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Publication Date

1986-02-01



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To be published in Journal of Physical Chemistry

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February 1986

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Electron Energy Loss and Thermal Desorption Spectroscopy of
Pyridine Adsorbed on Pt(111)

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Abstract

The chemisorption behavior of pyridine (NC_5H_5) on a Pt(111) surface has been examined using thermal desorption and electron energy loss spectroscopy as a function of adsorption temperature. The vibrational spectrum of pyridine adsorbed at room temperature on this surface shows intense loss peaks in the specular direction from vibrational modes which can be characterized as in-plane stretching and bending modes. This vibrational spectrum has been interpreted as the formation of an α -pyridyl species (NC_5H_4) on the surface. The pyridyl moiety is bonded to the platinum surface through the nitrogen and one of the α -carbon atoms with the pyridyl plane perpendicular to the metal surface. When pyridine is adsorbed at low temperature (120K), it bonds to the surface through both the nitrogen atom and the π and π^* orbitals of the pyridine ring. As the crystal is warmed to 260K, at saturation coverage, approximately 50% of the molecules desorb as molecular pyridine. The remaining pyridine molecules partially decompose on the surface to form an α -pyridyl fragment. The electron energy loss spectra of pyridine adsorbed at both low and room temperature is compared to the infrared spectra of two osmium cluster compounds $\text{Os}_3(\text{CO})_{11}(\text{NC}_5\text{H}_5)$, a pyridine complex and $\text{HOs}_3(\text{CO})_{10}(\text{NC}_5\text{H}_4)$, a pyridyl complex.

INTRODUCTION

The bonding of aromatic molecules on many different single crystal metals and different crystallographic planes has been a subject of considerable interest. Benzene has been the most studied aromatic molecule so far and by several different surface science techniques including photoelectron spectroscopy, low energy electron diffraction, thermal desorption spectroscopy (TDS) and electron energy loss spectroscopy (EELS). The adsorption geometry

of benzene on the low Miller index surface planes of many transition metals (Pt, Ni, Pd, Rh, Ir, Ru) has been established fairly conclusively.¹ The ring plane of benzene is oriented parallel to the surface plane. The π and π^* orbitals of benzene are considered to interact with the metal orbitals in forming the chemisorption bond. In the case of heteroaromatics there is the possibility of a similar interaction between the π and π^* orbitals of the adsorbate and the metal atoms. However, there is now the possibility of an interaction between the metal and the nonbonding lone pair of electrons localized on the heteroatom which can bond to the surface in a conventional sigma donor type interaction.

The importance of both of these interactions for heteroaromatics adsorbed on metal surfaces has been observed with pyridine on Ag(111) and Ni(100).² Both vibrational and electronic electron energy loss spectroscopy of pyridine adsorbed on Ag(111) have shown the conversion of a π bonded to nitrogen bonded pyridine as the pyridine coverage is increased. For coverages below 0.50 monolayers, pyridine adsorbs in a flat configuration, i.e. with the ring plane parallel to the surface plane in a geometry similar to that of benzene on the Ag(111) surface. For coverages above 0.5 monolayers there is a "phase transformation" to a nitrogen bonded species. Recently a study by DiNardo et al.² again using electron energy loss spectroscopy of pyridine adsorbed on Ni(100) showed the existence of two nitrogen bonded species. The surface chemistries of pyridine, deuterium labeled pyridines, and methyl substituted pyridines have been studied on the Ni(100) and Ni(111) surfaces.³

The IR and Raman spectra of pyridine and isotopically substituted pyridines have been well studied⁴⁻⁸ and normal coordinate analysis have been used to assign normal modes to the observed spectral features.⁵⁻⁷ The infrared spectra of pyridine complexed with a variety of mononuclear

transition metal complexes have been studied as a function of metal and coligands.⁹⁻¹⁰ In all of these metal complexes, pyridine is bonded to the metal atom through the nitrogen lone pair. There were only minor shifts, 15-25 cm^{-1} , in the vibrational frequencies of pyridine in these complexes relative to that of liquid pyridine. This is in contrast to benzene π complexes such as bis benzene chromium, $\text{Cr}(\text{C}_6\text{H}_6)_2$, and benzene chromium tricarbonyl, $\text{Cr}(\text{C}_6\text{H}_6)(\text{CO})_3$, where large shifts of 50-150 cm^{-1} from gas phase frequencies are observed for some modes.¹¹

We present here a vibrational study using electron energy loss spectroscopy of pyridine adsorbed on Pt(111) as a function of temperature. The complimentary technique of thermal desorption spectroscopy was also employed. The infrared spectrum of two prototype cluster compounds, $\text{Os}_3(\text{CO})_{11}(\text{NC}_5\text{H}_5)$ and $\text{Ho}_3(\text{CO})_{10}\text{NC}_5\text{H}_4$ (and their perdeuterated analogues) were also recorded for comparison. Triosmium and tricobalt cluster compounds have been extremely useful as prototype structural models in identifying adsorbed species on metal surfaces.^{12,13,14}

EXPERIMENTAL

All experiments were performed in a stainless steel ultra-high vacuum (UHV) chamber. The base pressure of the chamber was typically $<2 \times 10^{-10}$ torr. The chamber contained a 200 liter/sec ion pump and an auxillary titanium sublimation pump. The system was equipped with low energy electron diffraction (LEED) optics and Auger electron spectroscopy; the Auger electrons were analyzed with the four grid retarding field analyzer of the LEED optics. A Vacuum Generator quadrupole mass spectrometer was used to monitor individual masses during a thermal desorption spectroscopy experiment as well as background gases in the chamber.

The design of the Electron Energy Loss Spectrometer used in these experiments has been described elsewhere.¹⁵ Briefly the spectrometer consists of two 120° cylindrical sectors, a monochromator and an analyzer. An incident beam energy of 3 eV was used with an energy resolution of 50-80 cm⁻¹ (ΔE_{FWHM}). Typically the count rate in the elastic channel was 1 x 10⁵ counts/sec. For experiments done in the specular direction the angle of the incident and the scattered beam were both 60° with respect to the surface normal. Off-specular measurements were made by rotating the crystal about the vertical axis.

The platinum crystal was spark cut and oriented along the (111) direction using Laue X-ray back diffraction. The crystal, 1 cm in diameter, was then polished according to procedures described by Tsai *et al.*¹⁶ The crystal was spot-welded to two platinum wires 15 mil in diameter which mounted the crystal to a copper block. The crystal was cooled in a similar manner as described previously with minor modifications.¹⁷ Liquid nitrogen was passed through two copper cold fingers connected by copper braids 5" long and 1/4" in diameter to the copper block on which the sample was mounted. The crystal temperature, as monitored by a chromel-alumel thermocouple spot-welded to the bottom of the crystal, reached a minimum temperature of 120K. The sample was heated resistively through the same copper braids to override the cooling of the crystal. The copper block which mounted the crystal was mechanically attached to a stainless steel manipulator offset 2" from the center of the chamber so the sample could be rotated to any part of the chamber desired. A tilt mechanism permitted the crystal to be rotated on the crystal axis for off-specular electron energy loss measurements.

Pyridine (Mallinckrodt) and pyridine-d₅ (99.5% d) (Merck, Sharpe, and Dohme) were used without purification. The isotopically substituted

pyridines, 2,6-d² pyridine and 4-d¹ pyridine were prepared by the method of Bak et al. and as described elsewhere.^{3,18} The pyridine sample was degassed by several freeze-thaw cycles and then introduced into the vacuum chamber. All exposures are uncorrected for ion-gauge sensitivity.

To remove surface carbon, the Pt(111) crystal was heated to 1000K in the presence of oxygen at a pressure of 1×10^{-6} torr. After the crystal was free of carbon impurities (as shown by Auger spectra). Sulfur impurities that remained after this treatment can then be removed by heating the crystal to 800K in oxygen again at a pressure of 1×10^{-6} torr. Oxygen that remained on the surface after these cleaning treatments was removed by briefly heating the crystal to 1200K. Calcium was removed from the crystal surface by sputtering the surface with argon ions at an Ar pressure of 1×10^{-5} and an incident beam energy of 600 eV with the sample temperature at 300K. These procedures were repeated until the surface was determined to be clean by Auger electron spectroscopy. The lattice orientation was determined by low energy electron diffraction.

The tri-osmium cluster compounds were synthesized in a procedure slightly modified from that previously described by Yin et al.^{19,20} Standard vacuum and Schlenk line techniques were used in preparation of these compounds: reactions were carried out under an atmosphere of nitrogen. Using 300 mg of Os₃(CO)₁₂ (Strem Chemicals) dissolved in 100 mls of methylene chloride, excess of reagent grade pyridine (or perdeuterated pyridine), 1 ml, was transferred into the reaction flask. A solution of trimethylamine oxide (CH₃)₃NO (35 mg (CH₃)₃NO dissolved in 35 mls of CH₂Cl₂) was added dropwise to the osmium reaction mixture. The progress of the reaction was monitored by infrared spectroscopy in the 2100-1700 cm⁻¹ region where carbonyl absorption occurs. Trimethylamine oxide was added until the 2070 cm⁻¹ band corresponding to the

reactant, $\text{Os}_3(\text{CO})_{12}$ disappeared. The solvent was pumped off, the product $\text{Os}_3(\text{CO})_{11}(\text{NC}_5\text{H}_5)$ was redissolved in methylene chloride and pentane and then recrystallized. The solvent was decanted and evaporated off under vacuum.

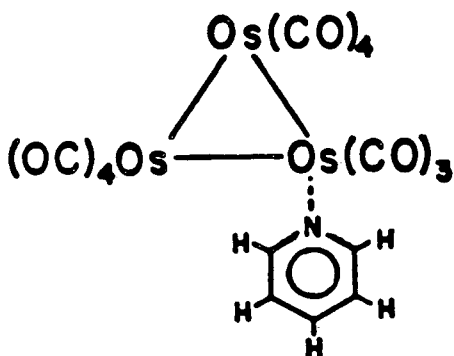
The product $\text{Os}_3(\text{CO})_{11}(\text{NC}_5\text{H}_5)$ (100 mg) was used in the preparation of the pyridyl complex, $\text{HOs}_3(\text{CO})_{10}(\text{NC}_5\text{H}_4)$. $\text{Os}_3(\text{CO})_{11}(\text{NC}_5\text{H}_5)$ was dissolved in 50 ml of n-octane containing a few drops of pyridine, and the mixture was refluxed for one hour. The pyridyl complex was isolated by chromatography using a silica column and pentane elutant. After the compound was collected, the solvent was evaporated.

