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Plasmon-Enhanced Monolayer MoS2 for Efficient and Sustainable Photocatalytic Water Splitting Platform

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#### UNIVERISTY OF CALIFORNIA

Los Angeles

Plasmon-Enhanced Monolayer MoS<sub>2</sub> for Efficient and Sustainable Photocatalytic

Water Splitting Platform

A dissertation submitted in partial satisfaction of the

requirements for the degree Doctor of Philosophy

in Materials Science and Engineering

by

Peiyi Ye

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#### ABSTRACT OF THE DISSERTATION

#### Plasmon-Enhanced Monolayer MoS<sub>2</sub> for Efficient and Sustainable Photocatalytic

Water Splitting Platform

by

Peiyi Ye

Doctor of Philosophy in Materials Science and Engineering University of California, Los Angeles, 2020 Professor Ya-Hong Xie, Chair

 $H_2$  has long been known to be one of the highest energy density fuels. Fuel cell technology has been actively pursued as an environmentally friendly power source for automobiles. The promise of the technology is hindered by the lacking of a sustainable way of producing  $H_2$ . Solar water-splitting via photo-electrochemical cells (PEC) is the most promising approach which converts the sustainable solar energy to the chemical energy store inside the  $H_2$  bonding. The materials that can be used in a PEC cell must fulfil a variety of thermodynamic and kinetic requirements to ensure good efficiency and durability. Single-layer  $MoS_2$  possess the corrected energy bandgap of 1.9 eV which allows for sufficient over-potential while still being capable of absorbing the majority of the solar spectrum. However, the limited optical absorbance from singlelayer  $MoS_2$  prevent it from widely used. This shortcoming of the single-layer thickness of  $MoS_2$  is overcome by superimposing the  $MoS_2$  with plasmonic surface that serves to amplify the enhanced electromagnetic field where the  $MoS_2$  is located, allowing singlelayer  $MoS_2$  to efficiently absorb Sun light thus producing H<sub>2</sub>.

In the first part of dissertation, a novel two-step chemical vapor deposition method is developed to consistently grow high coverage and exclusive single-layer MoS<sub>2</sub>. Up to 90% surface coverage and single-layer MoS<sub>2</sub> is successfully fabricated. This is a crucial step to conduct the follow-up experiments.

In the second part of dissertation, A facile one-pot synthetic approach for synthesis hollow Au nanoframes structure is reported for the first time. A growth mechanism has been revealed that involves a synergistic function of Ag and Br ions. The presence of  $Ag^+$  lead to observed self-limiting of Au film thickness whereas Au {111} facets are preferentially attacked by the presence of Br<sup>-</sup> in the reaction ambient. Combined simulation and experimental studies show strong plasmonic effect that the hybrid platform made of graphene/Au nanoframes is capable of detecting analytes at concentration levels down to  $10^{-9}$  M by using the surface-enhanced Raman spectroscopy (SERS) technique.

In the last part of the dissertation, the plasmonic effect generated from nanostructured metal surfaces is introduced to offset the small thickness from single-layer  $MoS_2$  and improve the overall absorption. It is demonstrated for the first time that using single-layer  $MoS_2$  as a well-defined nanospacer between Au-nanoparticles and Au-film (gap plasmon system). The field enhancement is known to be inversely proportional to this gap thickness. Hence reducing the gap thickness is important to achieve the highest possible field enhancement. In this work, it is demonstrated for the first time that using single-layer MoS<sub>2</sub> as a well-defined nanospacer between Au-nanoparticles (AuNPs) and Au-film, which could offer an extremely high localized electric field enhancement within the gap. The MoS<sub>2</sub> Raman intensity with the Surface Enhanced Raman Scattering (SERS) enhancement factor (EF) up to  $5 \times 10^6$  is obtained from the MoS<sub>2</sub>-Au gap plasmon system. A 5-fold increase in the photocurrent is obtained from the MoS<sub>2</sub> prepared under the same condition. Compared with individual metal nanoparticles commonly used to enhance thin-film photocatalytic process, gap-plasmon could theoretically produce 8 orders of magnitude higher SERS EF and precise control the hot spot location to superimpose where ultrathin materials locate thus using the higher incident energy available.

The dissertation of Peiyi Ye is approved.

Ali Mosleh

Dwight Streit

Xiaochun Li

Ya-Hong Xie, Committee Chair

University of California, Los Angeles

Dedicated to my family, teachers and friends

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#### Chapter 1 Introduction

#### 1.1 Motivation of the thesis

 $H_2$  has long been known to be one of the highest energy density fuels<sup>1</sup>. Fuel cell technology has been actively pursued as an environmentally friendly power source for automobiles<sup>2</sup>. The promise of the technology is hindered by the lacking of a sustainable way of producing H<sub>2</sub>. Solar water-splitting via photo-electrochemical cells (PEC)<sup>3</sup> is the most promising approach with intense research effort being dedicated to it. To date, only a handful of semiconductor materials meet these stringent requirements such as GaP<sup>4</sup>, TiO<sub>2</sub><sup>5-</sup> <sup>7</sup> and  $MoS_2^{8-10}$ . Among them, single-layer  $MoS_2$  possesses the largest chemical overpotential<sup>8-9</sup> with films thicker than a one-molecular layer having energy bandgap values too small for the intended purpose. MoS<sub>2</sub> is one member of the family of "van der Waals materials" (vdW), named after the nature of van der Waals bonding in between layers. The requirement of a one-molecular layer thickness makes the use of MoS<sub>2</sub> impractical because of the limited absorption of sunlight by the extremely thin film. Consequently, Most of the research efforts on solar water-splitting have been directed at using tandem semiconductor junctions with their own set of technological challenges that are yet to be overcome<sup>11-12</sup>. The thesis addresses the major roadblock that preventing the use of a single-layer MoS<sub>2</sub> for solar water-splitting by adapting surface plasmon resonance (SPR), a phenomenon of collective oscillation of conducting electrons at nanostructured metal surfaces when excited by light<sup>13</sup>. The primary shortcoming of the single molecular-layer thickness of MoS<sub>2</sub> is overcome by superimposing the MoS<sub>2</sub> over a plasmonic surface that serves to amplify the electromagnetic field where the MoS<sub>2</sub> locates, allowing single-layer MoS<sub>2</sub> to efficiently absorb sunlight thus producing  $H_2$ . This  $MoS_2$  coupled plasmonic platform proves the much higher  $H_2$  production, which help significant reduction in greenhouse gas emissions. It is an important and ingenious utilization of van der Waals materials and Plasmon resonance while chartering out the pathway for them to enter commercial-scale production.

This chapter introduces the background information, the properties and the common methods to grow transition metal dichalcogenide is introduced in section 1.2, the principle of surface plasmon resonance in section 1.3, the principle of the photo-electrochemical water-splitting platform is introduced in section 1.4, the outline of the thesis is lastly introduced in section 1.5.

#### 1.2 Transition Metal Dichalcogenides

The transition metal dichalcogenides (TMDs) are one type of van der Waals materials that contain one transition metal center (Mo, W) and six surrounded chalcogenide atoms (S, Se, Te) that sandwiching the transition metal atoms<sup>14</sup> (Figure 1.1a). Among them, the most famous and widely studied TMD materials belong to MoS<sub>2</sub>, it is a silvery black solid that naturally occurs as the mineral molybdenite<sup>15</sup>. It is widely utilized as an industrial lubricant due to its weak van der Waals interactions between layers, it is also relatively unreactive, unaffected by dilute acids and  $O_2^{16}$ . MoS<sub>2</sub> usually consists of a mixture of two major polytypes, with the 2-H form being more abundant, in 2-H crystal structure with a P63/mmc point group (Figure 1.1 b, c)<sup>14, 17</sup>. In 2-H form MoS<sub>2</sub> unit cell, each Mo center occupies a trigonal prismatic coordination sphere and is bound to six S, each S center is

pyramidal and is connected with 3 Mo.



Figure 1. 1 (a) Periodic table highlighting element available for the transition metals and chalcogenides, (b) Crystal structure of group 6 transition metal dichalcogenides; 2-H hexagonal trigonal prismatic crystal structure. A single layer is consisting of a chalcogen sandwiching a metal atom. (c)  $D_3h$  trigonal unit cell. Images adapted from reference<sup>17</sup>

The electronic properties of bulk MoS<sub>2</sub> is semiconducting with a 1.2 eV indirect band gap<sup>18</sup>. There is a considerable energy band shift from an indirect-to-direct band gap that is typically found within all group six TMD materials, the MoS<sub>2</sub> band structure shift while changing from bulk form to single-layer form can be seen in Figure 1.2. When the layer number gradually decreases till the separated two-dimension sheet, the energy of the valence bands decreases due to the antibonding feature of that orbital interactions<sup>19</sup>. On the other hand, the maximum of the valence band and the minimum of the conduction band do not change when the layers detach at the K-point. The energy states at the K-point consist of Mo-Mo  $d_{xy}$  and  $d_{dx}^2$ -y<sup>2</sup> exchanges, within the plane of the MoS<sub>2</sub> and it is independent of interlayer distance<sup>19</sup>. The gap between these two energy levels keeps constant around 1.9 eV and doesn't require momentum change<sup>19</sup>. In this case, when MoS<sub>2</sub> layer decrease to a single layer, there is a noticeable electronic structure shift from a 1.2 eV indirect band gap ( $\Gamma$ →K) to a 1.9 eV direct band gap (K→K).



Figure 1. 2 Energy band diagram in bulk, 4L, 2L and monolayer MoS2. The electronic signature shift from indirect to direct when thickness reaches to a monolayer. Image adapted from reference<sup>20</sup>.

Chemical vapor deposition (CVD) is an easy and most widely adapted bottom-up methods to grow two-dimension materials<sup>21</sup>. Starting from 2012, a lot of new CVD methods have been reported to grow single-layer MoS<sub>2</sub>. Atomically thin Mo film was deposited on a substrate followed by S annealing to form layer MoS<sub>2</sub> was reported<sup>15</sup>. Follow a similar procedure, MoO<sub>3</sub> was used as Mo precursor to deposited on a substrate followed by annealing sulfur vapor at a high temperature<sup>22</sup>. Another Mo precursor was also introduced by using (NH<sub>4</sub>)<sub>2</sub>MoS<sub>4</sub> solution on a substrate, followed by sulfur vapor annealing vapor<sup>23</sup>. From another paper, MoO<sub>3</sub> was first evaporated to form a MoO<sub>2</sub> template followed by treatment with sulfur vapor<sup>24</sup>. All of these methods above reports have two major drawbacks: first, the non-uniform MoS<sub>2</sub> thickness distribution along the substrate, second, the non-continuous MoS<sub>2</sub> film growth on the substrate. The performance

of electric and optoelectronic is significantly hindered by those drawbacks. In the second chapter of the dissertation a novel method to grow near continuously exclusive single-layer  $MoS_2$  is introduced that can be used as in a wider application including the photoelectrochemical cells in water splitting application that is described in Chapter 4.

#### 1.3 Surface Plasmonic Resonance

Surface plasmon resonance (SPR) is a resonant oscillation of conduction electrons at the interface between materials of negative and positive dielectric constants excited via incident light. SPR was first discovered by Kretschmann and Raether <sup>25</sup>. Since then, intensive attention has been paid to this field, the principles and related application can be found in numerous papers, for example, SPR principles from Davis, L.C. et al. at 1977<sup>26</sup>, SPR sensors from J. Homola. et al<sup>27-29</sup>, Surface-enhanced Raman scattering from A. Campion et al<sup>30</sup>.

