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Authors

Aad, G
Abbott, B
Abbott, DC
[et al.](#)

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Search for resonances decaying into a weak vector boson and a Higgs boson in the fully hadronic final state produced in proton–proton collisions at $\sqrt{s} = 13$ TeV with the ATLAS detector

The ATLAS Collaboration

A search for heavy resonances decaying into a W or Z boson and a Higgs boson produced in proton–proton collisions at the Large Hadron Collider at $\sqrt{s} = 13$ TeV is presented. The analysis utilizes the dominant $W \rightarrow q\bar{q}'$ or $Z \rightarrow q\bar{q}$ and $H \rightarrow b\bar{b}$ decays with substructure techniques applied to large-radius jets. A sample corresponding to an integrated luminosity of 139 fb^{-1} collected with the ATLAS detector is analyzed and no significant excess of data is observed over the background prediction. The results are interpreted in the context of the heavy vector triplet model with spin-1 W' and Z' bosons. Upper limits on the cross section are set for resonances with mass between 1.5 and 5.0 TeV, ranging from 6.8 to 0.53 fb for $W' \rightarrow WH$ and from 8.7 to 0.53 fb for $Z' \rightarrow ZH$ at the 95% confidence level.

1 Introduction

The search for physics beyond the Standard Model (SM) is a major focus of the physics program at the Large Hadron Collider (LHC). Since its discovery [1, 2], the Higgs boson has become a tool in this search. In particular, one may expect new heavy resonances to couple to Higgs bosons and weak vector bosons ($V = W$ or Z). Such resonances are expected to occur in a number of theories beyond the Standard Model. Theories that aim to solve the naturalness problem predict the existence of vector resonances as expected in composite Higgs models [3, 4], Little Higgs models [5], or models with extra dimensions [6, 7]. Theories with extended Higgs sectors predict scalar resonances as in two-Higgs-doublet models [8].

In this article, a search for WH and ZH resonances produced in proton–proton (pp) collisions at $\sqrt{s} = 13$ TeV is reported with a sample corresponding to an integrated luminosity of 139 fb^{-1} collected with the ATLAS detector during Run 2 of the LHC in 2015–2018. The search is designed for resonances with a mass of at least 1.5 TeV and with both the V and H bosons decaying hadronically in the modes $V \rightarrow q\bar{q}^{(\prime)}$ and $H \rightarrow b\bar{b}$, as shown in Figure 1. In this regime, the V and H bosons are produced with high transverse momentum (p_T), resulting in each boson being reconstructed as a single large-radius hadronic jet, and the invariant mass of this dijet system provides the final discriminating variable. Jet substructure techniques and b -tagging are then used to discriminate those jets from background jets originating from multijet, V +jets, and $t\bar{t}$ events – with QCD multijet events representing at least 85% of the total background. Due to difficulties in modeling the background from simulation, all background estimates are derived from the data.

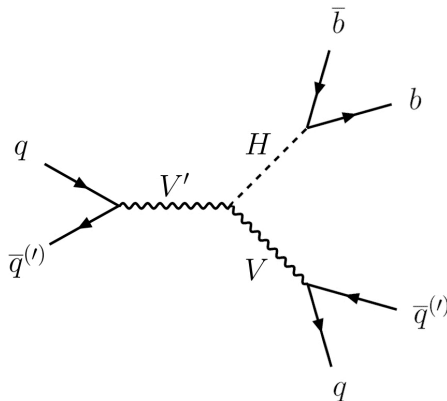


Figure 1: Feynman diagram for the production of a V' resonance with decay into a VH pair.

The results of the search are interpreted in the context of the heavy vector triplet (HVT) model [9], which is a simplified model providing a broad phenomenological framework for heavy resonances coupling to SM fermions and bosons. In this model, W' and Z' vector bosons interact with quarks and the Higgs field with coupling strength of g_q and g_H , respectively.¹ Coupling to the Higgs field gives rise to interactions with longitudinally polarized W and Z bosons. Two scenarios are considered as benchmarks for interpretation in this article. Model A corresponds to the choice $g_q = -0.55$ and $g_H = -0.56$, which reproduces the phenomenology of weakly coupled models based on an extended gauge symmetry [11]. Model B corresponds to $g_q = 0.14$ and $g_H = -2.9$, which implements a strongly coupled scenario as in composite Higgs models.

¹ Further details about the use of the HVT model in ATLAS analyses can be found in Ref. [10].

Previous searches for VH resonances have been carried out at $\sqrt{s} = 13$ TeV in the semileptonic final state ($\nu\nu bb$, $\ell\nu bb$, and $\ell\ell bb$) [12–14] and fully hadronic final state ($qqbb$) [15, 16]. The ATLAS and CMS collaborations report similar lower limits on the masses of W' and Z' bosons in these two sets of final states with about 36 fb^{-1} of integrated luminosity collected in 2015–2016. The strongest lower limit on the W' mass is set by CMS in the $\ell\nu bb$ channel [13] with a value of 2.9 TeV at the 95% confidence level (CL) in the context of HVT model B. For Z' bosons, the strongest lower limit on the mass is set by ATLAS in the combination of the $\nu\nu bb$ and $\ell\ell bb$ channels [12] with a value of 2.83 TeV at the 95% CL in HVT model B.

The results presented in this article differ from those previously published by ATLAS in the $qqbb$ channel [15] thanks to the following improvements. The integrated luminosity has increased by a factor of nearly four, an improved clustering algorithm combining measurements from the calorimeter and tracking systems is used to reconstruct V - and H -candidate jets, the b -tagging procedure used to identify H -candidate jets is performed on track-jets with a p_T -dependent radius that allows double b -tagging of H -candidate jets up to considerably higher p_T values, and the Higgs-candidate selection has been reoptimized with increased sensitivity.

2 ATLAS detector

The ATLAS experiment [17] at the LHC is a multipurpose particle detector with a forward–backward symmetric cylindrical geometry and a near 4π coverage in solid angle.² It consists of an inner detector (ID) for tracking surrounded by a thin superconducting solenoid providing a 2 T axial magnetic field, electromagnetic (EM) and hadronic calorimeters, and a muon spectrometer (MS). The inner detector covers the pseudorapidity range $|\eta| < 2.5$. It consists of silicon pixel, silicon microstrip, and transition radiation tracking detectors. An additional innermost pixel layer [18, 19] inserted at a radius of 3.3 cm has been used since 2015. Liquid-argon (LAr) sampling calorimeters provide EM energy measurements with high granularity. A hadronic scintillator-tile calorimeter covers the central pseudorapidity range ($|\eta| < 1.7$). The endcap and forward regions are instrumented with LAr calorimeters for both the EM and hadronic energy measurements up to $|\eta| = 4.9$. The muon spectrometer surrounds the calorimeters and features three large air-core toroidal superconducting magnet systems with eight coils each. The field integral of the toroids ranges between 2.0 and 6.0 Tm across most of the detector. The muon spectrometer includes a system of precision tracking chambers and fast detectors for triggering. A two-level trigger system [20] is used to select events. The first-level trigger is implemented in hardware and uses a subset of the detector information to reduce the accepted rate to at most 100 kHz. This is followed by a software-based trigger level that reduces the accepted event rate to 1 kHz on average.

3 Data and Monte Carlo simulation

The data sample for the analysis was collected by the ATLAS detector with high- p_T single-jet triggers utilizing the anti- k_t clustering algorithm [21] with a radius $R = 1.0$. The lowest unrescaled triggers were

² ATLAS uses a right-handed coordinate system with its origin at the nominal interaction point (IP) in the center of the detector and the z -axis along the beam pipe. The x -axis points from the IP to the center of the LHC ring, and the y -axis points upwards. Cylindrical coordinates (r, ϕ) are used in the transverse plane, ϕ being the azimuthal angle around the z -axis. The pseudorapidity is defined in terms of the polar angle θ as $\eta = -\ln \tan(\theta/2)$. The rapidity is defined relative to the beam axis as $y = (1/2) \ln[(E + p_z)/(E - p_z)]$. Angular distance is measured in units of $\Delta R \equiv \sqrt{(\Delta\eta)^2 + (\Delta\phi)^2}$.

used with the following p_T thresholds: 360 GeV in 2015, 420 GeV in 2016, and 460 GeV in 2017–2018. After requiring that the data were recorded with stable beam conditions and satisfied detector and data quality requirements, the integrated luminosity was measured to be 139 fb^{-1} using the methodology from Ref. [22].

The analysis relies on Monte Carlo (MC) samples to model signal events. Background MC samples are used only to optimize the signal event selection and to validate the data-driven background estimation method (Section 5).

Signal W' and Z' events for HVT model A were produced at leading-order (LO) precision in the strong coupling constant (α_s) with the MADGRAPH5_aMC@NLO 2.2.2 [23] event generator using the NNPDF23LO parton distribution function (PDF) set [24]. Separate generation of signal events for HVT model B is not required as both models A and B give rise to dijet mass peaks with a width that is dominated by the experimental resolution. The events were interfaced with PYTHIA 8.186 [25] for parton showering, hadronization, and the underlying event, and relied on the A14 set of tuned parameters [26]. Higgs boson decays to heavy-flavor final states $H \rightarrow b\bar{b}$ and $H \rightarrow c\bar{c}$ were included, corresponding to branching fractions of 58.2% and 2.9%, respectively [27].

Multijet events were produced with the PYTHIA 8.186 event generator, the NNPDF23LO PDF set, and the A14 tune. Samples of events with top-quark pairs were produced at next-to-leading order (NLO) with POWHEG-BOX [28] and the NNPDF30 NLO PDF set, interfaced with PYTHIA 8.183 and the A14 tune. The h_{damp} parameter was set to 1.5 times the top-quark mass [29]. Samples of W +jets and Z +jets events were produced with SHERPA 2.1.1 [30] and the CT10 PDF set [31] for up to two partons at NLO and up to four partons at LO. The cross sections used to normalize the multijet and the V +jets MC samples were computed with PYTHIA and SHERPA, respectively, and the top-quark pair cross section was taken to be $832_{-52}^{+46} \text{ pb}$ for a top-quark mass of 172.5 GeV. This value was calculated at next-to-next-to-leading order in α_s , including the resummation of next-to-next-to-leading logarithmic soft gluon terms, with Top++2.0 [32–38]. Other SM backgrounds originating from diboson production (VV) and weak vector-boson production in association with a Higgs boson (VH) are negligible and not considered.

For all MC samples, except those produced with SHERPA, b -hadron and c -hadron decays were handled by EVTGEN 1.2.0 [39]. Inelastic pp events generated using PYTHIA 8.186 with the A3 tune [40] and the NNPDF23LO PDF set were added to the hard-scattering interaction in such a way as to reproduce the effects of additional pp interactions (pileup) in each bunch crossing during data collection. The detector response was simulated with GEANT 4 [41, 42], and the events were processed with the same reconstruction software as for the data. Energy/momentum scale and efficiency corrections are applied to the results of the simulation to account for small differences between the simulation and the performance measured directly from the data.

4 Event reconstruction and selection

The analysis relies on the reconstruction of charged particles with $p_T > 500 \text{ MeV}$ in the inner detector to reconstruct pp collision vertices for each crossing of the LHC beams. The primary vertex is chosen to be the vertex with the largest $\sum p_T^2$ for the tracks associated with the vertex.

Jets are built from a combination of tracks and calibrated clusters of energy deposits in calorimeter cells [43], with the anti- k_r clustering algorithm using a radius parameter $R = 1.0$ as implemented in FASTJET [44]. The tracks are selected with the same requirements as in Ref. [45], except for the minimum p_T value, which

has been increased to 500 MeV. By combining calorimeter and tracking information, one benefits from both the better energy resolution of the calorimeter at high energy and the superior angular resolution for the tracks. This combination becomes highly beneficial at large jet p_T due to the small number of clusters produced and the limited angular resolution of the calorimeter. The resulting jets are referred to as Track-CaloCluster (TCC) jets. A detailed description of the algorithm can be found in Ref. [46] and its application to a search for high-mass diboson resonances is described in Ref. [47]. A trimming algorithm [48] is applied to minimize the impact of pileup. In this algorithm, the constituents of each jet are reclustered with the k_t algorithm [49] into smaller $R = 0.2$ subjets. Trimmed large- R jets are made up of constituents of those subjets with $p_T^{\text{subjet}} / p_T^{\text{jet}} > 0.05$, where p_T^{subjet} and p_T^{jet} are the transverse momenta of the subjet and original untrimmed jet, respectively. The energy and mass calibration of TCC jets is based on the simulation as described in Ref. [50]. As a result of the improved angular resolution of the energy distribution within the jet, the discrimination between signal W or Z jets and background QCD jets is noticeably improved. In addition to their masses, a powerful variable to discriminate between those jets is D_2 , defined as the ratio of three-point to two-point energy correlation functions that are based on the energies of the jet constituents and their pairwise angular separation [51, 52]. The D_2 variable exploits the two-body structure of the $V \rightarrow q\bar{q}^{(\prime)}$ decays, absent from typical QCD jets. Another variable that provides discrimination between signal and background jets is the number of tracks (n_{trk}) matched to the jets by ghost association [53]. This quantity is significantly higher for gluon-induced jets that are a component of the background than for quark-induced jets in signal events, due to the distinct energy scales involved and the different color factors for gluons and quarks.

To identify Higgs-boson jets, a separate collection of jets is built from tracks with the anti- k_t algorithm using a p_T -dependent radius $R = \rho/p_T$ [54], where the parameter ρ is set to 30 GeV and the radius is constrained to remain in the range between 0.02 and 0.4 [55]. The track-jets are assigned to specific large- R jets by ghost association with the original untrimmed large- R jets. The main advantage of using such variable-radius track-jets is that one can resolve the track-jets from $H \rightarrow b\bar{b}$ decays at high p_T and retain the ability to double b -tag the large- R Higgs-candidate jets.

Track-jets are tagged as likely to contain b -hadrons if they satisfy the selection criteria of the MV2c10 algorithm [56, 57] that takes advantage of the relatively long lifetime and large mass of b -hadrons. A working point corresponding to a b -tagging efficiency of 77% for true b -jets is used. For this efficiency value, rejection factors of 5 and 110 are obtained against c -quark and light-quark jets, with the efficiency and rejection factors determined in $t\bar{t}$ MC simulation [58].

