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July 31, 2017

Dear Editor:

Please find attached a review paper entitled "Biaxial mechanical properties of human vocal fold cover under vocal fold elongation", which I would like to be considered for publication in the Journal of the Acoustical Society of America.

No author names have been omitted. No part of the submitted article has been published or is under consideration for publication in any other journal.

Thank you for the consideration.

Sincerely,

Zhaoyan Zhang
Himadri Samajder
Jennifer L. Long

Department of Head and Neck Surgery
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Biaxial mechanical properties of human vocal fold cover under vocal fold elongation

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ABSTRACT

Mechanical properties of the human vocal fold cover layer were experimentally investigated in uniaxial and biaxial tensile tests. The results showed a coupling effect between the stress conditions along the anterior-posterior and transverse directions, with vocal fold elongation increasing vocal fold stiffness along both directions, thus allowing more efficient control of the fundamental frequency of voice through vocal fold elongation. This study also shows that vocal folds were nearly isotropic at resting conditions, thus a tendency to vibrate with incomplete glottal closure, but became increasingly anisotropic with increasing vocal fold elongation.

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Keywords: vocal fold stiffness; stiffness anisotropy; biaxial tests; pitch control

33 I. INTRODUCTION

34 The mechanical conditions within the human vocal folds play an important role in
35 determining the fundamental frequency of vocal fold vibration and the resulting voice quality. As
36 the vocal folds are postured for the production of different voice types, the mechanical conditions
37 are also expected to vary (Hirano, 1974). A better understanding of the mechanical conditions of
38 the vocal folds across different voice types would facilitate the development of biomaterials for
39 vocal fold injection or tissue engineered vocal fold replacement with similar material properties.
40 Computationally, a more quantitative characterization of the vocal fold mechanical properties
41 would provide data for developing material constitutive models, linear or nonlinear, for use in
42 computational models of phonation.

43 There have been many previous reports experimentally characterizing vocal fold
44 mechanical properties (e.g., Hirano and Kakita, 1985; Alipour-Haghighi and Titze, 1991; Haji et
45 al., 1992; Chan et al., 1999; Alipour and Vigmostad, 2012; Kelleher et al., 2013; Kazemirad et al.,
46 2014). While these studies have provided valuable contribution to our understanding of the vocal
47 fold mechanical properties, it still remains unclear how the anisotropic mechanical conditions vary
48 under different conditions of vocal fold posturing. For example, due to material nonlinearity,
49 elongating the vocal folds along the anterior-posterior (AP) direction increases vocal fold stiffness
50 along this direction. However, little is known about how vocal fold elongation affects the
51 mechanical conditions in the transverse plane (the plane perpendicular to the AP direction).
52 Previous studies have shown a cross-axis coupling effect such that AP elongation may also stiffen
53 the vocal folds in the transverse plane (e.g., Yin and Zhang, 2013). This cross-axis coupling, which
54 would allow more efficient control of the fundamental frequency (F_0) of voice through vocal fold
55 elongation, cannot be investigated in a uniaxial tensile test, as in most of the experimental studies

56 so far. Similarly, although there have been some recent studies toward quantifying stiffness
57 anisotropy of the vocal folds (Hirano and Kakita, 1985; Kelleher et al., 2013), these studies focused
58 mainly on the stiffness anisotropy at the resting state of the vocal folds, due to limitations of
59 uniaxial tests in evaluating biaxial mechanical properties at conditions of varying vocal fold
60 elongation. Maintaining a degree of stiffness anisotropy has been shown to play an important role
61 in achieving complete glottal closure during vocal fold vibration (Zhang, 2011; Xuan and Zhang,
62 2014).

63 The present study presents our first step toward quantification of the anisotropic
64 mechanical properties of the vocal folds in a physiologically more realistic setting. We will show
65 that the constraints along the AP ends of the vocal folds, induced by varying degrees of vocal fold
66 elongation, have a large effect on the mechanical properties within the transverse plane as well as
67 the degree of stiffness anisotropy of the vocal folds.