In general, there are two different types of surface plasmons, depending on the geometry of the metals shown in Figure 1.3a, the schematic of SPR on the interface between the thin metal film and a dielectric as well as on individual nanoparticles. SPR together with a metallic thin film of thickness can travel along the interface, which is called propagating SPR or surface plasmon polariton(SPP)<sup>31</sup>. The SPR associated with metal nanoparticles with the size between 10 to 200 nm is localized to the neighboring metal nanoparticle, which is named localized SPR (LSPR)<sup>32</sup>, schematically shown in Figure 1.3b.



Figure 1. 3 (a) Illustration of surface plasmon polariton (SPP) (b) Illustration of localized surface plasmon resonance (LSPR). (c) Typical dispersion curves of SPPs (red) and LSPs (blue) Adapted from reference<sup>33</sup>

LSPR is the collective electron charge oscillations within individual metal nanoparticles. They show enhanced near-field electromagnetic fields at the resonance frequency. This enhanced field is highly localized near the surface of metal nanoparticles which decreased rapidly away from the surface into the dielectric surrounding environment. The LSPR has a very high spatial resolution (subwavelength) due to the localization, limited only by the size of nanoparticles. SPPs have a continuous dispersion relation thus exists over a long range of frequency, however, LSPR only exists on a limited frequency range due to the additional constraints limited via its finite dimensions shown in Figure 1.3c. The particle's shape, size as well as dielectric functions of both the metal and the surrounding media are major parameters that control the spectral position of the SPR. Furthermore, SPPs cannot be directly coupled with propagating light while SPRs could.

A variety of metals, including commonly used gold, silver as well as some not commonly used copper, aluminum, sodium, indium, titanium, and chromium, can show SPR<sup>34</sup>. Compared with nanostructures made from other uncommon materials, gold and silver have been paid much more attention due to their unique optical properties. As shown in Figure 1.4, silver and gold nanostructures exhibit a wide range high quality factor or a dimensionless parameter that represents Q factor which is how underdamped a resonator is along the wavelength. In another word, Au and Ag SPR can be more easily for broadband resonant compared with other metal candidates. It is worth to mention that silver is too easy to oxidation in the air and most of the time it could be poisonous to biological samples, so usually gold is preferred for bio-sensing or used in a harsh environment. Recently, A wide range of applications have been realized via using plasmonic, such as high-resolution optical imaging below the diffraction limit, biodetection at single-molecule level, surface-enhanced Raman scattering, enhanced optical transmission through sub-wavelength apertures, etc. Undoubtedly, there will be more and more findings and new applications in the field of surface plasmon resonance in the near future.



Figure 1. 4 Quality factor (Q) of localized surface plasmon resonance between metal and air. Adapted from reference<sup>34</sup>

#### 1.4 Photoelectrochemical Water Splitting

A lot of previously published books<sup>35-37</sup> and review articles<sup>38-39</sup> have already discussed the very detailed working principles of PEC cell. However, the basic operation principle of the PEC cell is worth to introduce briefly here. The schematic in figure 1.5 showed a typical PEC cell with a two-electrode system, one electrode we called working electrode, which using an n-type semiconductor as photoanode. Incoming photons excite the semiconductor to generate electrons and holes pairs. The photogenerated electrons and holes pairs separate and move through the semiconductor in opposite directions; The holes participate the oxygen evolution reaction (OER) at the surface of the photoanode or working electrode. At the same time, the electrons are diffused towards the surface of the

counter electrode or counter electrode to participate the hydrogen evolution reaction (HER). It is also shown the minimum thermodynamic energy required for splitting water is 1.23 ev in Figure 1.5. A little overpotential is required to drive the kinetics of the HER and OER at the interface between solid and liquid. Minimizing these overpotentials is the key step to get the efficient catalysts for making highly efficient water splitting devices. Among all the available semiconductor, single-layer MoS<sub>2</sub> stands out as a great candidate for solar water-splitting. In addition to its corrosion resistance, single-layer MoS<sub>2</sub> is a direct bandgap semiconductor with ideal energy band structures: The conduction band edge (CB) is above that of the H<sup>+</sup>/H<sub>2</sub> reduction of H<sup>+</sup> in water, and the valence band edge (VB) is below that of H<sub>2</sub>O/O<sub>2</sub> potential<sup>20</sup>. The energy bandgap (E<sub>g</sub>) of 1.9eV allows for sufficient overpotential for H<sub>2</sub> production (1.23 eV)<sup>40</sup> while still being capable of absorbing almost all of the solar spectrum. The nature of direct bandgap means large oscillator strengths for electron-photon interaction and high quantum efficiency. The details on MoS<sub>2</sub>-based photoelectrochemical water splitting is studied in Chapter 4.



Figure 1. 5 Schematic and Band structure of an n-type semiconductor as photoanode within water splitting device. The processes include photon irradiation, electron–hole pair formation, charge transport, and interfacial reactions. Adapted from reference<sup>41</sup>

#### 1.5 Outline of the Thesis

Chapter 1 introduces the motivations of the thesis, basic principle and typical growth method for transitional metal dichalcogenides MoS<sub>2</sub>, the basic principle of surface plasmon resonance and photo-electrochemical water splitting platform.

Chapter 2 introduces a novel and controlled two-step method to grow  $MoS_2$  on carbon nanotube (CNT) substrate via typical chemical vapor deposition  $(CVD)^{42}$  setup, the growth method is significant in the following two aspects: first, direct growth of large quantity  $MoS_2$  nanoflakes on carbon nanotube sheet is realized for the first time; second, plasma activated CNT sheet is of critical importance to immobilize  $MoS_2$  nanoflakes. We reveal that plasma treatment of CNT sheet results in de-bundling and activation of carbon nanotube bundles, which promotes the deposition of  $MoS_2$  nanoflakes.

Chapter 3 presents a novel method for synthesizing well-defined Au nanoframes with tunable size on graphene substrate, which shows an extremely strong surface plasmon resonances via simulation and experimental study. A comprehensive investigation on structures and morphologies of graphene-Au nanoframes is studied to improve the electric field and the overall electric field enhancement nearby. A similar method is used for synthesizing Au nanoparticles for the solar water splitting application in Chapter 4

Chapter 4 demonstrates the feasibility by using single-layer  $MoS_2$  over a plasmonic surface for efficient water splitting applications. The primary shortcoming of the single molecular-layer thickness of  $MoS_2$  is overcome by superimposing the  $MoS_2$  over the plasmonic surface that serves to amplify the electromagnetic field where the  $MoS_2$  locates, allowing single-layer  $MoS_2$  to efficiently absorb Sunlight thus producing  $H_2$ .

Chapter 5 concludes the dissertation by providing each chapter's results with a section that proposes future work that can build on the results presented herein.

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Chapter 2 Growth of MoS<sub>2</sub> on Carbon Nanotube substrate via Chemical Vapor Deposition methods

#### Abstract

In this chapter, we demonstrate a novel two-step approach that is capable of growing uniform high-concentration MoS<sub>2</sub> nanoflakes on the carbon nanotube (CNT) sheet substrate. The CNT sheet is simply activated by the oxygen plasma to separate large CNT bundles and to create more active surface sites. A typical chemical vapor deposition (CVD) process is employed, enabling the direct growth of MoS<sub>2</sub> nanoflakes on the activated CNT substrate. Morphologies and structures of MoS<sub>2</sub>/CNT nanohybrid are characterized by SEM/EDS, high-resolution TEM (HRTEM) and Raman spectroscopy. Our results suggest that the plasma treatment of CNT sheet is critical to obtain well-dispersed large quantity of MoS<sub>2</sub> nanoflakes on CNT assemblies.

# 2.1 Introduction

Molybdenum disulfide (MoS<sub>2</sub>), one of the representative transition metal dichalcogenides (TMDs), has attracted increasing research interest in the field of twodimensional (2-D) nanomaterials<sup>1</sup>. The growth, structures, and emerging applications of  $MoS_2$  have been extensively investigated. In particularly,  $MoS_2$  and its composites have been widely employed as anode materials in lithium ion batteries (LIBs)<sup>2,3</sup>. Incorporation of  $MoS_2$  with highly conductive carbonaceous materials, such as CNTs and graphene<sup>4,5</sup>, can greatly improve the low-conductivity issue caused by  $MoS_2$  in LIBs<sup>6,7,8</sup>. Consequently, researchers have devoted great efforts to developing nanocomposites made of  $MoS_2$  with these carbon materials, particularly with CNTs<sup>8</sup>.

However, it is still challenging to grow high-concentration MoS<sub>2</sub> on pristine CNT platforms due to the chemical inertness of CNT walls. CNTs have to be functionalized to get sufficient active spots that can immobilize MoS<sub>2</sub><sup>9,10</sup>. Traditional functionalization of CNTs usually contains multiple steps, which is complicated, time-consuming, and costly <sup>11</sup>. Moreover, current research mainly focuses on using loose CNT powder as the support for MoS<sub>2</sub> <sup>8,9,10</sup>. To the best of our knowledge, reports on directly growing MoS<sub>2</sub> nanostructures on CNT assemblies, i.e. CNT yarns and sheets are very limited. CNT assemblies are the macro-form of individual CNTs, which transfer their excellent mechanical and physical properties from nanoscale to macroscopic scale<sup>12</sup>. Therefore, CNT macroscopic assemblies are the ideal platform to support foreign materials to make novel nanocomposites<sup>13</sup>.

In this study, we report a novel approach to directly grow high-concentration  $MoS_2$  nanoflakes on the CNT sheet with high uniformity. By using oxygen plasma, we not only increase effective surface areas of the CNT sheet via de-bundling process but also create more active sites to anchor  $MoS_2$ . The activated CNT sheet is then employed as the substrate and  $MoS_2$  nanoflakes can uniformly grow on it through a CVD process.

Compared with untreated CNT sheet, both uniformity and concentration of MoS<sub>2</sub> are significantly improved on the substrate of plasma treated CNT sheet.

#### 2.2 Experimental

# 2.2.1 Chemicals and Materials

CNT sheet was received from Nanocomp Technologies Inc. (Merrimack, NH), which was fabricated by a continuous CVD process. Sulfur (S, 99.998%) and molybdenum oxide (MoO<sub>3</sub>, 99.97%) were purchased from Sigma-Aldrich. The materials were used as received without any treatment or purification.

#### 2.2.2 Plasma treatment of CNT sheet.

CNT sheet was first cut into the dimension of  $10 \text{ cm} \times 2.54 \text{ cm}$ . The plasma treatment was performed by Atomflo<sup>TM</sup> plasma system (Surfx® Technologies LLC) with the power of 150 W, helium gas feeding rate at 30 L/min and oxygen at 0.5 L/min. Plasma scanning rate was 20 mm/s and the total exposure time was 90 s for one side. Double-side treatment was carried out. The plasma treatment is schematically shown as Figure 2.1, step 1.

#### 2.2.3 $MoS_2$ growth on the CNT sheet substrate.

We use traditional CVD method to grow  $MoS_2$  on the CNT sheet substrate. The equipment set-up is schematically shown in Figure 2.1 (step 2). In short, 10 mg  $MoO_3$  powder and 500 mg S powder were used as Mo and S precursors, respectively. The CNT sheet was put up-side down facing towards the  $MoO_3$  powder localized at the furnace center. The substrate temperature is set to be 850 °C, while S powder was put at the temperature window, where the T was set as 200 °C. The total growth time was 15min. During the growth, Ar was employed as the carrier gas with a flow rate of 100 sccm.



Figure 2. 1 Schematic illustration of two-step approach for MoS<sub>2</sub> growth on the CNT sheet.

