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Authors

Tracy, Craig A.
Widom, Harold

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Formulas for ASEP with Two-Sided Bernoulli Initial Condition

Craig A. Tracy · Harold Widom

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Abstract For the asymmetric simple exclusion process η_t on the integer lattice with two-sided Bernoulli initial condition, we derive exact formulas for the following quantities: (1) $\mathbb{P}(\eta_t(x) = 1)$, the probability that site x is occupied at time t ; (2) the correlation function $\mathbb{P}(\eta_t(x) = 1, \eta_0(0) = 1)$; (3) the distribution function for Q_t , the total flux across 0 at time t , and its exponential generating function.

Keywords Exclusion process · Bernoulli initial condition

1 Introduction

Since its introduction by Spitzer [11], the asymmetric simple exclusion process (ASEP) has attracted considerable attention in both the mathematics and physics literature due to the fact it is one of the simplest lattice models describing transport far from equilibrium [2, 7, 8]. Recall that ASEP on the integer lattice \mathbb{Z} is a continuous time Markov process η_t , where $\eta_t(x) = 1$ if $x \in \mathbb{Z}$ is occupied at time t , and $\eta_t(x) = 0$ if x is vacant at time t . Particles move on \mathbb{Z} according to two rules: (1) A particle at x waits an exponential time with parameter one (the particle’s “alarm clock”), and then chooses y with probability $p(x, y)$; (2) If y is vacant at that time it moves to y , while if y is occupied it remains at x and restarts its clock. The process is called “simple” since the jumps are restricted to nearest neighbors: $p(x, x + 1) = p$ and $p(x, x - 1) = q = 1 - p$. All clocks are independent of each other and one need not worry that two clocks ring simultaneously since this event has probability zero. A rigorous construction of this process can be found in Liggett [7]. We note that the process is called the totally asymmetric simple exclusion process (TASEP) if jumps occur

C.A. Tracy
Department of Mathematics, University of California, Davis, CA 95616, USA
e-mail: tracy@math.ucdavis.edu

H. Widom (✉)
Department of Mathematics, University of California, Santa Cruz, CA 95064, USA
e-mail: widom@ucsc.edu

only to the left ($q = 1$) or only to the right ($p = 1$); and if $p = q = 1/2$, the process is called the symmetric simple exclusion process (SSEP) [7, 8].

In previous work [14] the authors considered ASEP on the integer lattice \mathbb{Z} with *step Bernoulli initial condition*: initially a site in \mathbb{Z}^+ (the positive integers) is occupied with probability ρ while sites in \mathbb{Z}^- (the nonpositive integers) are unoccupied. Formulas were obtained for $\mathbb{P}(x_m(t) \leq x)$, the probability that at time t the m th particle from the left is at a site $\leq x$, with consequent asymptotic results for this probability. In this paper we consider the more general situation where initially a site in \mathbb{Z}^+ is occupied with probability ρ_+ while a site in \mathbb{Z}^- is occupied with probability ρ_- . This initial condition is called the *two-sided Bernoulli initial condition*. Now it makes no sense to speak of the m th particle from the left. Instead we find formulas for $\mathbb{P}(\eta_t(x) = 1)$, the probability that site x is occupied at time t , and the correlation function $\mathbb{P}(\eta_t(x) = 1, \eta_0(0) = 1)$. We also derive formulas for the generating function $\langle e^{\lambda Q_t} \rangle$, where Q_t is the total flux across 0 at time t .

For the special case of TASEP with two-sided Bernoulli initial condition much is already known including various limit theorems [1, 10]. (Also see [4] for the stationary case $\rho_- = \rho_+$.) A simplifying feature of TASEP is that it is a determinantal process [6] thus permitting the application of random matrix theory techniques. In [14] it was shown that the limit theorems in [1] for TASEP extend to ASEP for the special case of step Bernoulli initial condition ($\rho_- = 0, \rho_+ > 0$); and thus, these results [14] establish a KPZ universality theorem for ASEP. Whether the TASEP limit theorems of [1] for two-sided Bernoulli initial condition can be extended to ASEP from the formulas derived in this paper remains to be seen.

We shall assume throughout that $p, q \neq 0$. As in [14] we begin with a formula derived in [12] for $\mathbb{P}_Y(x_m(t) \leq x)$, the probability distribution for the m th particle from the left given an initial finite configuration Y ; then we average over all Y in a finite subset of \mathbb{Z} ; and then we let the subset of \mathbb{Z} become unbounded. The formula for $\mathbb{P}_Y(x_m(t) \leq x)$ involved multiple integrals over large contours, and this allowed Y to become unbounded on the right in [14]. Now, if Y is to become unbounded on left and right, we must sum over both large and small contours, so the formula for $\mathbb{P}_Y(x_m(t) \leq x)$ must first be recast. This is done in Sect. 2. In Sect. 3 we deduce the formula for $\mathbb{P}_Y(\eta_t(x) = 1)$, the probability that site x is occupied at time t given the initial configuration Y . In Sect. 4 we average over Y and take the limit to obtain the formula for $\mathbb{P}(\eta_t(x) = 1)$ for two-sided Bernoulli initial configuration on \mathbb{Z} . The formula for the correlation function requires only a small change. Finally, in Sect. 5 we deduce formulas for probability distribution of the total flux across the origin at time t and for its exponential generating function. For the special case of SSEP we connect this last result with a recent paper of Derrida and Gerschenfeld [3].

2 Finite Configuration, Large and Small Contours

We begin with Theorem 5.2 of [12] which gives a formula for $\mathbb{P}_Y(x_m(t) \leq x)$, the probability in ASEP when the initial configuration is a (deterministic) finite set Y . To state it we introduce some notation, which here will be slightly different.

First, we set $\tau = p/q$ and recall that the τ -binomial coefficient $\left[\begin{smallmatrix} N \\ n \end{smallmatrix} \right]_\tau$ is defined as

$$\frac{(1 - \tau^N)(1 - \tau^{N-1}) \cdots (1 - \tau^{N-n+1})}{(1 - \tau)(1 - \tau^2) \cdots (1 - \tau^n)}$$

when n is a positive integer, 1 when $n = 0$ (empty products are always defined to be 1), and 0 when n is a negative integer.

We define

$$\varepsilon(\xi) = p \xi^{-1} + q \xi - 1, \quad f(\xi_i, \xi_j) = \frac{\xi_j - \xi_i}{p + q \xi_i \xi_j - \xi_i},$$

and then

$$I(x, k, \xi) = I(x, k, \xi_1, \dots, \xi_k) = \prod_{i < j} f(\xi_i, \xi_j) \prod_i \frac{\xi_i^x e^{\varepsilon(\xi_i)t}}{1 - \xi_i},$$

$$\Delta I(x, k, \xi) = I(x, k, \xi) - I(x-1, k, \xi) = (1 - \prod_i \xi_i^{-1}) I(x, k, \xi).$$

All indices in the products run from 1 to k . Notice that I and ΔI depend on t , although it is not displayed in the notations.