The infrared spectra of both these compounds were recorded over the spectral range $400\text{--}3200\text{ cm}^{-1}$ using a Perkin Elmer 283 infrared spectrometer with a spectral resolution of $2\text{--}3\text{ cm}^{-1}$ and a frequency accuracy of $1\text{--}2\text{ cm}^{-1}$. Anhydrous KBr was used in the preparation of the KBr discs.

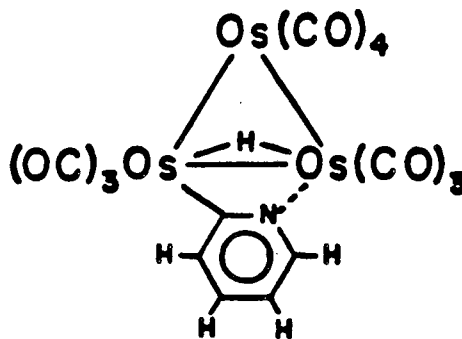
RESULTS

Infrared Spectra of $\text{Os}_3(\text{CO})_{11}(\text{NC}_5\text{H}_5)$ and $\text{HOs}_3(\text{CO})_{10}\text{NC}_5\text{H}_4$

We present the infrared spectra recorded for $\text{Os}_3(\text{CO})_{11}(\text{NC}_5\text{H}_5)$ (1a) and $\text{HOs}_3(\text{CO})_{10}(\text{NC}_5\text{H}_4)$ (1b) first so we may use these results when discussing the electron energy loss spectrum recorded for pyridine adsorbed on Pt(111).



1a



1b

Figure 1a shows the infrared spectra in the spectral range from 400 to 1650 cm^{-1} for the two osmium cluster compounds. The infrared spectra for the perdeuterated analogues are shown in Figure 1b. In both the pyridine and the pyridyl triosmium complexes the pyridine moiety is bonded to one of the metal atoms in the cluster through the nitrogen. In the pyridyl complex the pyridine moiety is additionally bonded to the metal cluster through the α carbon.

In Table I we list the low frequency absorption bands (400-610 cm^{-1}) for the pyridine and pyridyl cluster compounds along with the corresponding frequencies for $\text{Os}_3(\text{CO})_{12}$.²¹ The comparison of the low frequency bands observed in the IR spectra (Fig. 1) of the two prototype model compounds (1a and 1b), match reasonably well to the observed absorption bands in the IR spectra of $\text{Os}_3(\text{CO})_{12}$. Thus verifying that these absorptions are associated with the bending and stretching modes of the metal framework, the metal carbon stretching motion and the metal carbonyl bending motions. The spectral region from 600-1650 cm^{-1} is of more interest in this study, the absorption bands which occur at these frequencies are due to the vibrations of the pyridine and pyridyl fragments of the cluster. The assignment for the bands in this region are made by comparison to the vibrational spectrum and the assignment made for liquid pyridine (Table II). The frequencies of the absorption bands for these cluster compounds compare well with the vibrational frequencies for uncomplexed pyridine. The correspondence of the spectral features between pyridine and complexed pyridine (or pyridyl) permits the use of liquid phase vibrational assignments for guidance as we interpret the spectral observations. The average frequency shift from liquid pyridine frequencies is calculated to be 13 cm^{-1} for the pyridine complex (1a). As would be expected the frequencies of the vibrations of the pyridyl fragment do not correlate as well

to liquid pyridine frequencies with an average frequency shift of 27 cm^{-1} . However, qualitatively these infrared spectra confirm that the fingerprint regions for pyridine and pyridyl are reasonably similar if looked at with sufficiently low resolution, as is the case of a vibrational spectrum recorded using the technique of electron energy loss spectroscopy.

ROOM TEMPERATURE ADSORPTION OF PYRIDINE ON Pt(111)

I. Electron Energy Loss Spectroscopy

The electron energy loss spectrum recorded in the specular direction of pyridine adsorbed at room temperature onto platinum (the 111 surface) is shown in Figure 2a. Loss peaks listed in order of decreasing frequency are observed at 3080, 1570, 1450, 1230, 1150, 1010, and 740 cm^{-1} . Perdeuteropyridine, similarly studied (Fig. 2b), displays loss peaks in the specular direction, in order of decreasing frequency, at 2275, 1540, 1330, 1230, 1010, 860, and 680 cm^{-1} . Table III lists these frequencies with their liquid phase counterparts, again we feel that their correspondence to liquid phase pyridine permit the use of liquid phase vibrational assignments as a guide.

One of the most significant aspects to the spectrum in Figure 2 is the prominence of the loss at 3080 cm^{-1} . The frequency and the deuterium shift to 2275 cm^{-1} definitively identify the characteristic aromatic C-H stretching motion. Since C-H stretching modes are intrinsically in-plane vibrations, their intensity in the specular direction indicates that pyridine is not adsorbed with its skeletal plane parallel to the metal surface. This conclusion is strengthened by the high intensities observed in the spectral region $1000\text{-}1600 \text{ cm}^{-1}$ where C-C stretching and C-H in plane bending modes occur. It is also consistent with the striking difference from the corresponding spectra of benzene²² and toluene²³ on platinum (111) surfaces.

These substances, when adsorbed, have EELS spectra with extremely weak loss features both in the C-H aromatic stretching region and in the 1000-1600 cm^{-1} range.

The electron energy loss spectrum of two specifically labeled pyridines. 4d^1 pyridine and $2,6\text{d}^2$ pyridine adsorbed on Pt(111) at 300K have also been measured. In Figure 3 we show the EELS spectrum in the spectral region from 2000 to 3200 cm^{-1} for these two specifically labeled compounds along with the EELS spectrum for normal and perdeuterated pyridine in the same region for reference. There is a large difference in the intensities of the loss peaks associated with the C-D and C-H stretching motion between the two compounds 4d^1 pyridine and $2,6\text{d}^2$ pyridine. The intensity of the loss peak at 2275 cm^{-1} (the C-D stretching mode) is very weak in the specular direction for $2,6\text{d}^2$ pyridine. The intensity of this loss peak increases when the sample is tilted 10° off the specular angle. Clearly this loss peak is non-dipolar in nature. This is in contrast to 4d^1 pyridine where there is a significant dipole contribution to the intensity of the loss peak. The results from these specifically labeled compounds will be useful in determining the orientation and adsorption geometry of the pyridine ring.

II. Thermal Desorption Spectroscopy

The thermal desorption spectrum of pyridine adsorbed on Pt(111) at room temperature is shown in Figure 4. This figure shows the trace of the parent ion which corresponds to an AMU of 79. There is one molecular desorption peak in the thermal desorption curve, this peak has a temperature desorption rate maximum at approximately 473K. As the coverage is increased the intensity of the peak also increases until saturation coverage is reached.

The decomposition of pyridine as monitored by the thermal desorption

hydrogen curves for pyridine, d^5 pyridine, $4d^1$ pyridine, $2,6d^2$ pyridine, and $3,5d^2$ pyridine has been discussed previously.³ For normal pyridine (NC_5H_5) there are three hydrogen peaks in the thermal desorption spectrum. The desorption rate temperature maxima for these peaks occur at approximately 335K, 515K, and 620K. Through the use of the specifically labeled pyridines it was determined that the low temperature hydrogen (deuterium) peak was due to carbon-hydrogen bond cleavage at the α -carbon only. The regiospecific bond breaking of the α carbon-hydrogen bond led the authors to suggest the formation of an α -pyridyl intermediate, although no spectroscopic evidence was available at that time. The intensity ratio of the hydrogen desorption peaks was 0.5:2:2 for the 335K, 515K, and 620K peaks, respectively. The deviation from one of the low temperature peak was attributed to the observation of immediate hydrogen desorption upon pyridine adsorption. Our results on the decomposition of pyridine on Pt(111) is in qualitative agreement with these earlier results. Figure 7 shows H_2 desorption from low temperature pyridine adsorption, which will be discussed further in the next section. H_2 evolution for room temperature pyridine adsorption is similar except for the intensity of the first desorption maxima which is approximately one-half in the room temperature spectrum.

The amount of irreversibly bound pyridine can be calculated using the peak-to-peak height ratios of the carbon 272 eV Auger electron peak relative to the platinum 237 eV Auger electron peak. An Auger spectrum was taken after the crystal first was exposed to pyridine to reach saturation coverage which was then compared to the Auger spectrum taken after the thermal desorption experiment. The amount of carbon on the surface was calculated using equation (1).²⁴

$$\frac{I(272)}{I(237)} \times 0.64 = \theta_{\text{carbon}} \quad (1)$$

The amount of pyridine which irreversibly adsorbs was calculated to be approximately $60\% \pm 5\%$ for pyridine adsorbed at room temperature.

LOW TEMPERATURE ADSORPTION OF PYRIDINE ON Pt(111)

I. Electron Energy Loss Spectroscopy

The electron energy loss spectrum recorded for pyridine adsorbed on Pt(111) at 120K and saturation coverage is shown in Figure 5. This spectrum is substantially different than the spectrum for pyridine adsorbed at room temperature. Loss peaks are observed in order of decreasing frequency at: 3070, 1610, 1470, 1230, 1050, 840, and 660 cm^{-1} . And for perdeuteropyridine adsorbed on this platinum surface at 120K loss peaks are observed at 2270, 1570, 1330, 1250, 1010, 830, 580, and 290 cm^{-1} . For monolayer or submonolayer quantities of pyridine adsorbed on the Pt(111) surface at 120K we again find most of the modes observed in the specular direction can be characterized as primarily in-plane modes, as was seen in the room temperature spectrum. The relative intensities of the loss peaks differ greatly between spectra for room temperature and low temperature adsorption of pyridine. In addition to intensity differences between the EELS spectrum for pyridine adsorption at 120K and 300K we find no loss peak at 740 cm^{-1} in the low temperature EELS spectrum of normal pyridine. There are two new loss peaks observed in the low temperature spectrum at 660 cm^{-1} and 840 cm^{-1} . In Table III the vibrational frequencies for pyridine adsorbed on Pt(111) at 120K are compared to the vibrational frequencies for uncomplexed liquid or gas phase values of pyridine.

II. Thermal Desorption Spectroscopy

Figure 6 shows the thermal desorption spectrum of pyridine adsorbed at 120K, again monitoring the parent ion (AMU=79) as a function of coverage. At low coverages there are two molecular desorption peaks. As the coverage increases the intensity of the low temperature peak increases with coverage and the temperature of the rate maxima shifts to lower temperature. At low coverages (0.1 monolayers) $T_{\max} = 347\text{K}$, the maxima shift to a T_{\max} of 260K at saturation coverage. The high temperature peak with a $T_{\max} = 483\text{K}$ at $\theta=0.11$ shifts down to $T_{\max} = 450\text{K}$ at $\theta=4.0$. As the coverage increases a third desorption peak develops with a temperature desorption rate maximum at 190K. As the coverage is increased further, this is the only peak which increases in intensity.