Electrons are reconstructed by matching ID tracks to energy clusters in the EM calorimeter. The identification of electrons relies on a likelihood discriminant that takes the characteristic shape of electromagnetic showers into account [59]. Electrons are required to have $p_T > 7$ GeV and $|\eta| < 2.47$, and satisfy the *loose* identification criteria [59]. The associated tracks must have a transverse impact parameter significance $|d_0|/\sigma_{d_0} < 5$ relative to the beam axis and a longitudinal impact parameter $|z_0 \sin \theta| < 0.5$ mm relative to the primary vertex. Muons are reconstructed and identified by matching ID and MS tracks, and performing a global fit with all ID and MS measurements, taking the energy loss in the calorimeter into account [60]. Muons are required to have $p_T > 7$ GeV and $|\eta| < 2.5$, and satisfy the *loose* identification criteria [60]. The following track requirements are applied: $|d_0|/\sigma_{d_0} < 3$ and $|z_0 \sin \theta| < 0.5$ mm. Both the electrons and muons are required to satisfy loose isolation criteria [59, 60] based on the total transverse momentum of tracks surrounding the leptons within a cone of radius $\Delta R = \min(10/p_T^\ell [\text{GeV}], \Delta R_{\text{max}})$, where p_T^ℓ is the lepton p_T , and $\Delta R_{\text{max}} = 0.2$ for electrons and 0.3 for muons. The isolation criteria have an efficiency of 99% with negligible dependence on the lepton p_T value.

Events must satisfy the trigger requirements and contain a primary vertex. In addition, non-collision backgrounds originating from calorimeter noise, beam halo interactions or cosmic rays are suppressed by rejecting events that contain any $R = 0.4$ anti- k_t calorimeter jet failing to satisfy a set of quality criteria. These are based on the LAr pulse shape, the energy profile of the jet in different parts of the calorimeter, and track variables [61]. Events with one or more charged leptons (electrons or muons) are also rejected to retain orthogonality with other VH search channels.

The signal topology requires the presence of two large- R jets with $|\eta| < 2.0$ and $p_T > 200$ GeV. The leading (highest p_T) jet must satisfy $p_T > 500$ GeV. The invariant mass of the dijet system consisting of the two highest- p_T jets in the event (m_{JJ}) is required to be larger than 1.3 TeV. These kinematic requirements guarantee that the trigger is fully efficient. To suppress t -channel multijet production, the difference between the rapidities of the two leading jets must satisfy $|\Delta y| < 1.6$. Only the two leading- p_T jets are retained for further consideration.

As an initial step, the jet with the larger mass is taken to be the H -boson candidate while the other jet is taken to be the V -boson candidate. Discrimination between these V -boson or H -boson candidates and background jets relies on several properties of the large- R jets: mass, D_2 , and n_{trk} . An optimization procedure is applied to adjust the selection criteria involving those variables to maximize the significance of the resonance signal under study. In the case of V -boson jets, the selection is based on the three discriminating variables as developed in the search for heavy diboson resonances in the fully hadronic channel [47], with the exception that the n_{trk} selection is loosened slightly. In the case of the H -boson candidates, the selection criteria are optimized with regard to the jet mass and n_{trk} . The successful H -boson candidate has at least one associated track-jet and can be classified as either 1-tag or 2-tag, depending on the number of track-jets satisfying the b -tagging requirements. Only the two highest- p_T associated track-jets are considered for b -tagging. The variable D_2 provides little additional discrimination and is thus dropped.

The p_T -dependent jet mass windows for the $W(Z)$ -boson candidates that result from the optimization procedure vary from 80–100 (85–110) GeV for jets with p_T around 500 GeV to 55–130 (65–135) GeV for jets with p_T around 3000 GeV. The upper bounds on D_2 for the V -boson candidates vary approximately from 1.0 to 1.5, for the same p_T values. The upper cut on the n_{trk} variable varies between 25 (26) and 31 (29) for $W(Z)$ -boson candidates, with looser requirements for higher- p_T jets. Given the experimental jet mass resolution, no exclusive selection of W or Z bosons is performed and the WH and ZH final states are searched for independently. For H -boson candidates, the mass windows applied to events in the 1-tag (2-tag) category vary from 80–135 (95–150) GeV to 105–155 (100–170) GeV for jets with p_T of 500 and 3000 GeV, respectively. The H -boson candidates classified as 1(2)-tag are required to have an associated n_{trk} value below 32 (35) to 44 (55), looser at higher p_T .

The event selection defining the signal region (SR) is summarized in Table 1. The resulting signal acceptance times efficiency ($\mathcal{A} \times \varepsilon$) for events in each category is shown in Figure 2. In the 1-tag category, $\mathcal{A} \times \varepsilon$ rises from approximately 3% to 10% for WH resonance masses increasing from 1.5 to 5.0 TeV. In the 2-tag category, $\mathcal{A} \times \varepsilon$ remains essentially constant at 4% for WH resonances. The different trends for the 1-tag and 2-tag selections as a function of $m(V')$ are the result of a combination of effects including the p_T dependence of both V - and H -tagging as well as the signal to background ratio. The $\mathcal{A} \times \varepsilon$ values are about 0.5% lower for ZH resonances due to the smaller mass separation between Z and H bosons. The experimental mass resolution for resonance masses of 2 (4) TeV is 3.5% (2.6%).

Table 1: Event selection requirements and definition of the different regions used in the analysis. Events in the signal, control, and validation regions must satisfy the preselection requirements.

Preselection	Veto non- $qqqq$ channels: No e (μ) with $p_T > 7$ GeV and $ \eta < 2.47$ (2.5) Event kinematics: ≥ 2 large- R jets with $p_T > 200$ GeV and $ \eta < 2.0$ leading large- R jet with $p_T > 500$ GeV leading and subleading large- R jets with $m_{JJ} > 1.3$ TeV leading and subleading large- R jets with $ \Delta y < 1.6$
V/H assignment	V -boson (H -boson) candidate is large- R jet with lower (higher) mass
Signal region (SRWH / SRZH)	V and H bosons: W -boson candidate within W jet mass, D_2 , and n_{trk} windows Z -boson candidate within Z jet mass, D_2 , and n_{trk} windows H -boson candidate within H jet mass, n_{trk} windows, with 1 or 2 b -tagged track-jets
Control region (CR)	Fail both SRWH and SRZH Pass H -boson candidate n_{trk} $(m_V < 65$ GeV & $m_H < 70$ GeV) or $(m_V > 110$ GeV & $m_H > 150$ GeV) or $(m_V < 65$ GeV & $m_H > 150$ GeV)
Validation region (VR1A)	Fail both SRWH and SRZH Pass V -boson candidate n_{trk} Pass H -boson candidate n_{trk} $65 < m_V < 110$ GeV & $m_H > 150$ GeV
Validation region (VR1B)	Fail both SRWH and SRZH Fail V -boson candidate n_{trk} Pass H -boson candidate n_{trk} $65 < m_V < 110$ GeV & $m_H > 150$ GeV
Validation region (VR2A)	Fail both SRWH and SRZH Pass V -boson candidate n_{trk} Pass H -boson candidate n_{trk} $m_V < 65$ GeV & $70 < m_H < 150$ GeV
Validation region (VR2B)	Fail both SRWH and SRZH Fail V -boson candidate n_{trk} Pass H -boson candidate n_{trk} $m_V < 65$ GeV & $70 < m_H < 150$ GeV

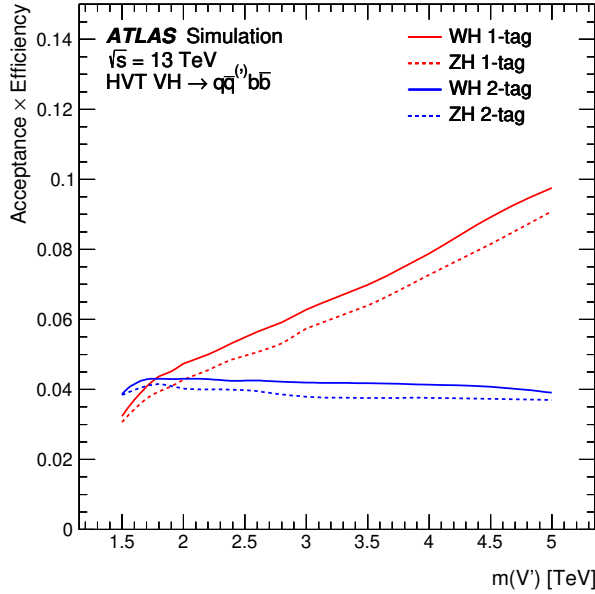


Figure 2: Signal acceptance times efficiency as a function of the resonance mass, for events in the WH (solid lines) and ZH (dashed lines) signal regions, in the 1-tag and 2-tag categories, with respect to the total number of generated events in each sample. The HVT MC samples include only $V' \rightarrow VH$ decays with $V \rightarrow q\bar{q}^{(\prime)}$ and $H \rightarrow b\bar{b}$ or $c\bar{c}$.

5 Background estimation and event yields

The dominant background in this search comes from multijet events, corresponding to at least 85% of the total background in the signal regions, where the remaining events come from $t\bar{t}$ and V +jets processes.

The totality of the background is estimated via a data-driven method that provides template m_{JJ} distributions for the WH and ZH final states in the 1-tag and 2-tag categories. These background m_{JJ} templates are obtained in three steps: (i) background templates for the WH and ZH final states are extracted from data in the 0-tag category, where the H -boson candidate has zero associated b -tagged jets; (ii) yield and shape corrections are derived from a control region and applied to these templates; and (iii) a rebinning and smoothing of the resulting m_{JJ} distributions is performed, to ensure robustness against statistical fluctuations. All steps are described in this section.

To define the 0-tag templates from which the background in the 1-tag and 2-tag categories is extracted, additional requirements are placed on 0-tag events, such that at least 1(2) variable-radius track-jet(s) is(are) associated with the H -boson jet, when estimating the background in the 1(2)-tag category. Events in the 0-tag category are not expected to directly describe either the yield or the m_{JJ} shape of the background in the 1-tag and 2-tag categories without further corrections: the b -tagging requirements enhance heavy-flavor components in the background and introduce kinematics-dependent effects.

Therefore, a control region (CR) is used to estimate the yield and shape corrections to the m_{JJ} distribution needed to extrapolate the 0-tag background events to the 1(2)-tag SRs. The CR has negligible contamination from signal. Validation regions (VR) with events that fail the SR selection are used to confirm the effectiveness of the background model and derive the associated systematic uncertainties. The definitions of the control and validation regions are shown in detail in Table 1 and illustrated in Figure 3. According to

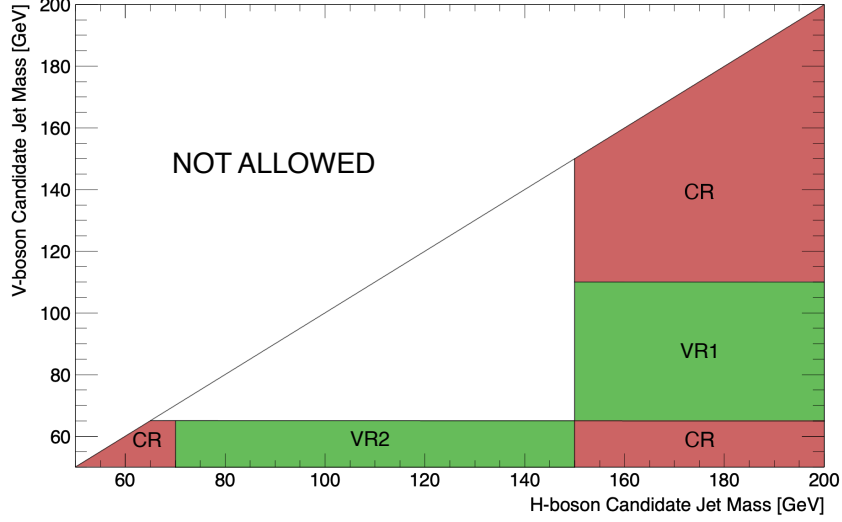


Figure 3: Illustration of control and validation regions, defined by the masses of the H -boson and V -boson candidates. The regions VR1 and VR2 are further split into two regions each, according to the n_{trk} requirement on the V -boson candidate. The same definitions are applied across the number of b -tags (0-, 1-, and 2-tag categories).

simulation, multijet processes are responsible for 85% to 99% of the background composition across the different control, validation, and signal regions. This variation and its impact on the background estimate are taken into account by the uncertainties described in Section 6.

To define the background normalization in the 1(2)-tag category, a normalization correction $\mu_{\text{CR}}^{1(2)\text{-tag}}$ is extracted from the ratio of 1(2)-tag to 0-tag yields in the control region:

$$\mu_{\text{CR}}^{1(2)\text{-tag}} = \frac{N_{\text{CR}}^{1(2)\text{-tag}}}{N_{\text{CR}}^{0\text{-tag}}},$$

where $N_{\text{CR}}^{0\text{-tag}}$ and $N_{\text{CR}}^{1(2)\text{-tag}}$ are the numbers of events observed in the 0-tag and 1(2)-tag CR event categories. The values of $\mu_{\text{CR}}^{1\text{-tag}}$ and $\mu_{\text{CR}}^{2\text{-tag}}$ are determined to be 0.160 ± 0.014 and 0.0167 ± 0.0028 , respectively, where the uncertainties are dominated by systematic effects discussed in Section 6. The difference in the corrections for 1-tag and 2-tag events can be understood based on studies of simulated multijet events. In 1-tag events, the two leading track-jets associated with Higgs candidate large- R jets are dominated by one true b -jet accompanied by one true light-jet, whereas in 2-tag events they are dominated by two true b -jets. This indicates that different processes are at work in 1-tag and 2-tag events. The expected number of background 1(2)-tag events in the SR is calculated by applying the $\mu_{\text{CR}}^{1(2)\text{-tag}}$ correction factor to the number of 0-tag events that pass all SR requirements except for the b -tagging.

To extract m_{JJ} background templates from the 0-tag category, a multidimensional kinematic reweighting [62] is performed using the control region events. A boosted decision tree (BDT) is used to perform the reweighting, by predicting the event weights needed to bring the shapes of kinematic distributions in the 0-tag and 1(2)-tag categories into agreement. The training is performed in the CR data and performed separately for 1-tag and 2-tag events. Variables that are sensitive to the presence of b -jets associated with the H -boson candidate and to the resulting kinematic differences are used to train the BDTs: the four-momenta of the two leading variable-radius track-jets and their angular separation, the transverse

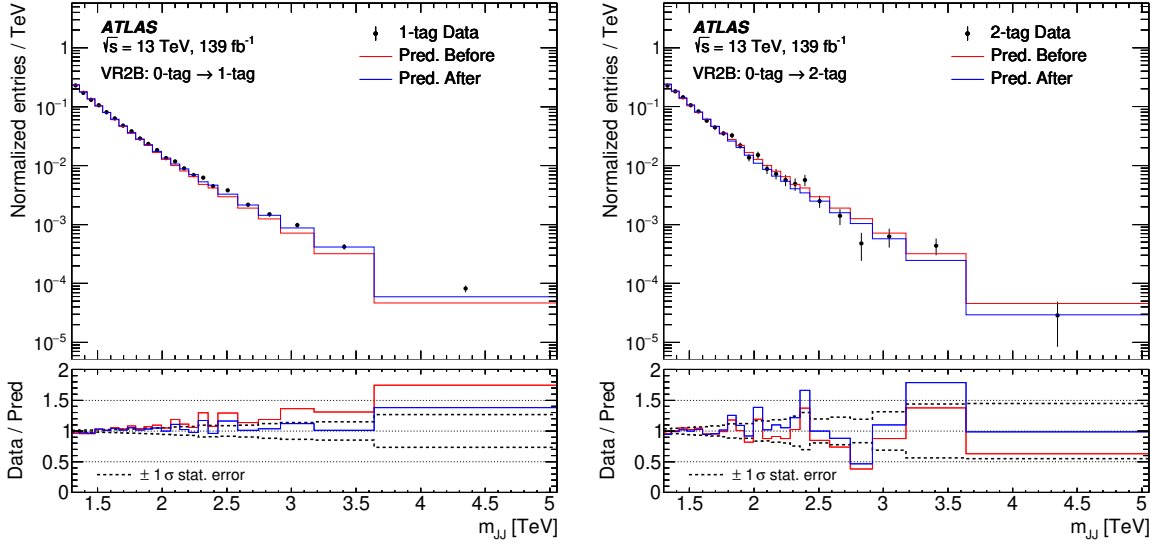


Figure 4: Dijet mass distributions in the 1-tag (left) and 2-tag (right) VR2B regions compared with the predicted background extracted from the 0-tag events (histograms) before and after BDT reweighting.

momenta of the H - and V -boson jets, and the number of tracks and variable-radius track-jets associated with the H -boson jet. The BDTs are built with 100 trees, a maximum depth of three layers, and a minimum of 500 events per leaf, with a learning rate set to 0.1.