Finally, given two sets of integers U and V we define

$$\sigma(U, V) = \#\{(u, v) : u \in U, v \in V, \text{ and } u \geq v\}.$$

(If $V = [1, N]$ ¹ and $U \subset [1, N]$ then $\sigma(U, V)$ equals the sum of the elements of U . Hence the notation.)

Theorem 5.2 of [12] states that when $q \neq 0$,

$$\mathbb{P}_Y(x_m(t) = x) = \sum_{k=1}^{|Y|} c_{m,k} \sum_{\substack{S \subset Y \\ |S|=k}} \tau^{\sigma(S, Y)} \int_{C_R} \dots \int_{C_R} \Delta I(x, k, \xi) \prod_i \xi_i^{-s_i} d\xi_1 \dots d\xi_k, \quad (1)$$

where

$$S = \{s_1, \dots, s_k\},$$

$$c_{m,k} = (-1)^m q^{k(k-1)/2} \tau^{m(m-1)/2-km} \begin{bmatrix} k-1 \\ m-1 \end{bmatrix}_{\tau}, \quad (2)$$

and C_R is the circle with center 0 and radius R , which is assumed so large that the denominators $p + q \xi_i \xi_j - \xi_i$ are nonzero on and outside the contours. The inner sum is taken over all subsets S of Y with $|S|=k$. Observe that $c_{m,k} = 0$ when $k < m$.

If we sum (1) on x from $-\infty$ to x , as we may do since $R > 1$, we obtain

$$\mathbb{P}_Y(x_m(t) \leq x) = \sum_{k \geq 1} c_{m,k} \sum_{\substack{S \subset Y \\ |S|=k}} \tau^{\sigma(S, Y)} \int_{C_R} \dots \int_{C_R} I(x, k, \xi) \prod_i \xi_i^{-s_i} d\xi_1 \dots d\xi_k. \quad (3)$$

Clearly (3) implies (1) immediately.

Because of the appearance of the factors $\prod_i \xi_i^{-s_i}$ in the integrands and because R is large we expect to be able to allow Y to be unbounded on the right but not on the left. So we want integrals over both large and small contours. Lemma 5.1 of [12] gives a formula which expresses an integral over a large contour as a sum of integrals over small contours, and we shall be able to apply it here.

¹This is short for the set $\{1, \dots, N\}$. We use this notation consistently.

We shall assume throughout that $Y = Y_- \cup Y_+$ where Y_- is to the left of Y_+ . Then each $S \subset Y$ is a union $S_- \cup S_+$ with $S_{\pm} \subset Y_{\pm}$. Set $k_{\pm} = |S_{\pm}|$, so $k = k_- + k_+$. Instead of using indices $i = 1, \dots, k$ for each S with $|S| = k$, we shall use k_- negative indices $-k_-, \dots, -1$ and k_+ positive indices $1, \dots, k_+$, with $s_i \in S_-$ if $i < 0$ and $s_i \in S_+$ if $i > 0$. In this notation (3) becomes

$$\begin{aligned} \mathbb{P}_Y(x_m(t) \leq x) &= \sum_{k_{\pm} \geq 0, k \geq 1} c_{m,k} \sum_{\substack{S_{\pm} \subset Y_{\pm} \\ |S_{\pm}| = k_{\pm}}} \tau^{\sigma(S_-, Y_-) + \sigma(S_+, Y_+)} \\ &\quad \times \int_{\mathcal{C}_R} \cdots \int_{\mathcal{C}_R} I(x, k, \xi) \prod_i \xi_i^{-s_i} \prod_i d\xi_i, \end{aligned} \quad (4)$$

where $k = k_- + k_+$. For the exponent of τ we used the bilinearity of σ and the obvious fact $\sigma(S_-, Y_+) = 0$.

We rephrase Lemma 5.1 of [12] to make it compatible with the present notation. Let A be a finite set of indices and let

$$g = g(\xi_i)_{i \in A}$$

be a function that is analytic for all $\xi_i \neq 0$. Assume that for $i < k$

$$g|_{\xi_i \rightarrow p/(1-q\xi_k)} = O(\xi_k^{-1})$$

as $\xi_k \rightarrow \infty$, uniformly when all the ξ_{ℓ} with $\ell \neq i, k$ are bounded and bounded away from zero. For $B \subset A$ denote by g_B the function obtained from g by setting all ξ_i with $i \notin B$ equal to 1. (In particular $g = g_A$.) Define

$$I_B(\xi) = \prod_{\substack{i < j \\ i, j \in B}} f(\xi_i, \xi_j) \frac{g_B(\xi)}{\prod_{i \in B} (1 - \xi_i)}.$$

Then when $p, q \neq 0$,

$$\begin{aligned} &\int_{\mathcal{C}_R^{|A|}} I_A(\xi) \prod_{i \in A} d\xi_i \\ &= \sum_{B \subset A} (-1)^{|A \setminus B|} \tau^{\sigma(B, A) - |A||B|} \frac{p^{|B|(|B|-1)/2}}{q^{|A|(|A|-1)/2}} \int_{\mathcal{C}_r^{|B|}} I_B(\xi) \prod_{i \in B} d\xi_i, \end{aligned} \quad (5)$$

where r is so small that all the zeros of the denominators lie outside \mathcal{C}_r . When B is empty the integral on the right side is interpreted as $g(1, \dots, 1)$.

We shall apply this to the integral on the right side of (4). In order to do this we change the set of indices in that integral from $[-k_-, -1] \cup [1, k_+]$ to $S_- \cup S_+$ in the obvious way. In the application of the lemma we take A to be the set S_- of negative indices, and g the integral over the remaining variables ξ_i with $i > 0$ (i.e., $i \in S_+$):

$$g(\xi) = \int_{\mathcal{C}_R^{k_+}} \prod_{\substack{i < j \\ j > 0}} f(\xi_i, \xi_j) \frac{\prod_i \xi_i^{x-s_i} e^{\varepsilon(\xi_i)t}}{\prod_{i>0} (1 - \xi_i)} \prod_{i>0} d\xi_i.$$

Then the integral on the right side of (4) equals the left side of (5).

Let us see that $g(\xi)$ satisfies the required conditions. Since R is arbitrarily large, $g(\xi)$ is analytic for $\xi_i \neq 0$ ($i < 0$). When $i < k < 0$ and we set $\xi_i = p/(1 - q\xi_k)$ the terms in the integral defining g that involve ξ_k are

$$\begin{aligned} & \prod_{j>0} f\left(\frac{p}{1-q\xi_k}, \xi_j\right) \cdot \prod_{j>0} f(\xi_k, \xi_j) \\ & \cdot \left(\frac{p}{1-q\xi_k}\right)^{x-s_i} \xi_k^{x-s_k} e^{[p/\xi_k + pq/(1-q\xi_k)]t}. \end{aligned}$$

As $\xi_k \rightarrow \infty$ the product of the f 's as well as the last factor are clearly bounded as $\xi_k \rightarrow \infty$. The product of the remaining two factors is $O(\xi_k^{s_i - s_k}) \rightarrow 0$ since $s_i < s_k$.