### 2.2.4 Characterizations

Scanning electron microscopy (SEM) images coupled with related energy dispersive X-ray (EDS) analysis of the samples were obtained from FEI Nova NanoSEM 230. Transmission electron microscopy (TEM) and high-resolution transmission electron microscopy (HRTEM) images were captured with FEI Titan S/TEM system at 200 keV. Raman spectroscopy analysis was obtained with a Renishaw In-Via Raman system with the laser length of 633 nm and laser power of 1mW at objective lens magnification of 50x and a grating spacing of 1200 *l*/mm.

# 2.3 Result and Discussion

Figure 2.2 (a) is the photograph of a typical spool of raw CNT sheet obtained from the supplier. From the microstructure of raw CNT sheet shown in Figure 2.2 (b) and (c), one can see the condensed network composed of CNT bundle. In these bundles, individual CNTs are bonded together by van de Waals interaction<sup>12</sup>. After plasma treatment (Figure 2.2 (d) and (e)), nanotubes are uniformly separated apart from original bundles, which is a so-called "de-bundling" process. The de-bundling would significantly enhance the effective surface area of CNT sheet, as schematically shown in Figure 2.2 (f). Another benefit from plasma treatment is the creation of more defective but active sites, which will be discussed in details shortly.



Figure 2. 2 (a) Photography of CNT sheet spool. (b) Low- and (c) High- magnification of SEM images of raw CNT sheet. (d) Low- and (e) High- magnification of SEM images of plasma activated CNT sheet. (f) De-bundling process induced by plasma treatment.

We now investigate the growth of  $MoS_2$  on the CNT sheet substrate. Firstly,  $MoS_2$  can grow on the raw CNT sheet. Their distribution on the CNT support, however, is quite nonuniform as shown in Figure 2.3 (a). In addition, they demonstrate very distinguishable shapes and sizes. These  $MoS_2$  nanostructures aggregate on the CNT bundles (Figure 2.3 (b)), creating large domains that are free of  $MoS_2$  (Figure 2.3 (d)). Such a poor dispersion of  $MoS_2$  can be identified from the EDS result in Figure 2.3 (i), where only a few particles have signals of S and Mo. Interestingly, we find a large ratio of  $MoS_2$  nanoflakes vertically aligned to the CNT surface (Figure 2.3 (c)). The measured thickness of a selected  $MoS_2$  nanoflake is ~ 55 nm.



Figure 2. 3 Microstructures and elemental analysis of  $MoS_2$  on CNT sheet. (a)-(d) SEM images of  $MoS_2$  grown on raw CNT sheet. (e)-(h) SEM images of  $MoS_2$  grown on plasma treated CNT sheet. (i) Elemental analysis of raw CTN sheet/  $MoS_2$ . (j) Elemental analysis of activated CTN sheet/  $MoS_2$ .

In contrast, a large quantity of  $MoS_2$  nanoflakes are uniformly deposited on the substrate of plasma treated CNT sheet, as illustrated in Figure 2.3 (e). It is of interest to note that  $MoS_2$  nanoflakes are successfully grafted on almost every individual CNTs and CNT bundles (Figure 2.3 (f)-(h)). The measured thicknesses of nanoflakes are in the range of 10-14 nm. It is reasonable to observe the thickness differences between  $MoS_2$  sheets on

different CNT substrates, as the incoming  $MoS_2$  molecules prefer to land on the asnucleated  $MoS_2$  islands instead of on inert CNT walls. With the treatment, anchoring sites on the CNT sheet for  $MoS_2$  molecules are enhanced significantly. As a result, MoS2 sheets are distributed more uniformly with thinner features. EDS result in Figure 2.3 (j) not only identifies nanoflakes are  $MoS_2$  but also confirms the uniform dispersion of these nanoflakes, as overall intensities of S and Mo signals are homogenously distributed on the entire substrate.  $MoS_2$  signal from the EDS spectrum of plasma treated sample is much more striking than that of raw CNT sheet/  $MoS_2$  (Figure 2.4), suggesting the significantly improved concentration of  $MoS_2$ .



Figure 2. 4 EDX spectra of raw CNT sheet-MoS<sub>2</sub> and plasma treated CNT sheet-MoS<sub>2</sub>.

Besides of de-bundling effect, another important change of the CNT sheet before and after plasma treatment is captured by Raman scattering.  $I_D / I_G$  (intensity of the D peak/ intensity of the G peak) characterizes the ratio of disorders of  $sp^2$  carbon network of crystalline carbon materials<sup>14</sup>. As shown in Figure 2.5 (a),  $I_D / I_G$  of CNT sheet increases from 0.06 to 0.58 after the plasma activation. The enhancement is attributed to the

introduction of oxidized species such as carboxyl and hydroxyl groups<sup>15</sup>. These functional groups make CNT surfaces more suitable to nucleate and immobilize external nano-species. In our study, de-bundling and oxidization from plasma treatment are the two key points that enable us to obtain high-concentration MoS<sub>2</sub> nanoflakes uniformly grown on the CNT sheet.

Figure 2.5 (b) presents the overall comparison of Raman scattering between raw CNT sheet deposited with MoS<sub>2</sub> ( $R_{CNT}$ - MoS<sub>2</sub>) and plasma treated CNT sheet deposited with MoS<sub>2</sub> ( $P_{CNT}$ - MoS<sub>2</sub>). Intensity of MoS<sub>2</sub> in the sample  $R_{CNT}$ - MoS<sub>2</sub> are barely identified (Figure 2.5 (c)), indicating very low concentration of MoS<sub>2</sub> grown successfully. In contrast, peaks at 373 cm<sup>-1</sup> and 402 cm<sup>-1</sup> of  $P_{CNT}$ - MoS<sub>2</sub> are ascribed to  $E_{2g}^1$  and  $A_{1g}$  modes of MoS<sub>2</sub>, respectively, which, once again, identify the existence of MoS<sub>2</sub>. Moreover, the signal intensity scattered from MoS<sub>2</sub> is comparable to the G peak of the CNT, elucidating the existence of extremely high concentration of MoS<sub>2</sub>. Interestingly, the ratios of I<sub>D</sub> / I<sub>G</sub> in the two samples are both enhanced after the grating of MoS<sub>2</sub> (Figure 2.5 (e)), which implies the deposition of MoS<sub>2</sub> on CNT support may degrade the crystallinity or integrity of the nanotubes.



Figure 2. 5 Raman spectroscopy characterizations of (a) D and G peaks of raw sheet and plasma activated sheet. (b) Raw sheet/MoS<sub>2</sub> ( $R_{CNT}$ - MoS<sub>2</sub>) and plasma activated sheet/MoS<sub>2</sub> ( $P_{CNT}$ - MoS<sub>2</sub>) in the frequency 100- 2400 cm<sup>-1</sup> (c) 280- 440 cm<sup>-1</sup> (d) 1050- 1800 cm<sup>-1</sup>.

Figure 2.6 (a) shows the representative TEM image of  $MoS_2$  nanoflakes deposited on the plasma treated CNT support, from which we observe two hexagonal  $MoS_2$ nanoflakes with lateral sizes of 135 nm and 157 nm, respectively. They are highly transparent under 200 keV electron beam because of the ultrathin nature. The edge as well as the corner of ultrathin  $MoS_2$  nanoflake can be observed from Figure 2.6 (b). A highmagnification image in Figure 2.6 (c) presents the atomic structure of  $MoS_2$  and the insert is its corresponding fast Fourier transformation (FFT) pattern. The typical six-fold symmetry is confirmed from the FFT, suggesting a single-crystalline nanoflake<sup>16</sup>. More importantly, the yellow dotted line highlights CNT-MoS<sub>2</sub> interface at the atomic scale.  $MoS_2$  is not perfectly straight but has a few layers impinged into CNT walls, which may cause the degradation of CNT crystallinity. This likely drives the increase of  $I_D / I_G$  as discussed in Figure 2.5 (d). Further analysis in Figure 2.6 (d) shows the lattice distance of  $MoS_2$  is 2.7 Å, which can be assigned to  $MoS_2\{100\}$  planes and corresponds to the six-fold spots of the FFT pattern<sup>17</sup>.



Figure 2. 6 TEM study of CNT-MoS2 interfaces. (a) Representative TEM image of  $P_{CNT}$ -MoS<sub>2</sub>. (b) MoS<sub>2</sub> edge and corner. (c) High-resolution TEM image of CNT-MoS<sub>2</sub> interface. (d) Lattice distance of MoS<sub>2</sub> captured by HRTEM.

# 2.4 Conclusion

In summary, we report a facile strategy to uniformly grow  $MoS_2$  nanoflakes on plasma activated CNT sheet. We demonstrate that plasma induced de-bundling as well as defective sites on the CNT sheet are two critical keys to obtain high-concentration and uniform  $MoS_2$  nanoflakes on the assembly of CNTs. This work may shed light on the controllable synthesis of  $MoS_2$ /carbon nanocomposites as well as their related applications in battery fields. 2.5 References

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17. Huang, X.; Zeng, Z.; Bao, S.; Wang, M.; Qi, X.; Fan, Z.; Zhang, H., Solutionphase epitaxial growth of noble metal nanostructures on dispersible single-layer molybdenum disulfide nanosheets. *Nature communications* 2013, *4* (1), 1-8. Chapter 3 One-pot Self-templated Growth of Gold Nano-frames For Enhanced SERS Performance

### Abstract

As one of the representative metallic hollow nanostructures, Au nanoframes have shown fascinating properties such as strong localized surface plasmon resonance (LSPR) associated with the emerging application as surface-enhanced Raman scattering (SERS) sensors. In this study, it is demonstrated, for the first time, a facile one-pot synthetic approach for hollow Au nanoframes by directly etching Au nanoplates, i.e. so-called self-templates. A novel growth mechanism has been revealed that involves a synergistic function of Ag and Br ions. The presence of  $Ag^+$  lead to observed self-limiting of Au film thickness whereas Au {111} facets are preferentially attacked by the presence of  $Br^-$  in the reaction ambient. Moreover, graphene is introduced to prevent/minimize aggregation during the formation of Au nanoframes. Combined simulation and experimental studies show that the hybrid platform made of graphene/Au nanoframes is capable of detecting analytes at concentration levels down to  $10^{-9}$  M by using the surface-enhanced Raman spectroscopy (SERS) technique.

### 3.1 Introduction

Hollow nanostructures have attracted increasing research interest recently owing to their unique characteristics<sup>1,2</sup>, such as large surface area, high fraction of active surface sites, low density, well-controlled dimensions, and multi-functional properties as nano-reactors<sup>3,4</sup>, catalysis<sup>5,6</sup>, and sensors<sup>7</sup>. A variety of strategies have been developed to fabricate hollow nanostructures, including top-down lithography<sup>8</sup>, nanowires self-coiling<sup>9</sup>, self-assembly processing<sup>10</sup>, galvanic replacement<sup>11-13</sup>, and chemical etching<sup>14-20</sup>. Particularly, to create metallic hollow structures on the nanoscale, chemical etching is the most commonly employed approach<sup>21</sup>. It has been witnessed that ultrathin Au nanoframes<sup>14,15</sup>, shape-controllable nanoframes of Au<sup>16</sup>, Pt<sup>17</sup>, and Pd<sup>18</sup>, complex trimetallic nanoframes such as Au@Pt-Ni<sup>19</sup> and AgAu@Pt<sup>20</sup> core-shell alloys, hierarchical 4H/fcc Ru nanotubes<sup>22</sup>, and so forth, are successfully synthesized via the template carving Such hollow metal nanostructures present tunable optical properties<sup>16</sup>, enhanced plasmonic and catalystic activies<sup>23</sup>, robust surface-enhanced Raman scattering (SERS) performance<sup>24</sup>, and increased electrocatalytic effects<sup>25</sup>.