In order to quantify the effectiveness of the reweighting, a binary classifier was trained to differentiate between the reweighted and target m_{JJ} distributions in the validation regions, and observed to classify them correctly at most 53% of the time, consistent with random guessing. The observed distributions of kinematic variables, including m_{JJ} , are found to be well described by the background model for 1-tag and 2-tag events in VR1A, VR1B, VR2A, and VR2B (defined in Table 1).

The modeling of the m_{JJ} distributions in the VR2B region is shown in Figure 4, for 1-tag and 2-tag events. A residual disagreement between the data and the expected background after reweighting is accounted for by a systematic uncertainty, as discussed in Section 6.

The numbers of 1-tag and 2-tag events observed in the control and validation regions are shown in Table 2 and compared with the predicted background yields.

After the normalization and reweighting corrections are applied to the events in the 0-tag category, the expected m_{JJ} background distributions in 1-tag and 2-tag categories are produced with a variable bin width that reflects the experimental mass resolution. Those distributions are then fit using a functional form that captures the smoothly falling behavior of the background:

$$f_{\text{Background}}(x) = e^{-p_0} (1-x)^{-p_1} x^{-p_2},$$

where $x = m_{JJ}/\sqrt{s}$ and p_0 , p_1 , and p_2 are the fit parameters. The results of these fits provide the background estimates that are used in the statistical analysis (described in Section 7) for the different signal regions.

Table 2: Data and estimated background yields for 1-tag and 2-tag events in the control and validation regions. The uncertainties correspond to the combination of statistical and systematic components. By construction, the uncertainties cover the differences between the observed and expected yields in the validation regions.

1-tag	Data	Background prediction
Control region (CR)	48668	--
Validation region (VR1A)	4440	4290 ± 380
Validation region (VR1B)	18361	18000 ± 1600
Validation region (VR2A)	11844	12600 ± 1100
Validation region (VR2B)	29436	31200 ± 2700
2-tag	Data	Background prediction
Control region (CR)	4976	--
Validation region (VR1A)	507	443 ± 74
Validation region (VR1B)	1922	1860 ± 310
Validation region (VR2A)	1337	1290 ± 220
Validation region (VR2B)	3065	3200 ± 540

6 Systematic uncertainties

Systematic uncertainties arise from several different sources: the data-driven background estimate, the modeling of experimental uncertainties affecting the signal, and the impact of signal theory uncertainties. At a resonance mass of 2 TeV, the background normalization and shape uncertainties dominate, while at a resonance mass of 4 TeV, the large-radius jet and boson tagging uncertainties dominate due to the small background contribution at high m_{JJ} .

The uncertainties affecting both the normalization and shape of the background predictions are determined from the validation regions. These uncertainties arise from limited sample sizes and from differences in the background composition in the various regions. The normalization uncertainty is taken to be the difference between the smallest and largest $\mu^{1(2)\text{-tag}}$ values obtained in any of the validation regions, resulting in a systematic error of 9% (17%) in the background estimate for the 1(2)-tag category. The uncertainty in the shape of the background m_{JJ} distribution is assessed from the ratio of data to prediction in the VR2B region, where the differences are the largest. This uncertainty is determined after smoothing the m_{JJ} distributions for both the data and the background prediction with the same functional form as described in Section 5. It results in changes to the background yield of approximately 5% at m_{JJ} of 2 TeV and up to 24% at 4 TeV.

An additional shape uncertainty is assigned to account for the choice of fitting function, assessed by fitting alternate empirical functions, amounting to a maximum uncertainty of 2% (14%) at an m_{JJ} value of 2 (4) TeV.

Experimental uncertainties related to MC simulation are applicable only to signal samples, and are divided into two categories: b -tagging and large- R jets. A set of b -tagging correction factors and corresponding uncertainties are applied as a function of p_T and η of the variable-radius track-jets to match the efficiencies for tagging b -jets measured in data, determined with $t\bar{t}$ events [58]. The uncertainties in the correction factors are extrapolated for track-jets with p_T larger than 400 GeV. An additional extrapolation uncertainty

is obtained by varying the inputs to the b -tagging algorithm according to their modeling uncertainties and by recomputing its efficiency in MC simulation [58].

Uncertainties in the p_T and mass scales of the large- R jet are determined with the R_{trk} method [63] adapted to the jet collection used in this article, relying on independent measurements by the calorimeter and the inner detector, and are of the order of 5% to 10% each. These uncertainties lead to shifts in m_{JJ} of the resonant signal peak as well as differences in the signal selection efficiency. The effects of resolution uncertainties on the p_T and mass measurements are estimated by degrading the resolution in MC simulation according to a Gaussian smearing of width 0.02 in $\sigma(p_T)/p_T$ and $0.20 \times \sigma(m)/m$ in $\sigma(m)/m$ [64, 65].

The H - and V -boson tagging techniques are assigned dedicated uncertainties to take into account the requirements on D_2 and n_{trk} . An MC efficiency correction factor of 0.92 ± 0.13 for V -boson tagging was determined in the search for heavy diboson resonances in the fully hadronic channel [47], by taking advantage of a control region in data that is enhanced in V +jets events. Given that the same method for V -tagging is used in this analysis, the same scale factor and uncertainty of 14% is assigned to the signal normalization.³ An efficiency correction and uncertainty in the n_{trk} requirement was also estimated in Ref. [66] for V -jets. This uncertainty is applied to H -tagging, with an additional component to cover topology differences, based on simulation studies, corresponding to a total 10% uncertainty. In particular, these studies compare the large-radius jet mass distributions between the data and the simulation, and the impact of the n_{trk} requirement on the data-to-MC agreement in the V -mass and H -mass regions.

Signal cross sections computed at leading order are used in the interpretation of the results. The impact of uncertainties in the PDF sets, initial- and final-state radiation, and multiparton interactions on the signal acceptance are included. Uncertainties related to the PDF sets are derived by applying the methodology outlined by the PDF4LHC group [24] and considering four additional PDF sets (CT14, MMHT2014, NNPDF3.0, and ATLAS-epWZ12), resulting in <1% uncertainties in the signal acceptance. An uncertainty due to choosing the A14 tune for the signal generation is estimated by varying the scales for initial- and final-state radiation, as well as multiparton interactions, and results in an uncertainty of 2% (3%) for WH (ZH) resonances.

Finally, an uncertainty in the Run 2 integrated luminosity of 1.7% [22] is considered, as obtained by the LUCID-2 detector [67] for the main luminosity measurements. The impact of the main systematic uncertainties on signal event yields is summarized in Table 3.

7 Results

The statistical analysis of the data is performed using a binned likelihood function, constructed from the m_{JJ} distributions in the 1-tag and 2-tag signal regions, using the procedure described in Ref. [1] and the RooStats framework [68]. The m_{JJ} histograms derived from MC simulation are used for the HVT W' and Z' processes, while the data-driven background estimates are used for the combined $t\bar{t}$, V +jets, and QCD multijet processes. The input m_{JJ} distribution bounds are [1.3, 6.0] TeV.

A test statistic based on the profile likelihood ratio [69] is used to test signal hypotheses, parameterized by the signal strength value, μ , acting as a scale factor on the predicted number of signal events for each

³ The only difference is in the n_{trk} requirement, which is looser in this analysis. Studies of relative signal efficiencies for the different n_{trk} requirements show that this approximation is well motivated.

Table 3: Systematic uncertainties affecting the signal event yields in the WH and ZH signal regions with 1-tag or 2-tag (denoted WH-1, WH-2, ZH-1, and ZH-2, respectively). The HVT model is used with resonance masses of 2 and 4 TeV.

Source	Signal (2 TeV)				Signal (4 TeV)			
	WH-1	WH-2	ZH-1	ZH-2	WH-1	WH-2	ZH-1	ZH-2
Jet energy scale	2.8%	3.1%	1.6%	2.1%	5.2%	8.4%	6.6%	8.0%
Jet mass scale	18%	18%	9.6%	14%	20%	17%	21%	18%
Jet mass resolution	43%	45%	37%	40%	22%	21%	21%	22%
Flavor tagging	17%	8.4%	16%	8.4%	26%	12%	26%	13%
H -boson tagging	10%	10%	10%	10%	10%	10%	10%	10%
V -boson tagging	13%	13%	13%	13%	13%	13%	13%	13%
Luminosity	1.7%	1.7%	1.7%	1.7%	1.7%	1.7%	1.7%	1.7%
MC statistical uncertainty	1.5%	1.4%	1.6%	1.4%	1.2%	1.4%	1.5%	1.7%

model assumption. The likelihood, L , is defined from the Poisson probability to observe N data events for a given signal s and background b expectation in each bin of the final discriminant:

$$L(\mu, \vec{\theta}) = \prod_{c=1}^{\text{categories}} \prod_{i=1}^{\text{bins}} \text{Pois} \left(N_{ci} | \mu s_{ci}(\vec{\theta}) + b_{ci}(\vec{\theta}) \right) \prod_{j \in \vec{\theta}} f_j(\theta_j),$$

where the index c represents the 1-tag or 2-tag event categories and i represents the bin in the m_{JJ} distribution. Nuisance parameters $\vec{\theta}$ are included in the likelihood function with Gaussian or log-normal constraint terms, $f_j(\theta_j)$. Those nuisance parameters which produce bin variations smaller than 1% from the nominal value are neglected.

Experimental uncertainties in the signal are fully correlated between the 1-tag and 2-tag signal regions, whereas background modeling uncertainties are kept independent. In order to avoid an overconstraining of the background modeling uncertainties in the high mass region due to the higher statistical power at low masses, the m_{JJ} shape uncertainties above and below 2.5 TeV are allowed to vary independently in the fit. The post-fit background expectation was found to be stable and independent of the particular choice of splitting point.

The fits are performed separately for the W' and Z' models, using data in the 1-tag and 2-tag regions from the SRWH and SRZH selections, respectively. The fit results are interpreted independently for the W' and Z' hypotheses – the WH and ZH signal regions are not orthogonal and have approximately 40% of events in common, in each category.

The pre- and post-fit m_{JJ} distributions in the signal region are shown in Figure 5, for signal resonances with a mass of 2 TeV. The numbers of data events in the signal regions are shown in Table 4, along with the predicted background and signal yields, post-fit. No events with m_{JJ} values above 5 TeV are selected.

A test of the background-only hypothesis is performed by setting μ equal to zero in the likelihood fit. Deviations from the background-only hypothesis are quantified by determining the local p_0 -value at each signal mass point. The largest deviation is observed in the fit to the WH signal regions and corresponds to a p_0 -value of 0.03 for a resonance mass of 2.8 TeV.

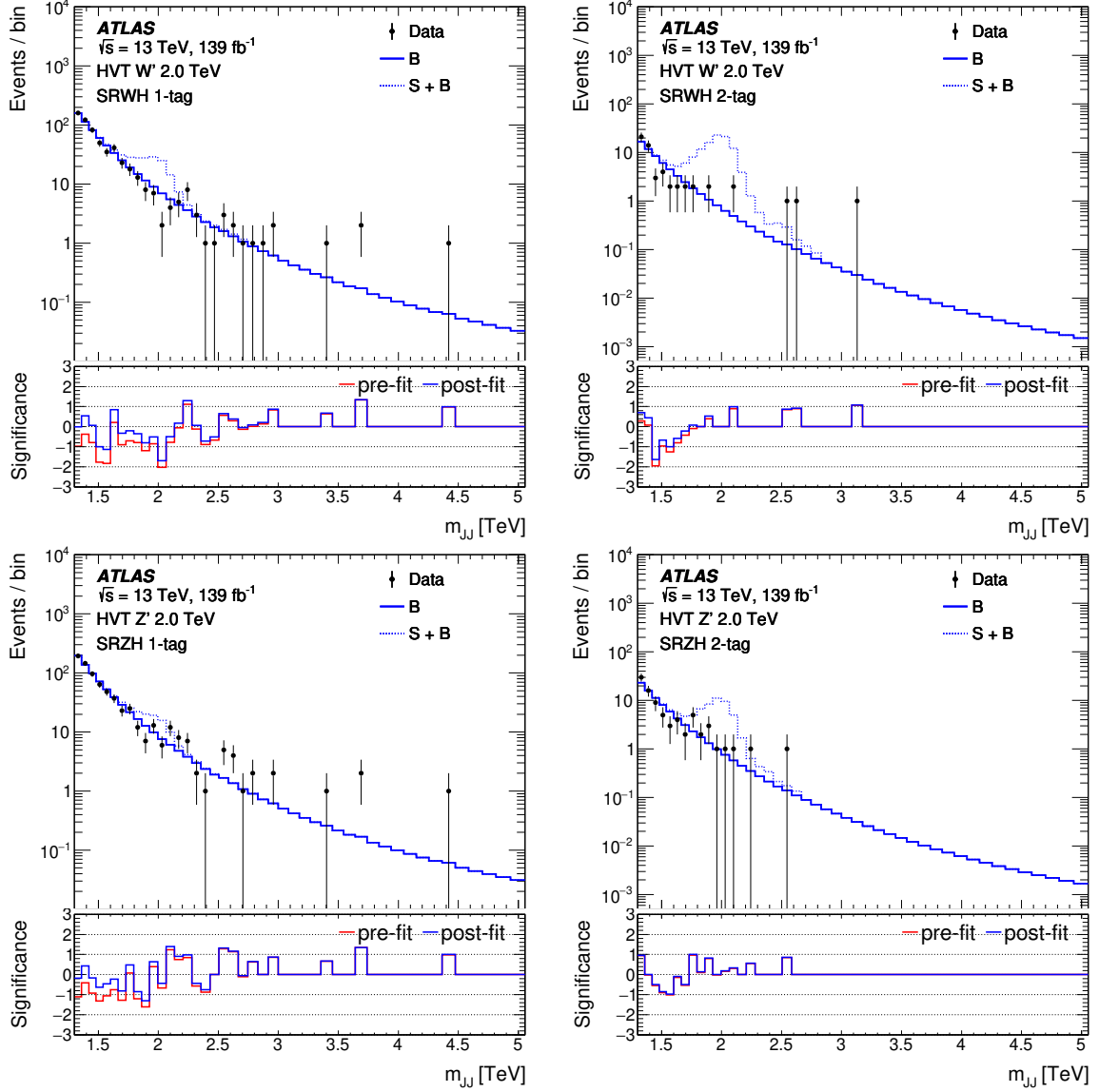


Figure 5: Dijet mass distributions in the WH (top) and ZH (bottom) signal regions, after the likelihood fit to events in the 1-tag (left) and 2-tag (right) categories. The black points correspond to data and the solid blue histogram to the post-fit background prediction. The WH and ZH signal regions are not orthogonal. The expected signal distributions for a V' boson with mass of 2 TeV are also shown (dashed histograms). The bin width varies and corresponds to the experimental mass resolution. Distributions of the significance of the observed deviations from the expected background are presented in the bottom panels before and after the fit. The significance calculation assumes Poisson probabilities and only accounts for statistical fluctuations.