The amount of pyridine which irreversibly adsorbs decreases when pyridine is adsorbed at low temperatures. The amount of decomposition is calculated using equation 1 to be 25%. This is approximately more than two-fold less than the amount of irreversible desorption when pyridine is adsorbed at room temperature. The thermal desorption hydrogen curve, for saturation coverage, is shown in Figure 7. Three peaks are observed in the spectrum with temperature maxima of 335K, 515K, and 620K.

DISCUSSION

Room Temperature Adsorption

The C_{2v} molecular symmetry of normal pyridine and d^5 pyridine breaks up the 27 fundamental modes of vibration into the irreducible representation as follows: $10a_1$, $3a_2$, $9b_1$, and $5b_2$ modes. The A_2 modes are Raman active only and the A_1 , B_1 , and B_2 modes are both infrared and Raman active. The A_2 and B_2 modes are out-of-plane modes and the 19 A_1 and B_1 modes are in-plane

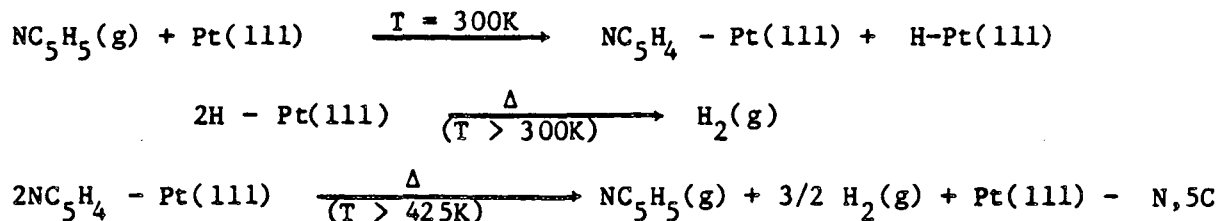
modes. The key question in this study is which modes will be active for pyridine adsorbed on the Pt(111) surface. To answer it, we take advantage of the fact that the infrared spectra recorded for the two cluster compounds $\text{Os}_3(\text{CO})_{11}(\text{NC}_5\text{H}_5)$ and $\text{HOs}_3(\text{CO})_{10}(\text{NC}_5\text{H}_4)$ were similar enough in the 600-1650 cm^{-1} region for the use of liquid pyridine absorption frequencies in the interpretation of orientation of the molecule on the surface regardless of whether it is pyridine or pyridyl on the surface. The EELS active modes will be dependent on the adsorption site and the orientation of the molecule on the surface. The surface normal dipole selection rule states that only modes which are totally symmetric with respect to all symmetry elements (A_1 modes) will be dipole active and therefore be strongly dependent on scattering angle. The intensity of those loss peaks will be at a maximum in the specular direction in the EELS spectrum if there is a significant dipole scattering contribution to the cross section of the loss peak. The symmetry of the pyridine-platinum complex can retain the C_{2v} symmetry of molecular pyridine or it can be reduced to C_s symmetry. If the pyridine molecule was nitrogen bonded to the surface and tilted at some angle away from the surface or lying in a flat configuration π -bonded parallel to the surface the only symmetry element would be the mirror plane perpendicular to the molecular plane that bisects the nitrogen atom and the C-H unit in the 4 position. If there was partial decomposition of the pyridine molecule to form a pyridyl fragment, (NC_5H_4) , the point group symmetry of the pyridyl fragment would again be C_s but now the molecular plane would be the mirror plane. In this case both the A_1 and B_1 irreducible representations of the C_{2v} point group would be lowered to the A' irreducible representation of the point group C_s and the vibrations of this representation would be "EELS allowed".

In Table III we show the frequencies for the normal modes of vibration of uncomplexed pyridine and compare these frequencies to the frequencies observed in the electron energy loss spectrum for pyridine adsorbed at 300K on the Pt(111) surface. Due to the poor resolution of EELS the frequencies of some of the modes have been used more than once in assigning the spectrum of adsorbed pyridine. Also some of the vibrations which are inherently weak transitions (i.e. weak absorption bands in the infrared spectrum of liquid phase pyridine) are not observed. There is some uncertainty in the assignment of the loss peak at 740 cm^{-1} in the pyridine EELS spectrum. This loss peak might be assigned to either of the strong absorptions at 700 cm^{-1} (ν_{26}) or 744 cm^{-1} (ν_{25}) for pyridine in the liquid phase which are assigned to out-of-plane motions of a pyridine. However, the corresponding loss peak in the NC_5D_5 spectrum (Fig. 2b) at 685 cm^{-1} displays a deuterium shift by a factor of only 1.08, much smaller than those displayed for liquid pyridine, 1.18 for ν_{25} and 1.33 for ν_{26} . Hence, the 740 cm^{-1} loss peak is assigned to the $B_1 \nu_{19}$ mode which implies that this mode shifts up in frequency by 13%. In the pyridyl molecular complex the ν_{19} mode is assigned to the band at 708 cm^{-1} and is therefore also shifted to higher frequency, here by 9%. It is unclear why this low frequency in-plane ring deformation would be more perturbed than other in-plane ring modes in the pyridyl structure. A detailed force-field coordinate analysis would perhaps answer this question.

As noted previously there are intense in-plane modes in the electron energy loss spectrum for pyridine adsorbed at 300K on the Pt(111) surface. As shown in Table III the modes observed in this spectrum are modes of A_1 and B_1 type symmetry, i.e. in-plane modes. These in-plane modes have high intensity at the specular angle, so they are the one with significant dipole derivatives perpendicular to the platinum surface.

The results for the specifically labeled pyridines support the proposal of the presence of a pyridyl fragment on the surface at room temperature. If only a dipolar contribution to the intensities of the loss peaks for the C-D stretching modes and the C-H stretching modes is considered in the specular direction, the expected intensity ratio for the C-H stretching mode relative to the C-D stretching mode depends upon the component of $\frac{du}{dr}$ perpendicular to the metallic surface and the fact that $(\frac{du}{dr})_{C-H} = 2(\frac{du}{dr})_{C-D}$. For a "stand up" geometry with the N-Pt bond perpendicular to the metallic surface, this ratio $\frac{I_{C-H}}{I_{C-D}}$ is predicted to be 4 for both the 2,6-dideutero and 4-monodeutero pyridines. For a geometry with one of the N-C bonds parallel to the metallic surface, the ratio $\frac{I_{C-H}}{I_{C-D}}$ is predicted to be infinity for the 2,6-dideuteropyridine (one of the deuterium atoms being lost to form an α C-Pt bond) and it would be 2 for the 4-monodeutero pyridine. The spectra in Figure 2 show for the 2,6-dideutero pyridine a very small C-D feature (at specular reflection) with a ratio $\frac{I_{C-H}}{I_{C-D}} = 10$. For the 4-monodeutero pyridine the experimental ratio is about 1.5. Thus the spectra distinctly favor the parallel N-C bond geometry with the loss of an α hydrogen to permit platinum bonding to both the nitrogen and one of the α carbon atoms.

In the case of room temperature adsorption of pyridine on Pt(111) the predominant reaction scheme proposed based on the evidence presented is:



An α -pyridyl species is formed upon pyridine adsorption on the Pt(111) surface. Atomic hydrogen on the surface desorbs as H_2 with a $T_{\text{max}} = 335\text{K}$.

The pyridyl structure was determined by electron energy loss spectroscopy to be stable up to a temperature of 425K. When the temperature is increased above 425K, disproportionation must occur, as revealed by the desorption of molecular pyridine (see Figure 4). Some pyridyl fragments undergo further decomposition. The available hydrogen from this process can then recombine with a remaining pyridyl fragment which then desorbs as molecular pyridine (AMU = 79). Although hydrogen evolution peaks at a somewhat higher temperature ($T_{\max} = 515\text{K}$ (see Figure 7)) then the pyridine desorption peak ($T = 473\text{K}$), the two processes, decomposition and desorption, begin at about the same temperature, near 425K.

Low Temperature Adsorption

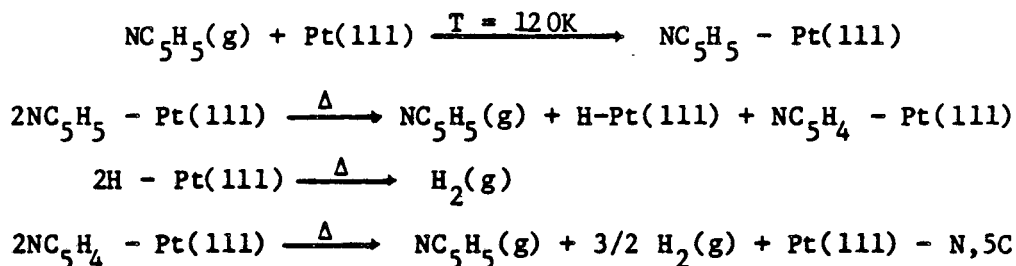
The assignment for low temperature monolayer or submonolayer pyridine covered surface is also given in Table III. The loss peak observed at 660 cm^{-1} is assigned to the ν_{19} ring deformation mode, close to the liquid phase value, 652 cm^{-1} , but distinctly different from 740 cm^{-1} in the room temperature spectrum. The loss peak at 840 cm^{-1} (which shifts to 580 cm^{-1} for d^5 pyridine) is assigned to the ν_{26} mode which corresponds to an in-phase out-of-plane bending motion. The increase of the C-H out-of-plane bending mode relative to the frequency of the free molecule is also observed for benzene adsorbed on Pt(111).²² The plane of the pyridine molecules is tilted from the surface plane as is suggested by the relative intensities of the in-plane modes. Although the pyridine ring is tilted away from the surface plane there is still an appreciable interaction with the surface metal d orbitals and the π and π^* orbitals of pyridine. This can be seen by the increase in frequency of the C-H out-of-plane bend from 700 cm^{-1} for liquid pyridine to 840 cm^{-1} for pyridine adsorbed on Pt(111) at 120K. At saturation coverage, a value of 52°

has been calculated for the angle between the pyridine ring and the Pt surface plane from Near Edge X-Ray Absorption Fine Structure data.²⁵ This is consistent with the electron energy loss spectroscopy results presented here. The bonding of pyridine to the surface at low temperatures appears to be through both the lone pair of electrons which are localized on the nitrogen atom and the delocalized molecular π orbitals on the ring. The lowest energy configuration is then a combination of these two different modes of bonding.

