

UC Berkeley

UC Berkeley Previously Published Works

Title

AGNs on the Move: A Search for Off-nuclear AGNs from Recoiling Supermassive Black Holes and Ongoing Galaxy Mergers with the Zwicky Transient Facility

Permalink

<https://escholarship.org/uc/item/3wr3q5cb>

Journal

The Astrophysical Journal, 913(2)

ISSN

0004-637X

Authors

Ward, Charlotte
Gezari, Suvi
Frederick, Sara
[et al.](#)

Publication Date

2021-06-01

DOI

10.3847/1538-4357/abf246

Peer reviewed

AGN on the move: A search for off-nuclear AGN from recoiling SMBHs and ongoing galaxy mergers with the Zwicky Transient Facility

CHARLOTTE WARD,¹ SUVI GEZARI,^{2,1} SARA FREDERICK,¹ ERICA HAMMERSTEIN,¹ PETER NUGENT,^{3,4}
SJOERT VAN VELZEN,^{1,5} ANDREW DRAKE,⁶ ABIGAIL GARCÍA-PÉREZ,⁷ IMMACULATE OYOO,⁸ ERIC C. BELLM,⁹
DMITRY A. DUEV,⁶ MATTHEW J. GRAHAM,⁶ MANSI M. KASLIWAL,⁶ STEPHEN KAYE,¹⁰ ASHISH A. MAHABAL,^{11,12}
FRANK J. MASCI,¹³ BEN RUSHOLME,¹³ MAAYANE T. SOUMAGNAC,^{3,14} AND LIN YAN¹⁰

¹*Department of Astronomy, University of Maryland, College Park, MD 20742, USA*

²*Space Telescope Science Institute, 3700 San Martin Dr., Baltimore, MD 21218, USA*

³*Lawrence Berkeley National Laboratory, 1 Cyclotron Road, Berkeley, CA 94720, USA*

⁴*Department of Astronomy, University of California, Berkeley, Berkeley, CA 94720, USA*

⁵*Center for Cosmology and Particle Physics, New York University, NY 10003*

⁶*Division of Physics, Mathematics, and Astronomy, California Institute of Technology, Pasadena, CA 91125, USA*

⁷*Instituto Nacional de Astrofísica, Óptica y Electrónica, Tonantzintla, Puebla 72840, México*

⁸*Prince George's Community College, Largo, MD 20774, USA*

⁹*DIRAC Institute, Department of Astronomy, University of Washington, 3910 15th Avenue NE, Seattle, WA 98195, USA*

¹⁰*Caltech Optical Observatories, California Institute of Technology, Pasadena, CA 91125, USA*

¹¹*Division of Physics, Mathematics and Astronomy, California Institute of Technology, Pasadena, CA 91125, USA*

¹²*Center for Data Driven Discovery, California Institute of Technology, Pasadena, CA 91125, USA*

¹³*IPAC, California Institute of Technology, 1200 E. California Blvd, Pasadena, CA 91125, USA*

¹⁴*Department of Particle Physics and Astrophysics, Weizmann Institute of Science, Rehovot 76100, Israel*

(Accepted to ApJ on March 24, 2021)

ABSTRACT

A supermassive black hole (SMBH) ejected from the potential well of its host galaxy via gravitational wave recoil carries important information about the mass ratio and spin alignment of the pre-merger SMBH binary. Such a recoiling SMBH may be detectable as an active galactic nucleus (AGN) broad line region offset by up to 10 kpc from a disturbed host galaxy. We describe a novel methodology using forward modeling with **The Tractor** to search for such offset AGN in a sample of 5493 optically variable AGN detected with the Zwicky Transient Facility (ZTF). We present the discovery of 9 AGN which may be spatially offset from their host galaxies and are candidates for recoiling SMBHs. Five of these offset AGN exhibit double-peaked broad Balmer lines which may arise from unobscured accretion disk emission and four show radio emission indicative of a relativistic jet. The fraction of double-peaked emitters in our spatially offset AGN sample is significantly larger than the 16% double-peaked emitter fraction observed for ZTF AGN overall. In our sample of variable AGN we also identified 52 merging galaxies, including a new spectroscopically confirmed dual AGN. Finally, we detected the dramatic rebrightening of SDSS1133, a previously discovered variable object and recoiling SMBH candidate, in ZTF. The flare was accompanied by the re-emergence of strong P-Cygni line features indicating that it may be an outbursting luminous blue variable star.

1. INTRODUCTION

Supermassive black holes (SMBHs) reside in the center of most galaxies (Kormendy & Richstone 1995; Ferrarese & Ford 2005). Galaxy growth via hierarchical

mergers therefore results in the formation of SMBH binaries. The time taken for these SMBH binaries to merge depends on the nature of their host galaxies. While binaries in gas-poor galaxies may stall at 1 pc separations (e.g. Milosavljević & Merritt 2001), SMBH binaries in gas-rich environments may merge on timescales of $10^6 - 10^7$ years (Escala et al. 2005).

A consequence of SMBH mergers in gas-rich environments may be the gravitational wave recoil of coalesced SMBHs after merger. In this process, the asymmetric emission of gravitational waves during SMBH coalescence imparts momentum to the coalesced black hole, ejecting it from the central potential well to wander about the galaxy halo for 10^6 to 10^9 years (Volonteri & Perna 2005; Campanelli et al. 2007; Loeb 2007; Volonteri & Madau 2008; Blecha & Loeb 2008). The ‘recoiling black hole’ is expected to carry broad line gas with it and continue to undergo regulated accretion, allowing it to be observable as an active galactic nucleus (AGN) spatially offset from the center of its host galaxy (Blecha & Loeb 2008). Simulations by Volonteri & Madau (2008), for example, show that an AGN with a 500 km/s kick velocity could still be accreting and observable as an off-center quasar at 30 kpc from its host center.

Other observable signatures of a recoil event include evidence of recent galaxy merging activity. Since the recoiling AGN may continue to accrete for a 10^6 year timescale after recoil (Blecha & Loeb 2008), tidal structures may still be visible while the recoiling SMBH is active and detectable. The ejected AGN could also leave behind a trail of feedback evidence in the form of enhanced H α emission leading to the galaxy centre (Loeb 2007).

Simulations show that the recoil velocity and maximum host-AGN spatial offset of the recoiling SMBH depends on the mass ratio and spin alignment of the black hole binary prior to merger (Campanelli et al. 2007; Blecha et al. 2016). Binaries with perfectly aligned spins can produce a maximum recoil velocity of only 200 km s $^{-1}$. A binary with misaligned spins can produce kicks up to 5000 km s $^{-1}$. A population of recoiling black holes would therefore provide strong constraints on the distribution of masses and spins in SMBH binaries. This would, in turn, inform simulations of SMBH spin alignment based on torques in the circumbinary gas disk (Bogdanović et al. 2007; Lodato & Facchini 2013) and stellar interactions during inspiral (Berczik et al. 2006).

A confirmed sample of recoiling SMBHs could be used to test the predictions of numerical relativity simulations on the fraction of massive black holes ejected via recoil from their host galaxies at different redshifts and the effects of this on observed black hole occupation fractions and the $M_{\text{BH}} - \sigma_*$ relation (Volonteri 2007; Volonteri et al. 2010; Blecha et al. 2011). Such a sample would also allow us to study the effect of displaced AGN feedback on the evolution of merger remnants, such as the expected increase in star formation rates and lengthening of the starburst phase (Blecha et al. 2011).

Despite the many motivations to search for recoiling SMBHs, only a few good candidates have been found to date. One such object is the radio-loud QSO 3C 186 which has an 11 kpc spatial offset from its host galaxy and a -2140 ± 390 km s $^{-1}$ velocity offset between the broad and narrow emission lines (Chiaberge et al. 2017). The tidal features of the host galaxy indicate recent merger activity. Integral field spectroscopy was performed to study the complex kinematics and determine if the velocity offset could result from a peculiar outflow (Chiaberge et al. 2018). The results were consistent with the recoiling SMBH scenario but final confirmation will require both James Webb Space Telescope IFU spectroscopy to map the H β region with 0.1” resolution and deep imaging from HST to rule out the presence of a second low mass galaxy (Chiaberge et al. 2018).

SMBH recoil may also be the origin of the variable object SDSS J113323.97+550415.8 (SDSS1133) (Koss et al. 2014). This object is 800 pc from the center of a low redshift dwarf galaxy and has displayed AGN-like stochastic variability over > 63 years. However, AGN-like variability can be mimicked by long-lived stellar transients (Burke et al. 2020) and giant stellar outbursts or supernovae such as SN2009ip (Maza et al. 2009) and UGC 2773 OT2009-1 (Boles 2009). Spectra of SDSS1133 show the presence of blue-shifted Balmer absorption lines and [Fe II] $\lambda 7155$ and [Ca II] $\lambda\lambda 7291, 7324$ forbidden emission lines which are highly unusual for an AGN. It is therefore possible that SDSS1133 is a luminous blue variable (LBV) star continuing to demonstrate non-terminal outbursts (Koss et al. 2014).

The lack of many recoiling SMBH candidates has motivated a number of systematic searches for offset AGN using different techniques. A Gaia analysis of a sample of low redshift, unobscured broad line AGN from SDSS showed that at least 99% were within 1 kpc of the host, 90% within 500 pc and 40% within 100 pc (Shen et al. 2019). This study used a technique called varstrometry (Hwang et al. 2020) to measure AGN-host spatial offsets via the astrometric jitter of the photocenter induced by the AGN flux variability, allowing them to rule out the existence of a substantial offset AGN population on 10 pc to 1 kpc scales at redshifts $0.3 < z < 0.8$.

From this study, it appears that unobscured, accreting, recoiling SMBHs with > 10 pc separations must be very rare at low redshifts, if they exist. This may be because SMBH spin alignment is always very efficient, inducing only small velocity recoils with small maximum separations. Recoiling black holes may also be more common in the early universe due to higher merger rates and lower galaxy masses. Cold gas inflow dur-

ing merger may increase the required escape velocity for many galaxies and gas drag could play a role in keeping recoiling black holes close to the galaxy center (Blecha et al. 2016). The level of accretion may be too small for recoiling SMBHs to be detectable and offset AGN may be frequently obscured by the gas environment induced by the merger (Shen et al. 2019).

While these results appear discouraging, another systematic search by Lena et al. (2014) undertook careful isophotal modeling of archival HST images of 14 nearby core elliptical galaxies and found that 10 of the 14 had small $\sim 1 - 10$ pc displacements between the AGN and host galaxy center, 6 of which were considered confident detections because the galaxy profiles were not asymmetric. 4 of the 6 galaxies showed alignment between the AGN-photocenter displacements and the radio jet axis. This correlation is predicted for gravitational recoil of SMBHs but may instead indicate that the spatial offset was induced by radio jet acceleration of the SMBH. This radio axis correlation would not be produced by interactions with massive perturbers or orbital motion prior to SMBH binary coalescence (Lena et al. 2014).

Other searches have used a multi-wavelength approach to successfully find offset AGN candidates. Skipper & Browne (2018) searched for radio-optical spatial offsets in a sample of 345 SDSS galaxies with nearby compact radio sources detected in the Cosmic-Lens All Sky Survey (CLASS) catalogue, finding 3 sources with offsets greater than 0.6 arcseconds. Condon et al. (2017) found one offset AGN candidate amongst 492 radio point sources from the NRAO VLA Sky Survey (VLASS; Condon et al. 1998) when crossmatched to extended sources in the Two Micron All-Sky Survey (2MASS) Extended Source Catalog (Jarrett et al. 2000).

A recent study by Reines et al. (2020) found that the majority of a sample of 13 radio AGN in dwarf galaxies were off-nuclear, likely because the lower escape velocities in dwarf galaxies make it easier for black holes to wander from the central potential. Kim et al. (2017) found a recoiling SMBH candidate in a systematic search for spatially offset X-ray AGN in a sample of 2542 sources with optical/near-infrared counterparts in archival HST images from the *Chandra* Source Catalogs - Sloan Digital Sky Survey Cross-match Catalog (Evans et al. 2010; Rots & Budavári 2011).

Even though these search strategies have yielded some recoiling SMBH candidates, there are many challenges to confirming the nature of these objects. Candidates with broad line gas at > 1000 km s $^{-1}$ velocities relative to narrow emission lines but no observable AGN-host spatial offset can often be explained by outflowing winds (Allen et al. 2015; Robinson et al. 2010), scattered broad

line emission from an SMBH binary (Robinson et al. 2010), asymmetric double-peaked emission from an elliptical accretion disk (Steinhardt et al. 2012), or two superposed AGN (Shields et al. 2009a,b).

For recoiling SMBH candidates with observable host-AGN spatial offsets, it must be ruled out that they are not QSOs with an additional undermassive host (Chiaberge et al. 2017). In this case, the true host galaxy may be on the lower end of the luminosity - SMBH mass scaling relation (McLure & Dunlop 2002) and very compact such that the extended galaxy emission around the QSO is not detectable, resulting in a false association with the brighter, offset companion galaxy in a merging system. This can occur when the AGN's host galaxy was tidally stripped as it merged with the larger galaxy (Bellovary et al. 2010) and has been proposed as the nature of the Jonker et al. (2010) recoiling SMBH candidate and the origin of a number of off-nuclear ultra-luminous X-ray sources such as HLX-1 (Farrell et al. 2009). Compact dwarf galaxies hosting SMBHs may be very common at low redshifts. Four ultra-compact dwarf galaxies with masses $M > 10^7 M_{\odot}$ in the Virgo and Fornax clusters have been shown to host SMBHs through analysis of their velocity dispersion and mass profiles (Seth et al. 2014; Ahn et al. 2017; Afanasiev et al. 2018; Ahn et al. 2018). These systems were likely produced through tidal stripping of a larger galaxy hosting an SMBH. Voggel et al. (2019) estimate that such stripped nuclei may host 8% to 32% of local SMBHs.

AGN in merging galaxies also have typical relative velocities of $10 - 400$ km s $^{-1}$ (Liu et al. 2018; Comerford et al. 2009; Comerford & Greene 2014) and these velocities are comparable to the predicted velocities of recoiling SMBHs from spin-aligned binaries. Comerford & Greene (2014) estimate that 4% to 8% of Type 2 AGNs are in galaxy mergers, so these systems may be quite common. Triple SMBH systems, in which a merger with a third galaxy occurs before the initial SMBH binary forms, can also be difficult to distinguish from recoiling SMBHs (Civano et al. 2010; Kalfountzou et al. 2017).

Offset AGN in merging galaxies are nonetheless important to find because they provide a way to study AGN fueling by galaxy merger triggered gas inflows (Surace et al. 1998; Canalizo & Stockton 2001; Treister et al. 2012). The increased incidence of galaxy mergers amongst X-ray selected AGN and increasing X-ray luminosity with decreasing AGN separation in dual AGN suggests that black hole accretion peaks during the merging process (Koss et al. 2012). Comerford & Greene (2014) also found that the fraction of AGN in galaxy mergers increases from 0.7% to 6% over the AGN bolometric luminosity range of $43 < \log(L_{bol})[\text{erg/s}] < 46$

which suggests that galaxy mergers trigger high luminosity AGNs. High-resolution hydrodynamical simulations by Van Wassenhove et al. (2012) predict that AGN triggering and the likelihood of dual AGN activity is strongest at $< 10 \text{ kpc}$ separations and that most merger-triggered AGN activity is non-simultaneous such that 90% of SMBHs in mergers appear as single or offset AGN instead dual AGN at $L_{\text{bol}} > 10^{44} \text{ erg s}^{-1}$ and separations $> 1 - 10 \text{ kpc}$. These predictions are supported by observations of the relative occurrence of offset vs dual AGN (e.g. Comerford et al. 2009). Discoveries of AGN in tidally stripped dwarf galaxies in mergers may also yield IMBH candidates which can be used to constrain models of BH seed formation in the early universe (e.g. Volonteri & Natarajan 2009; Reines & Comastri 2016).

Discoveries of AGN in mergers have occurred both serendipitously and in targeted searches. Binary AGN were found in X-ray imaging spectroscopy of ultraluminous infrared galaxy NGC 6240 (Komossa et al. 2003) and Mrk 739, a galaxy with two optically distinguishable bulges (Koss et al. 2011). A search for AGN companions to a sample of ultra-hard X-ray-selected AGNs from the all-sky Swift Burst Alert Telescope (BAT) survey with Chandra, XRT and XMM imaging combined with emission line diagnostics with SDSS and Gemini spectroscopy revealed 16 dual AGN (Koss et al. 2012).

Many searches for AGN in mergers on $< 10 \text{ kpc}$ scales have looked for double peaked narrow [O III] $\lambda 5007$ emission lines in large spectroscopic datasets such as SDSS (e.g. Xu & Komossa 2009; Wang et al. 2009; Liu et al. 2010; Smith et al. 2010). Comerford & Greene (2014) found 351 offset AGN candidates amongst a sample of 18,314 Type 2 AGNs by measuring velocity offsets between the forbidden and Balmer emission lines relative to the stellar absorption lines. Fu et al. (2011a) found 16 dual AGN candidates with high-resolution near-infrared images of 50 double-peaked [O III] $\lambda 5007$ AGNs and one of these was confirmed as a kpc scale binary AGN with high resolution radio images (Fu et al. 2011b).

Searches for dual AGN via X-ray, radio, and optical imaging suffer from different selection effects, and there are ongoing efforts to understand the affects of AGN obscuration by gas and dust during merger (Glikman et al. 2015; Kocevski et al. 2015; Koss et al. 2011). Koss et al. (2018) found that obscured luminous black holes, with X-ray emission but not visible broad $H\beta$ lines, were significantly more likely be in a later stage nuclear merger than a comparable sample of inactive galaxies.

The lack of a large sample of dual AGN and even a small sample of confirmed recoiling SMBHs motivates the development of new search strategies to find AGN offset from their host galaxies and from compan-

ion galaxies. A large transient survey which identifies offset AGN candidates via their variability provides one such approach.

In this paper we present a new method for a systematic search for offset AGN - both recoiling SMBHs and AGN in galaxy mergers - using the Zwicky Transient Facility (ZTF; Bellm et al. 2019; Graham et al. 2019; Dekany et al. 2020). In Section 2, we present our techniques for filtering ZTF transients to make a sample of 5493 optically variable AGN. We present a new version of **The Tractor** forward modeling software for confirmation of AGN-host spatial offsets and describe the candidate selection strategy used to obtain 9 offset AGN candidates and 52 AGN in mergers. In Section 3, we describe the multi-wavelength and spectroscopic properties of these new samples and present the rebrightening of the previously discovered recoiling SMBH candidate SDSS1133.

2. SAMPLE SELECTION

2.1. *The Zwicky Transient Facility*

ZTF is a wide-field optical transient survey of the northern sky which observes in g, r and i bands with an average 3 day cadence. It uses a 47 square degree field of view camera mounted on the Samuel Oschin 48-inch Schmidt telescope at Palomar Observatory. ZTF Phase 1 ran from November 2017 to October 2020. ZTF identifies variable sources by difference imaging, by which a change in flux is detected by subtracting a high-quality reference image from the data obtained at each epoch. ZTF provides a unique way to search for offset AGN via their optical variability by giving three main advantages for a systematic AGN search.