Table 4: Event yields for the data, predicted background, and signal in each of the signal regions. The signal corresponds to that expected for HVT model B with resonance masses of 2 and 4 TeV.

Region	Data	Background	Signal (2 TeV)	Signal (4 TeV)
SRWH 1-tag	598	612 ± 46	110	2.0
SRWH 2-tag	57	61 ± 13	100	1.0
SRZH 1-tag	717	725 ± 53	47	0.80
SRZH 2-tag	84	81 ± 17	44	0.42

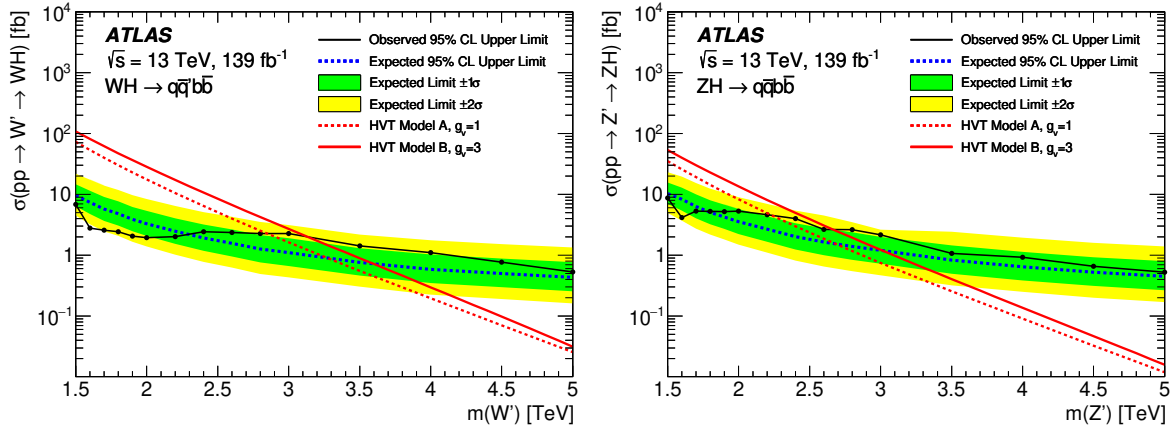


Figure 6: Observed and expected 95% CL upper limits on the cross section for $pp \rightarrow V' \rightarrow VH$ in the WH (left) and ZH (right) channels. The red solid (dashed) lines show the cross-section predictions as a function of the resonance mass in the context of HVT model B (A).

The data are used to set upper limits on the production cross section of new resonances. Exclusion limits are computed using the CL_s method [70], with a value of μ regarded as excluded at the 95% CL when the CL_s value is less than 5%. The cross-section limits are shown in Figure 6. The observed limits range from cross sections of 6.8 fb to 0.53 fb for WH and from 8.7 fb to 0.53 fb for ZH , corresponding to the exclusion of W' (Z') resonances up to a mass of 2.90 TeV (2.20 TeV) in the context of HVT model A and 3.20 TeV (2.65 TeV) in the context of HVT model B. The 2-tag category dominates the sensitivity of the search at low resonance mass while the impact of the 1-tag category increases at higher mass, surpassing the 2-tag category at masses above 3.9 TeV. As a test of the asymptotic approximation used in the statistical analysis, results are also obtained with ensembles of pseudoexperiments. The cross-section upper limits obtained in that case are looser by 10–20%, with a larger difference at high $m(V')$, and the mass limits are at most 0.05 TeV weaker.

These results can also be translated into exclusions in the $\{g_H, g_f\}$ plane, where g_f represents a universal coupling between the V' bosons and fermions. Here, g_q is taken to be equal to g_f . Figure 7 shows the 95% CL limits in this plane for several resonance masses.

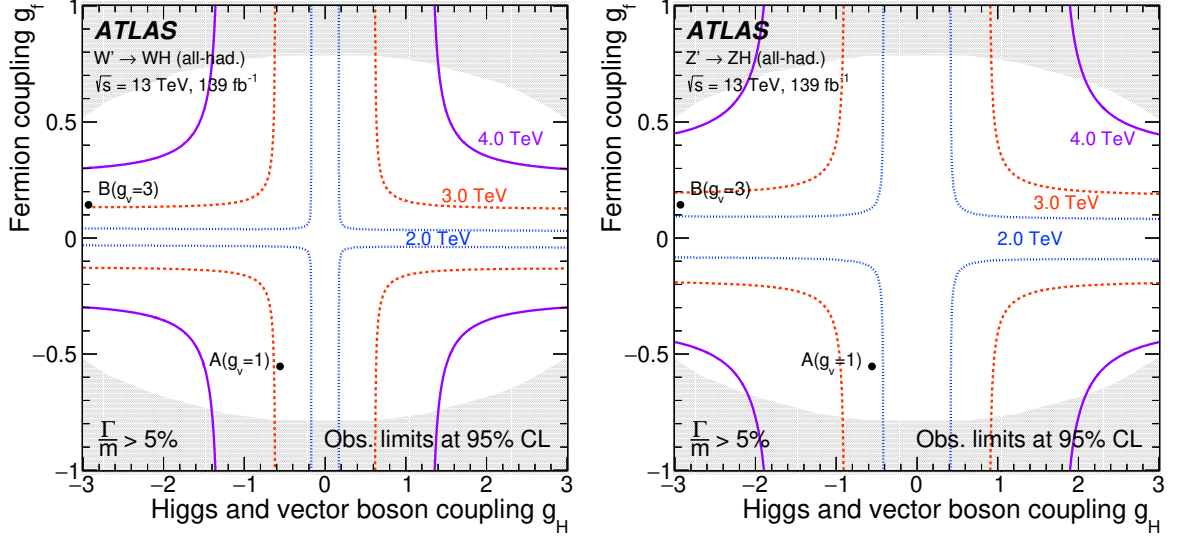


Figure 7: Limits at 95% CL in the g_f vs. g_H plane for resonance masses of 2, 3, and 4 TeV for the WH (left) and ZH (right) channels in the context of the HVT model. The coupling values corresponding to HVT models A and B are indicated by filled circles. The gray region indicates values of the couplings corresponding to V' resonances with Γ/m greater than 5%. In that region, the assumption that the V' signal m_{JJ} shape is dominated by the experimental resolution is no longer valid.

8 Conclusion

A search for heavy resonances decaying into a W or Z boson and a Higgs boson is reported. The results are based on a sample of pp collisions at $\sqrt{s} = 13$ TeV collected by the ATLAS detector at the Large Hadron Collider, corresponding to 139 fb^{-1} of integrated luminosity. The search exploits jet substructure techniques to study the fully hadronic $qqbb$ final state which results from the dominant decay modes of the W/Z and Higgs bosons. The main background contribution arises from multijet production, with a smaller contribution from top-quark pair and V +jets production. All background contributions are extracted directly from the data. Compared with the previously available results, the search benefits from an increased integrated luminosity, as well as from improvements in reconstruction and tagging of large- R jets and track-jets with p_T -dependent radius.

No significant excess of events is observed over the expected background and the upper limits set on the cross section for $pp \rightarrow W' \rightarrow WH$ and $pp \rightarrow Z' \rightarrow ZH$ range from 6.8 fb at $m(W') = 1.5$ TeV to 0.53 fb at $m(W') = 5.0$ TeV, and from 8.7 fb at $m(Z') = 1.5$ TeV to 0.53 fb at $m(Z') = 5.0$ TeV, at 95% CL. These results translate into lower limits on the mass of W' (Z') bosons of 2.90 TeV (2.20 TeV) in the context of the weakly coupled HVT model A and of 3.20 TeV (2.65 TeV) in the context of the strongly coupled HVT model B, at 95% CL.

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The ATLAS Collaboration

G. Aad¹⁰², B. Abbott¹²⁸, D.C. Abbott¹⁰³, A. Abed Abud³⁶, K. Abeling⁵³, D.K. Abhayasinghe⁹⁴, S.H. Abidi¹⁶⁷, O.S. AbouZeid⁴⁰, N.L. Abraham¹⁵⁶, H. Abramowicz¹⁶¹, H. Abreu¹⁶⁰, Y. Abulaiti⁶, B.S. Acharya^{67a,67b,n}, B. Achkar⁵³, L. Adam¹⁰⁰, C. Adam Bourdarios⁵, L. Adamczyk^{84a}, L. Adamek¹⁶⁷, J. Adelman¹²¹, M. Adersberger¹¹⁴, A. Adiguzel^{12c,ae}, S. Adorni⁵⁴, T. Adye¹⁴³, A.A. Affolder¹⁴⁵, Y. Afik¹⁶⁰, C. Agapopoulou⁶⁵, M.N. Agaras³⁸, A. Aggarwal¹¹⁹, C. Agheorghiesei^{27c}, J.A. Aguilar-Saavedra^{139f,139a,ad}, A. Ahmad³⁶, F. Ahmadov⁸⁰, W.S. Ahmed¹⁰⁴, X. Ai¹⁸, G. Aielli^{74a,74b}, S. Akatsuka⁸⁶, M. Akbiyik¹⁰⁰, T.P.A. Åkesson⁹⁷, E. Akilli⁵⁴, A.V. Akimov¹¹¹, K. Al Houry⁶⁵, G.L. Alberghi^{23b,23a}, J. Albert¹⁷⁶, M.J. Alconada Verzini¹⁶¹, S. Alderweireldt³⁶, M. Aleksa³⁶, I.N. Aleksandrov⁸⁰, C. Alexa^{27b}, T. Alexopoulos¹⁰, A. Alfonsi¹²⁰, F. Alfonsi^{23b,23a}, M. Alhroob¹²⁸, B. Ali¹⁴¹, S. Ali¹⁵⁸, M. Aliev¹⁶⁶, G. Alimonti^{69a}, C. Allaire³⁶, B.M.M. Allbrooke¹⁵⁶, B.W. Allen¹³¹, P.P. Allport²¹, A. Aloisio^{70a,70b}, F. Alonso⁸⁹, C. Alpigiani¹⁴⁸, E. Alunno Camelia^{74a,74b}, M. Alvarez Estevez⁹⁹, M.G. Alviggi^{70a,70b}, Y. Amaral Coutinho^{81b}, A. Ambler¹⁰⁴, L. Ambroz¹³⁴, C. Amelung²⁶, D. Amidei¹⁰⁶, S.P. Amor Dos Santos^{139a}, S. Amoroso⁴⁶, C.S. Amrouche⁵⁴, F. An⁷⁹, C. Anastopoulos¹⁴⁹, N. Andari¹⁴⁴, T. Andeen¹¹, J.K. Anders²⁰, S.Y. Andrean^{45a,45b}, A. Andreazza^{69a,69b}, V. Andrei^{61a}, C.R. Anelli¹⁷⁶, S. Angelidakis⁹, A. Angerami³⁹, A.V. Anisenkov^{122b,122a}, A. Annovi^{72a}, C. Antel⁵⁴, M.T. Anthony¹⁴⁹, E. Antipov¹²⁹, M. Antonelli⁵¹, D.J.A. Antrim¹⁷¹, F. Anulli^{73a}, M. Aoki⁸², J.A. Aparisi Pozo¹⁷⁴, M.A. Aparo¹⁵⁶, L. Aperio Bella⁴⁶, N. Aranzabal³⁶, V. Araujo Ferraz^{81a}, R. Araujo Pereira^{81b}, C. Arcangeletti⁵¹, A.T.H. Arce⁴⁹, F.A. Arduh⁸⁹, J-F. Arguin¹¹⁰, S. Argyropoulos⁵², J.-H. Arling⁴⁶, A.J. Armbruster³⁶, A. Armstrong¹⁷¹, O. Arnaez¹⁶⁷, H. Arnold¹²⁰, Z.P. Arrubarrena Tame¹¹⁴, G. Artoni¹³⁴, H. Asada¹¹⁷, K. Asai¹²⁶, S. Asai¹⁶³, T. Asawatavonvanich¹⁶⁵, N. Asbah⁵⁹, E.M. Asimakopoulou¹⁷², L. Asquith¹⁵⁶, J. Assahsah^{35e}, K. Assamagan²⁹, R. Astalos^{28a}, R.J. Atkin^{33a}, M. Atkinson¹⁷³, N.B. Atlay¹⁹, H. Atmani⁶⁵, K. Augsten¹⁴¹, V.A. Austrup¹⁸², G. Avolio³⁶, M.K. Ayoub^{15a}, G. Azeulos^{110,am}, H. Bachacou¹⁴⁴, K. Bachas¹⁶², M. Backes¹³⁴, F. Backman^{45a,45b}, P. Bagnaia^{73a,73b}, M. Bahmani⁸⁵, H. Bahrasemani¹⁵², A.J. Bailey¹⁷⁴, V.R. Bailey¹⁷³, J.T. Baines¹⁴³, C. Bakalis¹⁰, O.K. Baker¹⁸³, P.J. Bakker¹²⁰, E. Bakos¹⁶, D. Bakshi Gupta⁸, S. Balaji¹⁵⁷, R. Balasubramanian¹²⁰, E.M. Baldin^{122b,122a}, P. Balek¹⁸⁰, F. Balli¹⁴⁴, W.K. Balunas¹³⁴, J. Balz¹⁰⁰, E. Banas⁸⁵, M. Bandieramonte¹³⁸, A. Bandyopadhyay²⁴, Sw. Banerjee^{181,i}, L. Barak¹⁶¹, W.M. Barbe³⁸, E.L. Barberio¹⁰⁵, D. Barberis^{55b,55a}, M. Barbero¹⁰², G. Barbour⁹⁵, T. Barillari¹¹⁵, M-S. Barisits³⁶, J. Barkeloo¹³¹, T. Barklow¹⁵³, R. Barnea¹⁶⁰, B.M. Barnett¹⁴³, R.M. Barnett¹⁸, Z. Barnovska-Blenessy^{60a}, A. Baroncelli^{60a}, G. Barone²⁹, A.J. Barr¹³⁴, L. Barranco Navarro^{45a,45b}, F. Barreiro⁹⁹, J. Barreiro Guimarães da Costa^{15a}, U. Barron¹⁶¹, S. Barsov¹³⁷, F. Bartels^{61a}, R. Bartoldus¹⁵³, G. Bartolini¹⁰², A.E. Barton⁹⁰, P. Bartos^{28a}, A. Basalae⁴⁶, A. Basan¹⁰⁰, A. Bassalat^{65,aj}, M.J. Basso¹⁶⁷, R.L. Bates⁵⁷, S. Batlamous^{35f}, J.R. Batley³², B. Batool¹⁵¹, M. Battaglia¹⁴⁵, M. Baue^{73a,73b}, F. Bauer¹⁴⁴, P. Bauer²⁴, H.S. Bawa³¹, A. Bayirli^{12c}, J.B. Beacham⁴⁹, T. Beau¹³⁵, P.H. Beauchemin¹⁷⁰, F. Becherer⁵², P. Bechtel²⁴, H.C. Beck⁵³, H.P. Beck^{20,p}, K. Becker¹⁷⁸, C. Becot⁴⁶, A. Beddall^{12d}, A.J. Beddall^{12a}, V.A. Bednyakov⁸⁰, M. Bedognetti¹²⁰, C.P. Bee¹⁵⁵, T.A. Beermann¹⁸², M. Begalli^{81b}, M. Begel²⁹, A. Behera¹⁵⁵, J.K. Behr⁴⁶, F. Beisiegel²⁴, M. Belfkir⁵, A.S. Bell⁹⁵, G. Bella¹⁶¹, L. Bellagamba^{23b}, A. Bellerive³⁴, P. Bellos⁹, K. Beloborodov^{122b,122a}, K. Belotskiy¹¹², N.L. Belyaev¹¹², D. Benchekroun^{35a}, N. Benekos¹⁰, Y. Benhammou¹⁶¹, D.P. Benjamin⁶, M. Benoit²⁹, J.R. Bensinger²⁶, S. Bentvelsen¹²⁰, L. Beresford¹³⁴, M. Beretta⁵¹, D. Berge¹⁹, E. Bergeaas Kuutmann¹⁷², N. Berger⁵, B. Bergmann¹⁴¹, L.J. Bergsten²⁶, J. Beringer¹⁸, S. Berlendis⁷, G. Bernardi¹³⁵, C. Bernius¹⁵³, F.U. Bernlochner²⁴, T. Berry⁹⁴, P. Berta¹⁰⁰, A. Berthold⁴⁸, I.A. Bertram⁹⁰, O. Bessidskaia Bylund¹⁸², N. Besson¹⁴⁴, A. Bethani¹⁰¹, S. Bethke¹¹⁵, A. Betti⁴², A.J. Bevan⁹³, J. Beyer¹¹⁵, D.S. Bhattacharya¹⁷⁷, P. Bhattarai²⁶, V.S. Bhopatkar⁶, R. Bi¹³⁸, R.M. Bianchi¹³⁸, O. Biebel¹¹⁴, D. Biedermann¹⁹, R. Bielski³⁶, K. Bierwagen¹⁰⁰, N.V. Biesuz^{72a,72b}, M. Biglietti^{75a}, T.R.V. Billoud¹⁴¹, M. Bindi⁵³, A. Bingul^{12d},