The thermal desorption spectrum for pyridine adsorbed on Pt(111) at a temperature of 120K shows three molecular desorption peaks. The low temperature (190K) peak which begins to appear only at coverages greater than a monolayer is attributed to the desorption of multilayer pyridine. Using Redheads equation for first-order reaction kinetics,²⁶ and assuming a pre-exponential factor of 10^{13} sec^{-1} , for a temperature maximum of 190K the desorption energy was calculated to be 10.8 kcal/mole. The second molecular pyridine desorption peak has a temperature maximum which is coverage dependent (low coverages $T_{\text{max}} = 347\text{K}$; saturation coverage $T_{\text{max}} = 260\text{K}$), such behavior could be caused by a second order desorption process which follows second order reaction kinetics or a first order desorption process with an energy of activation that is coverage dependent. A first-order process with a coverage dependent energy of activation is the more probable cause of the variation of the temperature maximum as a function of coverage, due to the low mobility of the pyridine molecule at these temperatures. Again using a pre-exponential factor of 10^{13} sec^{-1} the energy of activation at low coverages is calculated to be 21.8 kcal/mole while at saturation coverage the energy of activation drops to 15.0 kcal/mole. The high temperature molecular desorption peak with a desorption rate maximum centered at $T=483\text{K}$ at $\Theta=0.11\text{L}$ which shifts to $T=450\text{K}$ at $\Theta=4.0\text{L}$ is the only molecular desorption peak for pyridine adsorbed at room

temperature. This peak discussed previously for room temperature results corresponds to the disproportionation of a NC_5H_4 fragment with a hydrogen atom which then desorbs as molecular pyridine.

For low temperature adsorption when the substrate is warmed the following reaction scheme is observed:



Initially the pyridine adsorbs on the surface through the nitrogen atom, which presumably acts as a conventional electron donor. Then as the sample is warmed some pyridine desorbs molecularly, with the temperature maximum of the desorption rate depending on the initial coverage. At saturation coverage approximately 50% of the pyridine desorbs from the surface at 260K. The remaining adsorbed pyridine molecules fragment to form the pyridyl-surface complex at $\sim T=300\text{K}$. On further heating the behavior is identical to that observed for pyridine adsorbed at room temperature. The recombination of a hydrogen atom and a NC_5H_4 fragment to desorb as molecular pyridine at temperatures greater than 400K. The amount of pyridine desorbing in this high temperature peak and the amount of decomposition is less when pyridine is adsorbed at low temperature. This supports the earlier conclusion that disproportionation of the pyridyl complex occurs when the temperature is raised above 400K.

The infrared spectra for the two cluster compounds provide additional supportive evidence for the proposed structure of pyridine adsorbed on Pt(111) at 120K and 300K. The absorption bands of the pyridine and pyridyl complexes

which shift the most in frequency are the ones that offer the best differentiation, but superimposed on that must be the certainty with which the assignment can be made. The second factor causes the bands at 1483 cm^{-1} , ν_5 , and 1474 cm^{-1} , ν_{14} , (frequencies for liquid pyridine) to be the first choice. For the pyridine complex (1a) $\Delta\nu$ ($\nu_{\text{complex}} - \nu_{\text{pyridine liquid}}$) is 7 cm^{-1} and 10 cm^{-1} for ν_5 and ν_{14} respectively, and for the pyridyl complex (1b) $\Delta\nu$ is -24 cm^{-1} and -21 cm^{-1} again for ν_5 and ν_{14} respectively. Now looking at the frequencies of the loss peaks in the EELS spectra we observe ν_5 and ν_{15} (unresolved in EELS spectrum), for pyridine adsorbed at 120K, to be at 1470 cm^{-1} for pyridine adsorbed at 300K these modes are observed at 1450 cm^{-1} . This loss peak is then 20 cm^{-1} higher in frequency for pyridine adsorbed at 120K relative to pyridine adsorbed at 300K. Following this comparison, with modes ν_4 (1583 cm^{-1}) and ν_{13} (1581 cm^{-1}), for the pyridine complex (1a) ν_4 increases by 30 cm^{-1} and ν_{13} decreases by 3 cm^{-1} relative to liquid pyridine, for the pyridyl complex ν_4 increases by 9 cm^{-1} and ν_{13} decreases by 32 cm^{-1} , again relative liquid pyridine frequencies. A net effect of an increase in frequency in this region for pyridine and a decrease in this region for pyridyl. The loss peak in the EELS spectrum at 1610 cm^{-1} for pyridine adsorbed at 120K has increased in frequency by $\sim 30\text{ cm}^{-1}$, and the loss peak for pyridine adsorbed at 300K has decreased in frequency, relative to liquid pyridine by 12 cm^{-1} .

The in-plane ring distortion mode ν_{19} at 652 cm^{-1} in pyridine shifts very little in the pyridine complex at 650 cm^{-1} . In the EELS spectrum for pyridine adsorbed at 120K this mode is assigned to the loss peak at 660 cm^{-1} which is similar unperturbed relative to liquid pyridine. In the pyridyl complex this mode shifts dramatically to 708 cm^{-1} , an increase of 56 cm^{-1} relative to the liquid pyridine value. As noted previously a similar behavior is observed for

pyridine adsorbed at room temperature, the ν_{19} mode has been assigned to the loss peak at 740 cm^{-1} . Lastly there is a large increase in the intensity of the adsorption in the perdeutero pyridine complex in the 800 cm^{-1} region shifted from the 1100 to 1200 cm^{-1} region in the perhydro pyridine cluster compound. A similar increase is observed in the 800 cm^{-1} region for low temperature perdeutero pyridine EELS spectrum.

CONCLUSIONS

The surface chemistry of pyridine on Pt(111) has been studied as a function of temperature. We have shown in the case of room temperature adsorption that there is partial decomposition of the pyridine molecule with an α -hydrogen abstracted, thereby forming a pyridyl surface species. For low temperature adsorption, pyridine adsorbs molecularly with the ring plane tilted toward but not parallel to the surface plane. A new reaction channel is seen for low temperature adsorption from the thermal desorption data. As the temperature is increased from 120K to approximately 260K (this temperature is dependent on the initial pyridine coverage) 50% of the pyridine desorbs as molecular pyridine. The remaining pyridine molecules form a pyridyl surface complex at about room temperature. The infrared spectra recorded for $\text{Os}_3(\text{CO})_{11}(\text{NC}_5\text{H}_5)$ and $\text{HOs}_3(\text{CO})_{10}(\text{NC}_5\text{N}_4)$ aided in the interpretation of the EELS data for the pyridine-platinum system.

Acknowledgement

This work was supported by the National Science Foundation Grant #CHE-83-07159 and by the Director, Office of Energy Research, Office of Basic Energy Sciences, Chemical Sciences Division of the U.S. Department of Energy under Contract No. DE-AC0376SF00098. We would like to thank Deiter Himmelreich for the guidance in preparing the Osmium cluster compounds. We would also like to thank G.C. Pimentel for his guidance in preparing this manuscript. Registry No. Pyridine 110-86-1: Pt. 7440-06-4

Table I
 Corresponding Frequencies of the Low Frequency Modes (400-610 cm^{-1}) for
 Tri-Osmium Complexes

$\text{Os}_3(\text{CO})_{12}^{\text{a}}$	$\text{Os}_3(\text{CO})_{11}(\text{NC}_5\text{H}_5)$	$\text{HOs}_3(\text{CO})_{10}(\text{NC}_5\text{H}_4)$
400 W ^b	405	403
410 S	417	412
429 M	435	439
448 W		451
462 S	460	
474 W	474	483
496 S	493	498
516 W	502	
534 W	528, 529	540
562 M	564	574
583 S	591	588
606 S	605	605

^a Reference: S.F.A. Kettle and P.L. Strangelin, *Inorganic Chemistry*, Vol. 18, 2749 (1979)

^b W=weak, M=medium, S=strong

Table II

Vibrational Assignment of $\text{Os}_3(\text{CO})_{11}(\text{NC}_5\text{H}_5)$ and $\text{HOs}_3(\text{CO})_{10}(\text{NC}_5\text{H}_4)$
(600 - 1650 cm^{-1})

Symmetry, Mode No.		Description	$\text{NC}_5\text{H}_5(\text{NC}_5\text{D}_5)^{\text{C}}$		$\text{Os}_3(\text{CO})_{11}(\text{NC}_5\text{H}_5)$ h^5 (d^5)		$\text{HOs}_3(\text{CO})_{10}(\text{NC}_5\text{H}_4)$ h^4 (d^4)	
A_1	4 (8a) ^B	C-C, C-N Stretch	1583	(1554)	1613	(1571)	1592	(1554)
	5 (19a)	C-C, C-N Stretch	1483	(1340)	1490	(1346)	1459	(1339)
	6 (9a)	C-H in-plane bend	1218	(882)	1223	(896)	1222	(869)
	7 (18a)	C-H in-plane bend	1072	(823)	1075	(838)	1082	(869)
	8 (12)	asymmetric ring breathing	1032	(1014)	1050	(1042)	1057	(1055)
	9 (1)	ring breathing	991	(963)	1019	(979)	1029	(1000)
	10 (6a)	in-plane ring distortion	601	(579)	639	A	680	A
B_2	23 (5)	o.o.p. ring distortion	1007	(828)	1019	(838)	1029	(869)
	24 (10b)	o.o.p. ring distortion	936	(765)		(783)		(738)
	25 (4)	C-H o.o.p. bend	744	(631)	762	(650)	758	(680)
	26 (11)	C-H in-phase o.o.p. bend	700	(526)	710	(540)	745,740	(549)
	27 (16b)	ring o.o.p. distortion	403	(367)		A		A
B_1	13 (8b)	C-C, C-N stretch	1581	(1546)	1578	(1564)	1549	(1519)
	14 (19b)	C-C, C-N stretch	1442	(1303)	1452	(1328)	1421	(1298)
	15 (14)	C-C ring stretch	1362	(1046)				
	16 (3)	C-C in-plane bend	1227	(1226)			1269	(1216)
	17 (15)	C-H in-plane bend	1143	(856)	1161	(843)	1162,1117	(825)
	18 (18b)	C-H in-plane bend	1079	(835)	1075	(838)	1082	(869)
	19 (6b)	in-plane ring distortion	652	(625)	650	A	708	632

A. Hidden under more intense absorption bands from low frequency vibrations not due to the pyridine moiety.

B. Wilson Mode Numbers for Benzene

C. Colson, S.D., J. Phys. Chem. 88, 6067 (1984).

Table III

Vibrational Frequencies and Assignment for the Electron Energy Loss Spectra of Pyridine and d⁵ Pyridine Adsorbed on Pt(111)