68 **II. METHOD**

69 Two cadaveric human larynges (L1 at age 20 and L2 at age 50, both male) were harvested
70 at autopsy and frozen at -80C until use. After thawing, entire membranous vocal folds were
71 resected, with care to exclude the cartilaginous vocal process. The vocal fold cover layer (the
72 epithelium and lamina propria) was dissected off the underlying muscle layer and used in the
73 tensile tests. An approximately square-shaped segment of the cover layer, with axes along the AP
74 and transverse direction (perpendicular to AP), were used in the experiments. The specimen was
75 mounted to a Bose ElectroForce 4-motor biaxial testing system using silk sutures placed at each
76 edge of the specimen so that the two axes of the biaxial testing system aligned approximately to
77 the AP and transverse direction of the specimens (Fig. 1). The mounting was adjusted so that the
78 sutures were neither under tension nor sagging noticeably. Suture marks were also placed in the

79 center region of the testing specimen and monitored during testing using a high-speed camera
80 (Phantom V711, Vision Research) mounted directly above the specimen.

81 Four experiments were performed for each specimen, two in a uniaxial setting and two in
82 a biaxial setting (Fig. 2c-2f). In the first two uniaxial experiments, the specimen was vibrated
83 along either the AP (Fig. 2c) or transverse direction (Fig. 2d), with the other edges of the specimen
84 unconstrained and free to move. The imposed vibration was a sinusoidal stretch at 1 Hz, from the
85 resting position to a maximum stretch of 60% of the in situ length of the specimen along the testing
86 direction. Results from these tests would provide data comparable to previous tensile tests. In the
87 third experiment, the specimen was stretched to and held at six given degrees of elongation (0-
88 50% in step of 10%) along the AP direction (Fig. 2e), and for each of the six AP elongation
89 conditions, the specimen was stretched sinusoidally at 1 Hz along the transverse direction, with a
90 maximum stretch of 60%. Data from these biaxial tests quantify vocal fold mechanical properties
91 along the transverse direction at different degrees of AP elongations, simulating conditions of
92 different combinations of thyroarytenoid and cricothyroid muscle activations. In the fourth
93 experiment (Fig. 2f), the specimen was held at the original length along the transverse direction
94 while a sinusoidal vibration of 1 Hz was imposed along the AP direction. Data from this
95 experiment allow evaluation of the AP stiffness when the vocal fold is subject to transverse
96 constraint, as in in vivo vocal fold conditions.

97 Each test lasted 50 seconds, and was followed by a short period of rest (less than one
98 minute) which was required to reprogram for the next experimental conditions, during which the
99 specimen was kept hydrated by dripping phosphate-buffered saline on the specimen. Force
100 relaxation was observed during the initial preparation stage, with the peak force of each cycle
101 decayed with time. No further preconditioning was applied during the experiment. In the data

102 reported below, force relaxation was generally small except for conditions of large vocal fold
103 elongation (greater than 30%) in which the peak force decayed by about 15% over 50 cycles of
104 vibration. Although this force relaxation may be further reduced by longer preconditioning, it was
105 not pursued in this study considering that preconditioning is not often required or occurs in human
106 voice communication. In other words, force relaxation may be a natural component of human
107 phonation.

108 For each test, the displacements and forces at each axis were measured by an accelerometer
109 included in the Bose ElectroForce system and a 250-gram load cell (Honeywell, OH), respectively,
110 at a sampling rate of 5000 Hz. The Bose ElectroForce system used a feedback control loop to
111 achieve accurate displacement control. This introduced some high-frequency noise in the force
112 data, which was filtered out by passing the force measurement through a low-pass filter with a cut-
113 off frequency of 100 Hz (which is much higher than the 1 Hz of vibration of the experiment). The
114 displacement and force data were then used to calculate the stress σ and strain ϵ using the initial
115 dimensions of the specimens. To calculate vocal fold stiffness, the stress-strain data were averaged
116 over loading and unloading and a tangent Young's modulus was calculated as the slope of the
117 resulting curve (i.e., $d\sigma/d\epsilon$) at different vocal fold deformations.

118 **III. RESULTS**

119 **A. Uniaxial tests: material properties at resting lengths**

120 Fig. 2a shows the stress-strain curves from the first two experiments with uniaxial tests
121 along the transverse direction (thick solid line) and the AP direction (thick dashed line). Similar
122 to previous experimental findings (e.g., Hirano and Kakita, 1985; Alipour and Vigmostad, 2012),
123 the stress-strain curves for both uniaxial tests include a roughly linear range at small strains, in
124 which the stress increased slowly with increasing strain, and a nonlinear range at large strains, in

125 which the stress rapidly increased with increasing strain. The linear range was much smaller for
126 the larynx L1 (strains less than approximately 0.2) than L2 for which the linear range extended up
127 to a strain of 0.3.