In the science of carving metal nanocrystals, chemical etching by corrosive chemicals such as hydrogen peroxide<sup>15</sup> and oxidative etching by the conjunction of oxygen and halogen (Cl<sup>-</sup>, Br<sup>-</sup>, and I<sup>-</sup>) ions<sup>18</sup>, <sup>26-28</sup> are realized at the expense of inside solid blocks, namely, sacrifice templates within the enclosed frames. These synthetic strategies are widely adopted to fabricate Au nanoframes, which is one of the representative hollow metal nanostructures that are studied intensively. For instance, Xue and co-authors obtained ultrathin Au nanoframes with tailorable widths in the range of 1.8 - 6.0 nm using Ag nanoplates as the sacrifice template<sup>14</sup>. Nie et al. reported the synthesis of circular and

triangular Au nanorings by excavating templates of Ag nanodisks and Ag nanoprisms, respectively<sup>16</sup>. In these approaches, however, materials of the sacrificial templates are different from the nanoframes, which cost extra steps for the preparation, causing difficulty in the controlling of the chemical synthesis due to the homogeneous nucleation of the coating material<sup>2</sup>. Moreover, these methods usually involve multiple cycles of metal deposition and etching<sup>29</sup>, which may significantly reduce the reproducibility and make it difficult to scale-up.

Herein, we demonstrate a one-pot strategy for the growth of Au nanoframes, Relying on the etching of Au nanoplates, i.e. self-templates. Well-defined Au nanoplates with tunable thickness are first synthesized. In the same solution, Au nanoplates can be gradually etched to Au nanoframes over time. With a systematically study, we reveal both Ag<sup>+</sup> and Br<sup>-</sup> ions are of critical importance to the synthesis. Moreover, we employ graphene template to reduce the aggregation of Au nanoframes. As a consequence, we show the hybrid platform made of graphene/Au nanoframes provides excellent SERS performance.

#### 3.2 Experimental

#### 3.2.1 Chemicals and Materials

Graphene template (AO-2) was purchased from Graphene Supermarket (Calverton, New York). Gold chloride hydrate (HAuCl<sub>4</sub> ·xH<sub>2</sub>O, 99.999%), silver nitrate (AgNO<sub>3</sub>,  $\geq$ 99.0 %), potassium bromide (KBr,  $\geq$  99.0 %), and potassium iodide (KI,  $\geq$  99.0 %) were purchased from Sigma-Aldrich, USA. L-ascorbic acid (> 99.0%) and Rhodamine 6G (99 %) were obtained from Fisher Scientific, USA.

# 3.2.2 Experimental Procedures

One-step synthesis of Au nanoframes: In a standard procedure, we first prepared 18 mL deionized (DI) water in a glass vial. Then, 85.5  $\mu$ M HAuCl<sub>4</sub> (60  $\mu$ L, 28.5 mM), 85.5  $\mu$ M KBr (50  $\mu$ L, 30 mM), and 125  $\mu$ M AgNO<sub>3</sub> (250  $\mu$ L, 10 mM) were added to the vial. After the mixture was homogeneous mixed with gently shaking, 500  $\mu$ L ascorbic acid (0.4 mM) was then added to the solution. The glass vial was capped with a PTFE lined cap and placed in an oven at 80 °C for 1 hours for the hydrothermal reduction. Subsequently, the glass vial was taken out from the oven and aged at room temperature for different time, i.e. 4 hrs, 24 hrs, 48 hrs and 96 hrs, to enable the etching.

In order to verify the roles of Ag<sup>+</sup> ions, more samples were prepared with 5  $\mu$ M, 20  $\mu$ M, 50  $\mu$ M, 400  $\mu$ M, and 800  $\mu$ M AgNO<sub>3</sub>, while keeping the concentrations of all other chemicals as well as the reaction condition the same. At the same time, to explore Br<sup>-</sup> ions function, a series of samples were prepared by only changing KBr concentrations as 0  $\mu$ L, 10  $\mu$ L, 200  $\mu$ L, and 500  $\mu$ L. Moreover, NaI providing an equivalent amount of I<sup>-</sup> ions was used to replace KBr in the synthesis, in order to comprehensively demonstrate the role of the Br<sup>-</sup> ions.

One-pot synthesis of Au nanoframes on graphene template: To improve the aggregation, we introduced graphene multilayers as the template for Au nanoframes. The synthesis was slightly modified from the previous section. First, 1 mg graphene powder is added into the glass vial containing 18 mL DI water and the mixture was ultrasonicated for 1 minute to obtain well-dispersed suspension. Thereafter, the same amount of chemicals, i.e. HAuCl<sub>4</sub> (60  $\mu$ L, 28.5 mM), KBr (50  $\mu$ L, 30 mM), AgNO<sub>3</sub> (250  $\mu$ L,10 mM), and ascorbic acid (500  $\mu$ L, 0.4 mM) were added to the vial. The mixture was capped closely

and then maintained in the oven at 80 °C for 1 hr. After that, the sample was taken out from the oven and placed at room temperature for 48 hr. Then the synthesized hybrid material was washed and collected by centrifugation for further characterization and other usage.

### 3.2.3 Characterizations

Surface-enhanced Raman scattering (SERS) performance: The synthesized graphene/Au nanoframes suspension was deposited on a silicon (Si) substrate with [100] orientation using a pipette. The Si supported hybrid was dried in the vacuum oven at 100  $^{\circ}$ C for 2 hours. The deposition process was repeated by 3 times to assure the substrate was totally covered with graphene/Au nanoframes. To test the SERS performance, a highly fluorescent dye material, Rhodamine 6G (R6G) aqueous solution (20 µL) was deposited on five prepared substrates, at the concentrations of 10<sup>-5</sup> M, 10<sup>-6</sup> M, 10<sup>-7</sup> M, 10<sup>-8</sup> M and 10<sup>-9</sup> M, respectively. The substrates were transferred to Raman equipment before the R6G solution was fully dried. Spots deposited with Au nanoframes can be observed from the optical microscope. When performing Raman scattering, we particularly check these spots and record the Raman signals of the analyte. Raman scattering spectra were obtained with a Renishaw In-Via Raman system with the laser length of 785 nm, expose time of 10 s, laser power of 1mW, objective lens magnification of 50 ×, and a grating spacing of 1200 *l*/mm.

Simulation of Plasmonic Properties of Au nanoframes: The near field electromagnetic field distribution was simulated via finite difference time domain simulation (Lumerical FDTD simulations). The perfect match layer (PML) was set as the absorptive boundary. The mesh size was set at  $0.1 \text{ nm} \times 0.1 \text{ nm} \times 0.1 \text{ nm}$  within the simulation area (plate/frame area). The edge lengths of Au nanoplates were appointed as 20 nm and 34.64 nm for

hexagonal and triangular shapes, respectively. Both of them had the thickness of 10 nm. The nanoframes had the same lateral sizes with the frame width of 10 nm. The dielectric function of Au was obtained from Johnson and Christy. Graphene was used as the substrate and its dielectric constant was taken from Falkovsky. The simulation result is the E/E0 distribution along the top surface of the plate/frame structure.

Characterizations: Top and tilt view scanning electron microscopy (SEM) images of samples with different *t* were obtained with FEI Nova NanoSEM 230. The system was coupled with an energy dispersive X-ray (SEM/EDX) detector (ThermoFisher Scientific), which can monitor spatial elemental maps of Au nanoplates, Au nanoframes and graphene supported Au nanoframes. Transmission electron microscopy (TEM) and high-resolution transmission electron microscopy (HRTEM) images were captured with FEI Titan S/TEM system at 300 keV. The X-ray diffraction (XRD) measurements were performed on a Bruker D1 diffractometer using Cu K $\alpha$ 1 radiation.

# 3.3 Result and Discussion

The details of the one-pot synthetic approach are revealed in the supporting information. We first investigate Au templates prior to the etching. At the time t = 4 h, well-defined two-dimensional Au nanoplates are obtained with a yield > 90%, as shown in Figure 3.1(a). It is of interest to note such a high yield of Au nanoplates is extremely difficult to obtain with seedless growth method unless a proper etchant is involved<sup>30</sup>. These nanoplates typically possess triangular and hexagonal shapes with a lateral size of a few microns. Figure 3.1(b) shows a 3-D simulation of two representative Au nanoplates, indicative of relatively smooth and flat surfaces. In addition,  $2\theta$ - $\omega$  XRD scan of the sample

in Figure 3.1(a) shows the synthesized nanoplates have face-centered cubic (FCC) structure, owing to the peaks observed at  $2\theta = 38.3^{\circ}$ , 44.5°, 82.1°, which belong to Au (111), Au (200), and Au (222), respectively (Figure 3.1(c))<sup>31,32</sup>. It is noteworthy to emphasize the presence of strikingly intensive Au (111) peak. The relative intensity of (111) / (200) is more than 20, indicating the Au nanoplates are predominantly (111) orientated<sup>33</sup>. Inset of Figure 3.1(c) is the elemental mapping from EDX measurement, confirming the nanoplates are pure gold. TEM selected area diffraction pattern of a selected Au nanoplate in Figure 3.1(d) shows diffraction spots of 1/3 (422) and (220) with six-fold symmetry, suggesting an FCC structure, which is consistent with the XRD result. This also indicates the diffracted facet, i.e. the basal plane of the plate is Au (111), which is the predominant crystallographic orientation. The measured lattice distance is 1.44 Å, corresponding to the plane distance of Au (220).<sup>34</sup>



Figure 3. 1 (a) A representative SEM image of Au nanoplates assembly. (b) 3-D simulation profile of two representative Au nanoplates. Simulation is obtained from the SEM image.

(c) X-ray diffraction (XRD)  $2\theta$ - $\omega$  scan obtained from the Au nanoplates. Inset is the element mapping from EDX. (d) Selected area electron diffraction pattern (SAED) of a nanoplate. (e) The atomic structure on the surface of Au nanoplate obtained from HRTEM.

We now focus on the transformation of Au nanoplates to Au nanoframes *via* the etching process. To begin, we point out that it is the Br<sup>-</sup> ion that plays the role of the etchant in our system. It is well established that halogen ions mainly impose two functional roles in the growth of noble metal nanocrystals: the capping agent<sup>35</sup> as well as the etchant<sup>28</sup>. As the capping agent, they would preferentially adsorb on certain facets of the nanocrystals and promote growth in directions other than the passivated facets. In our study, a proper amount of Br<sup>-</sup> ions can facilitate the formation of Au nanoplates, for example, where belt-like Au crystallites (Figure S3.1(a)) transform to Au nanoplates (Figure S3.1(b)) when using 15  $\mu$ M KBr during the synthesis. On the other hand, the etching effect from overdosed Br- ions causes severe damage to the Au nanoplates when more than 750  $\mu$ M KBr is used (Figures S3.2(c) and (d)).

It is thus very important to employ the appropriate concentration of Br<sup>-</sup> ions to obtain desirable Au nanoframes. Accordingly, not only high-yield Au nanoplates (the self-templates) can be synthesized (Figure3. 1(a)), but also well-defined hollow structured Au nanoframes (Figures 3.2(a)-(c)) are produced after etching at room temperature for 48 hrs. It is not surprising to note that the etching effect predominantly depends on the time *t*, as shown in Figure 3.2(d)-(g) with the t = 4, 24, 48, and 96 hrs, respectively. In particular, we observe flat Au nanoplates transform to concave nanoplates first (Figure 3.2(e) and Figure S3.2, 24 hrs), followed by the formation of center-hollowed Au nanoframes (Figure 3.2(f), 48 hrs). Interestingly, the thinner edge can be carved out, resulting in open frames when t extends to more than 96 hrs (Figure 3.2(g)). Selected examples of the etched Au nanoplates are shown in Figure 3.2(h), which suggests the etching is quite uniform at the very beginning but gradually becomes inhomogeneous. This is likely due to the non-uniformity of the thickness of Au nanoplates.