Firstly, ZTF provides an opportunity to find new AGN by applying light curve modeling techniques to new transients and identifying AGN-like stochastic variability. Light curve modeling of variable sources to find previously undiscovered AGN has been demonstrated for SDSS Stripe82 difference imaging and for transients in the Palomar Transient Factory (Baldassare et al. 2018, 2020).

Secondly, image subtractions containing the AGN can be used to locate the position of AGN-like variability relative to the host galaxy. Because spectroscopic surveys tend to have large plate sizes (for example, the SDSS spectroscopic plate size is $1.49''$), it can be difficult to confirm that the location of the AGN broad lines is associated with an offset point source. By detecting AGN-like variability from an offset point source, we can confirm the spatially offset nature of AGN candidates.

Finally, the sky coverage of ZTF allows us to search a very large area for offset AGN, which is important given

the apparent rarity of > 10 pc spatial offset recoiling SMBHs at low redshifts.

2.2. Selection of variable AGN in ZTF

We obtain our sample of variable AGN using data from the ZTF alert stream (Patterson et al. 2019). ZTF alerts contain details of the photometry and astrometry of single epoch transient detections (Masci et al. 2019). The ZTF pipeline produces approximately 100,000 alerts every night, so we implement a filter with the alert broker and analysis framework AMPEL (Alert Management, Photometry and Evaluation of Lightcurves; Nordin et al. 2019) to detect variable AGN amongst other transient phenomena.

Our method to filter out poor subtractions, moving sources and variable stars and find only extragalactic transients is similar to the approach used in the Tidal Disruption Event (TDE) filter of van Velzen et al. (2020). We apply a liberal cut of < 0.8 on the star-galaxy score (Tachibana & Miller 2018) to find transients associated with galaxies, and a cut of < 0.3 on the real-bogus score (Mahabal et al. 2019; Duev et al. 2019) to remove bogus transients. We remove objects in busy stellar fields by crossmatching to the Gaia and PanSTARRS catalogs with catsHTM (Soumagnac & Ofek 2018) to ensure that there are no more than 30 Gaia objects and 100 PanSTARRS objects within a $15''$ radius. We also require at least 3 significant detections > 0.01 days apart and a minimum flux increase of 2.5 magnitudes.

This filtering strategy primarily finds two kinds of common extragalactic transient: supernovae and variable AGN. To select AGN and remove supernovae within the AMPEL filter, we require that our transients either match an object in a series of AGN catalogs or have variability which is more characteristic of an AGN than a supernova. We use catsHTM and Extcats¹ look for a $2''$ crossmatch with objects in The Million Quasar Catalog (Flesch 2015), a machine learning based catalog of photometric AGN candidates (Brescia et al. 2015), and a catalog of 720,961 variable sources from the Palomar Transient Factory (PTF) and intermediate Palomar Transient Factory (iPTF) found between 2009 and 2016 which were not classified as a star (Miller et al. 2017) and had > 5 detections over > 24 hours.

For sources which do not have an AGN catalog crossmatch, we model their full ZTF light curve history within the AMPEL filter. We use the SNCOSMO supernova modeling tool (Barbary et al. 2016) to fit the ‘salt2’ SN Ia model to the g and r band light curves and extract

the reduced χ^2_{SN} goodness of fit for the best fit SN Ia model.

For comparison to the SN Ia goodness of fit, we implement the Butler & Bloom (2011) quasar modeling routine. This routine calculates the structure function for input light curves and compares this to the ensemble quasar structure function for Sloan Stripe 82 g- and r-band AGN light curves. The goodness of fit of the ensemble structure function model χ^2_Q gives a measure of how likely the ZTF light curve shows AGN-like variability.

The Butler & Bloom (2011) routine also calculates the reduced χ^2 for the null hypothesis that the source shows non-AGN like variability (such as from a variable star). We denote this as χ^2_{Q0} . For the purposes of separating AGN from supernovae, we found that AGN generally have much lower χ^2_{Q0} values than χ^2_{SN} , so this variability statistic was also effective at separating SN from AGN.

Our filter accepted any source where either the g- or r-band light curve had a χ^2_Q or χ^2_{Q0} value less than the SN Ia goodness of fit χ^2_{SN} . Based on tests with a sample of 111 spectroscopically classified supernovae with > 20 ZTF epochs, we determined that this method removes approximately 95% of SN Ia and 60% of Type II supernovae. With a sample of 166 spectroscopically confirmed AGN with 20 – 100 ZTF epochs, we were able to classify 80% correctly.

Objects which pass either the AGN crossmatch criteria or light curve fitting criteria pass the AMPEL filter and are pushed to the GROWTH Marshal science portal for arrangement of spectroscopic followup (Kasliwal et al. 2019). The AGN candidates are then confirmed either with existing SDSS spectroscopy, follow-up spectroscopic observations with the DeVeney spectrograph on the Lowell Discovery Telescope, or by their WISE color or variability history. To classify AGN based on their WISE W1-W2 color we use the criteria:

$$W1 - W2 > 0.662 \exp \{0.232(W2 - 13.97)^2\} \quad (1)$$

from Assef et al. (2013), and to classify AGN based on significant variability in their WISE light curve we require the χ^2 relative to a flat light curve to satisfy $\chi^2/\text{dof} > 10$.

We apply this procedure to ZTF alerts from a 2.5 year period between 2018-01-01 to 2020-07-06 and obtain a sample of 5493 AGN. This constitutes our final sample of strongly variable ZTF AGN with spectroscopic or WISE color/variability confirmation.

2.3. Selection of AGN spatially offset from their host galaxy

2.3.1. Image modeling with The Tractor

¹ <https://github.com/MatteoGiomi/extcats>

In order to model the positions of the variable AGN relative to their host galaxy, we apply **The Tractor** (Lang et al. 2016) to forward model the host galaxy profile and transient point source emission across ZTF images. **The Tractor** forward models in pixel space by parametrizing the astrometry solution, sky noise and point spread function of each image and modeling this simultaneously with the shape, flux and position of each source across images in multiple bands and surveys.

We developed a version of **The Tractor** to fit a host galaxy profile and an overlapping point source with a position in or around the host. The point source position and the host galaxy shape, flux and position are assumed constant over all epochs, while the point source flux may vary across single epochs. The version of **The Tractor** that we apply determines the best fitting one of two galaxy profile models: a pure de Vaucouleurs profile described by $I(r) = I_0 \exp(-7.67[(r/r_e)^{1/4}])$ and a pure exponential profile described by $I(r) = I_0 \exp(-1.68r/r_e)$. **The Tractor** is a highly advantageous tool for this analysis because it allows for forward modeling across images of different bands and different instruments. ZTF images taken when the transient is bright can be simultaneously modeled with higher resolution, deeper images from a different telescope to improve modeling of the host galaxy. Subtraction of the model from the coadded data can also help to reveal irregularities in galaxy structure in the residuals.

We model 3422 out of the sample of 5493 AGN with **The Tractor**. Due to limitations with archival storage of ZTF images on our filesystem and the availability of overlapping Legacy Survey (Dey et al. 2019) and PanSTARRS (Chambers et al. 2016) images for astrometric source matching, we do not model the remaining 2197. As the AGN are isotropically distributed in the sky, choosing to model a subset of the full sample with **The Tractor** does not introduce any biases with respect to the distribution of host-AGN offsets.

We remove 126 of the 3422 AGN because they are duplicates of existing AGN in the sample. The ZTF alerts for particular transients sometimes demonstrate such a large scatter in position that they are considered to be 2 or 3 different transients with separate ZTF names by the alert pipeline. As such, the alerts associated with a single transient can be distributed across 2 or 3 different transient objects. By applying an 8" cone search to all AGN which pass our AMPEL AGN filter, we find transients which are associated with one another and select the transient ID with the most alert packets containing real detections.

We model the sample of 3296 unique AGN by selecting the 30 ZTF g- and r-band images taken closest

to peak magnitude. This allows us to reach a median depth in AB magnitude of 22.4 in the r-band and 22.2 in the g band. We chose to model only 30 ZTF images due to computational and time constraints, but future work could model the whole sample with **The Tractor** to greater depths by modeling larger numbers of ZTF images and including higher resolution, deeper DECam imaging.

Of our 3296 AGN, 186 do not have sufficient S/N in the 30 ZTF images to model a host galaxy and point source with **The Tractor** above the limiting magnitude. This leaves 3110 objects with measured AGN-host offsets.

2.3.2. Determination of statistically significant spatial offsets

Since the ZTF camera has 1" pixels but much of the simulated recoiling black hole population is only observable at sub-arcsecond AGN-host spatial offsets, it is important to understand the positional accuracy that can be obtained from ZTF image subtractions.

In order to determine which AGN-host spatial offsets are statistically significant, we studied the distribution of offsets from the sample of 3110 objects with host galaxy and point source positions determined by simultaneously modeling 30 g- and r-band ZTF images with **The Tractor** (see Section 2.3.1). If the uncertainties in observed RA and Dec of the host galaxies are normally distributed with standard deviation σ_{ref} , and the uncertainties in observed RA and Dec of the variable point sources are normally distributed with standard deviation σ_{sci} , the radial distribution of spatial offsets between them will follow a Rayleigh distribution with $\sigma_{\text{R}}^2 = \sigma_{\text{ref}}^2 + \sigma_{\text{sci}}^2$.

In a study of radio AGN from the Cosmic-Lens All Sky Survey (CLASS) catalogue, Skipper & Browne (2018) found that the population of offsets between the radio AGN and optical galaxies in SDSS followed a mixture distribution consisting of a Rayleigh component and an exponential tail component, where the latter component may represent real AGN-host spatial offsets. Similarly, we find that the shape of our offset distribution from **Tractor** modeling of the ZTF AGN sample is described by the expected Rayleigh distribution at offsets $\lesssim 1''$ and at offsets $\gtrsim 1''$ the distribution is better described by a decaying exponential.

We therefore model our AGN-host spatial offset distributions by splitting the AGN sample into 3 different sub-samples based on the peak difference magnitude of the AGN. The ranges in peak-magnitude used to produce the sub-samples are 15-18, 18-19.5, and 19.5-23. These ranges were selected to ensure that each sample was large enough to model the Rayleigh and exponen-

tial tail components. The offset distributions of these sub-samples are shown in Figure 1a) c) and e). We fit a mixture distribution consisting of a Rayleigh component $\alpha(x)$ and an exponential component $\epsilon(x)$:

$$P_x = C\alpha(x) + (1 - C)\epsilon(x) \quad (2)$$

$$= C \frac{x}{\sigma_R^2} \exp\left(-\frac{x^2}{2\sigma_R^2}\right) + (1 - C)\tau e^{-\tau x} \quad (3)$$

for offsets x , mixture coefficient C , Rayleigh width σ_R and exponential decay parameter τ . We do this by directly minimizing the log likelihood between the model distribution and the data.

We first model the exponential decay parameter τ to fit only offsets $x > 1.0''$ so that the fit is not heavily biased by low-offset sources which dominate the distribution. We then fix the value of τ and fit the mixture distribution to the whole sample to find σ_R and C . The fits are shown in Figure 1a), c) and e), where we can see the Rayleigh component explaining the portion of the distribution which arises from positional uncertainty, and the exponential tail component showing the portion of the distribution which may contain physically real AGN-host offsets. For a given AGN-host offset x , we can calculate the probability that an offset $\geq x$ is drawn from the Rayleigh component as:

$$P_\alpha(x > R) = \frac{C \int_R^\infty \alpha(x)}{C \int_R^\infty \alpha(x) + (1 - C) \int_R^\infty \epsilon(x)} \quad (4)$$

This is shown as a function of offset x in Figure 1 b), d) and f). The probability function shows that as the offset increases, it is less likely to be explained by the Rayleigh distribution arising from positional uncertainties and more likely to be part of an exponential tail consisting of possibly real offsets. The spatial offset at which the probability of being drawn from the exponential component of the mixture distribution is 0.3% is marked by the dashed line in Figure 1 b), d) and f).

Using the probability functions to determine an offset cutoff for the 3 sub-samples, we determine 3σ offset cutoffs shown in the first column of Table 1. We select a cutoff of $0.511''$ for AGN with a peak magnitudes between 15 and 18, $0.773''$ for AGN with peak magnitudes between 18-19.5, and $0.976''$ for AGN with peak magnitudes between 19 and 23.

2.3.3. Matching of transient positions from *The Tractor* and ZTF difference images.

In order to confirm that the best-fit point source position from *The Tractor* modeling is consistent with

Table 1. Selected offset cutoffs

Peak magnitude	<i>Tractor</i> modeling	g-band subtraction	r-band of AGN	Number
19.5-23	0.732	0.946	0.976	64
18-19.5	0.605	1.009	0.773	164
15-18	0.574	0.959	0.551	23

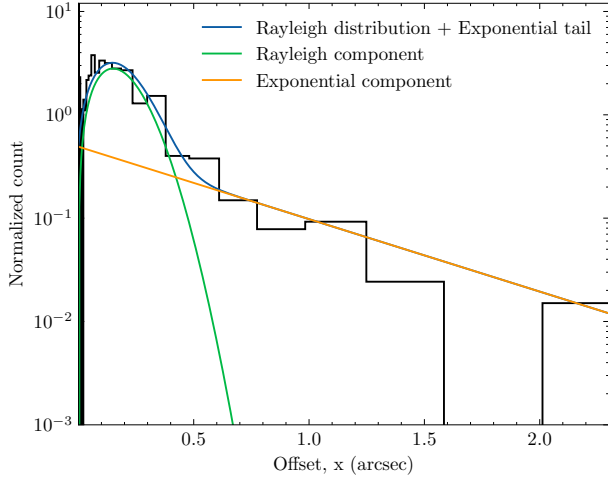
NOTE—Offset cutoffs (arcseconds) selected for AGN based on their peak magnitude. The first column shows the peak magnitude bin. The second column shows the $> 3\sigma$ cutoff for a significant AGN-host offset derived by *The Tractor*. The third and fourth columns show the uncertainty on the magnitude-weighted transient position derived from ZTF alert packets for r-band and g-band respectively. The fifth column shows the number of AGN which have $> 3\sigma$ *Tractor* offsets and match the the magnitude-weighted transient position using these cutoffs.

the position of the transient in the ZTF difference images, we calculate the magnitude-weighted position of the transient from the ZTF alert packets containing information about the position and magnitude of each single epoch difference image detection. The weights $1/\sigma_{\text{offset}}^2$ for the magnitude-weighted transient position are calculated using equation 3 from van Velzen et al. (2019):

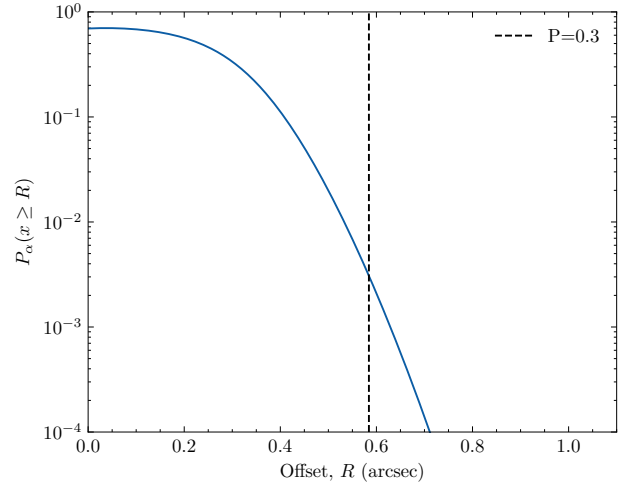
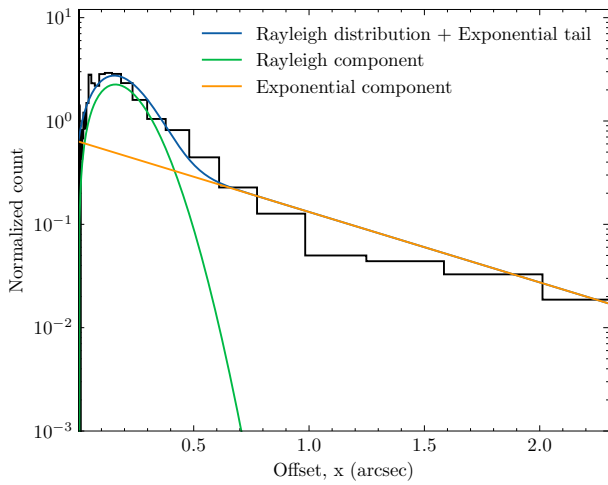
$$\sigma_{\text{offset}} = 0.24 + 0.04(m_{\text{diff}} - 20) \quad (5)$$

In order to determine the uncertainty in the magnitude-weighted transient position from the ZTF alert packets, we undertake the same offset distribution modeling procedure as we do for the *Tractor* AGN-host offsets. The modeling results for each magnitude binned sub-sample are shown in Figures 15 and 16 in the Appendix. As the distribution of magnitude-weighted offsets is substantially different for g-band images and r-band images, due to the differing contributions of the AGN towards the reference image in different bands, we find cutoffs for the two bands separately. Using the probability functions to determine an offset cutoff for the 3 sub-samples, we determine 3σ offset cutoffs shown in Table 1. When checking that the *Tractor* point source positions are consistent with the magnitude-weighted transient position from the alert packets, we require a match within these selected cutoffs.

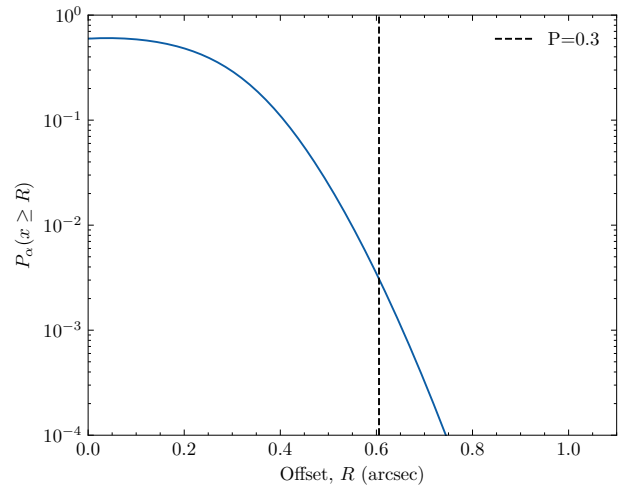
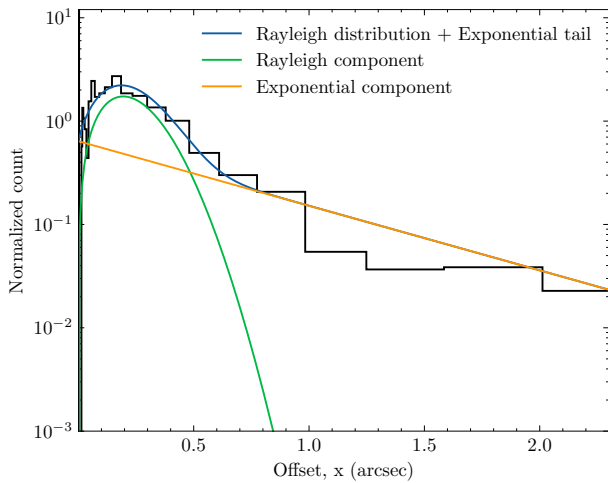
Figure 2 shows an example ZTF image subtraction for ZTF18aaxvmpg with the *Tractor* model overlaid. For this object, the *Tractor* point source position shown in the red contours was consistent with the ZTF transient seen in the image subtraction. This object was therefore considered to be an offset AGN candidate.