Symmetry, Mode No.	Description	NC ₅ H ₅ (NC ₅ D ₅) ^A		Pyridine-Pt(111)		Pyridine-Pt(111)		
				T=300K		T=120K		
				h ⁵	(d ⁵)	h ⁵	(d ⁵)	
A ₁	1 (2) ^B	C-H stretch	3094	(2302)	3080	(2275)	3070	(2270)
	2 (13)	C-H stretch	3072	(2276)	3080	(2275)	3070	(2270)
	3 (20A)	C-H stretch	3030	(2268)	3080	(2275)	3070	(2270)
	4 (8a)	C-C, C-N stretch	1583	(1554)	1570	(1540)	1610	(1570)
	5 (19a)	C-C, C-N stretch	1483	(1340)	1450	(1330)	1470	(1330)
	6 (9a)	C-H in-plane bend	1218	(882)	1150	(860)	1230	(830)
	7 (18a)	C-H in-plane bend	1072	(823)				
	8 (12)	asymmetric ring breathing	1032	(1014)	1010	(1010)	1050	(1010)
	9 (1)	ring breathing	991	(963)	1010	(1010)	1050	(1010)
	10 (6a)	in-plane ring distortion	601	(579)				
B ₂	23 (5)	o.o.p. ring distortion	1007	(828)				
	24 (10b)	o.o.p. ring distortion	936	(765)				
	25 (4)	C-H o.o.p. bend	744	(631)				
	26 (11)	C-H in-phase o.o.p. bend	700	(526)			840	(580)
	27 (16b)	ring o.o.p. distortion	403	(367)				
B ₁	11 (20b)	C-H stretch	3087	(2289)	3080	(2275)	3070	(2270)
	12 (7b)	C-H stretch	3042	(2256)	3080	(2275)	3070	(2270)
	13 (8b)	C-C, C-N stretch	1581	(1546)	1570	(1540)	1610	(1570)
	14 (19b)	C-C, C-N stretch	1442	(1303)	1450	(1330)	1470	(1330)
	15 (14)	C-C ring stretch	1362	(1046)				
	16 (3)	C-C in-plane bend	1227	(1226)	1230	(1230)	1230	(1250)
	17 (15)	C-H in-plane bend	1143	(856)	1150	(860)		
	18 (18b)	C-H in-plane bend	1079	(835)			1050	(830)
	19 (6b)	in-plane ring distortion	652	(625)	740	(685)	660	c

vPT-N (290)

A. Pyridine Frequencies and Assignments Taken From: Wiberg, K.B., Walters, V.A., Wong, K.N. and Colson, S.D., J. Phys. Chem. 88, 6067 (1984).

B. Wilson Mode #'s for Benzene

C. Loss Peak Not Observed at This Frequency, Most Probably Hidden Under More Intense Loss Peak at 580 cm⁻¹.

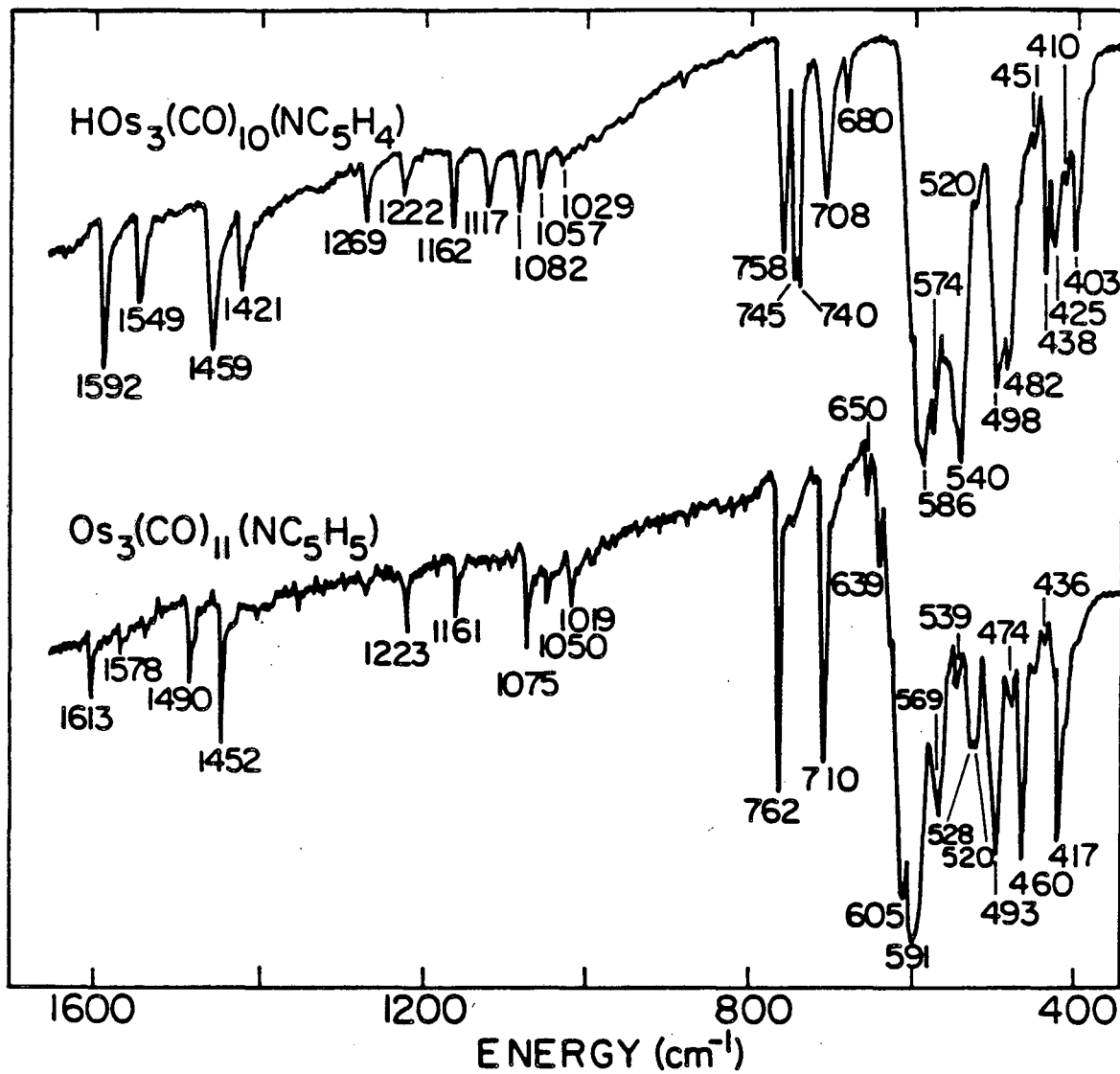
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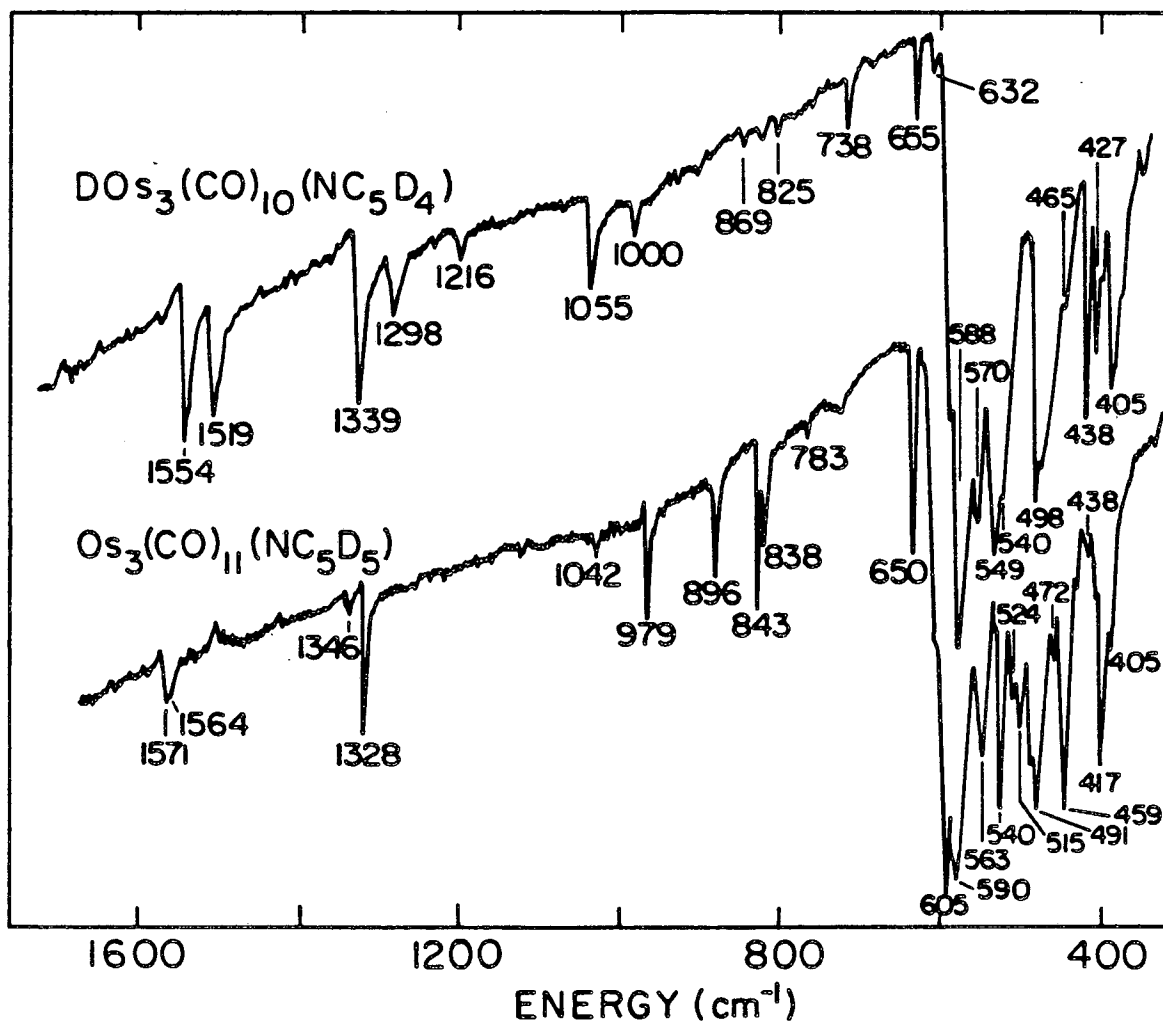
FIGURE CAPTIONS

- Fig. 1. The infrared spectrum of $\text{HOs}(\text{CO})_{10}\text{NC}_5\text{H}_4$ and $\text{Os}(\text{CO})_{11}\text{NC}_5\text{H}_5$ (and the perdeuterated compounds) in the region of $400\text{--}1600\text{ cm}^{-1}$. (a) perhydro compounds, (b) perdeutero compounds.
- Fig. 2. Electron energy loss spectrum of pyridine adsorbed on Pt(111) at room temperature and saturation coverage. (a) NC_5H_5 , (b) NC_5D_5 .
- Fig. 3. The carbon-hydrogen and carbon-deuterium stretching frequency region for specifically labeled pyridines, 4d^1 pyridine and $2,6\text{d}^2$ pyridine, with the normal and perdeuterated pyridine also shown for comparison.
- Fig. 4. Thermal desorption spectrum for pyridine adsorbed on Pt(111) at room temperature (monitoring the parent ion $\text{AMU} = 79$) for molecular desorption. The heating rate for all thermal desorption experiments was $25\text{--}30\text{ deg/sec}$.
- Fig. 5. Electron energy loss spectrum of pyridine (and perdeuterated pyridine) adsorbed on Pt(111) at 120K and a coverage of one monolayer. (a) NC_5H_5 , (b) NC_5D_5 .
- Fig. 6. Thermal desorption spectrum for reversible molecular desorption of pyridine adsorbed on Pt(111) at 120K as a function of coverage. Exposures noted are uncalibrated and in units of Langmuir's where $1\text{ Langmuir} = 1 \times 10^{-6}\text{ torr-sec}$.
- Fig. 7. Thermal desorption spectrum (monitoring $\text{AMU}=2$) for hydrogen evolution. Three hydrogen peaks are observed with temperature maxima of 335K , 515K , and 620K .