Figure 3. 2 (a)-(c) Representative SEM images of Au nanoframes from low magnification to high magnification. Au nanoplates at the etching time t (d) 4 hrs, (e) 24 hrs, (f) 48 hrs,

and (g) 96 hrs. (h) Four examples of etched Au nanoplates at different stages. (i) Dependence of Au nanoplates' thickness and the etching efficiency on Ag+ concentration.

In our synthesis, the Ag<sup>+</sup> ion plays a critical role as well in the formation of Au nanoframes. First, we observe well-defined Au nanoplates can be only synthesized with an appropriate amount of silver ions. Neither less than 20  $\mu$ M nor more than 800  $\mu$ M AgNO<sub>3</sub>, can be used to produce high-quality Au nanoplates (Figure S3.3). Second, the thickness of Au nanoplates is tunable from ~100 nm to ~10 nm, simply by varying the concentration of AgNO<sub>3</sub> in the solution from 20  $\mu$ M to 400  $\mu$ M (Figure 3.3). In addition, we find the etching efficiency is closely related to the original thickness of Au nanoplates, which increases with the decreasing thickness (Figure 3.2(i)). As a result, Au nanoplates with the thickness > 50 nm can be only partly carved in 48 hrs, leaving a pit hole with irregular shapes on the surface. In contrast, a complete hole or hollow structure from the surface to the bottom is created on nanoplates with the thickness < 50 nm can be geried. However, when the thickness is less than 10 nm, the etching becomes more difficult to control and a large number of open nanoframes are obtained.



Figure 3. 3 Thickness tunability of Au nanoplates by varying concentrations of AgNO<sub>3</sub>. (a) 20  $\mu$ M, (b) 50  $\mu$ M, (c) 125  $\mu$ M and (d) 400  $\mu$ M of AgNO<sub>3</sub>. The aging time of all samples shown here is 4 hrs.

According to the above results, we propose a growth mechanism of Au nanoframes that involves a synergistic effect of  $Ag^+$  and  $Br^-$  ions. In particular,  $Br^-$  ion serves as the etchant and dictates the etching rate, while  $Ag^+$  ion determines the total amount of gold template that is going to be etched up. Figure 3.4 is the schematic illustration of the etching mechanism that highlights individual roles of  $Ag^+$  and  $Br^-$  ions. In the shape-controlled synthesis of nanoparticles, it is widely known that  $Ag^+$  ions, even at a trace amount, can boost the growth of anisotropic gold nanostructures in a high-yield<sup>36-38</sup>, including Au nanorods, bipyramids, nanoprisms, ultrathin nanowires, and crystallographic high-index nanoparticles.

Thus far, however, there has been very limited report to reveal  $Ag^+$  ions are capable of tailoring the thicknesses of Au nanoplates. The origin of the anisotropy of gold nanostructures directed by  $Ag^+$  ions has been an elusive and controversial topic ever since it was discovered<sup>39</sup>. Silver ions ( $Ag^+$ ), complex pair of Ag-Br, and elemental silver ( $Ag^0$ ) that passivate certain facets of gold nanocrystals and promote the growth of anisotropic structures are all considered the plausible mechanisms<sup>40</sup>. Here, it is likely that free  $Ag^+$  ions or AgBr could adsorb on Au (111) facets, prevent gold atoms adding to (111) surface, and thus reduce the thickness of the plates accordingly. Nevertheless, the explicit mechanism of  $Ag^+$  ion that controls the thickness of Au nanoplates requires an in-depth study, which is out the scope of this research.



Figure 3. 4 Schematic illustration of the growth mechanism of Au nanoframes.  $Ag^+$  and  $Br^-$  ions show a synergistic function to obtain Au nanoframes.

To further investigate the etching effect, we choose to use Au nanoplates with the thickness of about 20 nm, which is realized by using 125  $\mu$ M AgNO<sub>3</sub>, as their etching is thorough but in a mild and controllable way. We find the etching initiates from the surface of Au nanoplates. As identified from Figures 3.1(d) and (e), crystallographic basal plane of the surface is Au (111). Previous studies show that Br ions preferentially adsorb on Au (111), causing the unbalanced aggregation of Br ions on Au (111) facets<sup>35</sup>. Meanwhile, it is well established that Au (111) facets possess planar defects such as stacking faults and twin boundaries, which result in the occurrence of the theoretically forbidden 1/3 (422) spots (Figure 3.1(d))<sup>34</sup>. Moreover, Au nanocrystals are observed to be unstable in the solution containing the halogen ions<sup>30,35</sup>. Thus, it is reasonable to speculate that Br<sup>-</sup> would favourably react with gold atoms on (111) planes given all the factors mentioned above. Aided by the oxygen solvated in the solution, the reaction can be listed as the following<sup>27</sup>:

$$O_2 + 2H_2O + 4e^- \rightarrow 4OH^-$$
  
Au + 4Br  $\rightarrow$  AuBr<sub>4</sub> + 3e<sup>-</sup>

Moreover, the edges of the nanoplate composed of (111) and (110) are occasionally found to be etched (Figure S3.4), indicating the etching process is not completely selective but preferable to a great extent. To verify such a role of  $Br^-$  in the etching, we use the same amount of I<sup>-</sup> to replace  $Br^-$  in the solution. As a consequence, we observe over-etched Au nanoplates (Figure S3.5) as I<sup>-</sup> is stronger than  $Br^-$  as an etchant. The synthesized Au nanoplates and Au nanoframes usually show aggregated morphology, which is undesirable for many applications, such as surface-enhanced Raman scattering (SERS). For instance, it is demonstrated that aggregation of Au nanoframes reduces, rather than increases, SERS effect due to the compromise of the inter- and intra-particle plasmonic fields<sup>41</sup>. To address this concern, we introduce the graphene multilayers in the solution and ultimately obtain graphene/Au nanoframes with improved dispersion, as shown in Figures 3.5(a) and (b). Besides, characterizations from XRD, SEM/EDX, and high-resolution TEM (Figures 3.5(c)-(i)) show that there are no changes in the crystalline orientations, atomic structures or element compositions of Au nanoframes when they are synthesized with graphene template. The improved dispersion is attributed to the presence of anchoring sites on graphene multilayers, which could immobilize Au nanocrystals from the nucleation stage and make the ultimate dispersion more homogenous<sup>35,42</sup>.



Figure 3. 5 Characterization of graphene grafted Au nanoframes. (a) Low and (b) high magnification SEM images of graphene supported Au nanoframes. Graphene and Au nanoframes are highlighted by white and red arrows, respectively in (b),

showing the improved dispersion of Au nanoframes. (c)  $2\theta \cdot \omega$  XRD scan of graphene/Au nanoframes. (d)-(f) EDX mapping of graphene/Au nanoframes with elements of C and Au captured. (g) A representative TEM image of Au nanoframes on graphene and (g) its selected area electron diffraction (SAED). (i) Atomic structure of Au nanoframes on the graphene template.

We now evaluate the SERS efficiency of the hybrid graphene/Au nanoframes platform. Prior to the measurement, we simulate the localized electrical field distributions around individual Au nanoframe in both triangular and hexagonal shapes. As illustrated in Figure 3.6(a), hot-spot regions of non-hollow Au nanoplates are limited in the tip area with extremely low intensity. On the other hand, hot-spot density is significantly higher in Au nanoframes as they are accumulated in areas along all the edges benefit from the hollow feature. As a consequence, the electrical field simulated from the frame is magnified significantly, which can enhance the Raman intensity by approximately  $(E/E_0)^4$ . We perform the Raman mapping of two different substrates (graphene/Au nanoframes vs. graphene/ Au nanoplates) decorated with 10<sup>-8</sup> M R6G and find dramatically high density of hotspots from the platform with Au nanoframes, as shown in Figures 3.6(b) and (c). Therefore, we are able to detect the analyte, i.e. R6G at a concentration as low as 10<sup>-9</sup> M using graphene/Au nanoframes as the SERS substrate (Figure 3.6(d)). Five signature peaks of the analyte at  $10^{-9}$  M can be still clearly detectable at 1182, 1306, 1362, 1506, and 1647  $\text{cm}^{-1}$  <sup>31</sup>. This SERS platform demonstrates outstanding reproducibility as well, as we are capable of probing R6G at 10<sup>-9</sup> M from arbitrarily selected nine spots containing Au nanoframes on the substrate (Figure 3.6(e)).



Figure 3. 6 (a) Simulation of localized electrical field enhancement of Au nanoplates and Au nanoframes. Raman hotspot mapping from the hybrid substrates of (b) graphene /Au nanoplates and (c) graphene/Au nanoframes with 10<sup>-8</sup> M analyte. (d) SERS detection of R6G at the concentration from 10<sup>-5</sup> to 10<sup>-9</sup> M using graphene/Au nanoframes hybrid platform. (e) Evaluation of SERS reproducibility from nine arbitrarily selected spots.

When aggregated Au nanoframes (without graphene) are used as the SERS substrate, the detection limit of the analyte is  $10^{-7}$  M (Figure S3.6). Their SERS sensitivity is quite close to those gained from Au nanoplates substrate.<sup>31</sup> Therefore, the detection limit from graphene/Au nanoframes is enhanced by two orders of magnitude compared with those substrates made of Au nanoframes and nanoplates. It has been reported graphene layers have shown multiple effects on SERS activities of plasmonic nanocrystals. First, there is the charge transfer between plasmonic Au nanostructures and graphene substrates<sup>33,43</sup>. The charge transfer contributes to the chemical enhancement (CM). Although weak, CM is one of the mechanisms responsible for the enhanced Raman scattering. The other mechanism is the electromagnetic enhancement (EM) that leads to the significant enhancement of localized surface plasmonic resonance (LSPR)<sup>31,44</sup>. Second, molecules of the analyte can have a more well-defined interaction with the hybrid SERS substrate due to the presence of oxygen-containing functional species on the graphene template. This will also improve the SERS performance<sup>45</sup>. Third, the benefit in terms of plasmonic activities from the unique hollow nanostructures disappears when Au nanoframes aggregates into an assembled entity, because there is actually no difference between aggregated Au nanoframes and Au nanoplates. It is demonstrated that aggregation of Au nanoframes reduces, rather than increases, SERS effect due to the compromise of the inter- and intra-particle plasmonic fields<sup>41</sup>. The dispersion of Au nanoframes is dramatically improved when graphene is introduced to the synthesis in this work. Accordingly, graphene template greatly enhances the SERS performance of Au nanoframes in our study.

# **3.4 Conclusions**

In summary, we demonstrate a facile one-pot route to synthesize Au nanoframes by the etching of self-templates. We demonstrate pivotal roles of both  $Ag^+$  and  $Br^-$  ions in obtaining Au nanoframes. In addition, graphene template effectively improves the aggregation issue of Au nanoframes and the hybridized graphene/Au nanoframes substrate exhibits high SERS sensitivity that is capable of probing the analyte down to  $10^{-9}$  M. Our approach greatly simplifies the tedious synthesis processes of Au nanoframes and proves the feasibility of synthesizing metallic nanoframes by etching self-templated nanoplates/nanosheets. This work may bring new insights to the synthesis of dimensioncontrollable Au nanocrystals and inspire the development of novel graphene-based nanocomposites.

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# 3.6 Supplementary Information



Figure S3. 1 Samples synthesized with different amount of KBr. (a) 0  $\mu$ L. (b) 15  $\mu$ M. (c) 300  $\mu$ M. (d) 750  $\mu$ M.