(a) AGN with a peak difference magnitude between 15 and 18.

(b) 3σ offset = $0.574''$.

(c) AGN with a peak difference magnitude between 18 and 19.5.

(d) 3σ offset = $0.605''$.

(e) AGN with a peak difference magnitude between 19.5 and 23.

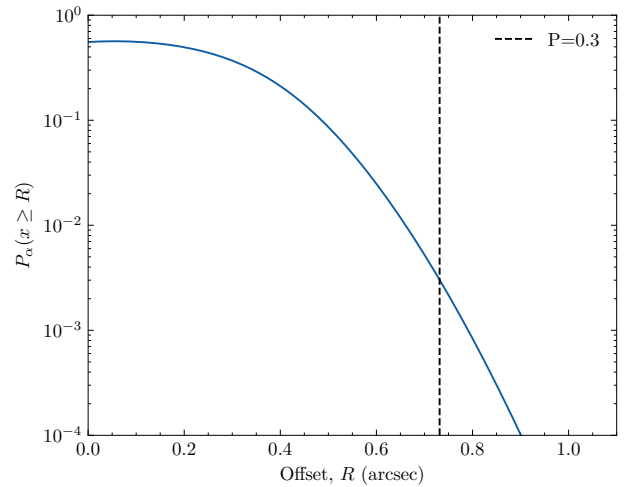
(f) 3σ offset = $0.732''$.

Figure 1. **Left:** Normalized histogram with logarithmic bins for AGN-host offsets obtained from *Tractor* modeling. The best-fit model of a mixture distribution with Rayleigh and exponential components are shown. **Right:** Probability that an offset greater than R is drawn from the Rayleigh component of the mixture distribution shown in (a) instead of the exponential component. The offset where this probability is 0.3% is shown with a dashed line.

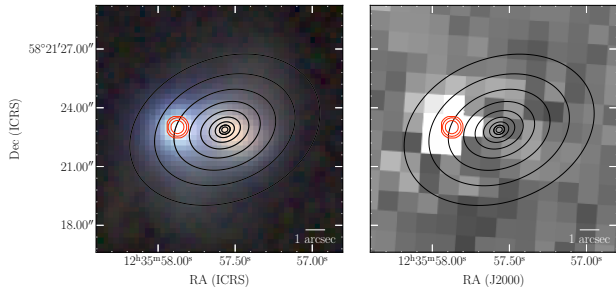


Figure 2. *Left:* Coadded g, r and z-band Legacy Survey images of ZTF18aaxvmpg. *Right:* ZTF image subtraction of ZTF18aaxvmpg when the AGN was close to peak magnitude. Overlaid contours show the best fit **Tractor** galaxy profile (black) and point source model (red) for a theoretical seeing of $1''$ derived from ZTF image modeling.

We find that 251 AGN have $> 3\sigma$ AGN-host offsets found by **Tractor** modeling which are consistent with the magnitude-weighted transient position from the alert packets. The breakdown of candidates from each peak magnitude bin is shown in the last column of Table 1.

2.4. Morphological classification of AGN hosts

In order to confirm the host-AGN offset found in ZTF images and classify the 251 offset AGN based on their morphology we undertook **Tractor** modeling with deeper, higher resolution images. For this task, we used archival images from the DESI Legacy Imaging Surveys (Dey et al. 2019). The combined DECam Legacy Survey, Mayall z-band Legacy Survey and Beijing-Arizona Sky Survey were taken between 2014 and 2019 and cover declinations from $-18^\circ < \delta < +84^\circ$, offering $0.262''/\text{pix}$ resolution and depths of 24.7, 23.9 and 23.0 for g, r and z bands respectively

We modeled the high resolution coadded g and r and z band Legacy Survey images of each system with a single galaxy profile and an offset point source. We then visually examined the images, **Tractor** models and residuals to determine if each offset AGN was well modeled as a point source or if there was excess unmodeled emission indicating the presence of a second host galaxy in the system which is centered on the AGN. We separated the sample into 5 categories. The number of objects in each category is shown in Table 2.

When the residuals of the galaxy and offset point source model showed a clear stellar bulge surrounding the offset AGN, it was considered likely that there are two galaxies in the system rather than one. If the residuals also showed morphological evidence of merging activity such as tidal structures, we considered the AGN to be part of a galaxy merger. These objects will be discussed further in Section 3.1.

Table 2. Classifications of offset AGN

Classification	Number
AGN in galaxy mergers	52
AGN offset from the stellar bulge of disturbed galaxy	9
AGN aligned with the stellar bulge of disturbed galaxy	21
AGN offset from an undisturbed galaxy	29
AGN without position confirmation in Legacy Survey modeling	140
Total	251

NOTE—The breakdown of the complete sample of 251 offset AGN into 5 morphology based classifications. The first row is the number of AGN which have extended galaxy emission around them and appear to be interacting and merging with a second galaxy. The second row shows the number of AGN which are not surrounded by a stellar bulge and appear to be spatially offset from the center of a galaxy with indications of recent merging activity. The third row shows the number of AGN in disturbed, post-merger systems where the stellar bulge is offset from the center of the extended galaxy profile and aligned with the AGN. The fourth row shows the number of AGN which appear to be point sources spatially offset from an undisturbed galaxy and are therefore more likely to be chance coincidences with background galaxies. The fifth row shows the number of AGN where the ZTF position could not be confirmed in archival Legacy Survey images.

When the system was well modeled by a single galaxy profile and offset point source and there were no tidal structures indicating recent merging activity it was considered likely to be chance coincidence of an AGN and an unrelated background galaxy. These AGN are discussed in Section 3.2.

For AGN which were well modeled by an offset point source and showed morphological evidence of recent merging activity in the host galaxy residuals, we considered that the AGN may be a candidate for a recoiling SMBH. We consider these objects as recoiling SMBH candidates because the recoiling AGN is expected to be visible for a period of 10^6 years after recoil while the host galaxy will still show evidence of previous merging activity in its morphology (Blecha et al. 2016). These objects are discussed in Section 3.3.

When the AGN was in a disturbed, irregularly-shaped host galaxy with a stellar bulge which was offset from the photometric center of the extended galaxy profile and the AGN was aligned with this stellar emission, we did not consider the AGN to be a recoiling SMBH candidate.

The remainder of the 251 objects did not have point source emission from an AGN present in the Legacy Survey images used for **Tractor** modeling. For these objects, confirmation of the spatial offset discovered in ZTF images requires follow-up with deeper, higher resolution imaging taken when the AGN emission is visible.

The distribution of physical offsets in kpc for the whole sample, the AGN in mergers, and the off-nuclear AGN is shown in Figure 3. The complete AGN sample shows a

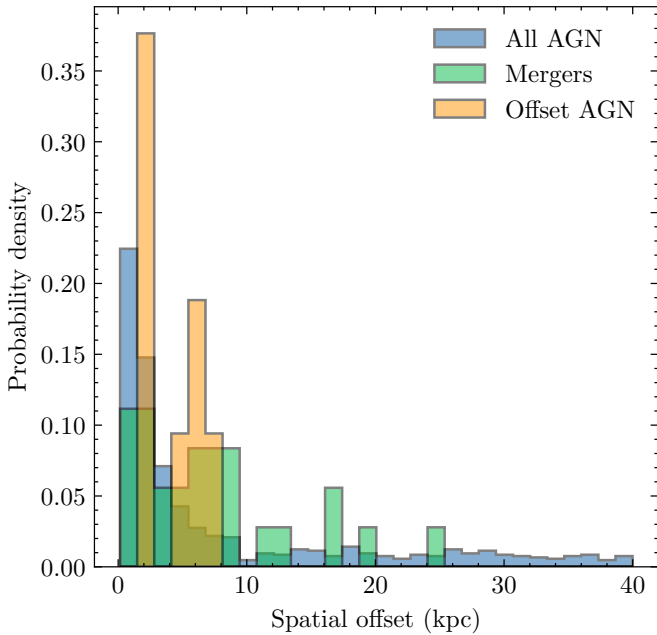


Figure 3. Distribution of physical spatial offsets between AGN and their closest galaxy for 3 samples: 898 AGN from the **Tractor** modeled sample with available spectroscopic redshifts, 27 of the 52 AGN in mergers with spectroscopic redshifts, and the 9 off-nuclear AGN. A tail of spatial offsets > 10 kpc can be seen in the blue histogram due to the presence of background galaxies.

tail extending beyond 40 kpc due to chance coincidences with background galaxies.

2.5. Spectroscopic analysis of ZTF broad line AGN

In order to model the distribution of broad line velocity offsets in the ZTF AGN sample as a whole, we modeled all 2422 ZTF AGN which had archival SDSS spectra and were classified as broad line AGN by the SDSS DR14 pipeline with Penalized Pixel Fitting (pPXF) (Cappellari & Emsellem 2004; Cappellari 2017). This method finds the velocity dispersion of stellar absorption lines using a large sample of high resolution templates of single stellar populations adjusted to match the spectral resolution of the input spectrum. We simultaneously fit the narrow $H\alpha$, $H\beta$, $H\gamma$, $H\delta$, $[S II] \lambda 6717, 6731$, $[N II] \lambda 6550, 6575$, $[O I] \lambda 6302, 6366$ and $[O III] \lambda 5007, 4959$ emission lines during template fitting. The emission line fluxes are each fit as free parameters but the line widths of the Balmer series are tied, as are the line widths of the forbidden lines.

In these fits the velocity of whichever broad $H\alpha$, $H\beta$, $H\gamma$, $H\delta$ lines were available within the SDSS spectroscopic wavelength range were allowed to vary up to 3000 km s^{-1} from the narrow lines. The velocity of each

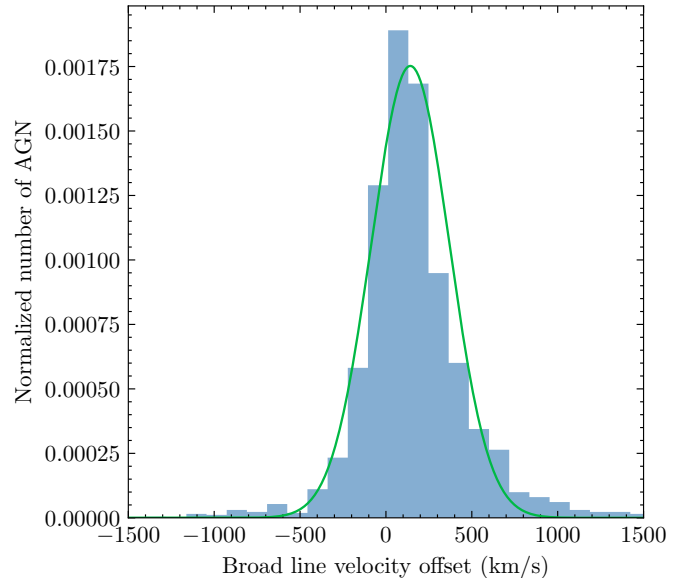


Figure 4. Distribution of Balmer broad line velocities relative to $[S II] \lambda 6717, 6731$, $[N II] \lambda 6550, 6575$, $[O I] \lambda 6302, 6366$ and $[O III] \lambda 5007, 4959$ emission line velocities found by fitting of archival SDSS spectra with pPXF. Velocities are shown for the 2422 AGN with broad Balmer lines in archival SDSS spectra out of the complete ZTF AGN sample of 5542 objects.

Balmer broad line was tied to the other Balmer broad lines.

We find that the ZTF AGN broad line velocities have a mean displacement of 143 km s^{-1} with a standard deviation of 126 km s^{-1} when the central component is fit with a Gaussian (Figure 4). Our distribution of broad line velocities for variable ZTF AGN is similar to the distribution of $H\beta$ broad line velocities found by Bonning et al. (2007) with a sample of 2598 SDSS AGN where they found a mean displacement of 100 km s^{-1} and a standard deviation of 212 km s^{-1} .

Our velocity distribution shows a substantial tail population with velocities up to $\pm 2500 \text{ km s}^{-1}$. The fraction of AGN f_ν with velocity magnitudes greater than 1000, 1500 and 2000 km s^{-1} are $f_{1000} = 0.025$, $f_{1500} = 0.009$ and $f_{2000} = 0.003$. These values are comparable to those found in the Bonning et al. (2007) sample, where they find fractions of $f_{1000} = 0.0035$, $f_{1500} = 0.0012$ and $f_{2000} = 0.0008$. It therefore appears that a variability selected AGN sample shows a broad line velocity distribution which is typical of a spectroscopically selected AGN sample.

3. RESULTS

3.1. AGN in galaxy mergers

52 of our spatially offset AGN were determined to be in a merger with a second galaxy based on the presence

Table 3. Galaxy merger summary

ZTF name	z	RA (hms)	Dec (dms)	Host-AGN offset arcseconds)	Host-AGN offset (kpc)	AGN flux/ companion flux	# AGN/ # spectra	Broad line shape
ZTF18aaxvmpg	0.212	12:35:57.810	58:21:21.726	1.296 ± 0.002	5.477 ± 0.008	21.25	1/2	C
ZTF18abamzru	-	17:23:27.486	42:21:22.683	5.231 ± 0.001	—	3.20	1/0	-
ZTF18aasxvyo	0.23635	14:17:56.535	25:43:21.506	5.068 ± 0.001	19.162 ± 0.004	24.24	1/2	-
ZTF18aaieguy	0.213	13:25:18.144	41:10:09.717	0.676 ± 0.007	2.88 ± 0.03	0.48	1/1	C
ZTF19aakjemw	0.147	15:50:07.939	27:28:10.996	5.435 ± 0.002	16.184 ± 0.006	1.16	1/2	2B
ZTF18aaifbku	0.067	13:2:58.854	16:24:27.806	0.985 ± 0.002	1.371 ± 0.003	1.75	1/1	C
ZTF18aampabj	0.216	16:52:58.864	44:48:45.540	3.118 ± 0.001	13.416 ± 0.004	15.31	1/2	C
ZTF19aaagygp	-	01:07:13.787	-11:36:2.998	3.969 ± 0.002	—	0.34	1/0	-
ZTF18abujubn	-	19:08:12.617	45:32:49.643	8.562 ± 0.002	—	25.97	1/1	MG
ZTF18acegbsb	0.037	09:04:36.964	55:36:02.772	9.027 ± 0.009	6.992 ± 0.007	2.14	1/2	RS
ZTF19aaozpdm	-	13:37:01.086	20:25:14.343	1.682 ± 0.013	—	0.88	1/0	-
ZTF18aacjltc	0.189	08:12:52.208	40:23:47.253	3.192 ± 0.001	12.091 ± 0.004	6.06	2/2	C
ZTF18abtpite	-	23:18:35.915	41:08:0.284	3.262 ± 0.007	—	0.41	1/0	-
ZTF18abvvrwx	-	22:04:07.944	-08:57:24.505	1.011 ± 0.003	—	5.87	1/0	-
ZTF19abaktpb	-	16:42:19.074	03:45:53.037	8.164 ± 0.008	—	8.5	1/0	-
ZTF18aaqjcxl	0.099	07:58:47.235	27:05:16.379	3.314 ± 0.001	6.712 ± 0.002	5.54	1/1	None
ZTF18abyoivl	-	00:22:52.018	08:24:0.757	0.391 ± 0.007	—	0.97	1/0	-
ZTF18aabdiug	0.062	12:31:52.060	45:04:43.273	0.572 ± 0.002	0.738 ± 0.003	11.13	1/2	C
ZTF19aaviuyv	-	18:56:20.579	37:12:36.076	2.651 ± 0.014	—	3.48	1/0	-
ZTF18aabxczq	0.063	10:38:33.425	46:58:06.741	0.413 ± 0.003	0.54 ± 0.004	1.69	1/2	C
ZTF18acvwlr	0.233	12:50:16.219	04:57:45.074	1.271 ± 0.005	5.887 ± 0.023	0.99	1/1	MG
ZTF19aasejqv	0.233	14:13:29.817	26:44:35.232	1.57 ± 0.009	7.261 ± 0.042	2.85	1/1	2B
ZTF18aazogyo	0.081	14:56:27.421	30:53:40.225	5.442 ± 0.001	9.123 ± 0.002	13.33	1/1	2B
ZTF18aceypyv	0.163	09:51:12.391	31:35:37.084	2.406 ± 0.003	7.913 ± 0.01	9.89	1/1	C
ZTF18acbweyd	0.189	10:20:38.565	24:37:12.421	4.24 ± 0.002	16.125 ± 0.008	11.47	1/1	C
ZTF18acablce	-	16:30:55.490	72:26:43.352	1.37 ± 0.008	—	1.77	1/0	-
ZTF18abhpvvr	-	00:38:33.041	41:28:53.681	3.595 ± 0.003	—	9.29	1/0	-
ZTF19abfqmjj	-	22:56:41.062	23:02:32.510	7.402 ± 0.003	—	48.85	1/0	-
ZTF18abmqwgr	-	20:27:53.382	14:08:50.604	1.539 ± 0.012	—	1.33	1/0	-
ZTF19aadgbih	0.196	12:46:33.522	45:34:21.773	0.381 ± 0.006	1.494 ± 0.024	1.82	1/1	C
ZTF19aalpfan	0.075	13:27:51.414	06:42:49.854	0.627 ± 0.003	0.964 ± 0.005	6.86	1/1	2B
ZTF18aawwfep	0.197	08:54:41.735	30:57:54.759	1.37 ± 0.003	5.424 ± 0.012	2.22	1/2	C
ZTF19aavxims	-	12:48:55.053	-06:59:54.802	4.735 ± 0.001	—	31.52	1/0	-
ZTF19aaaplct	-	14:57:28.940	08:34:22.879	2.835 ± 0.039	—	0.97	1/0	-
ZTF18aajnqqv	0.081	12:57:41.050	20:23:47.747	1.747 ± 0.004	2.901 ± 0.007	2.83	1/2	C
ZTF18abszfr	0.291	22:07:16.099	12:11:03.278	4.322 ± 0.002	24.646 ± 0.011	5.16	1/0	-
ZTF19abucbkt	0.16239	01:36:04.252	21:37:25.882	4.60 ± 0.01	12.935 ± 0.03	4.47	1/1	C
ZTF18adbhlyb	0.212	11:17:59.188	20:15:19.078	4.508 ± 0.001	19.075 ± 0.004	4.19	1/1	MG
ZTF18acxhoij	-	01:12:07.783	-21:04:28.682	1.099 ± 0.045	—	2.66	1/0	-
ZTF18acajwep	-	01:04:05.280	21:22:31.946	2.906 ± 0.001	—	7.08	1/0	-
ZTF19abipoqj	-	22:43:14.796	80:59:27.375	0.499 ± 0.005	—	0.57	1/0	-
ZTF19abpkooq	-	02:34:16.170	05:18:42.732	3.524 ± 0.001	—	8.38	1/0	-
ZTF18abztovy	-	08:29:24.624	34:50:45.655	1.719 ± 0.011	—	0.86	1/0	-
ZTF18acsllgd	-	03:45:45.495	22:23:58.156	1.041 ± 0.025	—	4.59	1/0	-
ZTF19aanxrki	0.114	15:32:27.165	04:19:22.283	3.557 ± 0.064	8.304 ± 0.149	1.29	1/2	C
ZTF18aamfuhc	0.086	13:42:34.214	19:13:34.184	4.845 ± 0.001	8.608 ± 0.002	28.39	1/1	MG
ZTF18aadwvyr	0.126	08:29:44.346	32:52:21.163	0.984 ± 0.003	2.53 ± 0.008	2.15	1/1	C
ZTF19abauzsd	0.285	15:54:32.681	21:43:48.220	0.483 ± 0.007	2.699 ± 0.039	1.03	1/1	RS
ZTF18abufbsq	-	23:46:15.513	12:47:07.733	0.891 ± 0.003	—	1.14	1/0	-
ZTF18abzuzrg	0.178	16:55:16.540	32:15:55.145	1.369 ± 0.007	4.908 ± 0.025	1.11	1/2	MG
ZTF18abtmcdb	-	01:20:12.473	07:12:58.251	0.535 ± 0.003	—	3.81	1/0	-
ZTF18aauhnby	0.09	12:04:15.954	56:02:58.100	1.092 ± 0.002	2.02 ± 0.004	1.39	1/1	C

NOTE—Summary of the properties of the 52 AGN in galaxy mergers. The host-AGN spatial offset and g-band AGN/host galaxy flux ratio is derived from *Tractor* modeling of Legacy Survey images. The 8th column lists the number of AGN in the system, confirmed by either spectroscopic narrow emission line ratios or the presence of WISE variability. The second number in this column indicates whether neither the AGN or host galaxy have a spectrum (0), only the AGN has a spectrum (1), or whether there is a spectrum of the host galaxy centroid available as well (2). The last column shows the classification of the shape of the broad line region where we adopt the scheme of *Strateva et al. (2003)*: prominent red shoulder (RS), prominent blue shoulder (BS), two prominent peaks (2P), two blended peaks (2B), complex multi-Gaussian structure (MS). We denote a Gaussian broad line with a ‘C’ and AGN with no broad line emission with ‘None’. If there are no SDSS spectra available for the source or the source does not have broad emission lines, we indicate this with a ‘-’.

of two galaxy nuclei and an interacting morphology visible in Legacy Survey images. The galaxy separations from *Tractor* modeling of Legacy Survey images ranged from $0.4 - 9''$, which for the 33 AGN with known redshifts, corresponded to physical separations of 0.54 to 24.65 kpc.