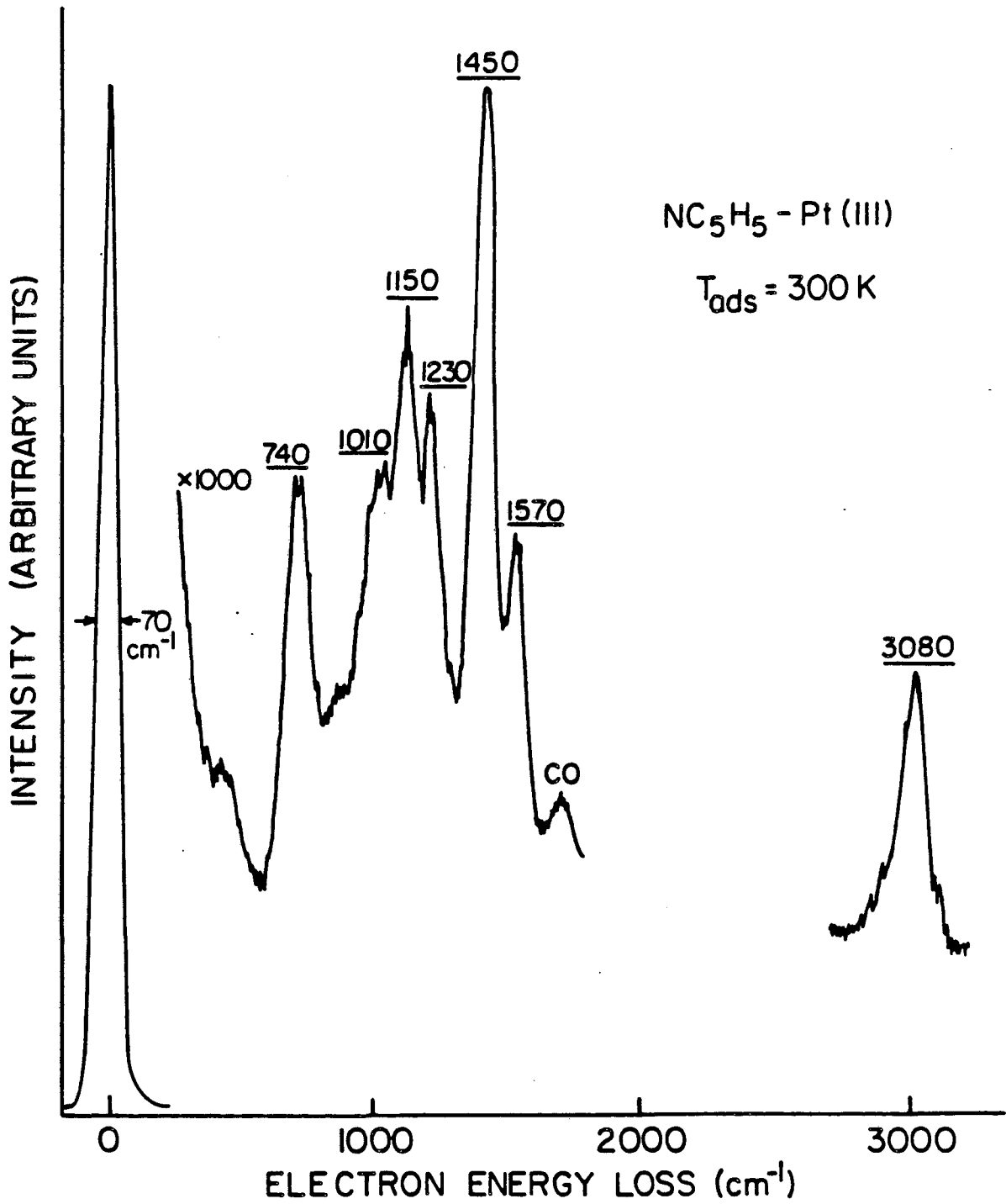
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Fig. 1a



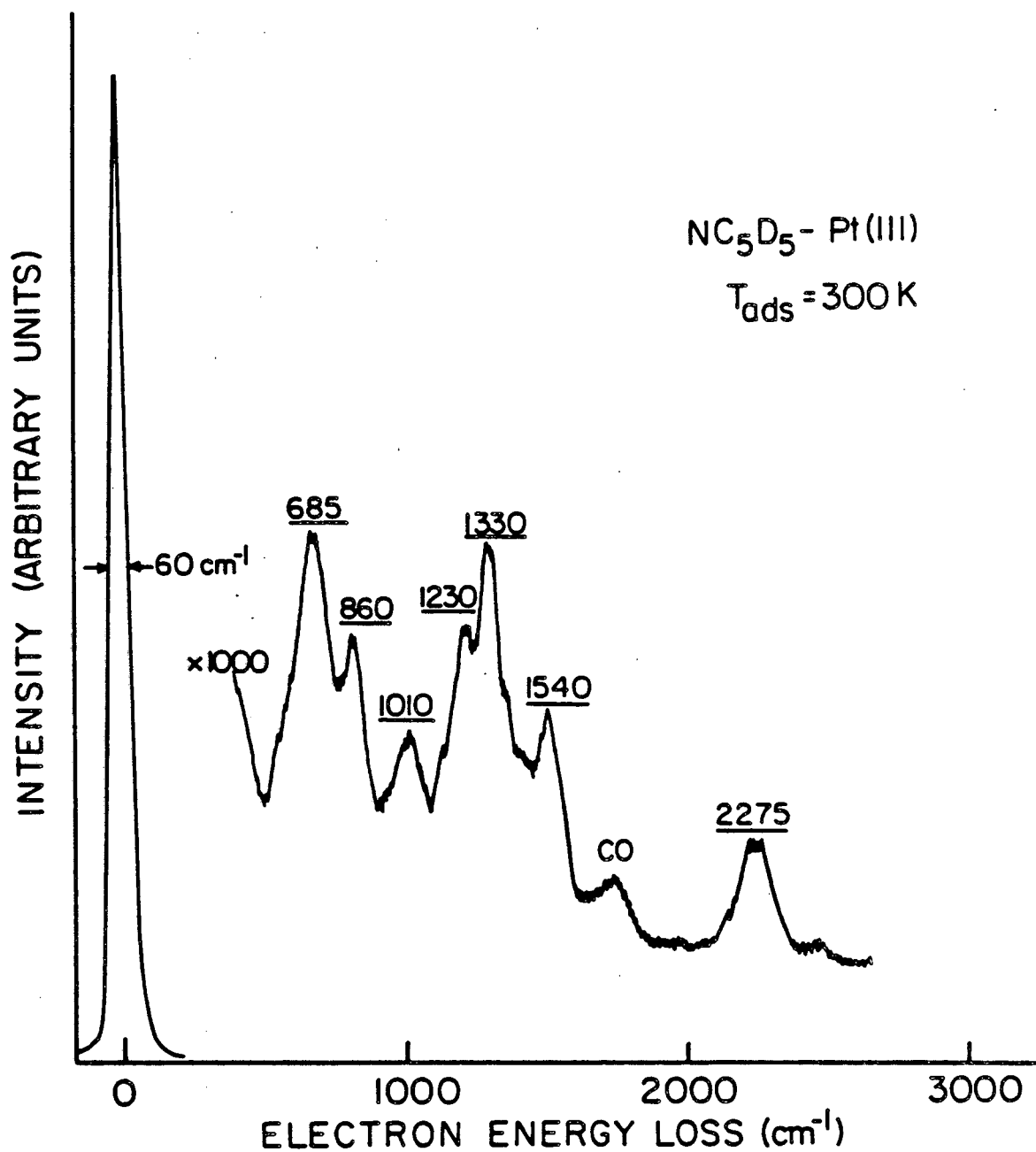
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Fig. 1b