So the hypothesis of the lemma holds. In its application² we replace $|A \setminus B| = k_- - |B|$ of the ξ_i with $i < 0$ by 1. Each replacement has the effect of reducing the number of variables in $g(\xi)$ by one and dividing by q^{k_+} , so in the integral over the B -variables we must divide by $q^{k_+(k_- - |B|)}$. Therefore to get the coefficients of the resulting integrals over large and small contours we must multiply the coefficient in (5) by $q^{-k_+(k_- - |B|)}$. The result is

$$\begin{aligned} & (-1)^{k_- + |B|} \tau^{\sigma(B, S_-) - k_- |B|} \frac{p^{|B|(|B|-1)/2}}{q^{k_-(k_- - 1)/2 + k_+(k_- - |B|)}} \\ & = (-1)^{k_- + |B|} \tau^{-\sigma(S_-, B)} \frac{p^{|B|(|B|+1)/2}}{q^{k_-(k_- - 1)/2 + k_+(k_- - |B|) + |B|}}. \end{aligned}$$

Here we used the identity

$$\sigma(U, V) + \sigma(V, U) = |U| |V| + |U \cap V| \quad (6)$$

to obtain $\sigma(B, S_-) = -\sigma(S_-, B) + (k_- + 1)|B|$. Thus the integral on the right side of (4) equals

$$\begin{aligned} & \sum_{B \subset S_-} (-1)^{k_- + |B|} \tau^{-\sigma(S_-, B)} \frac{p^{|B|(|B|+1)/2}}{q^{k_-(k_- - 1)/2 + k_+(k_- - |B|) + |B|}} \\ & \times \int_{C_r^{|B|}} \int_{C_R^{k_+}} I(x, |B| + k_+, \xi) \prod_{i \in B} \xi_i^{-s_i} \prod_{i>0} \xi_i^{-s_i} \prod_i d\xi_i. \end{aligned} \quad (7)$$

When $|B| + k_- = 0$ the integral is interpreted as 1.

If we multiply the coefficient in (7) by the coefficient of the integral in (4) and use

$$k(k-1)/2 = k_-(k_- - 1)/2 + k_+(k_+ - 1)/2 + k_- k_+$$

and the bilinearity of σ we get

$$\begin{aligned} & (-1)^{m+k_- + |B|} \tau^{m(m-1)/2 - km + \sigma(S_-, Y_- \setminus B) + \sigma(S_+, Y)} \\ & \times q^{k_+(k_+ - 1)/2 + (k_+ - 1)|B|} p^{|B|(|B|+1)/2} \begin{bmatrix} k-1 \\ m-1 \end{bmatrix}_\tau. \end{aligned}$$

²In the resulting integrals over C_r we must have rR large, since to maintain the analyticity of $g(\xi)$ for small ξ_i (with $i < 0$) we must have all the ξ_j (with $j > 0$) even larger.

Now instead of first summing over all $B \subset S_-$ and then over all $S_- \subset Y_-$, we reverse the order by first fixing B and then summing over all S_- satisfying $B \subset S_- \subset Y_-$. We change notation, so that in the end we get integrands like those in (4) with $s_i \in S_- \cup S_+$. To do this we interchange B and S_- , so the old k_- becomes $|B|$ and the new k_- is defined to be the new $|S_-|$. We label our indices so that $s_i \in S_-$ when $i < 0$ and $s_i \in S_+$ when $i > 0$. Then our formula becomes

$$\begin{aligned} \mathbb{P}_Y(x_m(t) \leq x) \\ = \sum_{k_\pm \geq 0} \sum_{\substack{S_\pm \subset Y_\pm \\ |S_\pm| = k_\pm}} c_{m, S_-, S_+} \int_{C_R^{k_+}} \int_{C_r^{k_-}} I(x, k, \xi) \prod_i \xi_i^{-s_i} d\xi_i \end{aligned} \quad (8)$$

where $k = k_- + k_+$ and

$$\begin{aligned} c_{m, S_-, S_+} &= (-1)^{m+k_-} \tau^{m(m-1)/2 - k_+ m + \sigma(S_+, Y)} q^{k_+(k_+-1)/2 + k_-(k_-+1)} p^{k_-(k_-+1)/2} \\ &\times \sum_{Y_- \supset B \supset S_-} (-1)^{|B|} \tau^{\sigma(B, Y_- \setminus S_-) - m|B|} \left[\begin{matrix} |B| + k_+ - 1 \\ m - 1 \end{matrix} \right]_\tau \\ &= (-1)^{m+k_-} \tau^{m(m-1)/2 - k_+ m + \sigma(S_+, Y) + k_-(k_-+1)/2} q^{k(k-1)/2} \\ &\times \sum_{Y_- \supset B \supset S_-} (-1)^{|B|} \tau^{\sigma(B, Y_- \setminus S_-) - m|B|} \left[\begin{matrix} |B| + k_+ - 1 \\ m - 1 \end{matrix} \right]_\tau. \end{aligned} \quad (9)$$

When $k = 0$ the integral in (8) is interpreted as 1. After our change of notation the restriction $k \geq 1$ in the sum in (4) becomes $|B| + k_+ \geq 1$.

To evaluate the sum in (9) we use³

$$\sum_{\substack{U \subset X \\ |U|=n}} \tau^{\sigma(U, X)} = \tau^{n(n+1)/2} \left[\begin{matrix} |X| \\ n \end{matrix} \right]_\tau. \quad (10)$$

We write the power of τ in the sum in (9) as

$$\sigma(B \setminus S_-, Y_- \setminus S_-) + \sigma(S_-, Y_- \setminus S_-),$$

and then the sum in (9) as

$$\begin{aligned} (-1)^{k_-} \tau^{\sigma(S_-, Y_- \setminus S_-) - mk_-} \sum_{Y_- \supset B \supset S_-} (-1)^{|B \setminus S_-|} \tau^{\sigma(B \setminus S_-, Y_- \setminus S_-) - m|B \setminus S_-|} \\ \times \left[\begin{matrix} |B \setminus S_-| + k - 1 \\ m - 1 \end{matrix} \right]_\tau. \end{aligned}$$

The restriction $|B| + k_+ \geq 1$ is the same as $|B \setminus S_-| + k \geq 1$. If we apply (10) with $U = B \setminus S_-$ and $X = Y_- \setminus S_-$, then the last sum may be written

$$\sum_{n=0}^{|Y_- \setminus S_-|} (-1)^n \tau^{n(n+1)/2 - mn} \left[\begin{matrix} |Y_- \setminus S_-| \\ n \end{matrix} \right]_\tau \left[\begin{matrix} n + k - 1 \\ m - 1 \end{matrix} \right]_\tau$$

³The left side equals $\sum_{1 \leq u_1 < \dots < u_n \leq |X|} \tau^{u_1 + \dots + u_n}$, which is the coefficient of z^n in the expansion of $\prod_{k=1}^{|X|} (1 + z\tau^k)$ which, by the τ -binomial theorem [9, p. 26], equals $\sum_{n=0}^{|X|} \tau^{n(n+1)/2} z^n \left[\begin{matrix} |X| \\ n \end{matrix} \right]_\tau$.