For 14 of these AGN with available SDSS spectra of both galaxies, only ZTF18aacjltc had narrow emission lines consistent with an AGN in both galaxies. The remaining 12 were single AGN, where the companion galaxy did not show narrow AGN emission line ratios. The observed fraction of dual vs offset AGN is consistent with the predictions of [Van Wassenhove et al. \(2012\)](#).

18 galaxy mergers had only one archival SDSS spectrum available, with the fiber centered on the variable AGN. Follow-up spectra of the companion galaxies will be required to determine if these mergers contain one or two AGN. 15 galaxy mergers have no archival SDSS spectra available, as the variable object was classified as an AGN based on optical ZTF and infrared WISE variability. Follow-up spectra of both the AGN and companion galaxy will be required to spectroscopically confirm the presence of one or two AGN.

Aside from appearing in a range of multi-wavelength AGN sample papers, a fraction of our sample of AGN in merging galaxies have been studied in more detail in the literature. ZTF18aasvxyo is a well studied X-ray bright BL Lac object ([Halpern et al. 1986](#)) and ZTF18aaqjcxl is also a known BL Lac ([Plotkin et al. 2008](#)). ZTF18acegbsb is a known double peaked emitter ([Strateva et al. 2003](#)) in a merger, which has been studied in the context of AGN photoionization of companion galaxies ([Keel et al. 2019](#)). ZTF18aamfuhc appeared in the same cross-ionization study. ZTF18abhpvvr is a known dual AGN ([Huang et al. 2014](#)) and ZTF18aawwfep and ZTF18aaajnqqv are also a known AGN pairs from [Liu et al. \(2011\)](#).

ZTF18aaifbku, an AGN in the double-lobed galaxy Mark 783, was imaged with the Karl G. Jansky Very Large Array (JVLA) at 5 GHz on September 6 2015 and showed radio emission from a compact core component and an extended component which was 26 kpc long ([Congiu et al. 2017](#)). The lack of jet emission led the authors to conclude that the radio emission was a relic of previous AGN activity before it entered a quiescent state. This AGN was first observed in ZTF on June 5 2018 and has shown continued variability to r-band magnitudes of 17.07 since then, suggesting that the AGN activity has turned on since on since the radio observations were made in 2015.

In order to determine if ZTF18aaxvmpg was a dual AGN, we undertook more detailed spectroscopic follow-

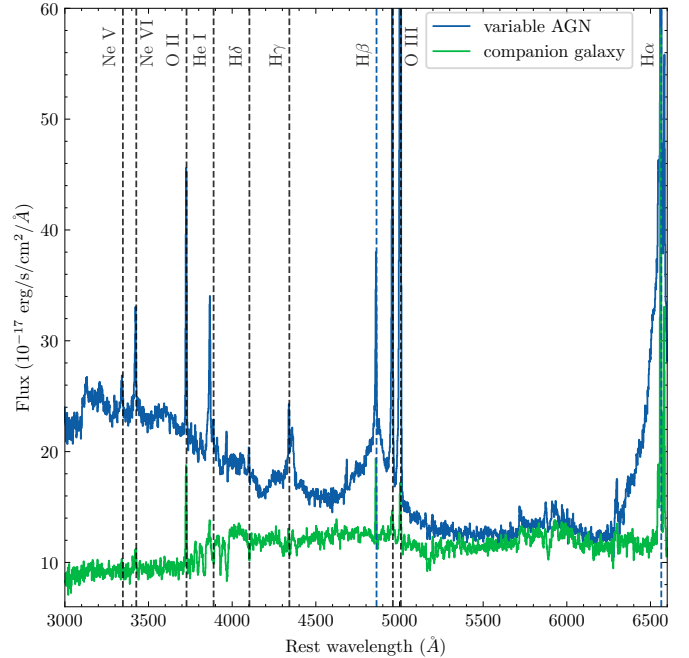


Figure 5. Spectrum of ZTF18aaxvmpg and its offset companion galaxy taken on 2019-10-29 with the DeVeny spectrograph on the Lowell Discovery Telescope. Broad Balmer emission features can be seen in the AGN spectrum but only narrow emission lines are visible from the companion galaxy.

up. The SDSS spectrum of ZTF18aaxvmpg taken on 2006-05-21 with the fiber centered on the AGN shows the presence of broad Balmer lines but the companion galaxy at 5.5 kpc from the AGN did not have an archival spectrum. We took a new spectrum of both the host galaxy and the AGN on 2019-10-29 with the DeVeny spectrograph on the Lowell Discovery Telescope using a $1.5''$ slit, central wavelength of 5700 \AA , a spectroscopic coverage of $3600\text{-}8000 \text{ \AA}$ and a total exposure time of 3200s. The spectrum of the AGN and the companion galaxy is shown in figure 5.

Using the best-fit line fluxes from pPXF modeling of the spectrum, we determine the AGN and host galaxy types using the $[\text{O III}] \lambda 5007$, $\text{H}\beta$, $[\text{N II}] \lambda 6573$ and $\text{H}\alpha$ line amplitudes. To separate starburst galaxies, we use the [Kauffmann et al. \(2003\)](#) condition:

$$\log([\text{O III}]/\text{H}\beta) < 0.61 / \log([\text{N II}]/\text{H}\alpha - 0.05) + 1.3 \quad (6)$$

The composite region is defined to be between the above condition and:

$$\log([\text{O III}]/\text{H}\beta) < 0.61 / \log([\text{N II}]/\text{H}\alpha - 0.47) + 1.19 \quad (7)$$

The emission line ratios and their classifications for the AGN and companion galaxy are shown in Figure 6. We find that the variable AGN has narrow line emission

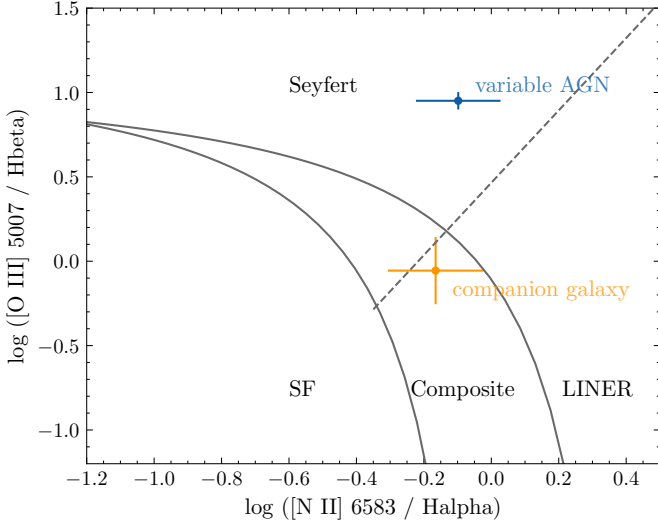


Figure 6. $[\text{O III}] \lambda 5007$, $\text{H}\beta$, $[\text{N II}] \lambda 6573$ and $\text{H}\alpha$ narrow emission line ratios from pPXF fitting of Deveny spectra taken of the variable AGN ZTF18aaxvmpg and its companion galaxy overlaid on the AGN classification regions. The ZTF AGN is classified as a Seyfert while the offset galaxy has a combination of star forming and LINER emission.

consistent with a Seyfert, while the host galaxy narrow line emission falls into the Composite/LINER category. Since LINER emission can be produced by either AGN or hot old stars on galactic scales we must distinguish between the two using a WHAN diagram, which classifies LINERs with a narrow $\text{H}\alpha$ equivalent width $> 3 \text{ \AA}$ and a $\log_{10}([\text{N II}] \lambda 6573 / \text{H}\alpha)$ flux ratio > -0.4 as AGN (Cid Fernandes et al. 2011; Mezcuca & Domínguez Sánchez 2020). With pPXF we measure the equivalent width of narrow $\text{H}\alpha$ to be $3.54 \pm 0.11 \text{ \AA}$ and the $\log_{10}([\text{N II}] \lambda 6573 / \text{H}\alpha)$ flux ratio to be -0.165 ± 0.097 . This places the spectrum in the ‘weak AGN’ class of the WHAN diagram, suggesting that the system does indeed host a second AGN. We therefore conclude that ZTF18aaxvmpg is part of a dual AGN.

6 of the 27 broad line AGN in mergers for which we have archival SDSS spectra showed double-peaked broad Balmer emission from an unobscured accretion disk, corresponding to 22% of the sample. The classifications of broad line region shapes of each AGN in a merger is shown in Table 3. It will be shown in Section 3.3.1 that 16% of the broad line ZTF AGN are double-peaked emitters. It is therefore possible that AGN in mergers are more likely to have double-peaked broad lines than normal ZTF AGN.

In order to determine the X-ray luminosities of the sample, we crossmatched with a $60''$ radius to the second ROSAT All Sky Survey catalog (Boller et al. 2016). 31 of the 52 AGN had X-ray detections. The X-ray fluxes

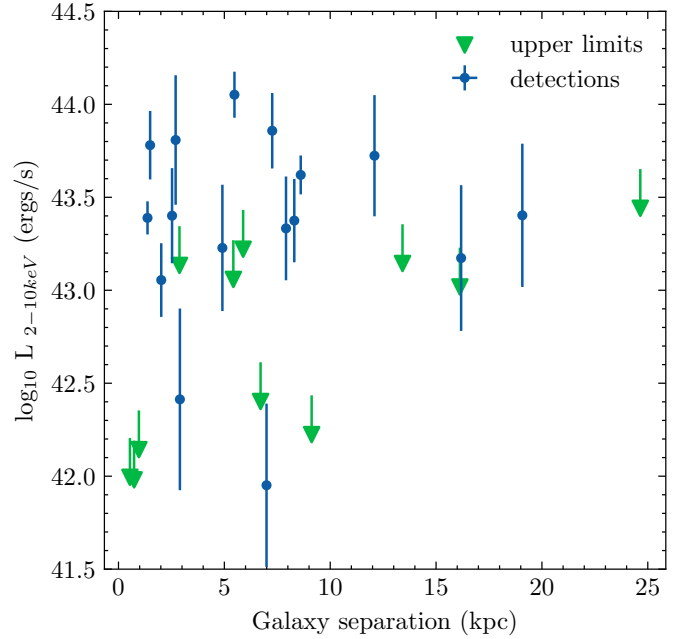


Figure 7. 2-10 keV luminosity for 27 AGN with known spectroscopic redshifts as a function of physical galaxy separation. 16 were detected in the second ROSAT All Sky Survey catalog and 11 have non-detections.

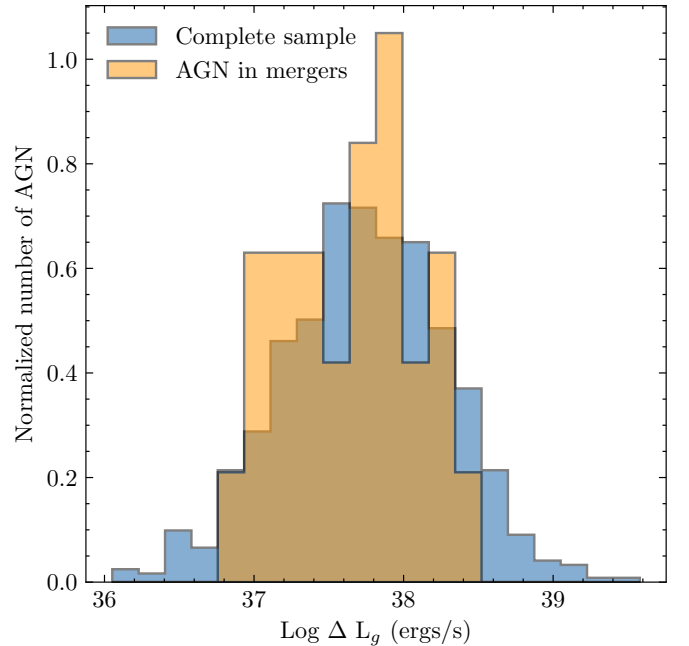


Figure 8. Peak luminosity change relative to the reference image amongst the 27 AGN in mergers and a larger sample of 689 variable AGN which had spectroscopic redshifts available. The larger sample was controlled to undergo the same quality cuts as the merger sample.

for the 31 AGN, along with luminosities for those with a known redshift, are shown in the appendix in Table 6. 9 of the AGN in mergers had 20cm radio detections in FIRST. 6 of the 10 AGN with detected radio emission were also X-ray bright. The smaller fraction of radio AGN compared to X-ray AGN may be due to delayed triggering of radio emission during the merger (Shabala et al. 2017; Skipper & Browne 2018).

The relationship between galaxy separation and X-ray luminosity is shown in Figure 7. We do not find a correlation between galaxy separation and X-ray luminosity in the range of 1-19 kpc separations represented by our sample. While Koss et al. (2012) found that the X-ray luminosity of AGN pairs decreased with increasing separation up to 90 kpc, they also did not find a strong correlation at < 20 kpc.

In order to compare the variability of the AGN in mergers with the complete AGN sample, we determined the maximum flux change between the ZTF reference image and single epoch science images for the AGN with spectroscopic redshifts, and determined the corresponding change in luminosity. The distribution of peak luminosity change is shown in Figure 8 for the merger sample and for a larger sample of 689 ZTF AGN, controlled for the quality cuts used to produce the merger sample. The two samples show a similar distribution of peak luminosity. Implementing a K-S test to compare the two samples confirms they are drawn from the same distribution with a p-value of 0.64.

3.2. Chance coincidences of AGN and background galaxies

In our sample of offset AGN we find 29 AGN which are offset from an undisturbed galaxy with a symmetrically shaped profile. While many of these objects are likely to be quasars coinciding with background galaxies, it is possible that a fraction of them are AGN with a real association to the spatially offset host. We therefore calculate the approximate number of chance coincidences with background galaxies that we expect for the sample of 3110 AGN which were modeled by **The Tractor**.

To determine the density of background galaxies, we use **Casjobs** to query the SDSS DR16 catalog to find the number of galaxies with g-band model magnitudes between 15 and 22.8. We choose a g-band limiting magnitude of 22.8 because it is the median effective depth of each set of 30 ZTF images used for **Tractor** modeling for the AGN sample. We note that the choice of the median limiting magnitude is a significant approximation because there is a wide distribution in effective depths for the **Tractor** models of the AGN. Future work should

take a more systematic approach to reach a consistent limiting magnitude for all objects.

We find that there are 119,364,394 galaxies in the SDSS survey area of $14,555 \text{ deg}^2$ within this magnitude range, corresponding to a density of 6.328×10^{-4} per arcsec². This means that we expect 1.789×10^{-2} background galaxies in a 3" radius circle around a given AGN. For our sample of 3110 AGN with **Tractor** modeled offsets, we therefore expect 56 AGN to have unassociated background galaxies within 3 arcseconds.

As we do not have an excess of AGN which are intrinsically offset from undisturbed host galaxies beyond the estimated number of chance coincidences, we do not have evidence that any of these objects could be recoiling SMBH candidates. The AGN falling under this category require spectra to confirm the host galaxy redshift and provide evidence that they are not chance coincidences.

3.3. AGN spatially offset from the center of disturbed galaxies

Our morphological classification scheme finds 9 AGN which are variable point sources spatially offset from a potential host galaxy and do not show evidence of a second stellar bulge around the AGN. The properties of the 9 offset AGN are summarized in Table 4. Their ZTF light curves are shown in Figure 9. We show the co-added g-, r- and z-band Legacy Survey images, the best-fit **Tractor** model and the residuals for each object in Figure 10. The residuals all show asymmetric, tidal structures indicative of previous merging activity. Spectra of the offset AGN are shown in Figure 11.

3.3.1. Broad line region properties

The details of the H α and H β broad line regions of the spatially offset AGN are shown in Figure 12. We do not have H α and H β spectra for ZTF18absvcae due to its higher redshift. 3 of the 8 AGN for which we do have spectra of the H α or H β lines (ZTF19aadgijf, ZTF18accptjn and ZTF19aadggaf) have Gaussian broad lines. The broad line velocity offsets determined from **pPXF** fitting are shown in Table 4. Because of the large distribution of velocity offsets in the overall ZTF AGN sample (Figure 4) and the fact that the ZTF AGN which have extreme $> 1500 \text{ km s}^{-1}$ velocities do not show any evidence of a spatial offset, we do not make any conclusions from the velocities of $200\text{--}600 \text{ km}^{-1}$ magnitude observed for these 3 AGN.