C. Bini^{73a,73b}, S. Biondi^{23b,23a}, C.J. Birch-sykes¹⁰¹, M. Birman¹⁸⁰, T. Bisanz³⁶, J.P. Biswal³, D. Biswas^{181,i}, A. Bitadze¹⁰¹, C. Bittrich⁴⁸, K. Bjørke¹³³, T. Blazek^{28a}, I. Bloch⁴⁶, C. Blocker²⁶, A. Blue⁵⁷, U. Blumenschein⁹³, G.J. Bobbink¹²⁰, V.S. Bobrovnikov^{122b,122a}, S.S. Bocchetta⁹⁷, D. Bogavac¹⁴, A.G. Bogdanchikov^{122b,122a}, C. Bohm^{45a}, V. Boisvert⁹⁴, P. Bokan^{172,53}, T. Bold^{84a}, A.E. Bolz^{61b}, M. Bomben¹³⁵, M. Bona⁹³, J.S. Bonilla¹³¹, M. Boonekamp¹⁴⁴, C.D. Booth⁹⁴, A.G. Borbély⁵⁷, H.M. Borecka-Bielska⁹¹, L.S. Borgna⁹⁵, A. Borisov¹²³, G. Borissov⁹⁰, D. Bortoletto¹³⁴, D. Boscherini^{23b}, M. Bosman¹⁴, J.D. Bossio Sola¹⁰⁴, K. Bouaouda^{35a}, J. Boudreau¹³⁸, E.V. Bouhova-Thacker⁹⁰, D. Boumediene³⁸, A. Boveia¹²⁷, J. Boyd³⁶, D. Boye^{33c}, I.R. Boyko⁸⁰, A.J. Bozson⁹⁴, J. Bracnik²¹, N. Brahim^{60d,60c}, G. Brandt¹⁸², O. Brandt³², F. Braren⁴⁶, B. Brau¹⁰³, J.E. Brau¹³¹, W.D. Breaden Madden⁵⁷, K. Brendlinger⁴⁶, R. Brenner¹⁶⁰, L. Brenner³⁶, R. Brenner¹⁷², S. Bressler¹⁸⁰, B. Brickwedde¹⁰⁰, D.L. Briglin²¹, D. Britton⁵⁷, D. Britzger¹¹⁵, I. Brock²⁴, R. Brock¹⁰⁷, G. Brooijmans³⁹, W.K. Brooks^{146d}, E. Brost²⁹, P.A. Bruckman de Renstrom⁸⁵, B. Brüers⁴⁶, D. Bruncko^{28b}, A. Bruni^{23b}, G. Bruni^{23b}, M. Bruschi^{23b}, N. Brusino^{73a,73b}, L. Bryngemark¹⁵³, T. Buanes¹⁷, Q. Buat¹⁵⁵, P. Buchholz¹⁵¹, A.G. Buckley⁵⁷, I.A. Budagov⁸⁰, M.K. Bugge¹³³, F. Bühner⁵², O. Bulekov¹¹², B.A. Bullard⁵⁹, T.J. Burch¹²¹, S. Burdin⁹¹, C.D. Burgard¹²⁰, A.M. Burger¹²⁹, B. Burghgrave⁸, J.T.P. Burr⁴⁶, C.D. Burton¹¹, J.C. Burzynski¹⁰³, V. Büscher¹⁰⁰, E. Buschmann⁵³, P.J. Bussey⁵⁷, J.M. Butler²⁵, C.M. Buttar⁵⁷, J.M. Butterworth⁹⁵, P. Butti³⁶, W. Buttinger³⁶, C.J. Buxo Vazquez¹⁰⁷, A. Buzatu¹⁵⁸, A.R. Buzykaev^{122b,122a}, G. Cabras^{23b,23a}, S. Cabrera Urbán¹⁷⁴, D. Caforio⁵⁶, H. Cai¹³⁸, V.M.M. Cairo¹⁵³, O. Cakir^{4a}, N. Calace³⁶, P. Calafiura¹⁸, G. Calderini¹³⁵, P. Calfayan⁶⁶, G. Callea⁵⁷, L.P. Caloba^{81b}, A. Caltabiano^{74a,74b}, S. Calvente Lopez⁹⁹, D. Calvet³⁸, S. Calvet³⁸, T.P. Calvet¹⁰², M. Calvetti^{72a,72b}, R. Camacho Toro¹³⁵, S. Camarda³⁶, D. Camarero Munoz⁹⁹, P. Camarri^{74a,74b}, M.T. Camerlingo^{75a,75b}, D. Cameron¹³³, C. Camincher³⁶, S. Campana³⁶, M. Campanelli⁹⁵, A. Camplani⁴⁰, V. Canale^{70a,70b}, A. Canesse¹⁰⁴, M. Cano Bret⁷⁸, J. Cantero¹²⁹, T. Cao¹⁶¹, Y. Cao¹⁷³, M.D.M. Capeans Garrido³⁶, M. Capua^{41b,41a}, R. Cardarelli^{74a}, F. Cardillo¹⁴⁹, G. Carducci^{41b,41a}, I. Carli¹⁴², T. Carli³⁶, G. Carlino^{70a}, B.T. Carlson¹³⁸, E.M. Carlson^{176,168a}, L. Carminati^{69a,69b}, R.M.D. Carney¹⁵³, S. Caron¹¹⁹, E. Carquin^{146d}, S. Carrá⁴⁶, G. Carratta^{23b,23a}, J.W.S. Carter¹⁶⁷, T.M. Carter⁵⁰, M.P. Casado^{14,f}, A.F. Casha¹⁶⁷, E.G. Castiglia¹⁸³, F.L. Castillo¹⁷⁴, L. Castillo Garcia¹⁴, V. Castillo Gimenez¹⁷⁴, N.F. Castro^{139a,139e}, A. Catinaccio³⁶, J.R. Catmore¹³³, A. Cattai³⁶, V. Cavaliere²⁹, V. Cavaioni^{72a,72b}, E. Celebi^{12b}, F. Celli¹³⁴, K. Cerny¹³⁰, A.S. Cerqueira^{81a}, A. Cerri¹⁵⁶, L. Cerrito^{74a,74b}, F. Cerutti¹⁸, A. Cervelli^{23b,23a}, S.A. Cetin^{12b}, Z. Chadi^{35a}, D. Chakraborty¹²¹, J. Chan¹⁸¹, W.S. Chan¹²⁰, W.Y. Chan⁹¹, J.D. Chapman³², B. Chargeishvili^{159b}, D.G. Charlton²¹, T.P. Charman⁹³, M. Chatterjee²⁰, C.C. Chau³⁴, S. Che¹²⁷, S. Chekanov⁶, S.V. Chekulaev^{168a}, G.A. Chelkov^{80,ah}, B. Chen⁷⁹, C. Chen^{60a}, C.H. Chen⁷⁹, H. Chen^{15c}, H. Chen²⁹, J. Chen^{60a}, J. Chen³⁹, J. Chen²⁶, S. Chen¹³⁶, S.J. Chen^{15c}, X. Chen^{15b}, Y. Chen^{60a}, Y-H. Chen⁴⁶, H.C. Cheng^{63a}, H.J. Cheng^{15a}, A. Cheplakov⁸⁰, E. Cheremushkina¹²³, R. Cherkaoui El Moursli^{35f}, E. Cheu⁷, K. Cheung⁶⁴, T.J.A. Chevalérias¹⁴⁴, L. Chevalier¹⁴⁴, V. Chiarella⁵¹, G. Chiarelli^{72a}, G. Chiodini^{68a}, A.S. Chisholm²¹, A. Chitan^{27b}, I. Chiu¹⁶³, Y.H. Chiu¹⁷⁶, M.V. Chizhov⁸⁰, K. Choi¹¹, A.R. Chomont^{73a,73b}, Y.S. Chow¹²⁰, L.D. Christopher^{33e}, M.C. Chu^{63a}, X. Chu^{15a,15d}, J. Chudoba¹⁴⁰, J.J. Chwastowski⁸⁵, L. Chytka¹³⁰, D. Cieri¹¹⁵, K.M. Ciesla⁸⁵, V. Cindro⁹², I.A. Cioară^{27b}, A. Ciocio¹⁸, F. Ciroto^{70a,70b}, Z.H. Citron^{180,j}, M. Citterio^{69a}, D.A. Ciubotaru^{27b}, B.M. Ciungu¹⁶⁷, A. Clark⁵⁴, M.R. Clark³⁹, P.J. Clark⁵⁰, S.E. Clawson¹⁰¹, C. Clement^{45a,45b}, Y. Coadou¹⁰², M. Cokal^{67a,67c}, A. Coccaro^{55b}, J. Cochran⁷⁹, R. Coelho Lopes De Sa¹⁰³, H. Cohen¹⁶¹, A.E.C. Coimbra³⁶, B. Cole³⁹, A.P. Colijn¹²⁰, J. Collot⁵⁸, P. Conde Muiño^{139a,139h}, S.H. Connell^{33c}, I.A. Connelly⁵⁷, S. Constantinescu^{27b}, F. Conventi^{70a,an}, A.M. Cooper-Sarkar¹³⁴, F. Cormier¹⁷⁵, K.J.R. Cormier¹⁶⁷, L.D. Corpe⁹⁵, M. Corradi^{73a,73b}, E.E. Corrigan⁹⁷, F. Corriveau^{104,ab}, M.J. Costa¹⁷⁴, F. Costanza⁵, D. Costanzo¹⁴⁹, G. Cowan⁹⁴, J.W. Cowley³², J. Crane¹⁰¹, K. Cranmer¹²⁵, R.A. Creager¹³⁶, S. Crépe-Renaudin⁵⁸, F. Crescioli¹³⁵, M. Cristinziani²⁴, V. Croft¹⁷⁰, G. Crosetti^{41b,41a}, A. Cueto⁵, T. Cuhadar Donszelmann¹⁷¹, H. Cui^{15a,15d}, A.R. Cukierman¹⁵³, W.R. Cunningham⁵⁷, S. Czekerda⁸⁵, P. Czodrowski³⁶,