when $k > 0$. When $k = 0$ the sum begins with $n = 1$. (Notice that all terms of the sum vanish unless $m \leq |Y_- \setminus S_-| + k$.) The sum from $n = 0$ is equal to⁴

$$(-1)^{|Y_- \setminus S_-|} \tau^{|Y_- \setminus S_-|(|Y_- \setminus S_-| - 2m + 1)/2} \left[\begin{matrix} k-1 \\ m - |Y_- \setminus S_-| - 1 \end{matrix} \right]_\tau.$$

When $k = 0$ the condition $m \leq |Y_- \setminus S_-| + k$ is the same as $m \leq |Y_- \setminus S_-|$, and then the sum from $n = 0$ equals zero. Therefore the sum from $n = 1$ equals minus the $n = 0$ term, which is

$$-\left[\begin{matrix} -1 \\ m-1 \end{matrix} \right]_\tau = (-1)^m \tau^{-m(m-1)/2}. \quad (11)$$

Using

$$|Y_- \setminus S_-|(|Y_- \setminus S_-| + 1)/2 = \sigma(Y_- \setminus S_-, Y_- \setminus S_-)$$

we see that when $k > 0$ the sum in (9) equals

$$(-1)^{|Y_-|} \tau^{\sigma(Y_-, Y_- \setminus S_-) - m|Y_-|} \left[\begin{matrix} k-1 \\ m - |Y_- \setminus S_-| - 1 \end{matrix} \right]_\tau$$

and so

$$\begin{aligned} c_{m, S_-, S_+} &= (-1)^{m+|Y_- \setminus S_-|} \tau^{m(m-1)/2 - k_+ m + \sigma(S_+, Y) + \sigma(Y_-, Y_- \setminus S_-) - m|Y_-| + k_- (k_- + 1)/2} \\ &\quad \times q^{k(k-1)/2} \left[\begin{matrix} k-1 \\ m - |Y_- \setminus S_-| - 1 \end{matrix} \right]_\tau. \end{aligned} \quad (12)$$

When $k = 0$ (i.e., $S_- = S_+ = \emptyset$) the sum in (9) equals $(-1)^m \tau^{-m(m-1)/2}$ when $m \leq |Y_-|$ and 0 otherwise, so $c_{m, \emptyset, \emptyset} = 1$ when $m \leq |Y_-|$ and 0 otherwise. Using (11) we find that this is exactly what we get if we substitute $k = 0$ into (12). Therefore (12) holds for all $k_\pm \geq 0$.

These are to be substituted into (8) to give $\mathbb{P}_Y(x_m(t) \leq x)$. To obtain the formula for $\mathbb{P}_Y(x_m(t) = x)$ we replace $I(x, k, \xi)$ in (8) by $\Delta I(x, k, \xi)$, obtaining

$$\mathbb{P}_Y(x_m(t) = x) = \sum_{k_\pm \geq 0} \sum_{\substack{S_\pm \subset Y_\pm \\ |S_\pm| = k_\pm}} c_{m, S_-, S_+} \int_{C_R^{k_+}} \int_{C_r^{k_-}} \Delta I(x, k, \xi) \prod_i \xi_i^{-s_i} d\xi_i. \quad (13)$$

Since the integral in (8) is interpreted as 1 for each x when $k = 0$, the integral here is interpreted as 0. Recall that in both formulas $s_i \in S_-$ when $i < 0$ and $s_i \in S_+$ when $i > 0$.

3 Formula for $\mathbb{P}_Y(\eta_x(t) = 1)$

To obtain $\mathbb{P}_Y(\eta_x(t) = 1)$, the probability that site x is occupied at time t given the initial configuration Y , we replace the coefficients in (13) by the sum of (12) over all m . The sum

⁴If we divide by the $n = 0$ term of the series the result is the τ -hypergeometric function ${}_2\phi_1(a, b; c; \tau, c/\tau)$ with $a = \tau^{-|Y_- \setminus S_-|}$, $b = \tau^k$, $c = \tau^{k-m+1}$. An application of Heine's q -Gauss identity [5, p. 14, (1.52)] with $q = \tau$ gives the stated equality. We "discovered" the equality by a Maple computation. Doron Zeilberger [15] showed us that it follows from the q -Gauss identity.

of those terms involving m is

$$\sum_m (-1)^m \tau^{m(m-1)/2 - (k_+ + |Y_-|)m} \begin{bmatrix} k-1 \\ m - |Y_- \setminus S_-| - 1 \end{bmatrix}_\tau,$$

and after making the substitution $m \rightarrow m + |Y_- \setminus S_-|$ this becomes

$$(-1)^{|Y_- \setminus S_-|} \tau^{|Y_- \setminus S_-|(|Y_- \setminus S_-| - 1)/2 - (k_+ + |Y_-|)|Y_- \setminus S_-|} \sum_m (-1)^m \tau^{m(m-1)/2 - mk} \begin{bmatrix} k-1 \\ m-1 \end{bmatrix}_\tau.$$

The factor outside the sum may be written

$$(-1)^{|Y_- \setminus S_-|} \tau^{-\sigma(Y_- \setminus S_-, Y_- \setminus S_-) - k |Y_- \setminus S_-|},$$

while another application of the τ -binomial theorem shows that when $k \geq 1$ the sum equals

$$(-1)^k \tau^{-k(k+1)/2} \prod_{j=1}^{k-1} (1 - \tau^j).$$

If we use

$$\sigma(Y_-, Y_- \setminus S_-) - \sigma(Y_- \setminus S_-, Y_- \setminus S_-) = -\sigma(Y_- \setminus S_-, S_-) + k_- |Y_- \setminus S_-|, \quad (14)$$

$$\sigma(S_+, Y) = \sigma(S_+, Y_+ \setminus S_+) + k_+(k_+ + 1)/2 + k_+ |Y_- \setminus S_-| + k_+ k_-, \quad (15)$$

which are obtained using the bilinearity of σ and (6), and multiply by the factors in (12) not involving m , we obtain for the sum over m

$$\begin{aligned} c_{S_-, S_+} &= (-1)^k \tau^{\sigma(S_+, Y_+ \setminus S_+) - \sigma(Y_- \setminus S_-, S_-)} q^{k(k-1)/2} \prod_{j=1}^{k-1} (1 - \tau^j) \\ &= -\tau^{\sigma(S_+, Y_+ \setminus S_+) - \sigma(Y_- \setminus S_-, S_-)} \prod_{j=1}^{k-1} (p^j - q^j). \end{aligned} \quad (16)$$

Thus, with these coefficients we have

$$\mathbb{P}_Y(\eta_x(t) = 1) = \sum_{k_\pm \geq 0} \sum_{\substack{S_\pm \subset Y_\pm \\ |S_\pm| = k_\pm}} c_{S_-, S_+} \int_{C_R^{k_+}} \int_{C_r^{k_-}} \Delta I(x, k, \xi) \prod_i \xi_i^{-s_i} d\xi_i, \quad (17)$$

where the integral is interpreted as 0 when $k = 0$.

4 $\mathbb{P}(\eta_x(t) = 1)$ for Bernoulli Initial Condition

We now derive a formula analogous to (17) for Bernoulli initial condition on \mathbb{Z} , with density ρ_- on $(-\infty, 0]$ and density ρ_+ on $[1, \infty)$. It will be a sum over k_\pm only. We shall begin with \mathbb{Z} replaced by $(-M, N]$, apply (17), and then let $M, N \rightarrow \infty$.