The remaining 5 of the 8 AGN (ZTF18aalsidi, ZTF19aautrth, ZTF18aaxmrom, ZTF19aayrjsx and ZTF18aaoeobb) for which we have spectra of the H α or H β lines have asymmetric Balmer broad line regions which are poorly fit by a Gaussian. Adopting the classification scheme of Strateva et al. (2003) we note that

Table 4. Properties of the offset AGN sample

ZTF name	z	RA (hms)	Dec (dms)	Host-AGN offset (arcseconds)	Host-AGN offset (kpc)	BLR velocity (km s ⁻¹)	X-ray?	Radio flux (mJy/beam)	R
ZTF19aaurth	0.208	16:30:41.964	30:36:2.448	1.076 ± 0.003	4.48 ± 0.012	BS+C	✓	< 0.98	< 1.56
ZTF19aadgijf	0.262	14:26:14.312	27:29:55.98	1.197 ± 0.005	6.195 ± 0.026	-462.0 ± 0.3	✓	< 1.24	< 2.05
ZTF18aaxmrom	0.347	16:09:11.257	17:56:16.271	1.057 ± 0.01	7.089 ± 0.067	RS+C	✓	444.48	1760
ZTF19aayrjsx	0.215	23:32:54.463	15:13:5.407	1.329 ± 0.003	5.695 ± 0.013	2P+C	✓	—	—
ZTF18aalsidi	0.348	15:53:57.736	47:52:32.015	0.901 ± 0.004	6.051 ± 0.027	RS+C	✓	< 0.88	< 10.32
ZTF18acctjn	0.214	22:12:17.117	3:50:40.531	0.648 ± 0.005	2.767 ± 0.021	344.3 ± 0.3	×	1.12	0.87
ZTF18absvcae	0.755	20:49:07.593	5:13:17.362	1.679 ± 0.004	22.035 ± 0.052	—	×	4.97	51.6
ZTF18aaoeobb	0.270	12:48:53.9	34:24:29.448	0.429 ± 0.004	2.284 ± 0.021	2B+C	✓	< 0.93	< 0.95
ZTF19aadggaf	0.266	13:42:06.57	5:05:23.898	0.488 ± 0.004	2.554 ± 0.021	608.2 ± 0.3	✓	3.80	2.89

NOTE—Spectroscopic and multi-wavelength properties of the sample of AGN spatially offset from disturbed host galaxies. The redshifts are based on the position of O III lines in the archival SDSS spectrum of the AGN. The spatial offsets are obtained from modeling of point source and galaxy profile positions in the Legacy Survey images with **Tractor**. For the 3 AGN with Gaussian broad lines the broad line velocities are found by modeling of a Gaussian Balmer broad line series with **pPXF**. For objects which are poorly modeled by a Gaussian and well-modeled by a double-peaked accretion disk model (Chen et al. 1989) we instead classify the shape of the broad line by adopting the scheme of Strateva et al. (2003): prominent red shoulder (RS), prominent blue shoulder (BS), two prominent peaks (2P), two blended peaks (2B). If a central Gaussian broad line is required in addition to the double-peaked accretion disk model we indicate this with ‘+C’. The X-ray detection column shows which objects have detections within 1.0’ of the AGN position in the ROSAT All Sky Survey catalog, and the radio detection column shows the 20cm flux density or upper limit if it was within the coverage of the FIRST radio survey (Helfand et al. 2015). The radio loudness is indicated in the last column.

Table 5. Best fit elliptical accretion disk parameters

Disk parameter	ZTF19aaurth	ZTF18aaxmrom	ZTF19aayrjsx	ZTF18aaoeobb	ZTF18aalsidi
Central component required	Yes	Yes	Yes	Yes	Yes
Inner radius (ϵ)	252	813	70	997	21
Outer radius (ϵ)	608	6800	7995	3841	2255
Inclination (degrees)	13.6	84.9	14.0	47.6	12.8
Turbulent broadening (c)	0.0035	0.0007	0.007	0.0019	0.0042
Ellipticity	0.46	0.54	0.91	0.17	0.93
Disk orientation (degrees)	18.2	36.6	39.0	31.6	357.1

NOTE—Best fit accretion disk parameters for the 5 double peaked emitters from modeling of the H α (ZTF19aadgijf, ZTF19aaurth, ZTF18aaxmrom and ZTF19aayrjsx) and H β (ZTF18aalsidi) double-peaked emission after subtracting a stellar continuum model derived from **pPXF**. Disk parameters were derived from a simple log likelihood minimization fit to illustrate that accretion disk emission models can describe the observed broad line shapes well.

ZTF18aalsidi and ZTF18aaxmrom have prominent red shoulders, ZTF18aaurth has a prominent blue shoulder, ZTF19aayrjsx has two prominent peaks which are red and blue-shifted, and ZTF18aaoeobb has two blended peaks which are red and blue-shifted. These classifications are displayed in Table 4. Such asymmetric structures can arise due to double-peaked emission from an unobscured, relativistic Keplerian accretion disk (Chen et al. 1989; Chen & Halpern 1989; Eracleous & Halpern 1994).

To determine if the velocity-offset peaks observed in the broad Balmer lines of ZTF19aaurth, ZTF18aaxmrom, ZTF19aayrjsx, ZTF18aaoeobb and ZTF18aalsidi can be accounted for by accretion disk emission, we model the broad lines with an elliptical accretion disk model where

an inner thick, hot ion torus illuminates a thin outer disk of ionized gas, which has a power law relation between emissivity and radius given by slope q (Chen et al. 1989; Strateva et al. 2003). The model depends on the inner and outer dimensionless gravitational radii of the disk ϵ_1 and ϵ_2 , local turbulent broadening parameter σ , azimuthal angle ϕ , inclination angle i where 0 degrees is face-on and 90 degrees is edge-on, ellipticity e and disk orientation ϕ_0 . We bound the inner radius to $0 < \epsilon_1 < 1000$, the outer radius to $0 < \epsilon_2 < 8000$, the emissivity slope to $0 < q < 5$ and the local broadening to $0 < \sigma < 0.01$, to ensure that we could find good fits to the data with disk shape parameters consistent with the known double-peaked emitter population from Strateva et al. (2003). All four objects required a cen-

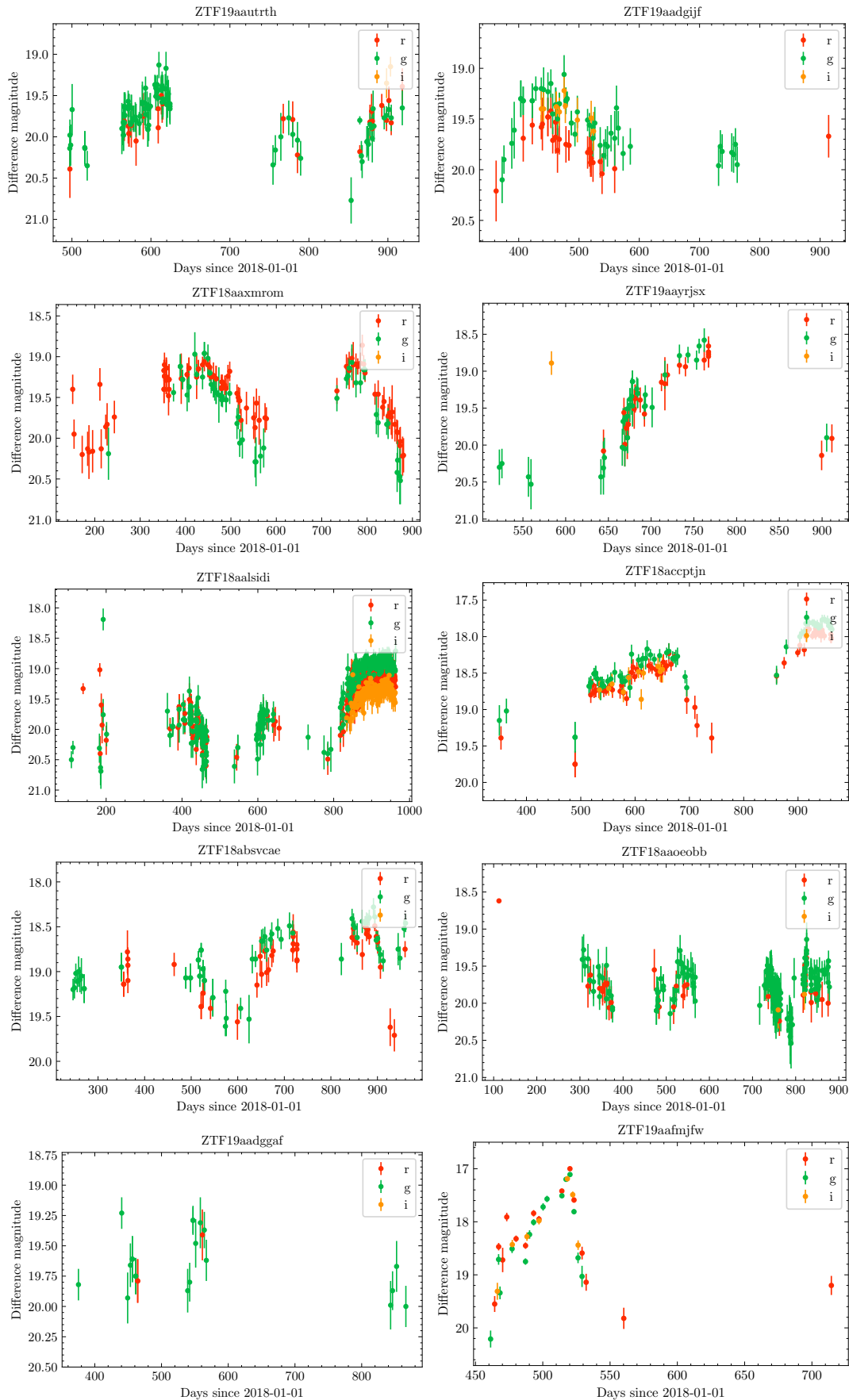


Figure 9. ZTF light curves in g, r and i-band for the sample of 9 offset AGN and the known recoiling SMBH candidate in SDSS1133 which rebrightened in ZTF (ZTF19aafmjfw). We show only the $> 3\sigma$ detections.

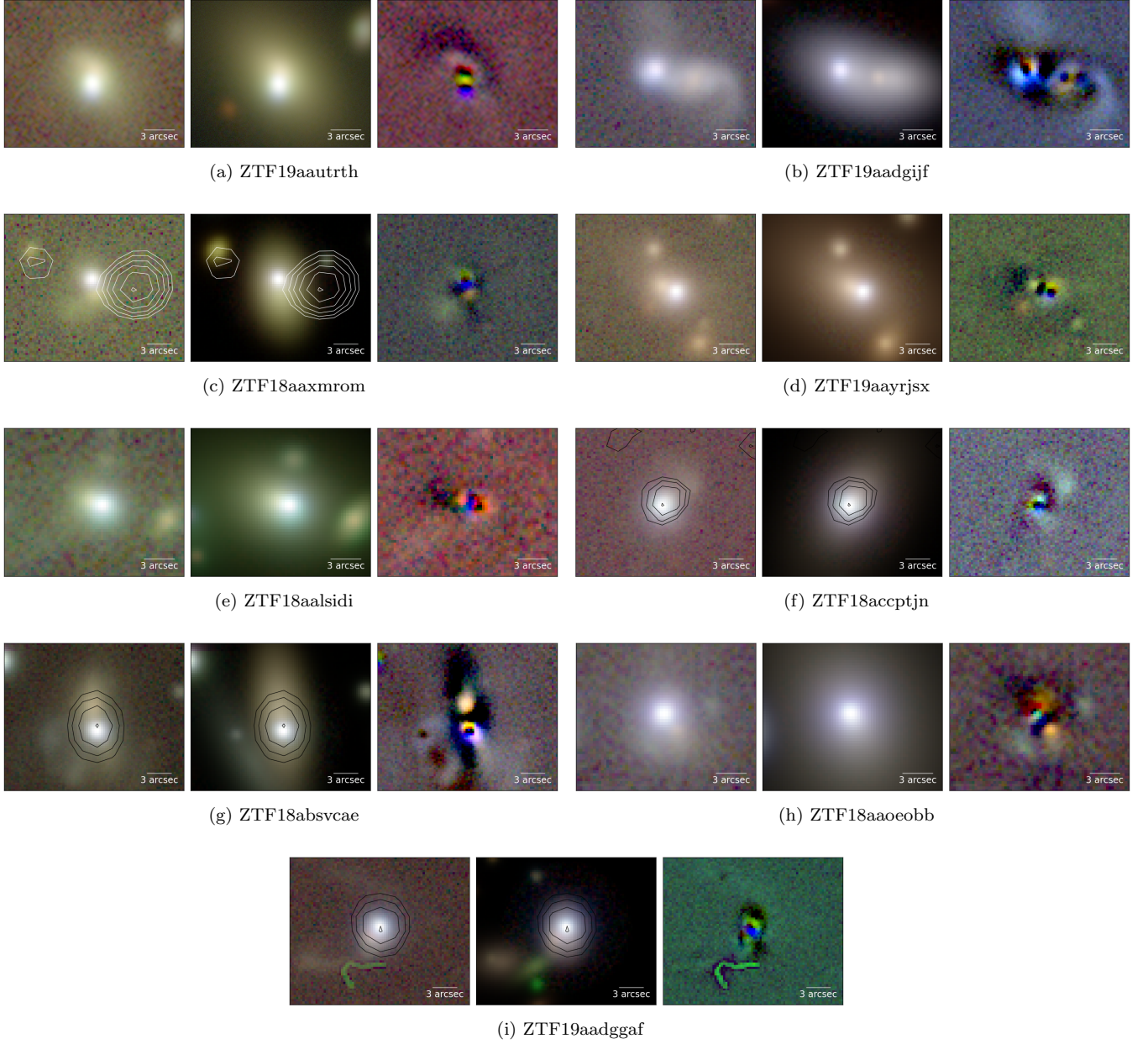


Figure 10. For each offset AGN: *Left*: Coadded g, r and z band Legacy Survey Images. *Middle*: Corresponding coadded **Tractor** model. *Right*: Image-model residuals showing galaxy tidal structures. For the 4 AGN with 20cm detections in FIRST (ZTF18aaxmrom, ZTF18acceptjn, ZTF18absvcae, ZTF19aadggaf) we overlay contours of the FIRST image.

tral Gaussian broad line in addition to the accretion disk model to produce a good fit to the data, which is not uncommon for double-peaked emitters (Strateva et al. 2003).

The best fit disk parameters from the $H\alpha$ and $H\beta$ fits of the 5 sources are shown in Table 5 and the models are shown in Figure 13. We note that our simple log likelihood minimization procedure for disk model fitting may not explore the whole parameter space and find the most optimal disk parameters. However, it serves to il-

lustrate that accretion disk emission models can account for the extra flux which is not well described by a Gaussian model. More rigorous fitting with MCMC could be used to better determine disk shape parameters and their uncertainties in the future.

In order to determine if the double-peaked emitter fraction of 63% amongst spatially offset AGN is substantially different from the entire variable ZTF AGN sample, we applied a multi-Gaussian fitting procedure to the 1923 variable AGN with archival SDSS spectra of

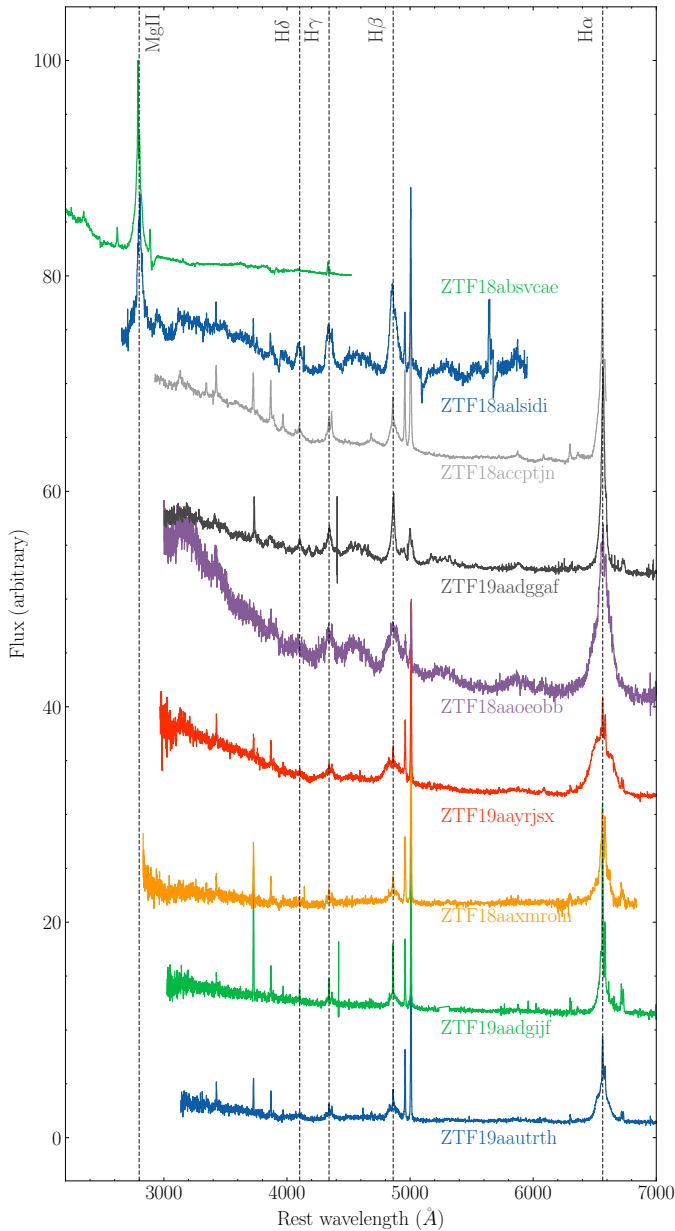


Figure 11. Spectra of the off-nuclear AGN, all from the SDSS archive except for ZTF18absvcae, ZTF18acctjn and ZTF18aalsidi which were observed with the Deveny spectrograph on LDT on 2020-09-13, 2020-09-15 and 2020-10-11 respectively. In all SDSS spectra, the fiber was centered on the AGN rather than the photometric center of the offset host galaxy. The Balmer series is shown with dotted lines.

the $H\alpha$ region which had broad Balmer lines to search for double-peaked emitters. We used the `pPXF` fitting procedure described in Section 2.4 to fit two different models. In the first model the spectrum was fit with a single broad Balmer line free to have a velocity offset to the narrow Balmer lines. This was fit at the same time as the narrow emission lines and stellar contin-

uum. In the second model the spectrum was fit with 3 broad lines: one with central velocity tied to the narrow Balmer lines, one with a velocity up to 6000 km s^{-1} relative to the narrow lines, and one with a velocity down to -6000 km s^{-1} . The widths of the 3 broad lines were not tied to each other.

To find double-peaked emitters from the fits of these two models, we required the χ^2 improvement from the multiple broad line model compared to the single broad line model to be > 250 . Then to be considered a double-peaked emitter candidate the 2 velocity offset $H\alpha$ broad lines were each required to have peak flux densities of $> 33\%$ of the narrow $H\alpha$ line and to have a velocity $> 500 \text{ km s}^{-1}$ from the narrow line velocity. As the effectiveness of this criteria depended on the relative brightness of narrow and broad lines, and the width of the broad line region, we visually inspected the 275 candidates found via this criteria and rejected 82, leaving 193 double-peaked emitters. We then visually inspected the remaining 1648 spectra to find any objects which may have missed by the criteria. We found 106 double-peaked spectra which were missed by the automatic classification scheme.

We therefore estimated that 299 of 1923, or 16%, of the variable AGN sample are double peaked emitters. This is much larger than the 3.6% fraction found for SDSS AGN by [Strateva et al. \(2003\)](#) with spectroscopic principal component analysis, suggesting that either variable AGN are more likely to be double-peaked emitters, or that our classification scheme is more likely to classify asymmetric broad line regions as double-peaked emitters than classification schemes used in previous studies. The 16% fraction of variable ZTF AGN with double-peaked broad Balmer lines is substantially smaller than 63% fraction seen in the spatially offset AGN sample.