Figure S3. 2 EDX mapping of Au nanoplates aged at 24 hrs.



Figure S3. 3 Au nanocrystals synthesized in the standard condition (80 °C 1hr and aged for 4 hrs at room temperature) with different amount of AgNO<sub>3</sub> solution. (a) 10  $\mu$ M AgNO<sub>3</sub>. (b) 800  $\mu$ M AgNO<sub>3</sub>.



Figure S3. 4 Examples of Au nanoplates with the etched edges synthesized in the standard condition and aged for 24 hrs.



Figure S3. 5 Au nanocrystals synthesized in the standard method and aged for 48 hrs except that the same amount of KI was used to replace KBr. Au nanoplates was not observed owing the stronger etching effect from  $I^-$  ions than  $Br^-$  ions.



Figure S3. 6 SERS substrate made from the aggregated Au nanoframes without graphene template and the Raman scattering of R6G using this substrate.

Chapter 4 Single-Layer MoS<sub>2</sub> Based Au Gap Plasmon Enhanced Photoelectrochemical Cells for Efficient Water Splitting

# Abstract

Metal nanoparticles and metal film coupling (Gap Plasmon) have been shown to possess distinguished plasmonic properties compared with other nano-metallic structures, the enhanced electric field is extremely localized within the narrow gap region when the gap thickness reaches 10 nm or less. The field enhancement is known to be inversely proportional to the gap thickness. Hence reducing the gap thickness is important to achieve the highest possible field enhancement. In this work, it is demonstrated for the first time that using single-layer  $MoS_2$  as a well-defined nanospacer between Au-nanoparticles (AuNPs) and Au-film, which could offer an extremely high localized electric field enhancement within the gap. The MoS<sub>2</sub> Raman intensity with the Surface Enhanced Raman Scattering (SERS) enhancement factor (EF) up to  $5 \times 10^6$  is obtained from the MoS<sub>2</sub>-Au gap plasmon system. A 5-fold increase in the photocurrent is obtained from the MoS<sub>2</sub>-Au gap plasmon as the working electrode compared to that from bare MoS<sub>2</sub> prepared under the same condition. Compared with individual metal nanoparticles commonly used to enhance thin-film photocatalytic process, gap-plasmon could theoretically produce 8 orders of magnitude higher SERS EF and precise control the hot spot location to superimpose where ultrathin materials locate thus using the higher incident energy available.

# 4.1 Introduction

 $H_2$  has long been known to be one of the highest energy density fuels<sup>1</sup> with water as the only by product. Fuel cell technology has been actively pursued as an environmentally friendly power source for automobiles<sup>2</sup>. The promise of the technology is hindered by the lacking of a sustainable way of producing  $H_2$ . Solar water-splitting via photoelectrochemical cells (PEC)<sup>3</sup> is the most promising approach with intense research effort being dedicated to it<sup>4-6</sup>. One bottleneck commonly identified for solar water splitting is the limited charge carrier diffusion length inside the semiconductor, typically below 20 nm<sup>7</sup>, electron and hole pairs that created inside the bulk semiconductor can hardly diffuse out to the surface which in result cannot contribute to the overall water splitting. Therefore, thin layers of semiconductors are very much preferred to increase the yield of hydrogen production<sup>7</sup>.

Another important factor for a semiconductor to be an effective PEC cell is its bandgap. The energy bands of the semiconductor (conduction band and valence band) need to be properly aligned with the redox potentials of water, with the energy bandgap ( $E_g$ ) being larger than the potential difference between the reduction of H<sub>2</sub> and the production of O<sub>2</sub>. In the meantime, the  $E_g$  needs to be sufficiently small in order to absorb much of the solar radiation available<sup>3</sup>. The preferred  $E_g$  range for PEC cell is typically reported to be between 1.6 ev to 2.4 ev<sup>8-9</sup>. To date, only a handful of semiconductor materials meet the requirement with the energy bands that being properly aligned with the redox potentials of water, such as BiVO<sub>4</sub><sup>10</sup>, GaP<sup>11</sup>, TiO<sub>2</sub><sup>12-14</sup>, Cu<sub>2</sub>O<sup>15</sup>, single-layer MoS<sub>2</sub><sup>16-17</sup> etc. Among them, single-layer MoS<sub>2</sub> possesses the correct chemical over-potential while with the bulk form MoS<sub>2</sub> having  $E_g$  values too small for the intended purpose<sup>18</sup>. Moreover, MoS<sub>2</sub> is an earth-abundant,

chemically stable, non-toxic material, which make it a perfect candidate for the photoelectrodes. However, single-layer MoS<sub>2</sub> still facing the problem of limited optical absorption cross-section due to the extreme thin property. Herein, it is demonstrated for the first time single-layer MoS<sub>2</sub> packed between AuNPs and Au film (referred MoS<sub>2</sub>-Au gap plasmon system), where extremely high electric field (hot spots) locates, similar gap plasmon structure studies have already demonstrated that gap thickness here plays a significant role on the electric field enhancement inside gap, the field enhancement is known to be inversely proportional to the gap thickness<sup>19</sup>. Hence reducing the gap thickness is important to achieve the highest possible field enhancement with gap thickness decreases. Considering the single-layer  $MoS_2$  is the thinnest  $MoS_2$  space layer that can be ever prepared, it is expected that using it as a well-defined sub-nanometer space layer will significantly enhance electric field within the gap. Though some studies have been done for plasmon enhanced single-layer MoS<sub>2</sub> water splitting from individual AuNPs<sup>16, 21</sup>, no previous report have been focused on the enhanced  $MoS_2$  photocatalysis from the gap plasmon system that using single layer MoS<sub>2</sub> as the space layer inside the gap where much higher field intensity is available.

This gap plasmon structure, is not limited to  $MoS_2$  but can easily be extended to other single-layer Van der Waals materials such as graphene, black phosphorus, hBN. Within the sub-nanometer thick gap plasmon structure, those significantly enhanced field intensity could enables a variety of interesting applications, such as surface-enhanced Raman scattering<sup>22</sup>.

## 4.2 Experimental

# 4.2.1 Sample Preparation.

Synthesis of Single-layer MoS<sub>2</sub>: We used reported chemical vapor deposition methods to grow MoS<sub>2</sub> on 300nm SiO<sub>2</sub> sheet substrate<sup>23</sup>. The equipment setup is schematically shown in Figure S4.1, the typical single-layer MoS<sub>2</sub> morphology and singlelayer MoS<sub>2</sub> surface converge shown in Figure S4.1 inset. In brief, 10 mg MoO<sub>3</sub> powder and 500 mg S powder were used as Mo and S precursors, respectively. The 300nm SiO<sub>2</sub> substrate was put up-side down facing towards the MoO<sub>3</sub> powder localized at the furnace centre. The substrate temperature is set to be 850 °C, while sulfur powder was put at the temperature window, where the temperature was set as 200 °C. The total growth time was 15min. During the growth, Ar was employed as the carrier gas with a flow rate of 100 sccm. The thickness of as-grown MoS<sub>2</sub> space layers before the transfer is measured via atomic force microscopy (AFM) images which is shown as shown in figure S4.2. The obtained average thickness is about 0.7nm, which corresponded with single layer MoS<sub>2</sub> thickness<sup>20</sup>.

Synthesis of Au nanoparticles: Gold chloride hydrate (HAuCl4·xH2O, 99.999% purity) and polyethyleneimine (PEI, branched, average molecular weight Mw ~25,000) are purchased from Sigma-Aldrich, USA. The chemicals are used without any further treatment or purification. The similar method reported in chapter 3 is adopted to synthesis AuNPs<sup>24</sup>. First, 220 mg of aqueous PEI solution (27 mg/mL) is mixed with 20 mL of deionized water in a glass vial. Then, 35  $\mu$ L of 40 mM HAuCl<sub>4</sub> solution is added and gently shaken for 1 minute. The glass vial is then capped and kept at 80 oC in the oven for hydrothermal reduction of gold precursor for 1 hr. PEI serves as the reducing agent as well as the stabilizer for Au nanoparticles in this standard synthesis procedure. At t = 0, the

solution appears light yellow and after 15 min, the color of the solution gradually changed to wine red. After completion of the reaction, the product, i.e. AuNPs, are used later.

Fabrication the MoS<sub>2</sub>-Au Gap Plasmon System: The 50 nm gold film was deposited on Si substrate via very slow electron-beam evaporate rate (0.1A s<sup>-1</sup>) after depositing 5 nm thick chromium as adhesion layer. The as-grown single-layer MoS<sub>2</sub> film was transferred to gold film using the reported PMMA-assisted method<sup>25</sup>. The as synthesized Au nanoparticle was deposited on as prepared MoS<sub>2</sub>-Au film substrate using a pipette. The hybrid substrate was dried in the vacuum oven at 100 °C for 2 hours. The deposition process was repeated by 3 times to assure the substrate was covered with enough AuNPs, the schematic of the MoS<sub>2</sub>-Au gap plasmon system is shown in figure 4.1 (d). The control experiments were carried out with 4 slightly different configurations shown in figure 4.1 to better elucidate the performance of MoS<sub>2</sub>-Au gap plasmon system



Figure 4. 1 (a)  $MoS_2$  on  $SiO_2$  substrate. (b)  $MoS_2$  on Au film substrate. (c)  $MoS_2$  on  $SiO_2$  substrate with AuNPs. (d)  $MoS_2$ -Au gap plasmon system

Surface Enhanced Raman Scattering Measurements: Raman spectroscopy analysis was obtained with a Renishaw In-Via Raman system, the of 633 nm laser length with power of 1mW was used, the objective lens magnification of 50x and a grating spacing of 1200 *l*/mm was used, MoS<sub>2</sub> Raman spectra was acquired from 4 different configurations as the control experiments.

Photocurrent Measurements: A xenon light source was used; optical power is measured to around 300 mW/cm<sup>2</sup>. To measure the photocurrent, a home-built electrochemical system with 3 electrodes configurations was used. MoS<sub>2</sub>-Au gap plasmon system (config 4) is used as the working electrode, Ag/AgCl electrode as reference electrode and Pt as counter electrode. All potentials are reported against standard Ag/AgCl voltage. The electrolyte was aqueous 0.1 M KH<sub>2</sub>PO<sub>4</sub> buffer at pH 7 with no additive. The voltammetry sweep was manual performed from 0 to 1.2 V at the step rate of 0.1v/sec immediately after the sample was fully immersed into the electrolyte. Two control experiments are carried out with MoS<sub>2</sub> on flat Au (config 2) and flat Au as the working electrode with rest setup remain the same.

FDTD Simulations: The near field electromagnetic field distribution was simulated via finite difference time domain simulation (Lumerical FDTD simulations). The perfect match layer (PML) was set as the absorptive boundary. The mesh size was set at 0.1 nm  $\times$  0.1 nm within the simulation area. The diameter of Au nanoparticles to was

appointed as 20 nm. the space layer of monolayer MoS<sub>2</sub> set as 1nm in thickness. The thickness of Au film underneath MoS<sub>2</sub> is set as 50 nm. The dielectric function of Au was obtained from Johnson and Christy<sup>26</sup>. The MoS<sub>2</sub> dielectric constant was taken from Laturia<sup>27</sup>. The simulation result shows the  $E/E_0$  distribution along the top surface of MoS<sub>2</sub> as well as the side view of the structure. The extremely strong near-field produced in the proposed MoS<sub>2</sub>-Au gap plasmon system (config 4) shown in Figure S4.4 which is due to couplings between the AuNPs and Au films<sup>28</sup>. The theoretical local SERS enhancement factor ( $E/E_0$ )<sup>4</sup> reaches up to 10<sup>8</sup>. Config 3 also shows some low near-field enhancement ( $E/E_0$ )<sup>4</sup> up to 10<sup>3</sup> mainly from AuNPs itself<sup>29</sup>, while config 1 and 2 didn't show enhanced field intensity near surface due to the non-metallic or flat surface under.