The probability for an initial configuration $Y = Y_- \cup Y_+$ is

$$\rho_-^{|Y_-|} (1 - \rho_-)^{M - |Y_-|} \rho_+^{|Y_+|} (1 - \rho_+)^{N - |Y_+|}. \quad (18)$$

To obtain the coefficient of the multiple integral corresponding to the indices k_{\pm} we have to multiply the above by c_{S_-, S_+} as given by (16), then sum over all $Y_{\pm} \supseteq S_{\pm}$, and finally multiply by $\prod \xi_i^{-s_i}$ and sum over all S_{\pm} with $|S_{\pm}| = k_{\pm}$. Because of the structure of c_{S_-, S_+} we need only compute the two limits

$$\lim_{N \rightarrow \infty} \sum_{\substack{S_+ \subset [1, N] \\ |S_+| = k_+}} \sum_{\substack{Y_+ \supset S_+ \\ Y_+ \subset [1, N]}} \rho_+^{|Y_+|} (1 - \rho_+)^{N - |Y_+|} \tau^{\sigma(S_+, Y_+ \setminus S_+)} \prod_{i > 0} \xi_i^{-s_i}, \quad (19)$$

where $S_+ = \{s_i, \dots, s_{k_+}\}$ with $s_1 < \dots < s_{k_+}$, and

$$\lim_{M \rightarrow \infty} \sum_{\substack{S_- \subset (-M, 0] \\ |S_-| = k_-}} \sum_{\substack{Y_- \supset S_- \\ Y_- \subset (-M, 0]}} \rho_-^{|Y_-|} (1 - \rho_-)^{M - |Y_-|} \tau^{-\sigma(Y_- \setminus S_-, S_-)} \prod_{i < 0} \xi_i^{-s_i}, \quad (20)$$

where $S_- = \{s_{-k_-}, \dots, s_{-1}\}$ with $s_{-k_-} < \dots < s_{-1}$, multiply them together, and then multiply by $\prod_{j=1}^{k_-} (p^j - q^j)$. In (19) the $|\xi_i|$ are large and in (20) the $|\xi_i|$ are small.

The limit (19) with the factor $\tau^{\sigma(S_+, Y_+)}$ instead of $\tau^{\sigma(S_+, Y_+ \setminus S_+)}$ was computed in [14].⁵ Using that result we find that (19) is equal to

$$\prod_{i=1}^{k_+} \frac{\rho_+}{\xi_i \cdots \xi_{k_+} - 1 + \rho_+ (1 - \tau^{k_+ - i + 1})}. \quad (21)$$

We required that these $|\xi_i|$ be sufficiently large to take the $N \rightarrow \infty$ limit.

For the limit (20) we set

$$\tilde{s}_i = -s_{-i} + 1, \quad \tilde{S}_+ = -S_- + 1, \quad \tilde{Y}_+ = -Y_- + 1,$$

(observe that $\tilde{s}_1 < \dots < \tilde{s}_{k_-}$) and use

$$\sigma(Y_- \setminus S_-, S_-) = \sigma(\tilde{S}_+, \tilde{Y}_+ \setminus \tilde{S}_+).$$

If we set $\tilde{\xi}_i = \xi_{-i}^{-1}$ ($i = 1, \dots, k_-$) then (20) becomes

$$\lim_{M \rightarrow \infty} \sum_{\substack{\tilde{S}_+ \subset [1, \dots, M] \\ |\tilde{S}_+| = k_-}} \sum_{\substack{\tilde{Y}_+ \supset \tilde{S}_+ \\ \tilde{Y}_+ \subset [1, \dots, M]}} \rho_-^{|\tilde{Y}_+|} (1 - \rho_-)^{M - |\tilde{Y}_+|} \tau^{-\sigma(\tilde{S}_+, \tilde{Y}_+ \setminus \tilde{S}_+)} \prod_{i > 0} \tilde{\xi}_i^{-\tilde{s}_i + 1}.$$

The limit formula (21) with an obvious modification shows that (20) equals

$$\prod_{i=1}^{k_-} \frac{\rho_- \tilde{\xi}_i}{\tilde{\xi}_i \cdots \tilde{\xi}_{k_-} - 1 + \rho_- (1 - \tau^{i - k_- - 1})} = \prod_{i=1}^{k_-} \frac{\rho_- \xi_{-i}^{-1}}{(\xi_{-i} \cdots \xi_{-k_-})^{-1} - 1 + \rho_- (1 - \tau^{i - k_- - 1})}.$$

We require that these $|\xi_i|$ be sufficiently small to take the $M \rightarrow \infty$ limit.

⁵The computation began with formula (5) of the cited paper and the limit is the displayed expression preceding identity (8).

Denote by $\varphi_{\pm}(k_{\pm}, \rho_{\pm}, \xi)$ the two limits,

$$\varphi_{+}(k_{+}, \rho_{+}, \xi) = \prod_{i=1}^{k_{+}} \frac{\rho_{+}}{\xi_i \cdots \xi_{k_{+}} - 1 + \rho_{+}(1 - \tau^{k_{+}-i+1})},$$

$$\varphi_{-}(k_{-}, \rho_{-}, \xi) = \prod_{i=1}^{k_{-}} \frac{\rho_{-} \xi_{-i}^{-1}}{(\xi_{-i} \cdots \xi_{-k_{-}})^{-1} - 1 + \rho_{-}(1 - \tau^{i-k_{-}-1})}.$$

If we recall (16) then we see that we have shown, formally, that for Bernoulli on \mathbb{Z} with densities ρ_{\pm}

$$\begin{aligned} \mathbb{P}(\eta_r(x) = 1) &= - \sum_{k_{\pm} \geq 0} \prod_{j=1}^{k-1} (p^j - q^j) \\ &\times \int_{C_R^{k_{+}}} \int_{C_r^{k_{-}}} \Delta I(x, k, \xi) \cdot \varphi_{-}(k_{-}, \rho_{-}, \xi) \varphi_{+}(k_{+}, \rho_{+}, \xi) \prod_i d\xi_i. \end{aligned} \quad (22)$$

The reason for the qualification “formally” is that we must justify the interchange of the limits $M, N \rightarrow \infty$ with the sum over k_{\pm} . To do this it is enough that we obtain a convergent series when the integrands in (16) with their coefficients are replaced by their upper bounds over M and N . Now not only was the limit (19) determined in [14], but an expression was found for finite N , namely

$$\rho_{+}^{k_{+}} \prod_{i=1}^{k_{+}} \xi_i^{-i} \sum_{\substack{t_i \geq 0 \\ \sum t_i \leq N-k_{+}}} \prod_{i=1}^{k_{+}} \left(\frac{1 - \rho_{+}(1 - \tau^{k_{+}-i+1})}{\xi_i \cdots \xi_{k_{+}}} \right)^{t_i}.$$

This gives the $N \rightarrow \infty$ limit if $R = |\xi_i|$ is sufficiently large (depending on τ) and also a uniform bound

$$A^{k_{+}^2} R^{-k_{+}^2/2+O(k_{+})}, \quad (23)$$

for some A and arbitrarily large R . Similarly for (20) we get a uniform bound

$$A^{k_{-}^2} r^{k_{-}^2/2-O(k_{-})} \quad (24)$$

for some A and arbitrarily small r .