3.3.2. Multi-wavelength properties

7 of the 9 offset AGN were detected in the ROSAT All Sky Survey and 4 AGN (ZTF18aaxmrom, ZTF18acctjn, ZTF18absvcae, ZTF19aadggaf) were detected at 20cm wavelengths in the FIRST radio survey, which uses the VLA to image a footprint largely coincident with SDSS to a detection sensitivity of 1 mJy ([Helfand et al. 2015](#)). The radio fluxes for these 4 AGN, as well as upper limits for another 4 which were within FIRST survey coverage but were not detected, are shown in Table 4. The contours of FIRST radio imaging are overlaid in Figure 10. ZTF18aaxmrom shows two radio lobes on either side of the galaxy which are likely the result of synchrotron emission from a jet. The other 3 AGN show radio point sources which coincide with the posi-

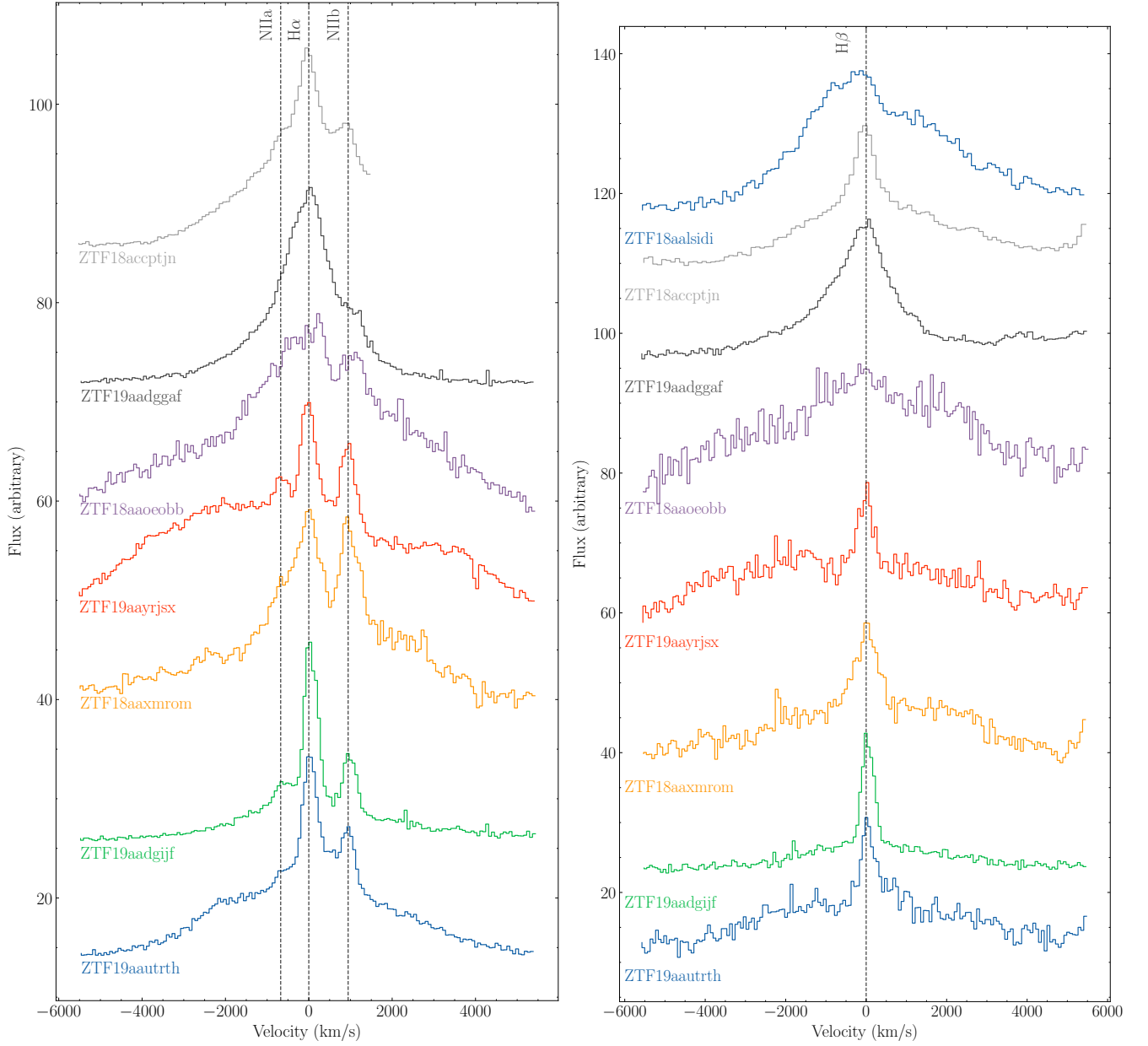


Figure 12. *Left:* $H\alpha$ broad line regions and *Right:* $H\beta$ broad line regions for the off-nuclear AGN candidates. 5 of the 8 off-nuclear AGN for which we have spectra of either the $H\alpha$ or $H\beta$ lines show asymmetric or double-peaked broad Balmer structures (ZTF18aaoeobb, ZTF19aayrsx, ZTF18aaxmrom, ZTF19aautrth, ZTF18aalsidi). The other three show standard Balmer broad lines (ZTF18accptjn, ZTF19aadggaf, ZTF19aadgijf).

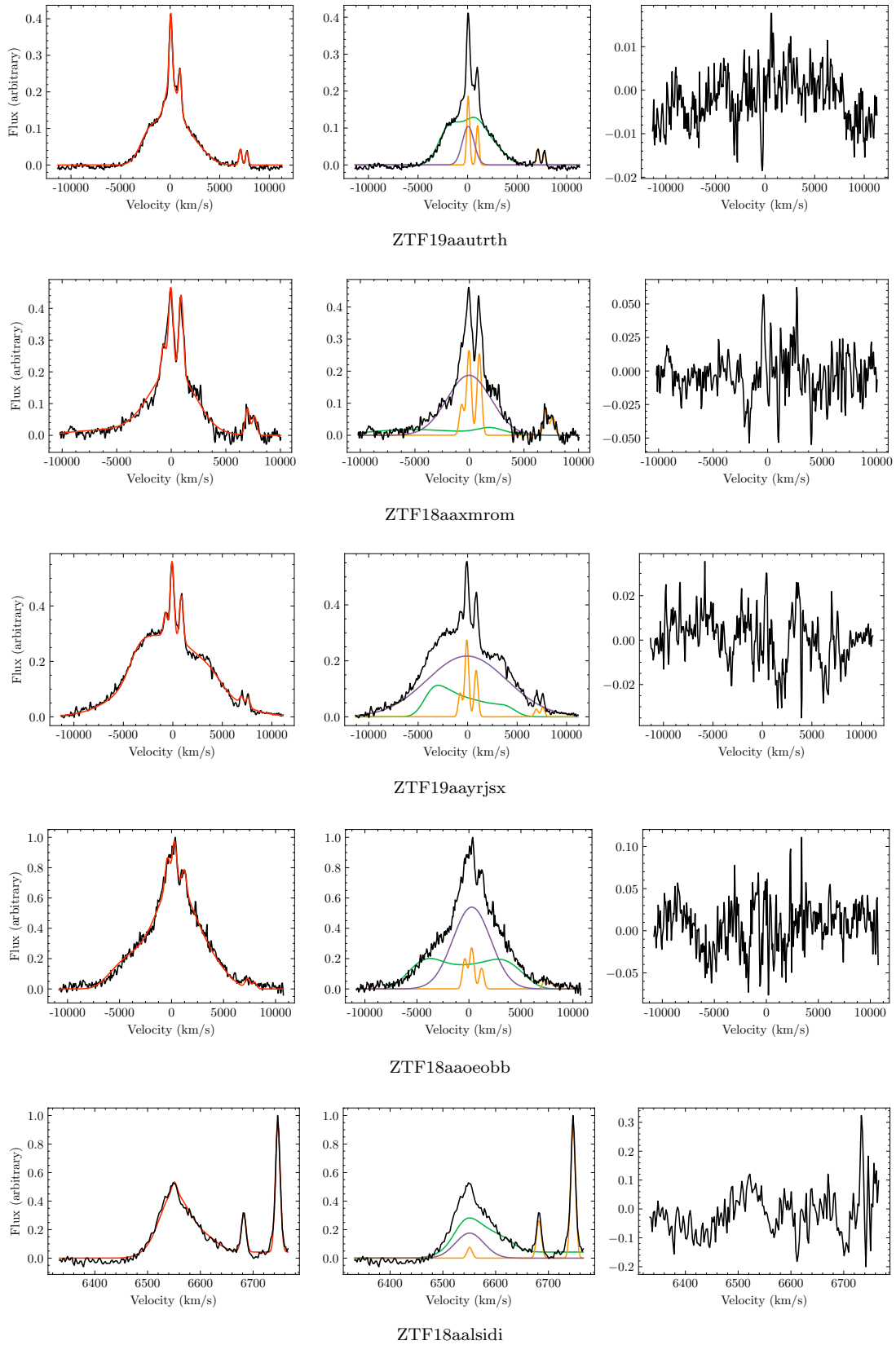


Figure 13. Best fit models of double-peaked $H\alpha$ accretion disk emission in off-nuclear AGN candidates ZTF19aaurth, ZTF18aaxmrom, ZTF19aayrjsx and ZTF18aaoebb and $H\beta$ disk emission in ZTF18aalsidi after subtracting a stellar continuum model derived from pPXF. On the left we show the continuum subtracted data in black and the best fit narrow emission line and accretion disk model in red. The middle plots show the separate narrow line (green), accretion disk (orange) and central broad line (purple) components of the fit. The rightmost plots show the error weighted flux residuals after subtracting the best fit model.

tion of the ZTF AGN. The radio loudness² is also shown in Table 4. Two AGN with radio emission are classified as radio-loud ($R > 10$, Kellermann et al. (1989)), 2 are radio-moderate, and 4 others have upper limits indicating that they are not radio loud.

The recoiling SMBH candidates from Chiaberge et al. (2017) and Lena et al. (2014) also demonstrate radio emission. As noted by Chiaberge et al. (2017), radio emission from a recoiling SMBH is not surprising given that the rapid spin needed to produce a relativistic jet can be produced by a binary black hole merger (Schnittman 2013; Hemberger et al. 2013) and that there is a link between radio loud AGNs and major galaxy mergers (Ivion et al. 2012; Chiaberge et al. 2015). Assuming that the radio jet axis matches the spin axis of the recoiling SMBH, we would also expect recoil velocity to be preferentially aligned with the radio jet (Lena et al. 2014).

3.3.3. Confirming the nature of the offset AGN

There are two main alternative explanations for the nature of these offset AGN. Firstly, they may be chance coincidences with disturbed background galaxies. Secondly, they may be AGN in mergers with compact or undermassive host galaxies such that extended emission around the AGN is very faint. Given the predicted frequency of SMBHs in tidally stripped nuclei which may appear as offset AGN (Voggel et al. 2019), it would not be surprising if a large fraction of our objects are in fact AGN in compact galaxies merging with a larger galaxy.

If the distribution of expected velocity and spatial offsets of the recoiling SMBH population follows the simulated distribution of Blecha et al. (2016), we would expect the velocity offsets of our sample to range from 300 to 3000 km s⁻¹ given their large spatial offsets. Spectroscopic fitting (Section 3.3.1) shows that the 3 AGN with Gaussian broad lines have velocity offsets ranging from 344-608 km/s (see Table 4) and the remainder of the sample with double-peaked emission have spectra consistent with a broad line region and accretion disk close to rest velocity (Figure 13). It is possible for large separation recoiling SMBHs to show small line of sight velocities, depending on their orbital trajectory at the time. However, if our objects are indeed all recoiling SMBHs, it is surprising that a fraction of the 9 do not have larger > 600 km/s line of sight velocities. The observed velocities therefore argue against the recoil hypothesis for these objects.

$$^2 R = \frac{L_{5\text{GHz}(\nu)}}{L_{4400\text{\AA}}(\nu)}$$

Much further work is therefore required to investigate the recoiling SMBH hypothesis for any of these offset AGN. In order to rule out that they are not AGN in mergers with undermassive host galaxies, Hubble Space Telescope (HST) IR imaging could be used to search for a second extended region of old stellar emission around the offset AGN. HST imaging would be well complemented by IFU observations to map the positions of the narrow line emission relative to the AGN broad line emission (Chiaberge et al. 2018).

Chandra X-ray imaging to search for a second obscured AGN would also be essential for these candidates (Comerford et al. 2017). If we observe just one X-ray point source, we can produce upper limits on the X-ray luminosity of a second AGN and show that they rule out the presence of an obscured AGN using the [O III] flux from the host center. This would strongly favor the recoiling black hole hypothesis.

As spatial offsets greater than 1.5 kpc are only possible in simulations where misaligned SMBH binary spins can occur (Blecha et al. 2016), if even a subset of our candidates could be confirmed as recoiling SMBHs, their extreme spatial offsets could demonstrate the existence of SMBH binaries with misaligned spins.

3.4. SDSS 1133 (ZTF19aafmjfw)

SDSS1133, a variable object and recoiling SMBH candidate discovered by Koss et al. (2014) was identified by our search pipeline during a dramatic flare in ZTF, labeled ZTF19aafmjfw. This object has a long history of variability and the alternative scenarios proposed for its origin include: an LBV which exploded as a type II supernova in 2001, an LBV which continues to exhibit giant eruptions, or an offset AGN with flaring and stochastic variability (Koss et al. 2014). It was first observed in 1950 with the 103aO DSS plate at a magnitude of 18.6 and was again observed at comparable magnitudes in 1994 and 1999. It was then discovered to be flaring to $g = 16.4$ in SDSS on 2001-12-18 and 2002-04-01 but had faded to $g = 18.7$ by 2003-03-09. It was observed at a minimum of $g = 20.18$ with PS1 on 2012-02-22 after which it brightened to $g = 19.3$ on 2014-01-20 (Koss et al. 2014).

We searched for this object in archival data from the Catalina Real-time Transient Survey (Drake et al. 2009) and found that this object demonstrated another small scale flaring event in 2014. After showing no detectable V-band activity between 2006-02-03 and 2014-05-12, it brightened by $V=0.22$ between 2014-04-28 and 2014-06-05. After a gap in observations it had faded again by 2014-12-27.

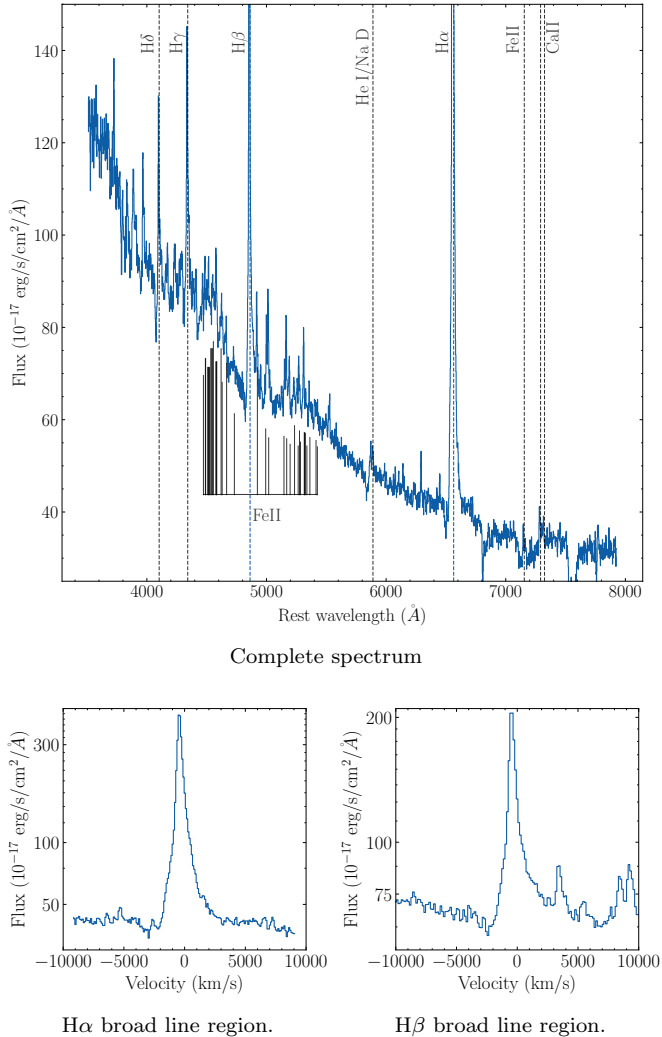


Figure 14. Spectrum of the flare ZTF19aafmjfw from SDSS1133 taken on 2019-05-29 with the Deveny spectrograph on the Lowell Discovery Telescope. Subplot (a) shows the complete spectrum, while (b) and (c) show close-ups of H α and H β in log scale to emphasize the P-cygni profile structure.

This object showed no evidence of variability in ZTF prior to 2019-04-07, when it became detectable at $m_g = 20.21$ on 2019-04-07, flared to $m_g = 17.00$ on 2019-06-05, and faded again by 2019-12-16. The ZTF flare ZTF19aafmjfw is shown in Figure 9. The 3 magnitude change in brightness seen in ZTF was of comparable scale to the SDSS flare in 2001. The ZTF rebrightening of this object suggests that the transient emission from 2001 to 2013 was not from a Type II n supernova with a variable star progenitor.

We obtained a spectrum of ZTF19aafmjfw with the Deveny spectrograph on LDT on 2019-05-29 when the object was bright. The spectrum showed the return of

broad H α and H β absorption lines blue-shifted over a 2000 to 8000 km s $^{-1}$ range (Figure 14), as had been observed in December 2013 and January 2014 (Koss et al. 2014). It also showed the return of [Fe II] $\lambda 7155$ and [Ca II] $\lambda\lambda 7291, 7324$ lines – features which have been seen only occasionally in AGN (e.g. Phillips (1976)) spectra, late-time supernovae spectra (e.g. Filippenko 1997; Pastorello et al. 2019) and outbursting LBVs (Solovyeva et al. 2019).

The 100 day timescale, 3-4 magnitude flux change and high velocity absorption lines make this outburst comparable to the transient SN2009ip, a supersonic stellar explosion from a $> 60M_{\odot}$ star (Smith et al. 2010; Foley et al. 2011). The presence of broad Balmer, He I and Na D lines and a very blue continuum also links ZTF19aafmjfw to SN2009ip. However, SN2009ip did not show the forest of Fe group lines and [Ca II] $\lambda\lambda 7291, 7324$ emission lines which we observe in ZTF19aafmjfw. Such Fe and Ca features were observed in the stellar outburst UGC 2773 OT2009-1, which is likely to be either a giant LBV eruption or extreme S Dor variability with circumstellar dust eruption (Smith et al. 2010; Foley et al. 2011). However, UGC 2773 OT2009-1 had a cooler temperature and weaker Balmer emission than SN2009ip, suggesting that it was a subsonic outburst rather than an explosion.