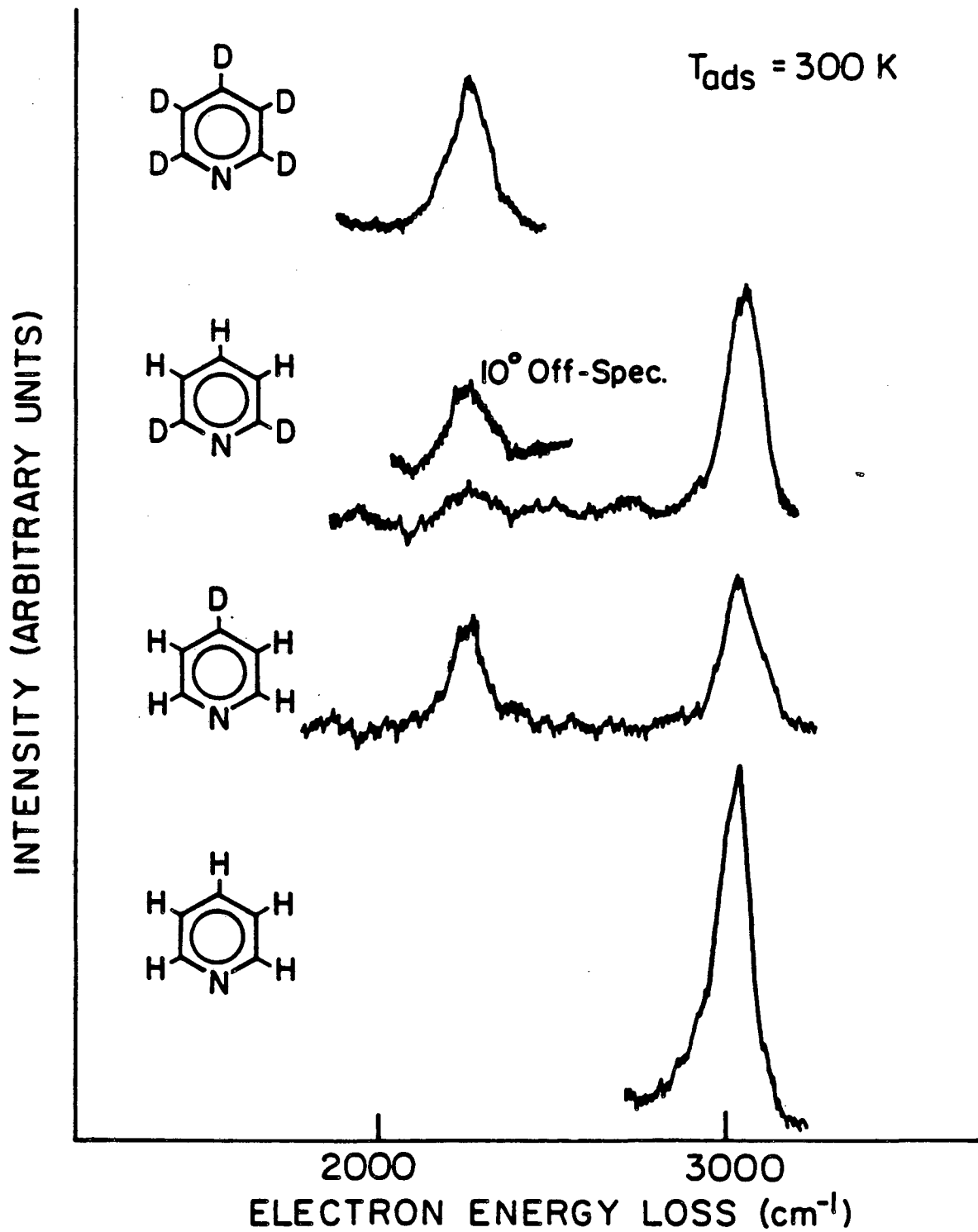
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Fig. 2a



XBL 8412-5929

Fig. 2b



XBL 8412-5930

Fig. 3

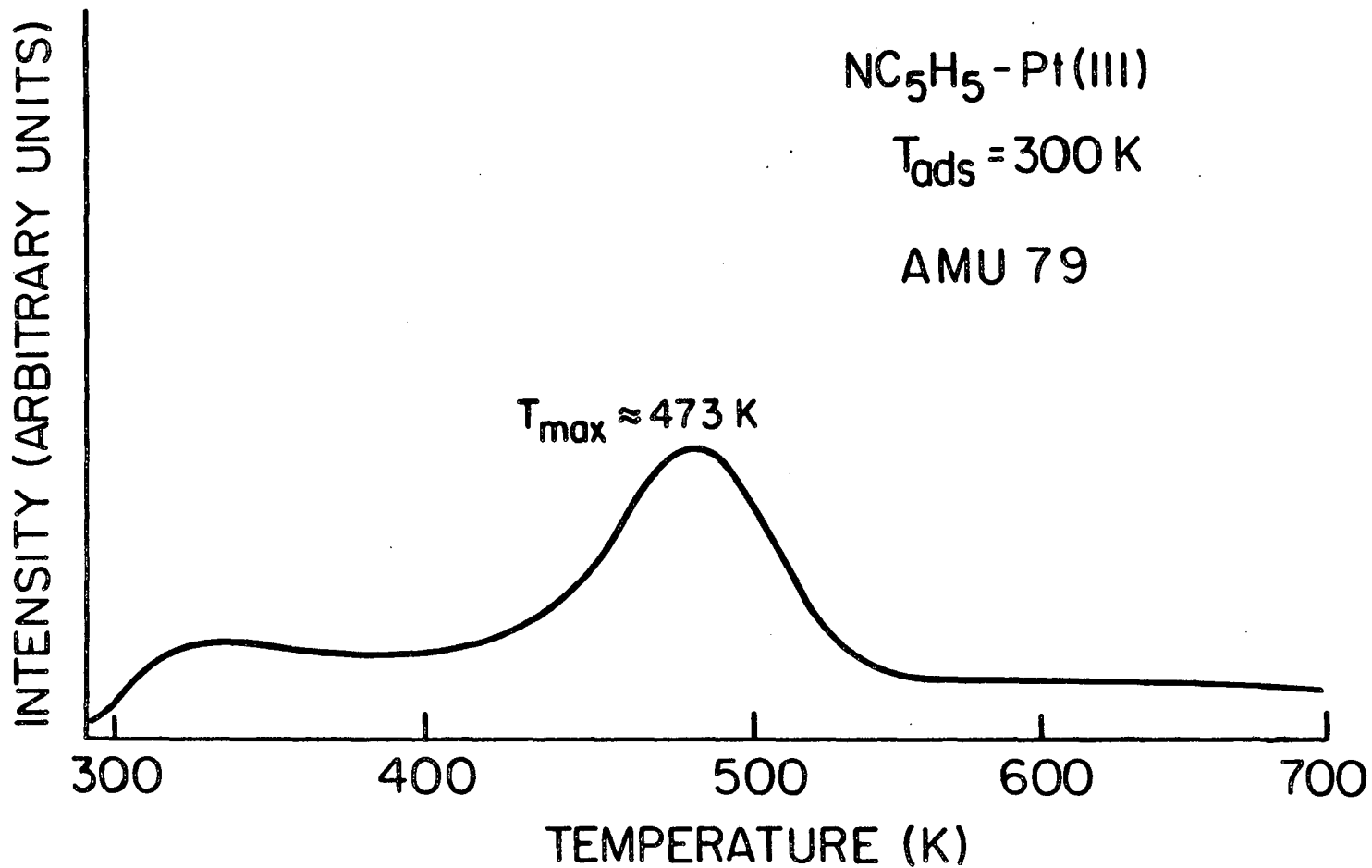
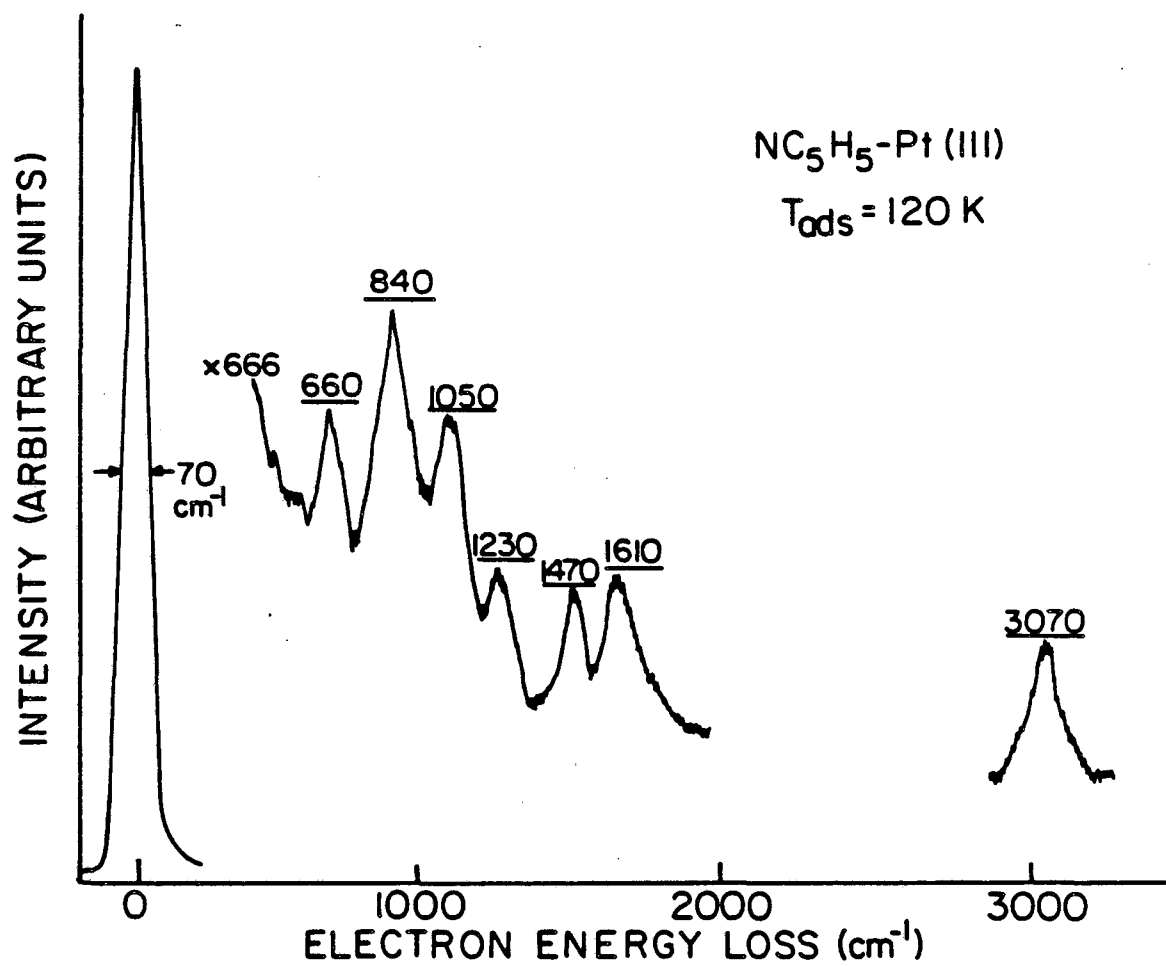