As for the factor $\Delta I(x, k, \xi)$ in the integrand, we have the easy bounds

$$\prod_{0 < i < j} |f(\xi_i, \xi_j)| \leq A^{k_{+}^2} R^{-k_{+}^2/2+O(k_{+})}, \quad (25)$$

$$\prod_{i < j < 0} |f(\xi_i, \xi_j)| \leq A^{k_{-}^2} r^{k_{-}^2/2-O(k_{-})}, \quad (26)$$

$$\prod_{\substack{i < 0 \\ j > 0}} |f(\xi_i, \xi_j)| \leq A^{k_{-}k_{+}} r^{-k_{-}k_{+}}. \quad (27)$$

(For the last we used $rR \gg 1$; see footnote 2.) The rest of the integrand is at most $A^{k_+ R + k_- r^{-1}} R^{O(k_+)} r^{-O(k_-)}$. If we use $k_- k_+ \leq k_-^2/2 + k_+^2/2$ and combine the above bounds we get the bound

$$A^{k_+^2 + k_-^2 + k_+ R + k_- r^{-1}} R^{-k_+^2 + O(k_+)} r^{k_-^2/2 - O(k_-)}.$$

If we take R sufficiently large and r sufficiently small (depending on this A) we get a convergent sum of integrals. So (22) is justified.

Remark For SSEP the only nonzero terms in (22) are those with $k_- = 0$, $k_+ = 1$ and $k_+ = 0$, $k_- = 1$, and we get a sum of two single integrals with ρ_{\pm} appearing only as coefficients:

$$\mathbb{P}(\eta_t(x) = 1) = \rho_+ \int_{C_R} \frac{\xi^{x-1} e^{\varepsilon(\xi)t}}{\xi - 1} d\xi + \rho_- \int_{C_r} \frac{\xi^{x-1} e^{\varepsilon(\xi)t}}{1 - \xi} d\xi. \quad (28)$$

Symmetrization. In [14], where $\rho_- = 0$, we symmetrized the integrand in (22), found that it was a determinant, and this led to a representation for $\mathbb{P}(x_m(t) \leq x)$ as an integral of Fredholm determinants. Although it does not lead to as nice an expression, we can do an analogous symmetrization here.

Write the integrand in (22) as

$$\left(1 - \prod_i \xi_i^{-1}\right) \prod_{\substack{i < 0 \\ j > 0}} f(\xi_i, \xi_j) \quad (29)$$

$$\times \prod_{i < j < 0} f(\xi_i, \xi_j) \prod_{i < 0} \frac{\xi_i^x e^{\varepsilon(\xi_i)t}}{1 - \xi_i} \varphi_-(k_-, \rho_-, \xi) \quad (30)$$

$$\times \prod_{0 < i < j} f(\xi_i, \xi_j) \prod_{i > 0} \frac{\xi_i^x e^{\varepsilon(\xi_i)t}}{1 - \xi_i} \varphi_+(k_+, \rho_+, \xi). \quad (31)$$

The factor (29) is symmetric separately in the ξ_i with $i < 0$ and the ξ_i with $i > 0$.

The factor (31) is a function of the ξ_i with $i > 0$. Its symmetrization, given by formula (9) of [14], is

$$\frac{1}{k_+!} q^{k_+(k_+-1)/2} \prod_{i \neq j} f(\xi_i, \xi_j) \prod_i \frac{\rho_+}{\xi_i - 1 + \rho_+(1-\tau)} \prod_i \frac{\xi_i^x e^{\varepsilon(\xi_i)t}}{1 - \xi_i}.$$

(All indices positive.) By the identity (3) of [13] this equals

$$\begin{aligned} & \frac{(-1)^{k_+}}{k_+!} p^{-k_+(k_+-1)/2} \det \left(\frac{1}{p + q \xi_i \xi_j - \xi_i} \right)_{1 \leq i, j \leq k_+} \prod_i \frac{\rho_+ (q \xi_i - p)}{\xi_i - 1 + \rho_+(1-\tau)} \xi_i^x e^{\varepsilon(\xi_i)t} \\ &= \frac{(-1)^{k_+}}{k_+!} p^{-k_+(k_+-1)/2} \det(K_+(\xi_i, \xi_j))_{1 \leq i, j \leq k_+}, \end{aligned}$$

where

$$K_+(\xi, \xi') = \frac{\rho_+ (q \xi - p)}{\xi - 1 + \rho_+(1-\tau)} \frac{\xi^x e^{\varepsilon(\xi)t}}{p + q \xi \xi' - \xi}.$$

For the factors

$$\prod_{i < j < 0} f(\xi_i, \xi_j) \varphi_-(k_-, \rho_-, \xi),$$

we write as before $\tilde{\xi}_i = \xi_{-i}^{-1}$, and this becomes

$$\begin{aligned} & \prod_{i > j > 0} \frac{\tilde{\xi}_i - \tilde{\xi}_j}{p\tilde{\xi}_i\tilde{\xi}_j + q - \tilde{\xi}_j} \prod_{i=1}^{k_-} \frac{\rho_- \tilde{\xi}_i}{\tilde{\xi}_i \cdots \tilde{\xi}_{k_+} - 1 + \rho_-(1 - \tau^{i-k_- - 1})} \\ &= \prod_{0 < i < j} \frac{\tilde{\xi}_j - \tilde{\xi}_i}{q + p\tilde{\xi}_i\tilde{\xi}_j - \tilde{\xi}_i} \prod_{i=1}^{k_-} \frac{\rho_- \tilde{\xi}_i}{\tilde{\xi}_i \cdots \tilde{\xi}_{k_+} - 1 + \rho_-(1 - \tau^{i-k_- - 1})}. \end{aligned}$$

The symmetrization of this is

$$\begin{aligned} & \frac{1}{k_-!} p^{k_-(k_- - 1)/2} \prod_{i \neq j} \frac{\tilde{\xi}_j - \tilde{\xi}_i}{q + p\tilde{\xi}_i\tilde{\xi}_j - \tilde{\xi}_i} \prod_i \frac{\rho_- \tilde{\xi}_i}{\tilde{\xi}_i - 1 + \rho_-(1 - \tau^{-1})} \\ & \quad (\text{all indices positive}) \\ &= \frac{1}{k_-!} p^{k_-(k_- - 1)/2} \prod_{i \neq j} \frac{\tilde{\xi}_j - \tilde{\xi}_i}{p + q\tilde{\xi}_i\tilde{\xi}_j - \tilde{\xi}_i} \prod_i \frac{\rho_- \tilde{\xi}_i^{-1}}{\tilde{\xi}_i^{-1} - 1 + \rho_-(1 - \tau^{-1})} \\ & \quad (\text{all indices negative}). \end{aligned}$$