M.M. Czurylo^{61b}, M.J. Da Cunha Sargedas De Sousa^{60b}, J.V. Da Fonseca Pinto^{81b}, C. Da Via¹⁰¹,
W. Dabrowski^{84a}, F. Dachs³⁶, T. Dado⁴⁷, S. Dahbi^{33e}, T. Dai¹⁰⁶, C. Dallapiccola¹⁰³, M. Dam⁴⁰,
G. D'amen²⁹, V. D'Amico^{75a,75b}, J. Damp¹⁰⁰, J.R. Dandoy¹³⁶, M.F. Daneri³⁰, M. Danninger¹⁵², V. Dao³⁶,
G. Darbo^{55b}, O. Dartsis⁵, A. Dattagupta¹³¹, T. Daubney⁴⁶, S. D'Auria^{69a,69b}, C. David^{168b}, T. Davidek¹⁴²,
D.R. Davis⁴⁹, I. Dawson¹⁴⁹, K. De⁸, R. De Asmundis^{70a}, M. De Beurs¹²⁰, S. De Castro^{23b,23a},
N. De Groot¹¹⁹, P. de Jong¹²⁰, H. De la Torre¹⁰⁷, A. De Maria^{15c}, D. De Pedis^{73a}, A. De Salvo^{73a},
U. De Sanctis^{74a,74b}, M. De Santis^{74a,74b}, A. De Santo¹⁵⁶, J.B. De Vivie De Regie⁶⁵, D.V. Dedovich⁸⁰,
A.M. Deiana⁴², J. Del Peso⁹⁹, Y. Delabat Diaz⁴⁶, D. Delgove⁶⁵, F. Deliot¹⁴⁴, C.M. Delitzsch⁷,
M. Della Pietra^{70a,70b}, D. Della Volpe⁵⁴, A. Dell'Acqua³⁶, L. Dell'Asta^{74a,74b}, M. Delmastro⁵,
C. Delporte⁶⁵, P.A. Delsart⁵⁸, D.A. DeMarco¹⁶⁷, S. Demers¹⁸³, M. Demichev⁸⁰, G. Demontigny¹¹⁰,
S.P. Denisov¹²³, L. D'Eramo¹²¹, D. Derendarz⁸⁵, J.E. Derkaoui^{35e}, F. Derue¹³⁵, P. Dervan⁹¹, K. Desch²⁴,
K. Dette¹⁶⁷, C. Deutsch²⁴, M.R. Devesa³⁰, P.O. Deviveiros³⁶, F.A. Di Bello^{73a,73b}, A. Di Ciaccio^{74a,74b},
L. Di Ciaccio⁵, W.K. Di Clemente¹³⁶, C. Di Donato^{70a,70b}, A. Di Girolamo³⁶, G. Di Gregorio^{72a,72b},
B. Di Micco^{75a,75b}, R. Di Nardo^{75a,75b}, K.F. Di Petrillo⁵⁹, R. Di Sipio¹⁶⁷, C. Diaconu¹⁰², F.A. Dias¹²⁰,
T. Dias Do Vale^{139a}, M.A. Diaz^{146a}, F.G. Diaz Capriles²⁴, J. Dickinson¹⁸, M. Didenko¹⁶⁶, E.B. Diehl¹⁰⁶,
J. Dietrich¹⁹, S. Díez Cornell⁴⁶, C. Diez Pardos¹⁵¹, A. Dimitrievska¹⁸, W. Ding^{15b}, J. Dingfelder²⁴,
S.J. Dittmeier^{61b}, F. Dittus³⁶, F. Djama¹⁰², T. Djobava^{159b}, J.I. Djuvslund¹⁷, M.A.B. Do Vale¹⁴⁷,
M. Dobre^{27b}, D. Dodsworth²⁶, C. Doglioni⁹⁷, J. Dolejsi¹⁴², Z. Dolezal¹⁴², M. Donadelli^{81c}, B. Dong^{60c},
J. Donini³⁸, A. D'onofrio^{15c}, M. D'Onofrio⁹¹, J. Dopke¹⁴³, A. Doria^{70a}, M.T. Dova⁸⁹, A.T. Doyle⁵⁷,
E. Drechsler¹⁵², E. Dreyer¹⁵², T. Dreyer⁵³, A.S. Drobac¹⁷⁰, D. Du^{60b}, T.A. du Pree¹²⁰, Y. Duan^{60d},
F. Dubinin¹¹¹, M. Dubovsky^{28a}, A. Dubreuil⁵⁴, E. Duchovni¹⁸⁰, G. Duckeck¹¹⁴, O.A. Ducu^{36,27b},
D. Duda¹¹⁵, A. Dudarev³⁶, A.C. Dudder¹⁰⁰, E.M. Duffield¹⁸, M. D'uffizi¹⁰¹, L. Dufflot⁶⁵, M. Dührssen³⁶,
C. Dülsen¹⁸², M. Dumancic¹⁸⁰, A.E. Dumitriu^{27b}, M. Dunford^{61a}, A. Duperrin¹⁰², H. Duran Yildiz^{4a},
M. Düren⁵⁶, A. Durglishvili^{159b}, D. Duschinger⁴⁸, B. Dutta⁴⁶, D. Duvnjak¹, G.I. Dyckes¹³⁶, M. Dyndal³⁶,
S. Dysch¹⁰¹, B.S. Dziedzic⁸⁵, M.G. Eggleston⁴⁹, T. Eifert⁸, G. Eigen¹⁷, K. Einsweiler¹⁸, T. Ekelof¹⁷²,
H. El Jarrari^{35f}, V. Ellajosyula¹⁷², M. Ellert¹⁷², F. Ellinghaus¹⁸², A.A. Elliot⁹³, N. Ellis³⁶, J. Elmsheuser²⁹,
M. Elsing³⁶, D. Emelianov¹⁴³, A. Emerman³⁹, Y. Enari¹⁶³, M.B. Epland⁴⁹, J. Erdmann⁴⁷, A. Ereditato²⁰,
P.A. Erland⁸⁵, M. Errenst¹⁸², M. Escalier⁶⁵, C. Escobar¹⁷⁴, O. Estrada Pastor¹⁷⁴, E. Etzion¹⁶¹,
G. Evans^{139a,139b}, H. Evans⁶⁶, M.O. Evans¹⁵⁶, A. Ezhilov¹³⁷, F. Fabbri⁵⁷, L. Fabbri^{23b,23a}, V. Fabiani¹¹⁹,
G. Facini¹⁷⁸, R.M. Fakhruddinov¹²³, S. Falciano^{73a}, P.J. Falke²⁴, S. Falke³⁶, J. Faltova¹⁴², Y. Fang^{15a},
Y. Fang^{15a}, G. Fanourakis⁴⁴, M. Fanti^{69a,69b}, M. Faraj^{67a,67c}, A. Farbin⁸, A. Farilla^{75a}, E.M. Farina^{71a,71b},
T. Farooque¹⁰⁷, S.M. Farrington⁵⁰, P. Farthouat³⁶, F. Fassi^{35f}, P. Fassnacht³⁶, D. Fassouliotis⁹,
M. Faucci Giannelli⁵⁰, W.J. Fawcett³², L. Fayard⁶⁵, O.L. Fedin^{137,o}, W. Fedorko¹⁷⁵, A. Fehr²⁰,
M. Feickert¹⁷³, L. Feligioni¹⁰², A. Fell¹⁴⁹, C. Feng^{60b}, M. Feng⁴⁹, M.J. Fenton¹⁷¹, A.B. Fenyuk¹²³,
S.W. Ferguson⁴³, J. Ferrando⁴⁶, A. Ferrante¹⁷³, A. Ferrari¹⁷², P. Ferrari¹²⁰, R. Ferrari^{71a},
D.E. Ferreira de Lima^{61b}, A. Ferrer¹⁷⁴, D. Ferrere⁵⁴, C. Ferretti¹⁰⁶, F. Fiedler¹⁰⁰, A. Filipčić⁹²,
F. Filthaut¹¹⁹, K.D. Finelli²⁵, M.C.N. Fiolhais^{139a,139c,a}, L. Fiorini¹⁷⁴, F. Fischer¹¹⁴, J. Fischer¹⁰⁰,
W.C. Fisher¹⁰⁷, T. Fitschen²¹, I. Fleck¹⁵¹, P. Fleischmann¹⁰⁶, T. Flick¹⁸², B.M. Flierl¹¹⁴, L. Flores¹³⁶,
L.R. Flores Castillo^{63a}, F.M. Follega^{76a,76b}, N. Fomin¹⁷, J.H. Foo¹⁶⁷, G.T. Forcolin^{76a,76b}, B.C. Forland⁶⁶,
A. Formica¹⁴⁴, F.A. Förster¹⁴, A.C. Forti¹⁰¹, E. Fortin¹⁰², M.G. Foti¹³⁴, D. Fournier⁶⁵, H. Fox⁹⁰,
P. Francavilla^{72a,72b}, S. Francescato^{73a,73b}, M. Franchini^{23b,23a}, S. Franchino^{61a}, D. Francis³⁶, L. Franco⁵,
L. Franconi²⁰, M. Franklin⁵⁹, G. Frattari^{73a,73b}, A.N. Fray⁹³, P.M. Freeman²¹, B. Freund¹¹⁰,
W.S. Freund^{81b}, E.M. Freundlich⁴⁷, D.C. Frizzell¹²⁸, D. Froidevaux³⁶, J.A. Frost¹³⁴, M. Fujimoto¹²⁶,
C. Fukunaga¹⁶⁴, E. Fullana Torregrosa¹⁷⁴, T. Fusayasu¹¹⁶, J. Fuster¹⁷⁴, A. Gabrielli^{23b,23a}, A. Gabrielli³⁶,
S. Gadatsch⁵⁴, P. Gadow¹¹⁵, G. Gagliardi^{55b,55a}, L.G. Gagnon¹¹⁰, G.E. Gallardo¹³⁴, E.J. Gallas¹³⁴,
B.J. Gallop¹⁴³, R. Gamboa Goni⁹³, K.K. Gan¹²⁷, S. Ganguly¹⁸⁰, J. Gao^{60a}, Y. Gao⁵⁰, Y.S. Gao^{31,1},
F.M. Garay Walls^{146a}, C. García¹⁷⁴, J.E. García Navarro¹⁷⁴, J.A. García Pascual^{15a}, C. Garcia-Argos⁵²,

M. Garcia-Sciveres¹⁸, R.W. Gardner³⁷, N. Garelli¹⁵³, S. Gargiulo⁵², C.A. Garner¹⁶⁷, V. Garonne¹³³,
S.J. Gasiorowski¹⁴⁸, P. Gaspar^{81b}, A. Gaudiello^{55b,55a}, G. Gaudio^{71a}, P. Gauzzi^{73a,73b}, I.L. Gavrilenko¹¹¹,
A. Gavriluk¹²⁴, C. Gay¹⁷⁵, G. Gaycken⁴⁶, E.N. Gazis¹⁰, A.A. Geanta^{27b}, C.M. Gee¹⁴⁵, C.N.P. Gee¹⁴³,
J. Geisen⁹⁷, M. Geisen¹⁰⁰, C. Gemme^{55b}, M.H. Genest⁵⁸, C. Geng¹⁰⁶, S. Gentile^{73a,73b}, S. George⁹⁴,
T. Geralis⁴⁴, L.O. Gerlach⁵³, P. Gessinger-Befurt¹⁰⁰, G. Gessner⁴⁷, S. Ghasemi¹⁵¹,
M. Ghasemi Bostanabad¹⁷⁶, M. Ghneimat¹⁵¹, A. Ghosh⁶⁵, A. Ghosh⁷⁸, B. Giacobbe^{23b}, S. Giagu^{73a,73b},
N. Giangiacomi^{23b,23a}, P. Giannetti^{72a}, A. Giannini^{70a,70b}, G. Giannini¹⁴, S.M. Gibson⁹⁴, M. Gignac¹⁴⁵,
D.T. Gil^{84b}, B.J. Gilbert³⁹, D. Gillberg³⁴, G. Gilles¹⁸², N.E.K. Gillwald⁴⁶, D.M. Gingrich^{3,am},
M.P. Giordani^{67a,67c}, P.F. Giraud¹⁴⁴, G. Giugliarelli^{67a,67c}, D. Giugni^{69a}, F. Giuli^{74a,74b}, S. Gkaitatzis¹⁶²,
I. Gkialas^{9,g}, E.L. Gkoukousis¹⁴, P. Gkountoumis¹⁰, L.K. Gladilin¹¹³, C. Glasman⁹⁹, J. Glatzer¹⁴,
P.C.F. Glaysher⁴⁶, A. Glazov⁴⁶, G.R. Gledhill¹³¹, I. Gnesi^{41b,b}, M. Goblirsch-Kolb²⁶, D. Godin¹¹⁰,
S. Goldfarb¹⁰⁵, T. Golling⁵⁴, D. Golubkov¹²³, A. Gomes^{139a,139b}, R. Goncalves Gama⁵³,
R. Gonçalves^{139a,139c}, G. Gonella¹³¹, L. Gonella²¹, A. Gongadze⁸⁰, F. Gonnella²¹, J.L. Gonski³⁹,
S. González de la Hoz¹⁷⁴, S. Gonzalez Fernandez¹⁴, R. Gonzalez Lopez⁹¹, C. Gonzalez Renteria¹⁸,
R. Gonzalez Suarez¹⁷², S. Gonzalez-Sevilla⁵⁴, G.R. Gonzalvo Rodriguez¹⁷⁴, L. Goossens³⁶,
N.A. Gorasia²¹, P.A. Gorbounov¹²⁴, H.A. Gordon²⁹, B. Gorini³⁶, E. Gorini^{68a,68b}, A. Gorišek⁹²,
A.T. Goshaw⁴⁹, M.I. Gostkin⁸⁰, C.A. Gottardo¹¹⁹, M. Gouighri^{35b}, A.G. Goussiou¹⁴⁸, N. Govender^{33c},
C. Goy⁵, I. Grabowska-Bold^{84a}, E.C. Graham⁹¹, J. Gramling¹⁷¹, E. Gramstad¹³³, S. Grancagnolo¹⁹,
M. Grandi¹⁵⁶, V. Gratchev¹³⁷, P.M. Gravila^{27f}, F.G. Gravili^{68a,68b}, C. Gray⁵⁷, H.M. Gray¹⁸, C. Grefe²⁴,
K. Gregersen⁹⁷, I.M. Gregor⁴⁶, P. Grenier¹⁵³, K. Grevtsov⁴⁶, C. Grieco¹⁴, N.A. Grieser¹²⁸, A.A. Grillo¹⁴⁵,
K. Grimm^{31,k}, S. Grinstein^{14,w}, J.-F. Grivaz⁶⁵, S. Groh¹⁰⁰, E. Gross¹⁸⁰, J. Grosse-Knetter⁵³, Z.J. Grout⁹⁵,
C. Grud¹⁰⁶, A. Grummer¹¹⁸, J.C. Grundy¹³⁴, L. Guan¹⁰⁶, W. Guan¹⁸¹, C. Gubbels¹⁷⁵, J. Guenther³⁶,
A. Guerguichon⁶⁵, J.G.R. Guerrero Rojas¹⁷⁴, F. Guescini¹¹⁵, D. Guest¹⁷¹, R. Gugel¹⁰⁰, A. Guida⁴⁶,
T. Guillemin⁵, S. Guindon³⁶, J. Guo^{60c}, W. Guo¹⁰⁶, Y. Guo^{60a}, Z. Guo¹⁰², R. Gupta⁴⁶, S. Gurbuz^{12c},
G. Gustavino¹²⁸, M. Guth⁵², P. Gutierrez¹²⁸, C. Gutschow⁹⁵, C. Guyot¹⁴⁴, C. Gwenlan¹³⁴,
C.B. Gwilliam⁹¹, E.S. Haaland¹³³, A. Haas¹²⁵, C. Haber¹⁸, H.K. Hadavand⁸, A. Hadeif^{60a}, M. Haleem¹⁷⁷,
J. Haley¹²⁹, J.J. Hall¹⁴⁹, G. Halladjian¹⁰⁷, G.D. Hallewell¹⁰², K. Hamano¹⁷⁶, H. Hamdaoui^{35f}, M. Hamer²⁴,
G.N. Hamity⁵⁰, K. Han^{60a,v}, L. Han^{15c}, L. Han^{60a}, S. Han¹⁸, Y.F. Han¹⁶⁷, K. Hanagaki^{82,t}, M. Hance¹⁴⁵,
D.M. Handl¹¹⁴, M.D. Hank³⁷, R. Hankache¹³⁵, E. Hansen⁹⁷, J.B. Hansen⁴⁰, J.D. Hansen⁴⁰,
M.C. Hansen²⁴, P.H. Hansen⁴⁰, E.C. Hanson¹⁰¹, K. Hara¹⁶⁹, T. Harenberg¹⁸², S. Harkusha¹⁰⁸,
P.F. Harrison¹⁷⁸, N.M. Hartman¹⁵³, N.M. Hartmann¹¹⁴, Y. Hasegawa¹⁵⁰, A. Hasib⁵⁰, S. Hassani¹⁴⁴,
S. Haug²⁰, R. Hauser¹⁰⁷, L.B. Havener³⁹, M. Havranek¹⁴¹, C.M. Hawkes²¹, R.J. Hawkings³⁶,
S. Hayashida¹¹⁷, D. Hayden¹⁰⁷, C. Hayes¹⁰⁶, R.L. Hayes¹⁷⁵, C.P. Hays¹³⁴, J.M. Hays⁹³, H.S. Hayward⁹¹,
S.J. Haywood¹⁴³, F. He^{60a}, Y. He¹⁶⁵, M.P. Heath⁵⁰, V. Hedberg⁹⁷, S. Heer²⁴, A.L. Heggelund¹³³,
C. Heidegger⁵², K.K. Heidegger⁵², W.D. Heidorn⁷⁹, J. Heilman³⁴, S. Heim⁴⁶, T. Heim¹⁸,
B. Heinemann^{46,ak}, J.G. Heinlein¹³⁶, J.J. Heinrich¹³¹, L. Heinrich³⁶, J. Hejbal¹⁴⁰, L. Helary⁴⁶, A. Held¹²⁵,
S. Hellesund¹³³, C.M. Helling¹⁴⁵, S. Hellman^{45a,45b}, C. Helsens³⁶, R.C.W. Henderson⁹⁰, Y. Heng¹⁸¹,
L. Henkelmann³², A.M. Henriques Correia³⁶, H. Herde²⁶, Y. Hernández Jiménez^{33e}, H. Herr¹⁰⁰,
M.G. Herrmann¹¹⁴, T. Herrmann⁴⁸, G. Herten⁵², R. Hertenberger¹¹⁴, L. Hervas³⁶, T.C. Herwig¹³⁶,
G.G. Hesketh⁹⁵, N.P. Hesse^{168a}, H. Hibi⁸³, S. Higashino⁸², E. Higón-Rodríguez¹⁷⁴, K. Hildebrand³⁷,
J.C. Hill³², K.K. Hill²⁹, K.H. Hiller⁴⁶, S.J. Hillier²¹, M. Hils⁴⁸, I. Hinchliffe¹⁸, F. Hinterkeuser²⁴,
M. Hirose¹³², S. Hirose¹⁶⁹, D. Hirschbuehl¹⁸², B. Hiti⁹², O. Hladik¹⁴⁰, J. Hobbs¹⁵⁵, N. Hod¹⁸⁰,
M.C. Hodgkinson¹⁴⁹, A. Hoecker³⁶, D. Hohn⁵², D. Hohov⁶⁵, T. Holm²⁴, T.R. Holmes³⁷, M. Holzbock¹¹⁵,
L.B.A.H. Hommels³², T.M. Hong¹³⁸, J.C. Honig⁵², A. Hönle¹¹⁵, B.H. Hooberman¹⁷³, W.H. Hopkins⁶,
Y. Horii¹¹⁷, P. Horn⁴⁸, L.A. Horyn³⁷, S. Hou¹⁵⁸, A. Houmada^{35a}, J. Howarth⁵⁷, J. Hoya⁸⁹,
M. Hrabovsky¹³⁰, J. Hrdinka⁷⁷, J. Hrivnac⁶⁵, A. Hrynevich¹⁰⁹, T. Hryn'ova⁵, P.J. Hsu⁶⁴, S.-C. Hsu¹⁴⁸,
Q. Hu²⁹, S. Hu^{60c}, Y.F. Hu^{15a,15d,ao}, D.P. Huang⁹⁵, X. Huang^{15c}, Y. Huang^{60a}, Y. Huang^{15a}, Z. Hubacek¹⁴¹,