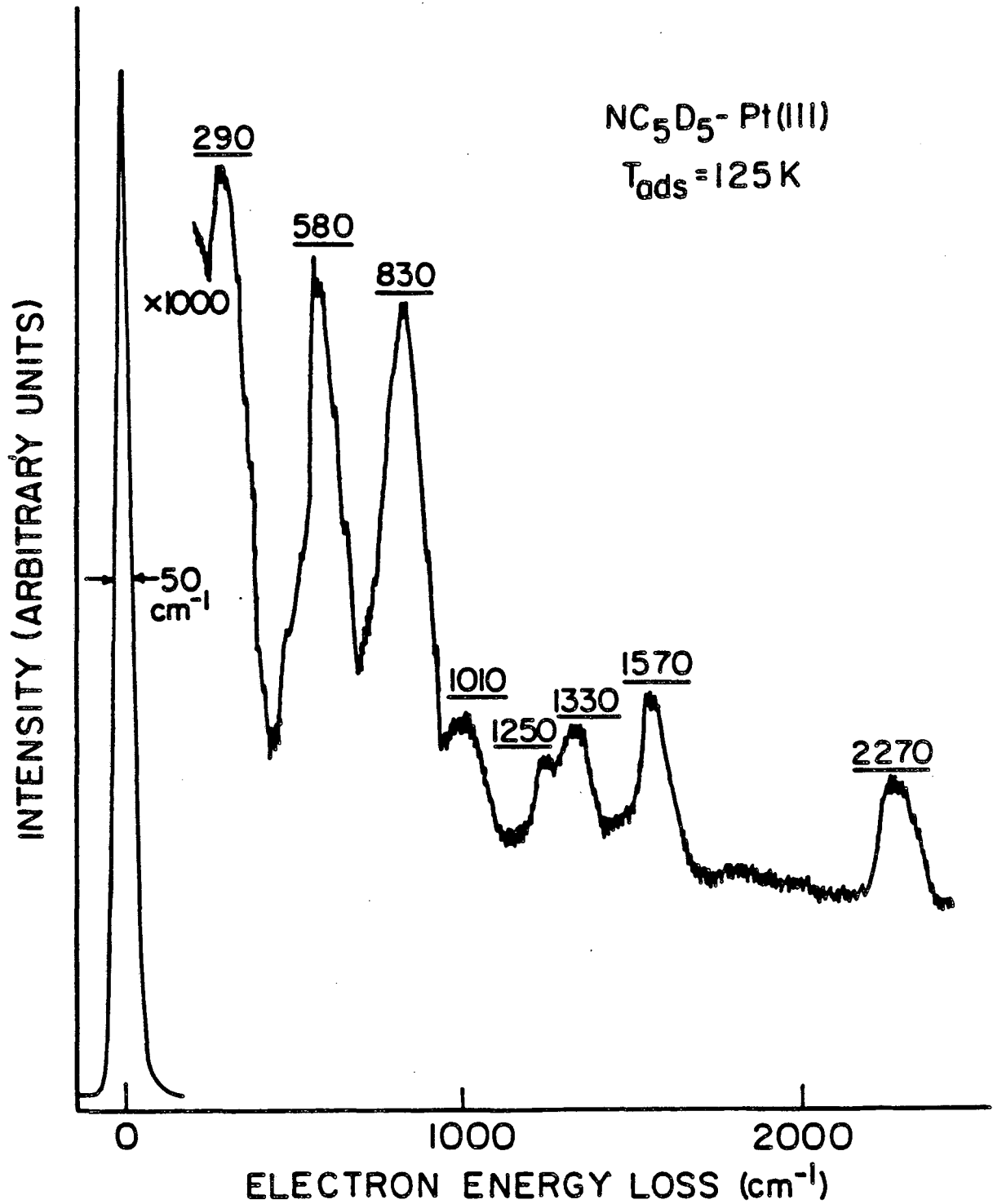
Fig. 4

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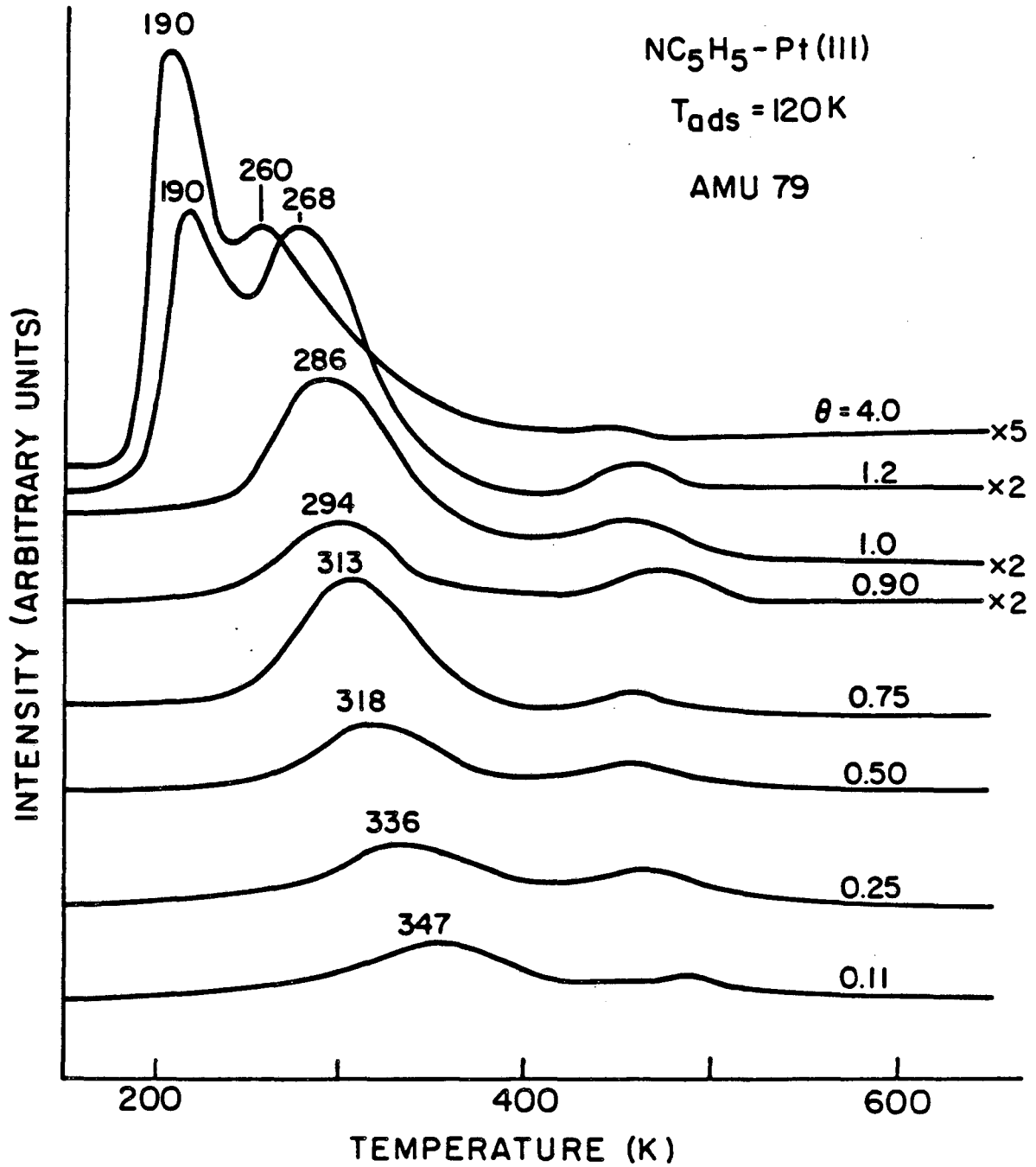
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Fig. 5a



XBL 8412-5933

Fig. 5b



XBL 8412-5935

Fig. 6

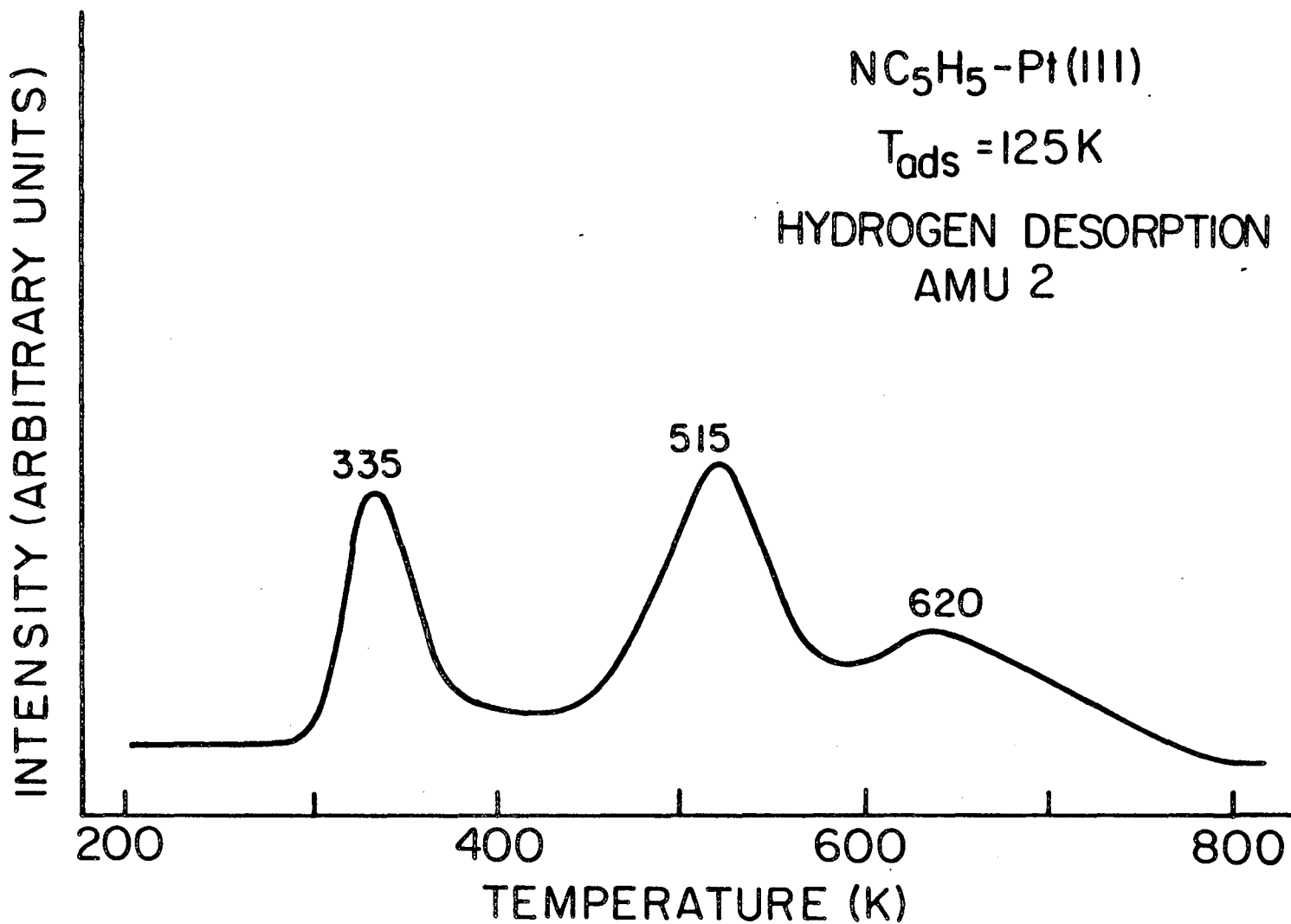


Fig. 7

XBL 8412-5936

This report was done with support from the Department of Energy. Any conclusions or opinions expressed in this report represent solely those of the author(s) and not necessarily those of The Regents of the University of California, the Lawrence Berkeley Laboratory or the Department of Energy.

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