As ZTF19aafmjfw has spectroscopic features in common with SN2009ip which could indicate a high velocity explosion (blueshifted Balmer absorption lines, bright and broad Balmer emission and a blue continuum) and other spectroscopic features in common with UGC 2773 OT2009-1 which could indicate a stellar outburst within a circumstellar dust envelope (FeII and CaII emission lines), the observed spectroscopic features may be due to a non-terminal, supersonic LBV outburst in a dusty circumstellar environment rather.

We note, however, that the presence of such a rare and massive star in a dwarf galaxy which does not show recent star formation in Keck adaptive optics and HST imaging is very unlikely (Koss et al. 2014), and while the observed combination of spectroscopic features would be unusual for a standard AGN, they may be induced by a recoil event. We therefore cannot yet completely rule out the recoil hypothesis in favor of an LBV or supernova imposter explanation.

4. CONCLUSIONS

We have described a novel search strategy for the discovery of spatially offset AGN from recoiling SMBHs and ongoing galaxy mergers. This strategy uses ZTF difference imaging to find variable AGN and Tractor forward modeling to determine the AGN position across

multiple ZTF epochs and in deep, high resolution Legacy Survey images.

We have found a sample of 52 AGN in galaxy mergers and a sample of 9 AGN which may be spatially offset from their host galaxies. 5 of the 8 offset AGN for which we have spectra of the $H\alpha$ or $H\beta$ broad line regions show irregularly shaped broad Balmer lines with velocity offset broad peaks. These structures may arise due to emission from an unobscured, relativistic Keplerian accretion disk around the AGN. The 63% fraction of double-peaked emitters in the offset AGN sample is much larger than the 16% observed for the whole ZTF AGN sample. The remaining 3 offset AGN with spectra have Gaussian broad lines with velocities ranging in magnitude from 200 to 600 km/s relative to the narrow emission lines.

Our search strategy detected the variable object and recoiling SMBH candidate SDSS1133 (Koss et al. 2014) when it re-brightened by 3 magnitudes in ZTF. Follow-up spectra showed the return of a blue continuum, blue-shifted absorption lines and [Fe II] λ 7155 and [Ca II] λ 7291, 7324 forbidden lines, suggesting that the source may be either a luminous blue variable star which continues to show non-terminal outbursts or a recoiling AGN which continues to show variability.

Further multi-wavelength follow-up is required to confirm that our recoiling SMBH candidates are not AGN with undermassive hosts in mergers or chance coincidences with disturbed background galaxies. If even a subset of our offset AGN candidates are confirmed to be recoiling black holes, their large spatial offsets could show that SMBH binaries with misaligned spins are able to form. Such binaries may be detectable at a later stage of evolution by the Laser Interferometer Space Antenna (LISA) and would provide strong constraints on models of SMBH binary formation. The success of our variability-based search strategy with ZTF suggests that future searches for offset AGN with the Vera Rubin Observatory (Ivezić et al. 2019) may yield large populations of recoiling SMBH candidates and AGN in mergers.

5. ACKNOWLEDGEMENTS

We thank the anonymous reviewer for their helpful feedback. S. Gezari is supported in part by NSF grant 1616566. This work relied on observations obtained with the Samuel Oschin Telescope 48-inch and the 60-inch Telescope at the Palomar Observatory as part of the Zwicky Transient Facility (ZTF) survey. ZTF is supported by the National Science Foundation under Grant No. AST-1440341 and a collaboration including Caltech, IPAC, the Weizmann Institute for Science, the Oskar Klein Center at Stockholm University, the University

of Maryland, the University of Washington, Deutsches Elektronen-Synchrotron and Humboldt University, Los Alamos National Laboratories, the TANGO Consortium of Taiwan, the University of Wisconsin at Milwaukee, and Lawrence Berkeley National Laboratories. Operations are conducted by COO, IPAC, and UW. SED Machine is based upon work supported by the National Science Foundation under Grant No. 1106171. This work was supported by the GROWTH project funded by the National Science Foundation under Grant No 1545949. MMK acknowledges generous support from the David and Lucille Packard Foundation.

This research used resources of the National Energy Research Scientific Computing Center, a DOE Office of Science User Facility supported by the Office of Science of the U.S. Department of Energy under Contract No. DE-AC02-05CH11231. P.E.N. acknowledges support from the DOE under grant DE-AC02-05CH11231, Analytical Modeling for Extreme-Scale Computing Environments.

These results made use of the Lowell Discovery Telescope (LDT) at Lowell Observatory. Lowell is a private, non-profit institution dedicated to astrophysical research and public appreciation of astronomy and operates the LDT in partnership with Boston University, the University of Maryland, the University of Toledo, Northern Arizona University and Yale University. The upgrade of the DeVeny optical spectrograph has been funded by a generous grant from John and Ginger Giovale.

This project used data obtained with the Dark Energy Camera (DECam), which was constructed by the Dark Energy Survey (DES) collaboration. Funding for the DES Projects has been provided by the U.S. Department of Energy, the U.S. National Science Foundation, the Ministry of Science and Education of Spain, the Science and Technology Facilities Council of the United Kingdom, the Higher Education Funding Council for England, the National Center for Supercomputing Applications at the University of Illinois at Urbana-Champaign, the Kavli Institute of Cosmological Physics at the University of Chicago, Center for Cosmology and Astro-Particle Physics at the Ohio State University, the Mitchell Institute for Fundamental Physics and Astronomy at Texas A&M University, Financiadora de Estudos e Projetos, Fundacao Carlos Chagas Filho de Amparo, Financiadora de Estudos e Projetos, Fundacao Carlos Chagas Filho de Amparo a Pesquisa do Estado do Rio de Janeiro, Conselho Nacional de Desenvolvimento Cientifico e Tecnol6gico and the Ministerio da Ciencia, Tecnologia e Inovacao, the Deutsche Forschungsgemeinschaft and the Collaborating Institu-

tions in the Dark Energy Survey. The Collaborating Institutions are Argonne National Laboratory, the University of California at Santa Cruz, the University of Cambridge, Centro de Investigaciones Energeticas, Medioambientales y Tecnologicas-Madrid, the University of Chicago, University College London, the DES-Brazil Consortium, the University of Edinburgh, the Eidgenossische Technische Hochschule (ETH) Zurich, Fermi National Accelerator Laboratory, the University of Illinois at Urbana-Champaign, the Institut de Ciencies de l'Espai (IEEC/CSIC), the Institut de Fisica d'Altes Energies, Lawrence Berkeley National Labora-

tory, the Ludwig-Maximilians Universitat Munchen and the associated Excellence Cluster Universe, the University of Michigan, the National Optical Astronomy Observatory, the University of Nottingham, the Ohio State University, the University of Pennsylvania, the University of Portsmouth, SLAC National Accelerator Laboratory, Stanford University, the University of Sussex, and Texas A&M University.

Software: Ampel (Nordin et al. 2019), Astropy (Robitaille et al. 2013; Price-Whelan et al. 2018), cat-sHTM (Soumagnac & Ofek 2018), extcats (github.com/MatteoGiomi/extcats), GROWTH Marshal (Kasliwal et al. 2019), The Tractor (Lang et al. 2016).

REFERENCES

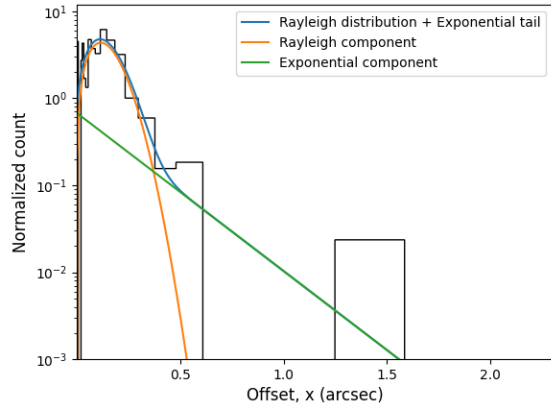
- Afanasiev, A. V., Chilingarian, I. V., Mieske, S., et al. 2018, *Monthly Notices of the Royal Astronomical Society*, 477, 4856
- Ahn, C. P., Seth, A. C., Brok, M. d., et al. 2017, *The Astrophysical Journal*, 839, 72
- Ahn, C. P., Seth, A. C., Cappellari, M., et al. 2018, *The Astrophysical Journal*, 858, 102
- Allen, J. T., Schaefer, A. L., Scott, N., et al. 2015, *Monthly Notices of the Royal Astronomical Society*, 451, 2780
- Assef, R. J., Stern, D., Kochanek, C. S., et al. 2013, *Astrophysical Journal*, 772, doi:10.1088/0004-637X/772/1/26
- Baldassare, V. F., Geha, M., & Greene, J. 2018, *The Astrophysical Journal*, 868, 152
- . 2020, *The Astrophysical Journal*, 896, 10
- Barbary, K., rbiswas4, Goldstein, D., et al. 2016, *sncosmo/sncosmo: v1.4.0*, , , doi:10.5281/zenodo.168220
- Bellm, E. C., Kulkarni, S. R., Graham, M. J., et al. 2019, *Publications of the Astronomical Society of the Pacific*, 131, 018002
- Bellovary, J. M., Governato, F., Quinn, T. R., et al. 2010, *Astrophysical Journal Letters*, 721, 148
- Berczik, P., Merritt, D., Spurzem, R., & Bischof, H.-P. 2006, *The Astrophysical Journal*, 642, L21
- Blecha, L., Cox, T. J., Loeb, A., & Hernquist, L. 2011, *Monthly Notices of the Royal Astronomical Society*, 412, 2154
- Blecha, L., & Loeb, A. 2008, *Monthly Notices of the Royal Astronomical Society*, doi:10.1111/j.1365-2966.2008.13790.x
- Blecha, L., Sijacki, D., Kelley, L. Z., et al. 2016, *Monthly Notices of the Royal Astronomical Society*, 456, 961
- Bogdanović, T., Reynolds, C. S., & Miller, M. C. 2007, *The Astrophysical Journal*, 661, L147
- Boles, T. 2009, *Central Bureau Electronic Telegrams*, 1931, 1
- Boller, T., Freyberg, M. J., Trümper, J., et al. 2016, *Astronomy & Astrophysics*, 588, A103
- Bonning, E. W., Shields, G. A., & Salviander, S. 2007, *The Astrophysical Journal*, 666, L13
- Brescia, M., Cavuoti, S., & Longo, G. 2015, *Monthly Notices of the Royal Astronomical Society*, 450, 3893
- Burke, C. J., Baldassare, V. F., Liu, X., et al. 2020, *The Astrophysical Journal*, 894, L5
- Butler, N. R., & Bloom, J. S. 2011, *Astronomical Journal*, 141, doi:10.1088/0004-6256/141/3/93
- Campanelli, M., Lousto, C. O., Zlochower, Y., & Merritt, D. 2007, *Physical Review Letters*, 98, 8
- Canalizo, G., & Stockton, A. 2001, *The Astrophysical Journal*, 555, 719
- Cappellari, M. 2017, *Monthly Notices of the Royal Astronomical Society*, 466, 798
- Cappellari, M., & Emsellem, E. 2004, *Publications of the Astronomical Society of the Pacific*, 116, 138
- Chambers, K. C., Magnier, E. A., Metcalfe, N., et al. 2016, *ArXiv e-prints*
- Chen, K., & Halpern, J. P. 1989, *The Astrophysical Journal*, 344, 115
- Chen, K., Halpern, J. P., & Filippenko, A. V. 1989, *The Astrophysical Journal*, 339, 742
- Chiaberge, M., Gilli, R., M. Lotz, J., & Norman, C. 2015, *The Astrophysical Journal*, 806, 147
- Chiaberge, M., Tremblay, G. R., Capetti, A., & Norman, C. 2018, *The Astrophysical Journal*, 861, 56
- Chiaberge, M., Ely, J. C., Meyer, E. T., et al. 2017, *Astronomy & Astrophysics*, 600, A57

- Cid Fernandes, R., Stasińska, G., Mateus, A., & Vale Asari, N. 2011, *Monthly Notices of the Royal Astronomical Society*, 413, 1687
- Civano, F., Elvis, M., Lanzuisi, G., et al. 2010, *Astrophysical Journal*, 717, 209
- Comerford, J. M., Barrows, R. S., Greene, J. E., & Pooley, D. 2017, *The Astrophysical Journal*, 847, 41
- Comerford, J. M., & Greene, J. E. 2014, *Astrophysical Journal*, 789, doi:10.1088/0004-637X/789/2/112
- Comerford, J. M., Gerke, B. F., Newman, J. A., et al. 2009, *Astrophysical Journal*, 698, 956
- Condon, J. J., Cotton, W. D., Greisen, E. W., et al. 1998, *The Astronomical Journal*, 115, 1693
- Condon, J. J., Darling, J., Kovalev, Y. Y., & Petrov, L. 2017, *The Astrophysical Journal*, 834, 184
- Congiu, E., Berton, M., Giroletti, M., et al. 2017, *Astronomy & Astrophysics*, 603, A32
- Dekany, R., Smith, R. M., Riddle, R., et al. 2020, *Publications of the Astronomical Society of the Pacific*, 132, 038001
- Dey, A., Schlegel, D. J., Lang, D., et al. 2019, *The Astronomical Journal*, 157, 168
- Drake, A. J., Djorgovski, S. G., Mahabal, A., et al. 2009, *Astrophysical Journal*, 696, 870
- Duev, D. A., Mahabal, A., Masci, F. J., et al. 2019, *Monthly Notices of the Royal Astronomical Society*, 489, 3582
- Eracleous, M., & Halpern, J. P. 1994, *The Astrophysical Journal Supplement Series*, 90, 1
- Escala, A., Larson, R. B., Coppi, P. S., & Mardones, D. 2005, *The Astrophysical Journal*, 630, 152
- Evans, I. N., Primini, F. A., Glotfelty, K. J., et al. 2010, *Astrophysical Journal, Supplement Series*, 189, 37
- Farrell, S. A., Webb, N. A., Barret, D., Godet, O., & Rodrigues, J. M. 2009, *Nature*, 460, 73
- Ferrarese, L., & Ford, H. 2005, *Space Science Reviews*, 116, 523
- Filippenko, A. V. 1997, *Annual Review of Astronomy and Astrophysics*, 35, 309
- Flesch, E. W. 2015, *Publications of the Astronomical Society of Australia*, 32, doi:10.1017/pasa.2015.10
- Foley, R. J., Berger, E., Fox, O., et al. 2011, *Astrophysical Journal*, 732, doi:10.1088/0004-637X/732/1/32
- Fu, H., Myers, A. D., Djorgovski, S. G., & Yan, L. 2011a, *Astrophysical Journal*, 733, 1
- Fu, H., Zhang, Z. Y., Assef, R. J., et al. 2011b, *Astrophysical Journal Letters*, 740, 1
- Glikman, E., Simmons, B., Maily, M., et al. 2015, *Astrophysical Journal*, 806, 218
- Graham, M. J., Kulkarni, S. R., Bellm, E. C., et al. 2019, *Publications of the Astronomical Society of the Pacific*, 131, 1
- Halpern, J. P., Impey, C. D., Bothun, G. D., et al. 1986, *The Astrophysical Journal*, 302, 711
- Helfand, D. J., White, R. L., & Becker, R. H. 2015, *The Astrophysical Journal*, 801, 26
- Hemberger, D. A., Lovelace, G., Lored, T. J., et al. 2013, *Physical Review D*, 88, 064014
- Huang, Y., Liu, X. W., Yuan, H. B., et al. 2014, *Monthly Notices of the Royal Astronomical Society*, 439, 2927
- Hwang, H.-C., Shen, Y., Zakamska, N., & Liu, X. 2020, *The Astrophysical Journal*, 888, 73
- Ivezić, v., Kahn, S. M., Tyson, J. A., et al. 2019, *The Astrophysical Journal*, 873, 111
- Ivison, R. J., Smail, I., Amblard, A., et al. 2012, *Monthly Notices of the Royal Astronomical Society*, 425, 1320
- Jarrett, T. H., Chester, T., Cutri, R., et al. 2000, *The Astronomical Journal*, 119, 2498
- Jonker, P. G., Torres, M. A., Fabian, A. C., et al. 2010, *Monthly Notices of the Royal Astronomical Society*, 407, 645
- Kalfountzou, E., Lleo, M. S., & Trichas, M. 2017, *The Astrophysical Journal*, 851, L15
- Kasliwal, M. M., Cannella, C., Bagdasaryan, A., et al. 2019, *Publications of the Astronomical Society of the Pacific*, 131, 038003
- Kauffmann, G., Heckman, T. M., Tremonti, C., et al. 2003, *Monthly Notices of the Royal Astronomical Society*, 346, 1055
- Keel, W. C., Bennert, V. N., Pancoast, A., et al. 2019, *Monthly Notices of the Royal Astronomical Society*, 483, 4847
- Kellermann, K. I., Sramek, R., Schmidt, M., Shaffer, D. B., & Green, R. 1989, *The Astronomical Journal*, 98, 1195
- Kim, D.-C., Yoon, I., Privon, G. C., et al. 2017, *The Astrophysical Journal*, 840, 71
- Kocevski, D. D., Brightman, M., Nandra, K., et al. 2015, *Astrophysical Journal*, 814, 104
- Komossa, S., Burwitz, V., Hasinger, G., et al. 2003, *The Astrophysical Journal*, 582, L15
- Kormendy, J., & Richstone, D. 1995, *Annual Review of Astronomy and Astrophysics*, 33, 581
- Koss, M., Mushotzky, R., Treister, E., et al. 2012, *The Astrophysical Journal*, 746, L22
- . 2011, *Astrophysical Journal Letters*, 735, doi:10.1088/2041-8205/735/2/L42
- Koss, M., Blecha, L., Mushotzky, R., et al. 2014, *Monthly Notices of the Royal Astronomical Society*, 445, 515