128 As in previous studies, the vocal fold was significantly stiffer along the AP direction than
129 the transverse direction. The derived tangent Young's moduli are shown in Fig. 2b. For both
130 larynges, the tangent Young's modulus was larger along the AP direction than the transverse
131 direction. At the resting state (unstretched in either direction), the tangent Young's modulus for
132 L1 was 3.73 kPa and 1.17 kPa along the AP and transverse direction, respectively, with a stiffness
133 anisotropy ratio of 3.2. For L2, the tangent Young's modulus was 1.05 kPa and 0.60 kPa along
134 the AP and transverse direction, respectively, with a ratio of 1.8. Note that because the AP and
135 transverse uniaxial tests provided information about the mechanical properties along two different
136 directions, estimation of the AP and transverse stiffnesses at one same vocal fold state was only
137 possible for the resting unstretched vocal fold state. In other words, uniaxial tests do not allow
138 estimation of the degree of stiffness anisotropy at conditions of varying degree of vocal fold
139 elongation.

140 **B. Biaxial tests: effect of vocal fold elongation**

141 Fig. 2a also shows data obtained in two biaxial tests, when the vocal fold specimen was
142 vibrated along one direction while the specimen edges at the other direction, unlike being free in
143 uniaxial test, were fixed at the initial length (thin solid and dashed lines). This fixed condition at
144 the two edges induced a coupling effect between the transverse and AP directions, and increases
145 the stress along the direction of vibration, compared to that in a uniaxial test. In other words,
146 greater force was required in the biaxial tests to produce the same amount of vocal fold deformation
147 as in a uniaxial test. More importantly, the slope of the stress-strain curve was larger in the biaxial

148 setting, indicating an increase in vocal fold stiffness. Note that this coupling effect appears to be
149 larger when the AP direction was constrained as compared to that when the transverse direction
150 was constrained.

151 For biaxial experiments with the AP direction held at different degrees of elongation (Fig.
152 1e), the tangent Young's modulus along the transverse direction was extracted and plotted in Fig.
153 3 as a function of transverse strain, for different degrees of AP elongation. The bottom half of Fig.
154 3 shows the same data but focused in a smaller range of transverse strains which is more likely to
155 occur during normal phonation. For both larynges, the coupling effect significantly increased the
156 transverse Young's moduli. This magnitude of the coupling effect was larynx specific. For L1,
157 increasing the AP elongation from 0 to 0.33 led to the transverse Young's modulus at $\epsilon_t=0$ being
158 almost doubled. For L2, the transverse Young's modulus at $\epsilon_t=0$ increased by almost four times
159 as the AP elongation increased from 0 to 0.45. For both larynges, this increase gradually plateaued
160 as the degree of AP elongation approached the nonlinear range (AP strain below about 0.2 for L1
161 and 0.3 for L2).

162 **C. Stiffness anisotropy**

163 Quantifying the degree of stiffness anisotropy requires data of the AP and transverse
164 tangent Young's moduli measured around the same vocal fold state. In the present study, the two
165 tangent Young's moduli data were available only for conditions of zero transverse strain (i.e., the
166 specimen length along the transverse direction was kept constant at the initial value) at different
167 degrees of AP elongation. For these conditions, an anisotropy ratio was calculated as the ratio
168 between the AP and transverse tangent Young's moduli, and is shown in Fig. 4 as a function of
169 AP elongation, for larynx L2 (solid line). Fig. 4 shows that at zero AP elongation, the degree of
170 stiffness anisotropy was about 1.5, indicating a nearly isotropic mechanical behavior. The

171 anisotropy ratio increased significantly with increasing AP elongation, increased to 4.8 and 55 for
172 an AP elongation of 18% and 45% AP elongation, respectively.