Hence the symmetrization of (30) is (all indices negative)

$$\begin{aligned} & \frac{1}{k_-!} p^{k_-(k_- - 1)/2} \prod_{i \neq j} f(\xi_i, \xi_j) \prod_i \frac{\rho_- \xi_i^{-1}}{\xi_i^{-1} - 1 + \rho_-(1 - \tau^{-1})} \prod_i \frac{\xi_i^{x-1} e^{\varepsilon(\xi_i)t}}{1 - \xi_i} \\ &= \frac{(-1)^{k_-}}{k_-!} q^{-k_-(k_- - 1)/2} \det \left(\frac{1}{p + q\xi_i\xi_j - \xi_i} \right) \prod_i \frac{\rho_- (q - p\xi_i^{-1})}{\xi_i^{-1} - 1 + \rho_-(1 - \tau^{-1})} \xi_i^x e^{\varepsilon(\xi_i)t} \\ &= \frac{(-1)^{k_-}}{k_-!} q^{-k_-(k_- - 1)/2} \det(K_-(\xi_i, \xi_j))_{-k_- \leq i, j \leq -1}, \end{aligned}$$

where

$$K_-(\xi, \xi') = \frac{\rho_- (q - p\xi^{-1})}{\xi^{-1} - 1 + \rho_-(1 - \tau^{-1})} \frac{\xi^x e^{\varepsilon(\xi)t}}{p + q\xi\xi' - \xi}.$$

So our formula may be written

$$\begin{aligned} \mathbb{P}(\eta_t(x) = 1) &= \sum_{\substack{k_-, k_+ \geq 0 \\ k_- + k_+ \geq 1}} \frac{1}{k_-! k_+!} p^{-k_+(k_+ - 1)/2} q^{-k_-(k_- - 1)/2} \prod_{j=1}^{k_- + k_+ - 1} (q^j - p^j) \\ &\times \int_{C_R^{k_+}} \int_{C_r^{k_-}} (1 - \prod_i \xi_i^{-1}) \prod_{\substack{i < 0 \\ j > 0}} f(\xi_i, \xi_j) \end{aligned}$$

$$\begin{aligned} & \times \det(K_-(\xi_i, \xi_j))_{-k_- \leq i, j \leq -1} \cdot \det(K_+(\xi_i, \xi_j))_{1 \leq i, j \leq k_+} \\ & \cdot \prod_{i<0} d\xi_i \cdot \prod_{i>0} d\xi_i. \end{aligned} \quad (32)$$

Remark 1 Suppose the change of density occurs at y rather than zero. Then to use the computations of (19) and (20) we would make the substitutions $Y_\pm = \tilde{Y}_\pm + y$, $S_\pm = \tilde{S}_\pm + y$, $s_i = \tilde{s}_i + y$ and apply those computations to \tilde{Y}_\pm , etc. The result in the end is that each ξ_i^x in the formula for $\Delta I(x, k, \xi)$ becomes ξ_i^{x+y} . In case $\rho_- = \rho_+$ it makes no difference what y is so the formulas we got were independent of x , as they should be.

Remark 2 To obtain formulas for the correlation function $\mathbb{P}(\eta_t(x) = 1, \eta_0(0) = 1)$ we write it as

$$\mathbb{P}(\eta_t(x) = 1) - \mathbb{P}(\eta_t(x) = 1, \eta_0(0) = 0). \quad (33)$$

The first probability we know. For the second we modify (18) by multiplying by $1 - \rho_-$, the probability that site 0 is initially unoccupied, and take only those Y_- that are contained in $(-\infty, -1]$. To use the preceding computation we use the one-one correspondence between subsets $S_- \subset (-\infty, -1]$ and subsets $S'_- \subset (-\infty, 0]$ given by $S'_- = S_- + 1$. Then $\sigma(Y_- \setminus S_-, S_-) = \sigma(Y'_- \setminus S'_-, S'_-)$ and, with obvious notation, $\prod \xi_i^{-s_i} = \prod \xi_i \times \prod \xi_i^{-s'_i}$. It follows that for the second probability in (33) we multiply the integrands in (22) and (32) by $(1 - \rho_-) \prod_{i<0} \xi_i$. Therefore for (33) itself we multiply the integrands by

$$1 - (1 - \rho_-) \prod_{i<0} \xi_i.$$

5 Formula for the Total Flux

Suppose $\max Y_- \leq 0 < \min Y_+$. When $t = 0$ there are $|Y_-|$ particles ≤ 0 , so Q_t , the total flux to the left across 0 at time t , is the number of particles ≤ 0 at time t minus $|Y_-|$. Thus $\mathbb{P}_Y(Q_t \geq m) = \mathbb{P}_Y(x_{m+|Y_-|}(t) \leq 0)$. Therefore in (8) and (12) we replace m by $m + |Y_-|$ and set $x = 0$. If we use (14) and (15) and replace m by $m + |Y_-|$ (so the new m may be negative) then (12) becomes

$$\begin{aligned} c_{m+|Y_-|, S_-, S_+} &= (-1)^{m+k_-} \tau^{\sigma(S_+, Y_+ \setminus S_+) - \sigma(Y_- \setminus S_-, S_-) + m(m-1)/2 + k_+(k_++1)/2 - mk_+} \\ &\times q^{k(k-1)/2} \left[\begin{matrix} k-1 \\ m+k_- - 1 \end{matrix} \right]_\tau. \end{aligned}$$

Using this, comparing with (16), and following the argument of the last section we obtain the formulas

$$\begin{aligned} \mathbb{P}(Q_t \geq m) &= \sum_{k_\pm \geq 0} (-1)^{m+k_-} \tau^{m(m-1)/2 + k_+(k_++1)/2 - mk_+} q^{k(k-1)/2} \left[\begin{matrix} k-1 \\ m+k_- - 1 \end{matrix} \right]_\tau \\ &\times \int_{C_R^{k_+}} \int_{C_r^{k_-}} I(0, k, \xi) \cdot \varphi_-(k_-, \rho_-, \xi) \varphi_+(k_+, \rho_+, \xi) \prod_i d\xi_i, \end{aligned}$$

$$\begin{aligned} \mathbb{P}(Q_t \geq m) &= \sum_{k_{\pm} \geq 0} \frac{(-1)^{m+k_+}}{k_-! k_+!} \tau^{m(m-1)/2-(m-1)k_+} q^{k_-k_+} \left[\begin{matrix} k-1 \\ m+k_- - 1 \end{matrix} \right]_{\tau} \\ &\times \int_{\mathcal{C}_R^{k_+}} \int_{\mathcal{C}_r^{k_-}} \prod_{\substack{i < 0 \\ j > 0}} f(\xi_i, \xi_j) \cdot \det(K_-(\xi_i, \xi_j))_{-k_- \leq i, j \leq -1} \\ &\cdot \det(K_+(\xi_i, \xi_j))_{1 \leq i, j \leq k_+} \cdot \prod_i d\xi_i, \end{aligned} \quad (34)$$

where in the expressions for $K_{\pm}(\xi_i, \xi_j)$ we set $x = 0$.

For $\mathbb{P}(Q_t = m)$ we subtract from these what we get by replacing m by $m + 1$. To get a formula for $\langle e^{\lambda Q_t} \rangle$, the expected value of $e^{\lambda Q_t}$, we then multiply by $e^{\lambda m}$ and sum over all m .