F. Hubaut¹⁰², M. Huebner²⁴, F. Huegging²⁴, T.B. Huffman¹³⁴, M. Huhtinen³⁶, R. Hulsken⁵⁸,
 R.F.H. Hunter³⁴, P. Huo¹⁵⁵, N. Huseynov^{80,ac}, J. Huston¹⁰⁷, J. Huth⁵⁹, R. Hyneman¹⁵³, S. Hyrych^{28a},
 G. Iacobucci⁵⁴, G. Iakovidis²⁹, I. Ibragimov¹⁵¹, L. Iconomidou-Fayard⁶⁵, P. Iengo³⁶, R. Ignazzi⁴⁰,
 O. Igonkina^{120,y,*}, R. Iguchi¹⁶³, T. Iizawa⁵⁴, Y. Ikegami⁸², M. Ikeno⁸², N. Ilic^{119,167,ab}, F. Iltzsche⁴⁸,
 H. Imam^{35a}, G. Introzzi^{71a,71b}, M. Iodice^{75a}, K. Iordanidou^{168a}, V. Ippolito^{73a,73b}, M.F. Isacson¹⁷²,
 M. Ishino¹⁶³, W. Islam¹²⁹, C. Issever^{19,46}, S. Istin¹⁶⁰, J.M. Iturbe Ponce^{63a}, R. Iuppa^{76a,76b}, A. Ivina¹⁸⁰,
 J.M. Izen⁴³, V. Izzo^{70a}, P. Jacka¹⁴⁰, P. Jackson¹, R.M. Jacobs⁴⁶, B.P. Jaeger¹⁵², V. Jain², G. Jäkel¹⁸²,
 K.B. Jakobi¹⁰⁰, K. Jakobs⁵², T. Jakoubek¹⁸⁰, J. Jamieson⁵⁷, K.W. Janas^{84a}, R. Jansky⁵⁴, M. Janus⁵³,
 P.A. Janus^{84a}, G. Jarlskog⁹⁷, A.E. Jaspan⁹¹, N. Javadov^{80,ac}, T. Javûrek³⁶, M. Javurkova¹⁰³, F. Jeanneau¹⁴⁴,
 L. Jeanty¹³¹, J. Jejelava^{159a}, P. Jenni^{52,c}, N. Jeong⁴⁶, S. Jézéquel⁵, H. Ji¹⁸¹, J. Jia¹⁵⁵, Z. Jia^{15c}, H. Jiang⁷⁹,
 Y. Jiang^{60a}, Z. Jiang¹⁵³, S. Jiggins⁵², F.A. Jimenez Morales³⁸, J. Jimenez Pena¹¹⁵, S. Jin^{15c}, A. Jinaru^{27b},
 O. Jinnouchi¹⁶⁵, H. Jivan^{33e}, P. Johansson¹⁴⁹, K.A. Johns⁷, C.A. Johnson⁶⁶, E. Jones¹⁷⁸, R.W.L. Jones⁹⁰,
 S.D. Jones¹⁵⁶, T.J. Jones⁹¹, J. Jongmanns^{61a}, J. Jovicevic³⁶, X. Ju¹⁸, J.J. Junggeburth¹¹⁵,
 A. Juste Rozas^{14,w}, A. Kaczmarzka⁸⁵, M. Kado^{73a,73b}, H. Kagan¹²⁷, M. Kagan¹⁵³, A. Kahn³⁹, C. Kahra¹⁰⁰,
 T. Kaji¹⁷⁹, E. Kajomovitz¹⁶⁰, C.W. Kalderon²⁹, A. Kaluza¹⁰⁰, A. Kamenshchikov¹²³, M. Kaneda¹⁶³,
 N.J. Kang¹⁴⁵, S. Kang⁷⁹, Y. Kano¹¹⁷, J. Kanzaki⁸², L.S. Kaplan¹⁸¹, D. Kar^{33e}, K. Karava¹³⁴,
 M.J. Kareem^{168b}, I. Karkani¹⁶², S.N. Karpov⁸⁰, Z.M. Karpova⁸⁰, V. Kartvelishvili⁹⁰, A.N. Karyukhin¹²³,
 E. Kasimi¹⁶², A. Kastanas^{45a,45b}, C. Kato^{60d}, J. Katzy⁴⁶, K. Kawade¹⁵⁰, K. Kawagoe⁸⁸, T. Kawaguchi¹¹⁷,
 T. Kawamoto¹⁴⁴, G. Kawamura⁵³, E.F. Kay¹⁷⁶, S. Kazakos¹⁴, V.F. Kazanin^{122b,122a}, J.M. Keaveney^{33a},
 R. Keeler¹⁷⁶, J.S. Keller³⁴, E. Kellermann⁹⁷, D. Kelsey¹⁵⁶, J.J. Kempster²¹, J. Kendrick²¹, K.E. Kennedy³⁹,
 O. Kepka¹⁴⁰, S. Kersten¹⁸², B.P. Kerševan⁹², S. Ketabchi Haghighat¹⁶⁷, M. Khader¹⁷³, F. Khalil-Zada¹³,
 M. Khandoga¹⁴⁴, A. Khanov¹²⁹, A.G. Kharlamov^{122b,122a}, T. Kharlamova^{122b,122a}, E.E. Khoda¹⁷⁵,
 A. Khodinov¹⁶⁶, T.J. Khoo⁵⁴, G. Khorauli¹⁷⁷, E. Khramov⁸⁰, J. Khubua^{159b}, S. Kido⁸³, M. Kiehn³⁶,
 E. Kim¹⁶⁵, Y.K. Kim³⁷, N. Kimura⁹⁵, A. Kirchhoff⁵³, D. Kirchmeier⁴⁸, J. Kirk¹⁴³, A.E. Kiryunin¹¹⁵,
 T. Kishimoto¹⁶³, D.P. Kisliuk¹⁶⁷, V. Kitali⁴⁶, C. Kitsaki¹⁰, O. Kivernyk²⁴, T. Klapdor-Kleingrothaus⁵²,
 M. Klassen^{61a}, C. Klein³⁴, M.H. Klein¹⁰⁶, M. Klein⁹¹, U. Klein⁹¹, K. Kleinknecht¹⁰⁰, P. Klimek¹²¹,
 A. Klimentov²⁹, T. Klingl²⁴, T. Klioutchnikova³⁶, F.F. Klitzner¹¹⁴, P. Kluit¹²⁰, S. Kluth¹¹⁵, E. Kneringer⁷⁷,
 E.B.F.G. Knoop¹⁰², A. Knue⁵², D. Kobayashi⁸⁸, M. Kobel⁴⁸, M. Kocian¹⁵³, T. Kodama¹⁶³, P. Kodys¹⁴²,
 D.M. Koeck¹⁵⁶, P.T. Koenig²⁴, T. Koffas³⁴, N.M. Köhler³⁶, M. Kolb¹⁴⁴, I. Koletsou⁵, T. Komarek¹³⁰,
 T. Kondo⁸², K. Köneke⁵², A.X.Y. Kong¹, A.C. König¹¹⁹, T. Kono¹²⁶, V. Konstantinides⁹⁵,
 N. Konstantinidis⁹⁵, B. Konya⁹⁷, R. Kopeliansky⁶⁶, S. Koperny^{84a}, K. Korcyl⁸⁵, K. Kordas¹⁶²,
 G. Koren¹⁶¹, A. Korn⁹⁵, I. Korolkov¹⁴, E.V. Korolkova¹⁴⁹, N. Korotkova¹¹³, O. Kortner¹¹⁵, S. Kortner¹¹⁵,
 V.V. Kostyukhin^{149,166}, A. Kotsokchagia⁶⁵, A. Kotwal⁴⁹, A. Koulouris¹⁰,
 A. Kourkoumeli-Charalampidi^{71a,71b}, C. Kourkoumelis⁹, E. Kourlitis⁶, V. Kouskoura²⁹, R. Kowalewski¹⁷⁶,
 W. Kozanecki¹⁰¹, A.S. Kozhin¹²³, V.A. Kramarenko¹¹³, G. Kramberger⁹², D. Krasnopevtsev^{60a},
 M.W. Krasny¹³⁵, A. Krasznahorkay³⁶, D. Krauss¹¹⁵, J.A. Kremer¹⁰⁰, J. Kretzschmar⁹¹, P. Krieger¹⁶⁷,
 F. Krieter¹¹⁴, A. Krishnan^{61b}, M. Krivos¹⁴², K. Krizka¹⁸, K. Kroeninger⁴⁷, H. Kroha¹¹⁵, J. Kroll¹⁴⁰,
 J. Kroll¹³⁶, K.S. Krowpman¹⁰⁷, U. Kruchonak⁸⁰, H. Krüger²⁴, N. Krumnack⁷⁹, M.C. Kruse⁴⁹,
 J.A. Krzysiak⁸⁵, A. Kubota¹⁶⁵, O. Kuchinskaia¹⁶⁶, S. Kудay^{4b}, D. Kuechler⁴⁶, J.T. Kuechler⁴⁶, S. Kuehn³⁶,
 T. Kuhl⁴⁶, V. Kukhtin⁸⁰, Y. Kulchitsky^{108,af}, S. Kuleshov^{146b}, Y.P. Kulinich¹⁷³, M. Kuna⁵⁸, A. Kupco¹⁴⁰,
 T. Kupfer⁴⁷, O. Kuprash⁵², H. Kurashige⁸³, L.L. Kurchaninov^{168a}, Y.A. Kurochkin¹⁰⁸, A. Kurova¹¹²,
 M.G. Kurth^{15a,15d}, E.S. Kuwertz³⁶, M. Kuze¹⁶⁵, A.K. Kvam¹⁴⁸, J. Kvita¹³⁰, T. Kwan¹⁰⁴, F. La Ruffa^{41b,41a},
 C. Lacasta¹⁷⁴, F. Lacava^{73a,73b}, D.P.J. Lack¹⁰¹, H. Lacker¹⁹, D. Lacour¹³⁵, E. Ladygin⁸⁰, R. Lafaye⁵,
 B. Laforge¹³⁵, T. Lagouri^{146c}, S. Lai⁵³, I.K. Lakomic^{84a}, J.E. Lambert¹²⁸, S. Lammers⁶⁶, W. Lampl⁷,
 C. Lampoudis¹⁶², E. Lançon²⁹, U. Landgraf⁵², M.P.J. Landon⁹³, M.C. Lanfermann⁵⁴, V.S. Lang⁵²,
 J.C. Lange⁵³, R.J. Langenberg¹⁰³, A.J. Lankford¹⁷¹, F. Lanni²⁹, K. Lantzscht²⁴, A. Lanza^{71a},
 A. Lapertosa^{55b,55a}, J.F. Laporte¹⁴⁴, T. Lari^{69a}, F. Lasagni Manghi^{23b,23a}, M. Lassnig³⁶, V. Latonova¹⁴⁰,