173 Unfortunately, the fourth experiment was not performed for L1 so that the AP Young's
174 modulus data for conditions of zero transverse strain were not available for L1. However, an
175 estimation of the anisotropy ratio for L1 can be obtained as the ratio between the AP Young's
176 moduli measured in a uniaxial test (i.e., with the transverse edges unconstrained) and the transverse
177 Young's moduli as zero transverse strain. Because the AP Young's modulus measured in a
178 uniaxial test was expected to be smaller than that at zero transverse strain (Fig. 2), this
179 approximation would lead to an underestimation of the anisotropy ratio. Such estimations are
180 shown in Fig. 4 for both larynges. Similar trends as to that in L2 can be observed: the degree and
181 stiffness anisotropy was generally low at resting states and increased with increasing AP
182 elongation.

183 **IV. DISCUSSION AND CONCLUSION**

184 This study shows that in a physiologically more realistic condition with the vocal fold
185 subject to constraint along both the AP and transverse directions, the mechanical stress within the
186 vocal fold was higher than that measured in a uniaxial tensile test. In particular, the boundary
187 conditions at the AP ends of the vocal folds can significantly increase both the stress and stiffness
188 along the transverse direction (the transverse tangent Young's modulus increased by four times
189 for L2). This coupling effect between the AP and transverse directions, specifically the increase
190 in the transverse stiffness accompanying vocal fold elongation, allows a more efficient control of
191 the fundamental frequency through vocal fold elongation, especially at low-pitch voice production.
192 The effect of this cross-axis coupling on F0 control can be estimated by examining the effect of
193 vocal fold elongation on vocal fold eigenfrequencies, which are important determinants of F0.

194 Using the data obtained for L2 and a vocal fold geometry similar to that used in Zhang (2016), an
195 eigenvalue analysis showed that a vocal fold elongation of 30%, which increased AP stiffness from
196 1.5 kPa to 20 kPa and the transverse stiffness from 1.5 kPa to 4.5 kPa, would increase the first in
197 vacuo eigenfrequency of the vocal fold from 31 Hz to 69 Hz, which is an extra 46% increase
198 compared to an increase from 31 Hz to 57 Hz if the accompanying increase in the transverse
199 stiffness was not accounted for.

200 This study also shows that the degree of stiffness anisotropy was generally low at resting
201 conditions and increased significantly with increasing vocal fold AP elongation. It is interesting
202 to note that L2 was almost isotropic at the resting state whereas L1 was more anisotropic. This
203 inherent difference in the stiffness anisotropy may explain the observation that in excised larynx
204 experiments some vocal folds are able to vibrate with complete closure whereas others cannot
205 (Isshiki, 1989). Our previous studies show that vocal fold models of nearly isotropic mechanical
206 properties tend to vibrate with incomplete glottal closure (Xuan and Zhang, 2014). Thus, vocal
207 folds with nearly isotropic mechanical properties, such as L2 in the present study, are likely to
208 vibrate with incomplete glottal closure, whereas vocal folds with a higher degree of anisotropy,
209 such as L1 in the present study, may be able to vibrate with better glottal closure.

210 The two larynges in this study exhibited large difference in their mechanical behaviors,
211 with L1 having a smaller linear range and higher degree of stiffness anisotropy compared to L2
212 (Fig. 2). It may be tempting to attribute such differences to the age difference (age 20 for L1 vs.
213 age 50 for L2), among other factors. However, due to the small sample size of this study, this
214 potential age effect would require further investigation in future studies.

215

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220

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249

250 Figure Captions:

251

252 Fig. 1: (Color online) The experimental setup for the biaxial tests (a) and tissue mounting (b). (c)-
253 (f) shows the four experimental protocols used in this study, including two uniaxial settings and
254 two biaxial setting.

255

256 Fig. 2. (Color online) (a) The stress-strain curves obtained from uniaxial and biaxial tests for L1.
257 (b) The derived tangent Young's moduli derived from the uniaxial tests for larynges L1 and L2.

258

259 Fig. 3. (Color online) The transverse tangent Young's modulus as a function of the transverse
260 strain for different conditions of AP elongation. The bottom panels show the same data but for a
261 smaller range of transverse strain that is more likely to occur in normal phonation.

262

263 Fig. 4. (Color online) The stiffness anisotropy ratio, calculated as the ratio between the AP and
264 transverse tangent Young's moduli, as a function of AP vocal fold elongation. Solid line: L2. The
265 dotted and dashed line are an underestimation of the anisotropy ratio for L1 and L2, respectively.

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