy that has been almost entirely transformed to martensite by cooling it to a temperature at or below its respective martensite - finish temperature  $M_f$ . In the absence of an external and appreciable internal stresses, all the martensite variants that are crystallographically equivalent would form with equal probability. Therefore, no net macroscopic strain as measured from the zero - point would be expected and hence no recovery stresses would develop during the reverse transformation to the htp. In order to realize a recovery stress, the material must acquire a macroscopic strain relative to the zero - point. In terms of the microstructural features, this corresponds to a certain degree of directionality in the distribution of the martensite variants through-out the microstructure. It has been shown experimentally that any strain within the recoverable range is accommodated by re-orientation of a volume fraction of the thermally induced martensite [7 - 10]. This occurs by preferential growth of the most favorably oriented martensite variants. It is then expected that the recovery stress,  $\sigma_r$ , would be proportional to the volume fraction of re-oriented martensite ( $V_M$ ). This determines the functional dependence of  $\sigma_r$  on strain  $\epsilon$ .

Upon heating to within the  $A_s - A_f$  temperature range, a certain volume fraction of martensite tends to revert to the htp. The driving force for such reversion arises from the chemical free energy difference between the two respective structures, the elastic strain energy stored in the lattice during the transformation to martensite and energy associated with reversible defects, e.g., twins. As the temperature increases within the range  $A_s - A_f$ , the driving force increases and in turn the volume fraction of martensite that tends to revert to the htp ( $V_{htp}$ ).

It follows that  $\sigma_r$  increases as  $V_{\text{htp}}$  increases with temperature. This determines the functional dependence of  $\sigma_r$  on temperature within the range  $A_s - A_f$ .

It can be concluded from the above discussion that:

$$\sigma_r \propto V_{\text{htp}}(T) \cdot V_M(\epsilon) \dots\dots\dots(1)$$

$\sigma_r$  would be limited by the minimum stress required to initiate irreversible plastic deformation in the martensite structure and/or the htp. The yield stresses of the htp are, however, expected to be lower than those of the martensite due to: (i) the lower symmetry of the martensite structure and (ii) the finer grain size of the martensite phase. Therefore, the proportionality constant in equation (1) can be taken to be the yield stress of the htp,  $\sigma_Y^{\text{htp}}$ , that is consistent with the state of temperature and strain and hence,

$$\sigma_r = \sigma_Y^{\text{htp}} \cdot V_{\text{htp}}(T) \cdot V_M(\epsilon) \dots\dots\dots(2)$$

From dilatometric measurements on a Ti - 50.4 at pct. Ni alloy it appeared that  $V_{\text{htp}}$  increases exponentially with temperature within the range  $A_s - A_f$ .  $V_{\text{htp}}(T)$  could be expressed as:

$$V_{\text{htp}}(T) = 1 - \exp^{-K \frac{(T - A_s)}{A_f - T}} \dots\dots\dots(3)$$

where K is a constant which by analogy to chemical reactions, can be considered as a reaction rate constant. It is to be noted that the type of equation (3) has been found empirically to describe the kinetics of a wide variety of reactions in metals (11).

The probability that re-orientation of an existing martensite occurs under an external stress is expected to be proportional to the number N of martensite variants that are crystallographically equivalent. If as has been assumed all the variants exist with equal probability, the interfaces contained in the microstructure (twin - boundaries in the case of Ni - Ti [7] assume all the possible orientations with respect to the applied stress. Therefore some martensite

variants would be unaffected and others would be affected to different extents depending on their relative orientation. It then follows that  $V_M$  would never achieve unity for any finite strain. A relation which gives  $V_M$  as a function of  $\epsilon$  must then satisfy the following conditions:

- $V_M = 0$  for  $\epsilon = 0$ , and
- $V_M$  increases with  $\epsilon$  (as measured from the zero - point) but never achieves unity.

These conditions can be satisfied by an exponential function of the form:

$$V_M = 1 - \exp^{-C\epsilon} \dots\dots\dots(4)$$

where C is a constant that can be taken to be the number N of crystallographically equivalent martensite variants.

$\sigma_r (T, \epsilon)$  can then be expressed as:

$$\sigma_r (T, \epsilon) = \sigma_Y^{htp} (1 - \exp^{-K(\frac{T - A_S}{A_f - T})}) (1 - \exp^{-N\epsilon}) \dots\dots\dots(5)$$

The characteristic features of the recovery stress as depicted from equation (5) are:

- (i) For a given strain  $\epsilon$ ,  $\sigma_r$  is sensitive to temperature only within the range  $A_S - A_f$ .
- (ii)  $\sigma_r$  achieves its maximum value of:  
 $\sigma_r \text{ max.} = \sigma_Y^{htp} (1 - \exp^{-N\epsilon}) \dots\dots\dots(6)$   
 at the  $A_f$  temperature where  $V_{htp} = 1$ .
- (iii) For temperatures higher than the  $A_f$ ,  $\sigma_r$  becomes sensitive only to strain.

These features are in agreement with the experimental observations [5, 6]. Experimental work is now in progress in order to compare the measured and predicted values of the recovery stresses developed by a Ni - Ti alloy.

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