Characterizations: Top view scanning electron microscopy (SEM) images of samples were obtained with FEI Nova NanoSEM 230. Transmission electron microscopy (TEM) and high-resolution transmission electron microscopy (HRTEM) images were captured with FEI Titan S/TEM system at 60 keV. Raman spectra are measured using Renishaw inVia microscope under 633 nm laser excitation. The laser power used is ~1 mW. The grating used is 1800 l/mm and the objective lens used is 50x.

## 4.3 Results and Discussion

#### 4.3.1 MoS<sub>2</sub>-Au gap plasmon system

The  $MoS_2$ -Au gap plasmon system (Figure 1(d)) is made of a three-layer structure consisting of 20nm diameter AuNPs sitting on top of near continuous single-layer  $MoS_2$ 

with another 50nm thick Au film on bottom Si substrate, low and high-magnification of SEM images are utilized to observe the surface morphology of different stages related to preparing MoS<sub>2</sub>-Au gap plasmon system, three stages include the structure of very slowly deposited (0.1 A/sec) 50nm Au film (Figure 4.2 (a), (d)), transferred single-layer MoS<sub>2</sub> on 50nm Au film(Figure 4.2 (b), (e)), and finally MoS<sub>2</sub>-Au gap plasmon system after AuNPs deposition(Figure 4.2 (c), (f)). The bare Au film is observed in figure 4.2 (a), (d), it is worth to mention that some surface roughness on the Au film is observed under highmagnification of SEM image shown in figure 4.2 (d), it is expect that surface roughness is contributed to the later slight enhanced Raman intensity from  $MoS_2$ . It is observed in figure 4.2 (b) that  $MoS_2$  surface coverage on Au film to be more than 80% which is similar as the coverage before transfer. It is also clearly observed in figure 4.2 (d) the edge of singlelayer MoS<sub>2</sub> resulting from the different surface morphology between MoS<sub>2</sub> and the uncovered Au film. The similar roughed morphology of Au film under MoS<sub>2</sub> can also be clearly observed due to the extremely thin property of the single-layer MoS<sub>2</sub>. It is found in figure 4.2 (f) that the AuNPs deposited on  $MoS_2$  could be clearly visualized especially on the MoS<sub>2</sub> films as MoS<sub>2</sub> is a well-proven diffusion barrier<sup>16</sup> which could avoid AuNPs to migrate through. The AuNPs can barely be visualized from Au film without placing singlelayer  $MoS_2$  film in between due to the metal migration. High coverage of single-layer  $MoS_2$ on Au film significantly avoids subsequent deposited AuNPs migrating to the bottom Au film. The thickness of as-prepared  $MoS_2$  space layers is measured via atomic force microscopy (AFM) images which shown in figure S4.2. The thickness obtained is about 0.7 nm, corresponded with single layer MoS<sub>2</sub> profile<sup>20</sup>. It confirms that the thinnest available MoS<sub>2</sub> as space layer in gap plasmon system is successfully being prepared. This geometrically proves that single-layer  $MoS_2$  could be effectively used as a well-defined spacer layer inside gap plasmon system. The AuNPs size and coverage is clearly visualized from HRTEM image in figure S4.3 (a). It is also obtained that the average AuNPs diameter is around 20 nm and the average separation distance between nearby AuNPs is around 7 nm, the corresponding statistic histogram of AuNPs size and separation distance is shown in figure S4.3 (b) and (c). The very high coverage of AuNPs is very beneficial to the entire  $MoS_2$  surface to make full use of incident light above across the entire  $MoS_2$  film.



Figure 4. 2 Low-magnification of SEM images of structures used in (a) Au film (b)  $MoS_2$ on Au film substrate. (c)  $MoS_2$ -Au gap plasmon system. High-magnification of SEM images of structures (d) a zoomed-in image of Au film. (e) a zoomed-in image of  $MoS_2$  on Au film substrate. (f) a zoomed-in image of  $MoS_2$ -Au gap plasmon system.

It is shown in Figure 4.3 (a) the high-resolution TEM image of single-layer MoS<sub>2</sub> films after AuNPs were synthesized and deposited on. It is observed that many AuNPs packed on top of MoS<sub>2</sub> film. The lattice spacing of 2.7A with six-fold symmetry corresponds to the MoS<sub>2</sub> (100) plane<sup>30</sup> and the clear inter-planar spacing of 2.4A corresponds to Au (111) plane<sup>24</sup>, respectively. These results are in agreement with the selected area electron diffraction (SEAD) pattern shown in Figure 4.3 (b). The diffraction spots can be divided into two categories, which are assigned to AuNPs and MoS<sub>2</sub> film correspond. The diffraction ring of Au (111) planes is observed as shown in yellow circle. The bright spots circled in white square can be indexed to (100) diffraction spots of MoS<sub>2</sub><sup>30</sup>, respectively. Furthermore, the element distribution of Au-MoS<sub>2</sub> composite were investigated by energy-dispersive X-ray spectrometry (EDX) shown in Figure 4.3 c-f. The element Mo, S and Au distribution can be observed. It shows the co-existence of Mo, S and Au and relative uniform distribution along the selective area.



Figure 4. 3 (a) HRTEM image of AuNPs directly deposited on  $MoS_2$ , showing the distinguishable lattice parameter between Au and  $MoS_2$ . (b) the selected area electron diffraction (SEAD) pattern of  $MoS_2$ -Au compound (c) selected area TEM image (d) EDX mapping of Mo element (e)EDX mapping of S element (f) EDX mapping of Au element

# 4.3.2 Enhanced SERS in MoS<sub>2</sub>-Au plasmon gap System

The field enhancement is first studied by detecting Raman intensity from  $MoS_2$  from 4 different configurations shown in figure 4.1. Since Raman spectroscopy is a widely used tool to analysis the optical response towards monolayer  $MoS_2^{31}$  and its SERS intensity directly correlates with the strength of the localized electrical field that is enhanced by the

surface plasmon from neighboring AuNPs or Au films<sup>32-34</sup>, so MoS<sub>2</sub> Raman intensity is the best benchmark for investigating the near electrical field enhancement. The 633nm laser is used to for all the SERS measurement to match the plasmon resonance of the as-prepared AuNPs. As shown in figure 4.4, config 1 represent the typical Raman spectra for monolayer MoS<sub>2</sub> layer on SiO<sub>2</sub>/Si with its 2 major peaks, E<sub>2g</sub> (385cm<sup>-1</sup>) and A<sub>1g</sub>(405cm<sup>-1</sup>) can be identified but show a very low signal-to-noise ratio due to the typical low Raman activity, the E<sub>2g</sub> peak represents the vibration mode of Mo S atoms in the basal plane, whereas the A<sub>1g</sub> peak represents the S atoms out of plane vibration modes. Here, 20cm<sup>-1</sup> Raman shift difference between A1g and E2g confirms the single-layer MoS2 property<sup>31</sup>, this finding corresponds with our AFM result before. Another 460 cm<sup>-1</sup> can also be distinguishable which is original from the LA and LA' photon at the M point<sup>35</sup>, apart from these, no other MoS<sub>2</sub> related peak can be observed from config 1. Raman intensity of the MoS<sub>2</sub> from Config 2, 3, 4 all show various enhanced intensity compared to normal Raman spectra from config 1, it can be clearly seen that S/N ratios of Raman spectra from config 3 and 4 are much higher than those from config 1 and 2 mostly due to surface plasmon coupling between neighboring nano units, this finding corresponds with our FDTD simulation results before (Figure S4.4). The spectra from config 4 shows the most enriched spectral information as expected due to the extremely strong surface plasmon resonance from gap plasmon system. The newly exposed peaks at 528 cm<sup>-1</sup>, 573 cm<sup>-1</sup> and 645 cm<sup>-1</sup> are assigned to MoS<sub>2</sub> E<sub>g</sub>(M)+TA'(M), A<sub>1g</sub>(M)+TA'(M), and A<sub>1g</sub>(M)+LA'(M)<sup>35</sup>, respectively. The spectra from config 3 shows some enhancement due to the partilly resonace from neighboring AuNPs as was confirmed simulation result in figure S4.4c. The sepctra from config 2 shows slight enhancement possibly due to the factor the as-prepared Au film is

not atomiclly flat which renders some surface plasmoc resonance towards the MoS<sub>2</sub> layer on top.

To quantitatively describe the spectra enhacement, SERS enhancemcemnt factors (EF) are adopted as a figure of merit to represent the boosted raman signal, we employ the following equation:

$$EF = (I_{SERS \times} A_{Raman}) / (I_{Raman \times} A_{hotspot})$$

where  $I_{SERS}$  and  $I_{Raman}$  represent the SERS and normal Raman intensities of the analyte acquired under identical conditions.  $A_{Raman}$  is the area of the laser excitation spot on flat region. We used 50× lens with NA 0.55 for 633 nm laser, and the calculated laser spot size is ~1400nm<sup>35</sup>. We take  $A_{Raman}$  as a 1400×1400 nm area and take  $A_{hotspot}$  as a 10×10 nm area. The SERS EF for  $A_{1g}$  band of  $MoS_2$  is calculated to be ~  $4x10^5$  and ~ $9x10^5$  for config 2 and config 3 as to normal Raman from config 1, respectively. Remarkably, For the MoS<sub>2</sub>-Au gap plasmon system (config 4), the typical SERS EF is about  $\sim 5x10^6$  compared with normal Raman from config 1, this enhanced ratio is significantly higher than just simple linear superposition from two separate structure with MoS<sub>2</sub>, i.e. Au film under MoS<sub>2</sub> (SPP from the Au film) and AuNPs on MoS<sub>2</sub> (LSP from the AuNPs). The significant enhancement is promoted due to the strong plasmonic couplings between the AuNP-Au film coupling with only 0.7nm thick single-layer  $MoS_2$  as the space layer. It should be mentioned, to our best knowledge, it is first time surface enhanced Raman study on gap plasmon system with single-layer  $MoS_2$  as the spacer layer that is ever been reported and it is one of highest SERS EF ever reported so far for plasmon-enhanced  $MoS_2^{35-37}$ .



Figure 4. 4 (a) Normal Raman spectra of  $MoS_2$  from config 1. (b) SERS spectra of  $MoS_2$  on Au film from config 2. (c) SERS spectra of  $MoS_2$  with AuNPs deposited on top from config 3. (d)  $MoS_2$ -Au gap plasmon system from config 4. The excitation wavelength is 633 nm.