If $\mathbb{P}(Q_t \geq m) = p_m$ then $\langle e^{\lambda Q_t} \rangle = \sum_{m=-\infty}^{\infty} (p_m - p_{m+1}) e^{\lambda m}$. The series converges for all λ and represents an entire function of λ .⁶ Therefore if we find a formula that holds when $\Re \lambda > 0$, and the formula represents an entire function of λ , it will hold in general. When $\Re \lambda > 0$ we can write the series as the difference of two convergent series in the obvious way. Changing the summation variable in the second series and combining gives $(1 - e^{-\lambda}) \sum_{m=-\infty}^{\infty} p_m e^{\lambda m}$. This is what we do below.

The factor that involves m in the first formula is

$$(-1)^m \tau^{m(m-1)/2-mk_+} \left[\begin{matrix} k-1 \\ m+k_- - 1 \end{matrix} \right]_{\tau}.$$

When we make the replacement, subtract, multiply by $e^{\lambda m}$, and sum over m we get

$$(1 - e^{-\lambda}) \sum_{m=-k_- + 1}^{k_+} (-1)^m \tau^{m(m-1)/2-mk_+} e^{\lambda m} \left[\begin{matrix} k-1 \\ m+k_- - 1 \end{matrix} \right]_{\tau}$$

and another application of the τ -binomial theorem shows that this equals

$$\begin{aligned} &(-1)^{k_+} (1 - e^{-\lambda}) e^{\lambda k_+} \tau^{-k_+(k_++1)/2} \prod_{j=1}^{k-1} (1 - e^{-\lambda} \tau^j) \\ &= (-1)^{k_+} e^{\lambda k_+} \tau^{-k_+(k_++1)/2} \prod_{j=0}^{k-1} (1 - e^{-\lambda} \tau^j). \end{aligned}$$

Therefore

$$\begin{aligned} \langle e^{\lambda Q_t} \rangle &= \sum_{k_{\pm} \geq 0} (-1)^k q^{k(k-1)/2} e^{\lambda k_+} \prod_{j=0}^{k-1} (1 - e^{-\lambda} \tau^j) \\ &\times \int_{\mathcal{C}_R^{k_+}} \int_{\mathcal{C}_r^{k_-}} I(0, k, \xi) \cdot \varphi_-(k_-, \rho_-, \xi) \varphi_+(k_+, \rho_+, \xi) \prod_i d\xi_i. \end{aligned}$$

⁶For positive m , if $Q_t \geq m$ then the particle initially the m th to the right of 0 must have moved at least m steps to the left at time t . The probability of this is less than the probability that a free particle would so move, and this probability is $O(e^{-m} \log m + O(m))$ as $m \rightarrow +\infty$. Similarly $\mathbb{P}(Q_t \geq m) = 1 - O(e^{-|m|} \log |m| + O(|m|))$ as $m \rightarrow -\infty$.

To justify the interchange of the sums over m and k_{\pm} , and to show that the sum represents an entire function of λ , we use estimates analogous to those used in Sect. 4. For the integral (34) we have the bound $A^{k^2+k_+R+k_-r^{-1}} R^{-k_+^2+O(k_+)} r^{k_-^2/2-O(k_-)}$ and for its coefficient another bound A^{k^2} for all m , since $|m| \leq k$ when the coefficient is nonzero. And $|e^{\lambda m}| \leq e^{k|\lambda|}$ for such m . If R is large enough and r small enough the sum (over k_{\pm} and m) of the bounds is finite and uniformly bounded in bounded λ -sets. This justifies the computation.

In the case of SSEP ($\tau = 1$) our formula simplifies since

$$\varphi_{\pm}(k_{\pm}, \rho_{\pm}, \xi) = \rho_{\pm}^{k_{\pm}} \varphi_{\pm}(k_{\pm}, 1, \xi),$$

and we can write it as

$$\begin{aligned} \langle e^{\lambda Q_I} \rangle &= \sum_{k_{\pm} \geq 0} (-1)^k 2^{-k(k-1)/2} \rho_-^{k_-} \rho_+^{k_+} (1 - e^{-\lambda})^{k_-} (e^{\lambda} - 1)^{k_+} \\ &\quad \times \int_{C_R^{k_+}} \int_{C_r^{k_-}} I(0, k, \xi) \cdot \varphi_-(k_-, 1, \xi) \varphi_+(k_+, 1, \xi) \prod_i d\xi_i. \end{aligned}$$

In particular, this is a function of the two quantities $\rho_- (1 - e^{-\lambda})$ and $\rho_+ (e^{\lambda} - 1)$.

This fact was derived in [3] in a different way. The authors went further to show that it is a function of a particular combination of these quantities. This allowed them to reduce to the case when $\rho_- = 0$ and to obtain, for general ρ_{\pm} , a series for the generating function involving integrals only over small contours.

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References

1. Ben Arous, G., Corwin, I.: Current fluctuations for TASEP: A proof of the Prähofer-Spohn conjecture. [arXiv:0905.2993](https://arxiv.org/abs/0905.2993)
2. Derrida, B.: An exactly soluble non-equilibrium system: the asymmetric simple exclusion process. *Phys. Rep.* **301**, 65–83 (1998)
3. Derrida, B., Gerschenfeld, A.: Current fluctuations of the one dimensional symmetric simple exclusion process with step initial condition. *J. Stat. Phys.* **136**, 1–15 (2009)
4. Ferrari, P.L., Spohn, H.: Scaling limit for the space-time covariance of the stationary totally asymmetric simple exclusion process. *Commun. Math. Phys.* **265**, 1–44 (2006)
5. Gasper, G., Rahman, M.: Basic Hypergeometric Series. Cambridge University Press, Cambridge (2004)
6. Johansson, K.: Shape fluctuations and random matrices. *Commun. Math. Phys.* **209**, 437–476 (2000)
7. Liggett, T.M.: Interacting Particle Systems. Springer, Berlin (2005) [Reprint of the 1985 edition]
8. Liggett, T.M.: Stochastic Interacting Systems: Contact, Voter and Exclusion Processes. Springer, Berlin (1999)
9. Macdonald, I.G.: Symmetric Functions and Hall Polynomials. Clarendon, Oxford (1995)
10. Prähofer, M., Spohn, H.: Current fluctuations for the totally asymmetric simple exclusion process. *Prog. Probab.* **51**, 185–204 (2002)
11. Spitzer, F.: Interaction of Markov processes. *Adv. Math.* **5**, 246–290 (1970)
12. Tracy, C.A., Widom, H.: Integral formulas for the asymmetric simple exclusion process. *Commun. Math. Phys.* **279**, 815–844 (2008)

13. Tracy, C.A., Widom, H.: A Fredholm determinant representation in ASEP. *J. Stat. Phys.* **132**, 291–300 (2008)
14. Tracy, C.A., Widom, H.: On ASEP with step Bernoulli initial condition. *J. Stat. Phys.* **137**, 825–838 (2009)
15. Zeilberger, Z.: Private communication, Nov. 2009