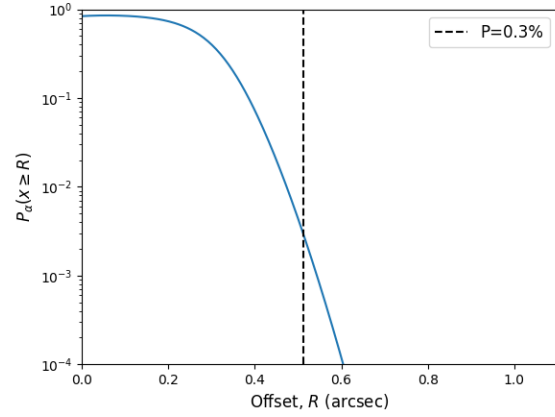
- Koss, M. J., Blecha, L., Bernhard, P., et al. 2018, *Nature*, 563, 214
- Lang, D., Hogg, D. W., & Mykytyn, D. 2016, *The Tractor: Probabilistic astronomical source detection and measurement*, *Astrophysics Source Code Library*, ,
- Lena, D., Robinson, A., Marconi, A., et al. 2014, *Astrophysical Journal*, 795, doi:10.1088/0004-637X/795/2/146
- Liu, X., Guo, H., Shen, Y., Greene, J. E., & Strauss, M. A. 2018, *The Astrophysical Journal*, 862, 29
- Liu, X., Shen, Y., Strauss, M. A., & Greene, J. E. 2010, *Astrophysical Journal*, 708, 427
- Liu, X., Shen, Y., Strauss, M. A., & Hao, L. 2011, *The Astrophysical Journal*, 737, 101
- Lodato, G., & Facchini, S. 2013, *Monthly Notices of the Royal Astronomical Society*, 433, 2157
- Loeb, A. 2007, *Phys. Rev. Lett.*, 99, 41103
- Mahabal, A., Rebbapragada, U., Walters, R., et al. 2019, *Publications of the Astronomical Society of the Pacific*, 131, 038002
- Masci, F. J., Laher, R. R., Rusholme, B., et al. 2019, *Publications of the Astronomical Society of the Pacific*, 131, 1
- Maza, J., Hamuy, M., Antezana, R., et al. 2009, *Central Bureau Electronic Telegrams*, 1928, 1
- McLure, R. J., & Dunlop, J. S. 2002, *Monthly Notices of the Royal Astronomical Society*, 331, 795
- Mezcua, M., & Domínguez Sánchez, H. 2020, *The Astrophysical Journal*, 898, L30
- Miller, A. A., Kulkarni, M. K., Cao, Y., et al. 2017, *The Astronomical Journal*, 153, 73
- Milosavljević, M., & Merritt, D. 2001, *The Astrophysical Journal*, 563, 34
- Nordin, J., Brinnel, V., van Santen, J., et al. 2019, *Astronomy & Astrophysics*, 631, A147
- Pastorello, A., Reguitti, A., Morales-Garoffolo, A., et al. 2019, *Astronomy and Astrophysics*, 628, 1
- Patterson, M. T., Bellm, E. C., Rusholme, B., et al. 2019, *Publications of the Astronomical Society of the Pacific*, 131, 18001
- Phillips, M. M. 1976, *The Astrophysical Journal*, 208, 37
- Plotkin, R. M., Anderson, S. F., Hall, P. B., et al. 2008, *The Astronomical Journal*, 135, 2453
- Price-Whelan, A. M., Sipőcz, B. M., Günther, H. M., et al. 2018, *The Astronomical Journal*, 156, 123
- Reines, A. E., & Comastri, A. 2016, *Publications of the Astronomical Society of Australia*, 33, doi:10.1017/pasa.2016.46
- Reines, A. E., Condon, J. J., Darling, J., & Greene, J. E. 2020, *The Astrophysical Journal*, 888, 36
- Robinson, A., Young, S., Axon, D. J., Kharb, P., & Smith, J. E. 2010, *Astrophysical Journal Letters*, 717, 122
- Robitaille, T. P., Tollerud, E. J., Greenfield, P., et al. 2013, *Astronomy and Astrophysics*, 558, 1
- Rots, A. H., & Budavári, T. 2011, *Astrophysical Journal, Supplement Series*, 192, 1
- Schnittman, J. D. 2013, *Classical and Quantum Gravity*, 30, 244007
- Seth, A. C., Van Den Bosch, R., Mieske, S., et al. 2014, *Nature*, 513, 398
- Shabala, S. S., Deller, A., Kaviraj, S., et al. 2017, *Monthly Notices of the Royal Astronomical Society*, 464, 4706
- Shen, Y., Hwang, H.-c., Zakamska, N., & Liu, X. 2019, *The Astrophysical Journal*, 885, L4
- Shields, G. A., Bonning, E. W., & Salviander, S. 2009a, *Astrophysical Journal*, 696, 1367
- Shields, G. A., Rosario, D. J., Smith, K. L., et al. 2009b, *Astrophysical Journal*, 707, 936
- Skipper, C. J., & Browne, I. W. A. 2018, *Monthly Notices of the Royal Astronomical Society*, 475, 5179
- Smith, K. L., Shields, G. A., Bonning, E. W., et al. 2010, *Astrophysical Journal*, 716, 866
- Solovyeva, Y., Vinokurov, A., Fabrika, S., et al. 2019, *Monthly Notices of the Royal Astronomical Society: Letters*, 484, L24
- Soumagnac, M. T., & Ofek, E. O. 2018, *Publications of the Astronomical Society of the Pacific*, 130, 1
- Steinhardt, C. L., Schramm, M., Silverman, J. D., et al. 2012, *Astrophysical Journal*, 759, 1
- Strateva, I. V., Strauss, M. A., Hao, L., et al. 2003, *The Astronomical Journal*, 126, 1720
- Surace, J. A., Sanders, D. B., Vacca, W. D., Veilleux, S., & Mazzarella, J. M. 1998, *The Astrophysical Journal*, 492, 116
- Tachibana, Y., & Miller, A. A. 2018, *Publications of the Astronomical Society of the Pacific*, 130, 128001
- Treister, E., Schawinski, K., Urry, C. M., & Simmons, B. D. 2012, *Astrophysical Journal Letters*, 758, doi:10.1088/2041-8205/758/2/L39
- van Velzen, S., Gezari, S., Cenko, S. B., et al. 2019, *The Astrophysical Journal*, 872, 198
- van Velzen, S., Gezari, S., Hammerstein, E., et al. 2020, *ArXiv e-prints*
- Van Wassenhove, S., Volonteri, M., Mayer, L., et al. 2012, *Astrophysical Journal Letters*, 748, 5
- Voggel, K. T., Seth, A. C., Baumgardt, H., et al. 2019, *The Astrophysical Journal*, 871, 159
- Volonteri, M. 2007, *The Astrophysical Journal*, 663, L5
- Volonteri, M., Gültekin, K., & Dotti, M. 2010, *Monthly Notices of the Royal Astronomical Society*, 404, 2143

- Volonteri, M., & Madau, P. 2008, *The Astrophysical Journal*, 687, L57
- Volonteri, M., & Natarajan, P. 2009, *Monthly Notices of the Royal Astronomical Society*, 400, 1911
- Volonteri, M., & Perna, R. 2005, *Monthly Notices of the Royal Astronomical Society*, 358, 913
- Wang, J. M., Chen, Y. M., Hu, C., et al. 2009, *Astrophysical Journal*, 705, 76
- Xu, D., & Komossa, S. 2009, *Astrophysical Journal*, 705, 20

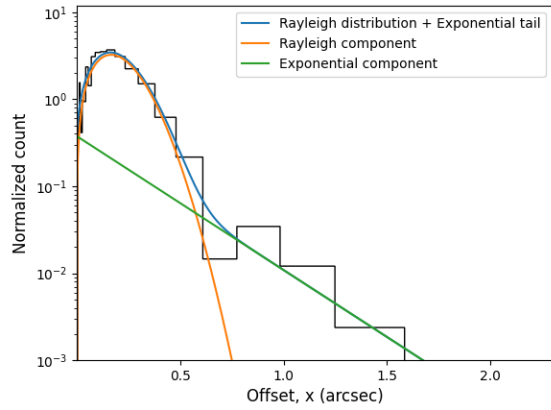
APPENDIX



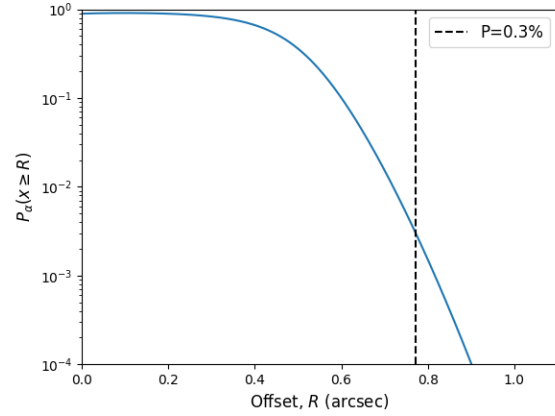
(a) AGN with a peak difference magnitude between 15 and 18.



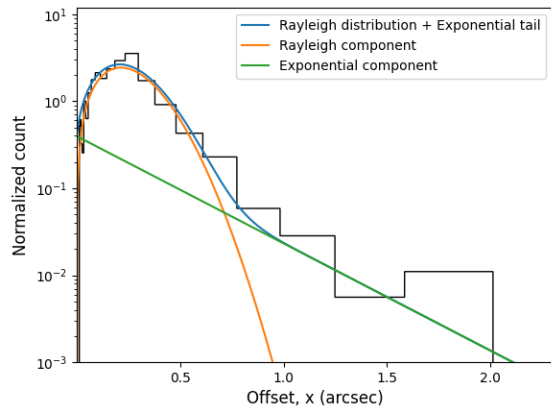
(b) 3σ offset = $0.551''$.



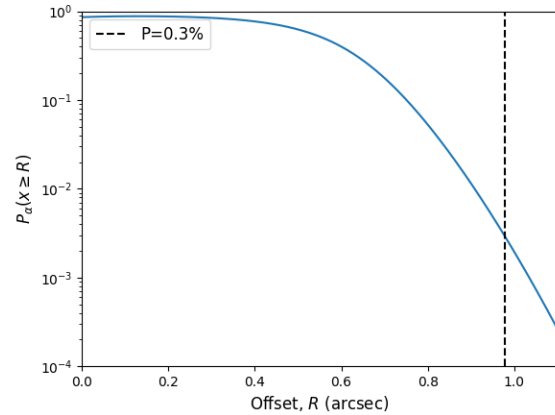
(c) AGN with a peak difference magnitude between 18 and 19.5.



(d) 3σ offset = $0.773''$.

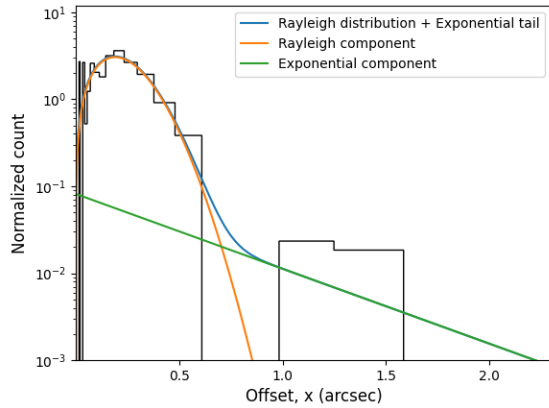


(e) AGN with a peak difference magnitude between 19.5 and 23.

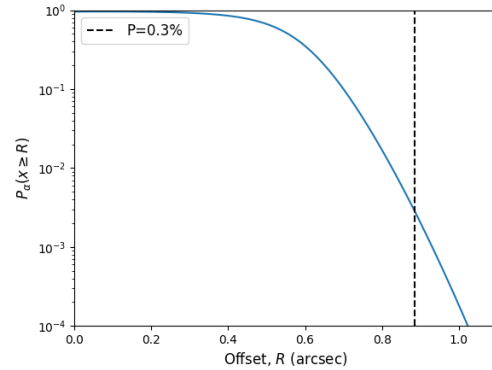
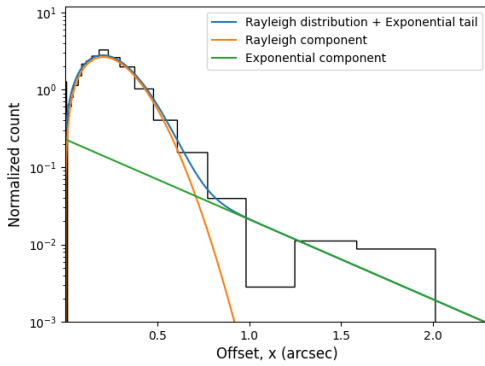


(f) 3σ offset = $0.976''$.

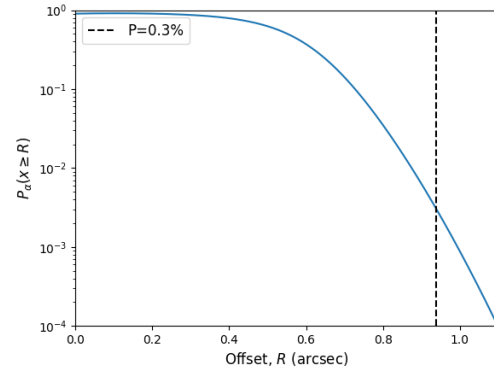
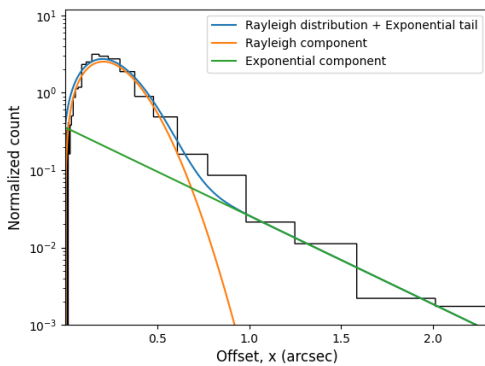
Figure 15. **Left:** Normalized histogram with logarithmic bins for magnitude weighted AGN-host offsets obtained from r-band ZTF difference images. The best-fit model of a mixture distribution with Rayleigh and exponential components are shown. **Right:** Probability that an offset greater than R is drawn from the Rayleigh component of the mixture distribution shown in (a) instead of the exponential component. The offset where this probability is 0.3% is shown with a dashed line.



(a) AGN with a peak difference magnitude between 15 and 18.

(b) 3σ offset = $0.959''$.

(c) AGN with a peak difference magnitude between 18 and 19.5

(d) 3σ offset = $1.009''$.

(e) AGN with a peak difference magnitude between 19.5 and 23.

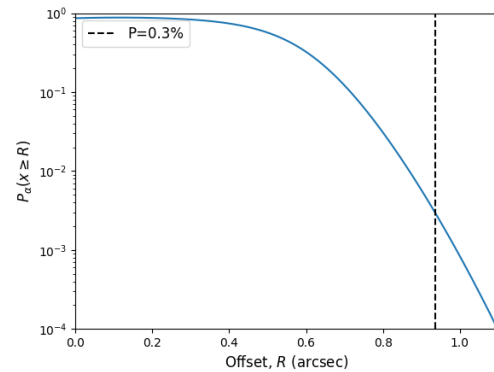
(f) 3σ offset = $0.946''$.

Figure 16. **Left:** Normalized histogram with logarithmic bins for magnitude weighted AGN-host offsets obtained from g-band ZTF difference images. The best-fit model of a mixture distribution with Rayleigh and exponential components are shown. **Right:** Probability that an offset greater than R is drawn from the Rayleigh component of the mixture distribution shown in (a) instead of the exponential component. The offset where this probability is 0.3% is shown with a dashed line.

Table 6. Galaxy merger X-ray and radio detections

ZTF name	ROSAT flux (cts/s)	$L_{2-10\text{keV}}$ ($\times 10^{44}$ ergs/s)	FIRST flux ($\times 10^{24}$ ergs/cm ² /s)	$L_{5\text{GHz}}$ ($\times 10^{55}$ ergs/s)
ZTF18aaxvmpg	0.196 ± 0.022	1.126 ± 0.129	< 1.39	< 18.552
ZTF18abamzru	< 0.05	—	< 1.33	—
ZTF18aasxvyo	1.728 ± 0.06	—	24.3 ± 1.36	—
ZTF18aaieguy	< 0.05	—	< 1.54	< 20.944
ZTF19aakjemw	0.031 ± 0.011	0.149 ± 0.054	< 1.46	< 8.666
ZTF18aaifbku	0.302 ± 0.025	0.245 ± 0.02	185.3 ± 1.5	208.297 ± 1.686
ZTF18aampabj	< 0.05	—	24.1 ± 1.34	335.637 ± 18.662
ZTF19aaagygp	0.04 ± 0.012	—	—	—
ZTF18abujubn	< 0.05	—	—	—
ZTF18acegbsb	0.034 ± 0.014	0.009 ± 0.004	12.8 ± 0.97	4.206 ± 0.319
ZTF19aaozpdm	< 0.05	—	< 1.39	—
ZTF18aacjltc	0.051 ± 0.015	0.529 ± 0.159	< 1.76	< 18.156
ZTF18abtpite	0.024 ± 0.01	—	—	—
ZTF18abvvrxu	0.163 ± 0.03	—	< 1.42	—
ZTF19abaktpb	< 0.05	—	—	—
ZTF18aaqjcxl	< 0.05	—	< 1.47	< 3.706
ZTF18abyoivl	< 0.05	—	< 1.1	—
ZTF18aabdiug	< 0.05	—	37.4 ± 1.43	35.732 ± 1.366
ZTF19aaviuyv	0.051 ± 0.01	—	—	—
ZTF18aabxczq	< 0.05	—	< 1.45	< 1.43
ZTF18acvwlrf	< 0.05	—	11.6 ± 1.51	192.822 ± 25.1
ZTF19aasejqv	0.086 ± 0.016	0.721 ± 0.135	< 1.34	< 22.161
ZTF18aazogyo	< 0.05	—	< 1.44	< 2.41
ZTF18aceypvy	0.047 ± 0.012	0.215 ± 0.055	< 1.39	< 10.349
ZTF18acbweyd	< 0.05	—	< 1.36	< 14.168
ZTF18acablce	0.022 ± 0.005	—	—	—
ZTF18abhpvvr	0.072 ± 0.014	—	—	—
ZTF19abfqmjj	< 0.05	—	—	—
ZTF18abmqwgr	0.077 ± 0.015	—	—	—
ZTF19aadgbih	0.092 ± 0.016	0.603 ± 0.102	24.9 ± 1.5	279.113 ± 16.814
ZTF19aalpfan	< 0.05	—	< 1.48	< 2.056
ZTF18aawwfep	< 0.05	—	< 1.39	< 15.897
ZTF19aavxims	0.041 ± 0.015	—	< 1.52	—
ZTF19aaaplct	0.165 ± 0.025	—	< 1.48	—
ZTF18aajnqqv	0.02 ± 0.009	0.026 ± 0.012	36.5 ± 1.47	59.825 ± 2.409
ZTF18abszfur	< 0.05	—	< 1.5	< 41.386
ZTF19abucbkt	< 0.05	—	—	—
ZTF18adbhlyb	0.036 ± 0.013	0.253 ± 0.09	< 1.34	< 17.931
ZTF18acxhoij	0.086 ± 0.016	—	—	—
ZTF18acajwep	< 0.05	—	—	—
ZTF19abipoqj	< 0.05	—	< 1.05	—
ZTF19abpkouu	0.043 ± 0.016	—	< 1.08	—
ZTF18abztovy	< 0.05	—	< 1.54	—
ZTF18acsllgd	0.031 ± 0.012	—	—	—
ZTF19aanxrki	0.068 ± 0.014	0.237 ± 0.049	< 1.44	< 4.96
ZTF18aamfuhc	0.341 ± 0.033	0.417 ± 0.04	< 1.42	< 2.693
ZTF18aadwvyr	0.06 ± 0.014	0.252 ± 0.059	< 1.02	< 4.353
ZTF19abauzsd	0.027 ± 0.009	0.643 ± 0.206	< 1.5	< 39.366
ZTF18abufbsq	0.048 ± 0.013	—	21.5 ± 1.39	—
ZTF18abzuzrg	0.028 ± 0.009	0.169 ± 0.053	< 1.45	< 13.173
ZTF18abtmcdb	0.032 ± 0.01	—	< 1.11	—
ZTF18aauhnby	0.1 ± 0.018	0.113 ± 0.021	< 1.53	< 3.168

NOTE—X-ray and radio properties of the 52 ZTF AGN in galaxy mergers. We show the 2-10keV flux in cts/s for the AGN with detections in the second ROSAT All-Sky Survey and conversions to luminosity for those with a spectroscopically confirmed redshift. We also the 20cm flux from the FIRST survey, including upper limits where available, and the corresponding luminosities.