## 4.3.3 Photocurrent Measurements

It is illustrated in Figure 4.5 (a), the electron hole pairs are generated when shining the light towards working electrode. The generated electrons were transferred from the Au-MoS<sub>2</sub> working electrode through the Au substrate towards the counter electrode Pt, where hydrogen ion was reduced to generate hydrogen gas. On the working electrode, the left

behind holes on the Au-MoS<sub>2</sub> surface oxidize water molecules to generate oxygen. The linear sweep voltammograms show that the working electrode made from config 2 (pure MoS<sub>2</sub> on Au film) exhibited a photocurrent of 110  $\mu$ A/cm<sup>2</sup> at 0.8 V under visible light illumination (Figure 4.5 (b)) similar to the previous reported monolayer  $MoS_2$  on FTO substrate photocurrent values<sup>38</sup>. The same measurement is carried out using config 4 as working electrode (MoS<sub>2</sub>-Au gap plasmon system), with the gap plasmon resonance within the MoS<sub>2</sub> the photocurrent can boost the photocurrent up to 590  $\mu$ A/cm<sup>2</sup> at 0.8 V (Figure 4.5 (b)), a 5-fold increase in photocurrent has been achieved for gap plasmon system. Here the gap plasmon system plays a few important roles here, first, AuNPs coupling with the Au film separated by only single-layer  $MoS_2$  distance generates an extremely high localized electric field(Figure S4.4 (d)), which facilitate more electron-hole pairs generation and separation in  $MoS_2^{38}$ . Second, Au films severs as the electron sink to trap the generated electron move away from  $MoS_2$  and to the counter electrode<sup>39</sup>, third, AuNPs as a hot plasmon-excited electrons provider for the direct water reduction towards the counter electrode<sup>39-42</sup>. The switched ON and OFF photocurrent of working electrodes made from config 2 and 4 is shown in Figure 4.5 (c). It is observed that when light is on, the photocurrent increases immediately, and the on photocurrent is significantly larger than when the light is off, the on and off ratio reaches to about 100 when using config 4 (gap plasmon) as the working electrode. These results show that single-layer MoS<sub>2</sub> thin film exhibits extreme sensitive towards photo activity, the normalized-on photocurrent ratio from config 4 and config 2 is around 5, which is consistent with the photocurrent from config 4 and config 2 when potential vs Ag/AgCl is at 0.8v in Figure 4.5 (b).



Figure 4. 5 (a) Schematic of MoS<sub>2</sub>-Au gap plasmon system as photoanode and commercial Pt as photocathode respectively (b) Linear sweep voltammograms for PECs with different working electrodes, config 2 as working electrodes shows in blue, config 4 as working electrodes shows in red. (c) Normalized photocurrent cycles with potential at 0.8 V for PEC with different working electrodes config 2 as working electrodes shows in blue, config 4 as working electrodes shows in blue, config 4 as working electrodes config 2 as working electrodes shows in blue, config 4 as working electrodes config 2 as working electrodes shows in blue, config 4 as working electrodes config 2 as working electrodes shows in blue, config 4 as working electrodes config 2 as working electrodes shows in blue, config 4 as working electrodes config 2 as working electrodes shows in blue, config 4 as working electrodes shows in red.

## 4.4 Conclusion

In conclusion, single-layer MoS<sub>2</sub>-Au gap plasmon system with the enhanced photoelectrochemical water-splitting activity is demonstrated for the first time. The use of sub-nanometer thick MoS<sub>2</sub> as space layer enables gap plasmonic coupling throughout the MoS<sub>2</sub> with the significantly higher field enhanced. A 5-fold increase in photocurrent is achieved through MoS<sub>2</sub>-Au gap plasmon as the working electrode. With the systematically study, it is revealed that single-layer MoS<sub>2</sub>-Au gap plasmon system can theoretically generate up to 8 orders of magnitudes SERS EF locally confirmed via FDTD simulation. The MoS<sub>2</sub> Raman intensity is tremendously enhanced from the MoS<sub>2</sub>-Au gap plasmon system, the SERS EF up to  $5x10^6$  is obtained, one of the highest values that is ever reported to plasmon-enhanced MoS<sub>2</sub> Raman study. Our results provide a plausible solution to enhance the light trapping and photocatalytic response towards the ultrathin MoS<sub>2</sub>, solving the problem that trade-off the small diffusion length of carrier and the low light absorption.

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# 4.6 Supplementary Information



Figure S4. 1 Schematic illustration of chemical vapor deposition method for  $MoS_2$  grow on SiO<sub>2</sub> substrate. Inset SEM image shows the nearly continuous monolayer  $MoS_2$  (up to 90% coverage)



Figure S4. 2 AFM image of the obtained from CVD growth  $MoS_2$  samples on Si/SiO<sub>2</sub> substrate. The height profile obtains from the edge  $MoS_2$  crystal. The height measured 0.7nm corresponds with the monolayer  $MoS_2$  profiles.



Figure S4. 3 (a)high-resolution TEM image of as-fabricated AuNPs deposited on TEM grids. (b) AuNPs diameter distribution histogram, (c) AuNPs interparticle distance distribution histogram.



Figure S4. 4 Electrical Field enhancement ( $E/E_0$ ) distribution side view and top view from 4 different configurations
## 5.1 Conclusions

This dissertation demonstrated the feasibility of the gap plasmon enhanced monolayer  $MoS_2$  based efficient and sustainable photocatalytic water splitting platform. The following is a summary of the dissertation work presented:

In Chapter 2, we introduce a novel strategy of growing MoS<sub>2</sub> nanoflakes on CNT sheet that shows significance in the following two aspects: first, the direct growth of large quantity MoS<sub>2</sub> nanoflakes on carbon nanotube sheet is realized. To the best of our knowledge, this is the first time to grow large quantities MoS<sub>2</sub> nanoflakes directly on the substrate of the CNT sheet. Second, plasma-activated CNT sheet is of critical importance to immobilize MoS<sub>2</sub> nanoflakes. We reveal that plasma treatment of the CNT sheet results in de-bundling and activation of carbon nanotube bundles, which promotes the deposition of MoS<sub>2</sub> nanoflakes.

In Chapter 3, we present a novel way to synthesis Au nanoframes, which have been shown to possess distinguished plasmonic properties comparing to other metal nanocrystal geometry. However, current approaches to obtain these nanoframes usually involve multistep treatment. In this work, we demonstrate a facile one-step growth-etching approach of self-templated hollow Au nanoframes, which shows significant merits in the following aspects: First, A facile one-step route to synthesize Au nanoframes based on the etching of self-templates is realized for the first time. High-quality Au nanoplates with tunable thickness are first synthesized, from which Au nanoframes can be evolved with enough etching time. We find a synergistic effect from  $Ag^+$  and  $Br^-$  ions as the key factors to the formation of predominantly Au nanoframes. Second, the dispersion issue of Au nanoframes is significantly improved using graphene as a template. The improved dispersion is attributed to the presence of anchoring sites on graphene, which could immobilize Au nanocrystals from the nucleation stage and make the ultimate dispersion more homogenous. Third. Graphene/Au nanoframe hybrid platform shows great potential in surface enhanced Raman spectroscopy (SERS) that is highly relevant to the application such as biosensing. The hybrid platform is capable of rendering single molecule sensitivity proven via both simulation and experiment results.

In Chapter 4, we present using single-layer  $MoS_2$  as the space layer inside gapplasmon system to achieved enhanced photoelectrochemical water splitting. The enhancement in this work is due to plasmon resonances between AuNPs and Au films. A 5-fold increase in the photocurrent was achieved across  $MoS_2$ -Au gap plasmon working electrode. With the systematic study, we reveal that  $MoS_2$ -Au gap plasmon system can theoretically increase nearly 8 orders of magnitude's enhancement factor via FDTD simulation. The  $MoS_2$  Raman intensity is tremendously enhanced from the  $MoS_2$ -Au gap plasmon system, the SERS EF up to  $5x10^6$  is obtained, one of the highest values that is ever reported to plasmon-enhanced  $MoS_2$  Raman study. Our results provide a plausible solution to enhance the light trapping and photocatalytic response towards the ultrathin  $MoS_2$ , solving the problem that trade-off the small diffusion length of carrier and the low light absorption.

## 5.2 Future Work

For the platform to be sustainable in nature as well as the broadband resonance within the visible light region, replacing precious metal Au employed for the plasmonic substrate with other more environmentally rich elements like Ag, Cu, Al is investigated. The preliminary study using FDTD modeling of replacing bottom Au film with sustainable Ag, Cu, Al is simulated shown in figure 5.1, the reason we only replace the bottom film instead of the AuNPs is that the Ag, Cu, Al is not corrosion-resistant, by placing metal under continuous MoS<sub>2</sub> film, MoS<sub>2</sub> as a well-known diffusion barrier can prevent O<sub>2</sub> penetrate through to corrode metal film underneath<sup>1</sup>.

It is observed that AuNPs and Au films coupling only resonance while incident wavelength at 650 nm and 596 nm, significantly less hot spot and hot spot density are observed while further decrease the incident wavelength in the simulation result in figure 5.1 row one. The result corresponds with the quality value of Au shown in figure 1.4. It is observed that Cu films show strong resonance with AuNPs not only at 650 nm and 596 nm as well as at 513 nm, which means Cu have broader band resonance with AuNPs, however, further decrease incident wavelength to 450 nm, the resonance also decrease dramatically observed from simulation results in row two. When using Al or Ag as bottom metal film in gap plasmon system, we find strong resonance on all four wavelengths that was simulated. Strong hot spots and high hot spots density can be clearly observed under 450 nm wavelength shown in figure 5.1



Figure 5. 1 EM field distributions for the AuNP coupling with various metal films separated by a 1 nm MoS<sub>2</sub> spacer at 450 nm, 513 nm, 596nm and 785 nm, respectively.

In order to better compare which metal film resonate with AuNPs more efficiently, we need to consider both hot spot and hot spot density, here the quality factor is defined as equation below

$$Q = \iint_{x,y=-40}^{x,y=40} E^2 dx dy$$

This self-defined quality factor represents the total stored energy of the resonator along the surface of the  $MoS_2$  under the cross-section area of the Au nanoparticles to by considering both hot spot and hot spot density, the data is extracted from simulation results.

Q value comparison with different metal under different wavelength is obtained in figure 5.2. From the plot we can draw the conclusion that Ag film can store the most available light energy from 450nm to 650nm, which proves Ag film resonates more with AuNPs in a broader band resonance between 450nm to 650nm. This finding corresponds with figure 1.4 that Ag shows the highest Q factor from 400nm and up.



Figure 5. 2 Quality factor as a function of wavelength for different metal films and AuNPs coupling.

However, the substitution may bring the potential problem of Ag corrosion. Here, MoS<sub>2</sub> has already been proven to be a very effective diffusion barrier for O<sub>2</sub> penetration, the strong sulfur-oxygen interaction makes it difficult for oxygen molecule to penetrate through even for single-layer MoS<sub>2</sub>, the diffusion energy barrier for O<sub>2</sub> molecular to penetrate through is calculated to be 13.94 eV<sup>1</sup>, this diffusion energy barrier is significantly larger than the reported barrier for suspended graphene (5.98 eV) which is widely known as a good corrosion protective barrier<sup>1-3</sup>. Ag plasmonic substrate will be completely protected if we have continuously monolayer MoS<sub>2</sub> covered on top to prevent potential corrosion or oxidation. This event may raise another potential issue that is a method to fabricate continuously exclusive single-layer MoS<sub>2</sub>, current growth recipe can hardly control for the perfect exclusive continuous single-layer MoS<sub>2</sub> growth, the typical result is either near continuously single-layer MoS<sub>2</sub> film or sometimes slightly more than one single-layer MoS<sub>2</sub>, it is relative hard to control the perfect timing to grow exact the 100% coverage single-layer MoS<sub>2</sub>. We will address this issue by first over-grow MoS<sub>2</sub> to be slightly more than one single layer. We propose to use an energy-band gap dependent, selflimiting etching process for produce exact single layer MoS<sub>2</sub>. A selective photo-activated etching recipe need to be developed. The physics basis for the selective photo-activation is the narrower  $E_g$  of multiple layer  $MoS_2$  compared to that of single layer  $MoS_2$ . By employing an etchant that can barely etches  $MoS_2$  in the absence of light, we expect that illuminating the surface of MoS2 with patches of different thickness as obtained from typical chemical vapor deposition (CVD) growths<sup>4</sup> using photons with energy precisely controlled to be just below the  $E_g$  of single layer  $MoS_2$  could provide the etch rate differential needed to obtain single layer MoS<sub>2</sub> covering the entire substrate surface.

## 5.3 References

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