T.S. Lau^{63a}, A. Laudrain¹⁰⁰, A. Laurier³⁴, M. Lavorgna^{70a,70b}, S.D. Lawlor⁹⁴, M. Lazzaroni^{69a,69b}, B. Le¹⁰¹, E. Le Guirriec¹⁰², A. Lebedev⁷⁹, M. LeBlanc⁷, T. LeCompte⁶, F. Ledroit-Guillon⁵⁸, A.C.A. Lee⁹⁵, C.A. Lee²⁹, G.R. Lee¹⁷, L. Lee⁵⁹, S.C. Lee¹⁵⁸, S. Lee⁷⁹, B. Lefebvre^{168a}, H.P. Lefebvre⁹⁴, M. Lefebvre¹⁷⁶, C. Leggett¹⁸, K. Lehmann¹⁵², N. Lehmann²⁰, G. Lehmann Miotto³⁶, W.A. Leight⁴⁶, A. Leisos^{162,u}, M.A.L. Leite^{81c}, C.E. Leitgeb¹¹⁴, R. Leitner¹⁴², D. Lellouch^{180,*}, K.J.C. Leney⁴², T. Lenz²⁴, S. Leone^{72a}, C. Leonidopoulos⁵⁰, A. Leopold¹³⁵, C. Leroy¹¹⁰, R. Les¹⁰⁷, C.G. Lester³², M. Levchenko¹³⁷, J. Levêque⁵, D. Levin¹⁰⁶, L.J. Levinson¹⁸⁰, D.J. Lewis²¹, B. Li^{15b}, B. Li¹⁰⁶, C-Q. Li^{60c,60d}, F. Li^{60c}, H. Li^{60a}, H. Li^{60b}, J. Li^{60c}, K. Li¹⁴⁸, L. Li^{60c}, M. Li^{15a,15d}, Q. Li^{15a,15d}, Q.Y. Li^{60a}, S. Li^{60d,60c}, X. Li⁴⁶, Y. Li⁴⁶, Z. Li^{60b}, Z. Li¹³⁴, Z. Li¹⁰⁴, Z. Liang^{15a}, M. Liberatore⁴⁶, B. Liberti^{74a}, A. Liblong¹⁶⁷, K. Lie^{63c}, S. Lim²⁹, C.Y. Lin³², K. Lin¹⁰⁷, R.A. Linck⁶⁶, R.E. Lindley⁷, J.H. Lindon²¹, A. Linss⁴⁶, A.L. Lioni⁵⁴, E. Lipeles¹³⁶, A. Lipniacka¹⁷, T.M. Liss^{173,al}, A. Lister¹⁷⁵, J.D. Little⁸, B. Liu⁷⁹, B.X. Liu¹⁵², H.B. Liu²⁹, J.B. Liu^{60a}, J.K.K. Liu³⁷, K. Liu^{60d,60c}, M. Liu^{60a}, M.Y. Liu^{60a}, P. Liu^{15a}, X. Liu^{60a}, Y. Liu⁴⁶, Y. Liu^{15a,15d}, Y.L. Liu¹⁰⁶, Y.W. Liu^{60a}, M. Livan^{71a,71b}, A. Lleres⁵⁸, J. Llorente Merino¹⁵², S.L. Lloyd⁹³, C.Y. Lo^{63b}, E.M. Lobodzinska⁴⁶, P. Loch⁷, S. Loffredo^{74a,74b}, T. Lohse¹⁹, K. Lohwasser¹⁴⁹, M. Lokajicek¹⁴⁰, J.D. Long¹⁷³, R.E. Long⁹⁰, I. Longarini^{73a,73b}, L. Longo³⁶, K.A. Looper¹²⁷, I. Lopez Paz¹⁰¹, A. Lopez Solis¹⁴⁹, J. Lorenz¹¹⁴, N. Lorenzo Martinez⁵, A.M. Lory¹¹⁴, P.J. Lösel¹¹⁴, A. Lösle⁵², X. Lou^{45a,45b}, X. Lou^{15a}, A. Lounis⁶⁵, J. Love⁶, P.A. Love⁹⁰, J.J. Lozano Bahilo¹⁷⁴, M. Lu^{60a}, Y.J. Lu⁶⁴, H.J. Lubatti¹⁴⁸, C. Luci^{73a,73b}, F.L. Lucio Alves^{15c}, A. Lucotte⁵⁸, F. Luehring⁶⁶, I. Luise¹³⁵, L. Luminari^{73a}, B. Lund-Jensen¹⁵⁴, M.S. Lutz¹⁶¹, D. Lynn²⁹, H. Lyons⁹¹, R. Lysak¹⁴⁰, E. Lytken⁹⁷, F. Lyu^{15a}, V. Lyubushkin⁸⁰, T. Lyubushkina⁸⁰, H. Ma²⁹, L.L. Ma^{60b}, Y. Ma⁹⁵, D.M. Mac Donell¹⁷⁶, G. Maccarrone⁵¹, A. Macchiolo¹¹⁵, C.M. Macdonald¹⁴⁹, J.C. MacDonald¹⁴⁹, J. Machado Miguens¹³⁶, D. Madaffari¹⁷⁴, R. Madar³⁸, W.F. Mader⁴⁸, M. Madugoda Ralalage Don¹²⁹, N. Madysa⁴⁸, J. Maeda⁸³, T. Maeno²⁹, M. Maerker⁴⁸, V. Magerl⁵², N. Magini⁷⁹, J. Magro^{67a,67c,q}, D.J. Mahon³⁹, C. Maidantchik^{81b}, T. Maier¹¹⁴, A. Maio^{139a,139b,139d}, K. Maj^{84a}, O. Majersky^{28a}, S. Majewski¹³¹, Y. Makida⁸², N. Makovec⁶⁵, B. Malaescu¹³⁵, Pa. Malecki⁸⁵, V.P. Maleev¹³⁷, F. Malek⁵⁸, D. Malito^{41b,41a}, U. Mallik⁷⁸, D. Malon⁶, C. Malone³², S. Maltezos¹⁰, S. Malyukov⁸⁰, J. Mamuzic¹⁷⁴, G. Mancini^{70a,70b}, I. Mandić⁹², L. Manhaes de Andrade Filho^{81a}, I.M. Maniatis¹⁶², J. Manjarres Ramos⁴⁸, K.H. Mankinen⁹⁷, A. Mann¹¹⁴, A. Manousos⁷⁷, B. Mansoulie¹⁴⁴, I. Manthos¹⁶², S. Manzoni¹²⁰, A. Marantis¹⁶², G. Marceca³⁰, L. Marchese¹³⁴, G. Marchiori¹³⁵, M. Marcisovsky¹⁴⁰, L. Maroccia^{74a,74b}, C. Marcon⁹⁷, M. Marjanovic¹²⁸, Z. Marshall¹⁸, M.U.F. Martensson¹⁷², S. Marti-Garcia¹⁷⁴, C.B. Martin¹²⁷, T.A. Martin¹⁷⁸, V.J. Martin⁵⁰, B. Martin dit Latour¹⁷, L. Martinelli^{75a,75b}, M. Martinez^{14,w}, P. Martinez Agullo¹⁷⁴, V.I. Martinez Outschoorn¹⁰³, S. Martin-Haugh¹⁴³, V.S. Martoiu^{27b}, A.C. Martyniuk⁹⁵, A. Marzin³⁶, S.R. Maschek¹¹⁵, L. Masetti¹⁰⁰, T. Mashimo¹⁶³, R. Mashinistov¹¹¹, J. Masik¹⁰¹, A.L. Maslennikov^{122b,122a}, L. Massa^{23b,23a}, P. Massarotti^{70a,70b}, P. Mastrandrea^{72a,72b}, A. Mastroberardino^{41b,41a}, T. Masubuchi¹⁶³, D. Matakias²⁹, A. Matic¹¹⁴, N. Matsuzawa¹⁶³, P. Mättig²⁴, J. Maurer^{27b}, B. Maček⁹², D.A. Maximov^{122b,122a}, R. Mazini¹⁵⁸, I. Maznas¹⁶², S.M. Mazza¹⁴⁵, J.P. Mc Gowan¹⁰⁴, S.P. Mc Kee¹⁰⁶, T.G. McCarthy¹¹⁵, W.P. McCormack¹⁸, E.F. McDonald¹⁰⁵, A.E. McDougall¹²⁰, J.A. Mcfayden¹⁸, G. Mchedlidze^{159b}, M.A. McKay⁴², K.D. McLean¹⁷⁶, S.J. McMahon¹⁴³, P.C. McNamara¹⁰⁵, C.J. McNicol¹⁷⁸, R.A. McPherson^{176,ab}, J.E. Mdhluli^{33e}, Z.A. Meadows¹⁰³, S. Meehan³⁶, T. Megy³⁸, S. Mehlhase¹¹⁴, A. Mehta⁹¹, B. Meirose⁴³, D. Melini¹⁶⁰, B.R. Mellado Garcia^{33e}, J.D. Mellenthin⁵³, M. Melo^{28a}, F. Meloni⁴⁶, A. Melzer²⁴, E.D. Mendes Gouveia^{139a,139e}, A.M. Mendes Jacques Da Costa²¹, L. Meng³⁶, X.T. Meng¹⁰⁶, S. Menke¹¹⁵, E. Meoni^{41b,41a}, S. Mergelmeyer¹⁹, S.A.M. Merkt¹³⁸, C. Merlassino¹³⁴, P. Mermoud⁵⁴, L. Merola^{70a,70b}, C. Meroni^{69a}, G. Merz¹⁰⁶, O. Meshkov^{113,111}, J.K.R. Meshreki¹⁵¹, J. Metcalfe⁶, A.S. Mete⁶, C. Meyer⁶⁶, J-P. Meyer¹⁴⁴, M. Michetti¹⁹, R.P. Middleton¹⁴³, L. Mijović⁵⁰, G. Mikenberg¹⁸⁰, M. Mikestikova¹⁴⁰, M. Mikuž⁹², H. Mildner¹⁴⁹, A. Milic¹⁶⁷, C.D. Milke⁴², D.W. Miller³⁷, A. Milov¹⁸⁰, D.A. Milstead^{45a,45b}, R.A. Mina¹⁵³, A.A. Minaenko¹²³, I.A. Minashvili^{159b}, A.I. Mincer¹²⁵, B. Mindur^{84a}, M. Mineev⁸⁰, Y. Minegishi¹⁶³, Y. Mino⁸⁶, L.M. Mir¹⁴, M. Mironova¹³⁴,

K.P. Mistry¹³⁶, T. Mitani¹⁷⁹, J. Mitrevski¹¹⁴, V.A. Mitsou¹⁷⁴, M. Mittal^{60c}, O. Miu¹⁶⁷, A. Miucci²⁰,
 P.S. Miyagawa⁹³, A. Mizukami⁸², J.U. Mjörnmark⁹⁷, T. Mkrtychyan^{61a}, M. Mlynarikova¹⁴², T. Moa^{45a,45b},
 S. Mobius⁵³, K. Mochizuki¹¹⁰, P. Mogg¹¹⁴, S. Mohapatra³⁹, R. Moles-Valls²⁴, K. Mönig⁴⁶, E. Monnier¹⁰²,
 A. Montalbano¹⁵², J. Montejo Berlingen³⁶, M. Montella⁹⁵, F. Monticelli⁸⁹, S. Monzani^{69a}, N. Morange⁶⁵,
 A.L. Moreira De Carvalho^{139a}, D. Moreno^{22a}, M. Moreno Llácer¹⁷⁴, C. Moreno Martinez¹⁴,
 P. Morettini^{55b}, M. Morgenstern¹⁶⁰, S. Morgenstern⁴⁸, D. Mori¹⁵², M. Morii⁵⁹, M. Morinaga¹⁷⁹,
 V. Morisbak¹³³, A.K. Morley³⁶, G. Mornacchi³⁶, A.P. Morris⁹⁵, L. Morvaj¹⁵⁵, P. Moschovakos³⁶,
 B. Moser¹²⁰, M. Mosidze^{159b}, T. Moskalets¹⁴⁴, P. Moskvitina¹¹⁹, J. Moss^{31,m}, E.J.W. Moyse¹⁰³,
 S. Muanza¹⁰², J. Mueller¹³⁸, R.S.P. Mueller¹¹⁴, D. Muenstermann⁹⁰, G.A. Mullier⁹⁷, D.P. Mungo^{69a,69b},
 J.L. Munoz Martinez¹⁴, F.J. Munoz Sanchez¹⁰¹, P. Murin^{28b}, W.J. Murray^{178,143}, A. Murrone^{69a,69b},
 J.M. Muse¹²⁸, M. Muškinja¹⁸, C. Mwewa^{33a}, A.G. Myagkov^{123,ah}, A.A. Myers¹³⁸, G. Myers⁶⁶, J. Myers¹³¹,
 M. Myska¹⁴¹, B.P. Nachman¹⁸, O. Nackenhorst⁴⁷, A.Nag Nag⁴⁸, K. Nagai¹³⁴, K. Nagano⁸², Y. Nagasaka⁶²,
 J.L. Nagle²⁹, E. Nagy¹⁰², A.M. Nairz³⁶, Y. Nakahama¹¹⁷, K. Nakamura⁸², T. Nakamura¹⁶³, H. Nanjo¹³²,
 F. Napolitano^{61a}, R.F. Naranjo Garcia⁴⁶, R. Narayan⁴², I. Naryshkin¹³⁷, M. Naseri³⁴, T. Naumann⁴⁶,
 G. Navarro^{22a}, P.Y. Nechaeva¹¹¹, F. Nechansky⁴⁶, T.J. Neep²¹, A. Negri^{71a,71b}, M. Negrini^{23b}, C. Nellist¹¹⁹,
 C. Nelson¹⁰⁴, M.E. Nelson^{45a,45b}, S. Nemecek¹⁴⁰, M. Nessi^{36,e}, M.S. Neubauer¹⁷³, F. Neuhaus¹⁰⁰,
 M. Neumann¹⁸², R. Newhouse¹⁷⁵, P.R. Newman²¹, C.W. Ng¹³⁸, Y.S. Ng¹⁹, Y.W.Y. Ng¹⁷¹, B. Ngair^{35f},
 H.D.N. Nguyen¹⁰², T. Nguyen Manh¹¹⁰, E. Nibigira³⁸, R.B. Nickerson¹³⁴, R. Nicolaidou¹⁴⁴,
 D.S. Nielsen⁴⁰, J. Nielsen¹⁴⁵, M. Niemeyer⁵³, N. Nikiforou¹¹, V. Nikolaenko^{123,ah}, I. Nikolic-Audit¹³⁵,
 K. Nikolopoulos²¹, P. Nilsson²⁹, H.R. Nindhito⁵⁴, A. Nisati^{73a}, N. Nishu^{60c}, R. Nisius¹¹⁵, I. Nitsche⁴⁷,
 T. Nitta¹⁷⁹, T. Nobe¹⁶³, D.L. Noel³², Y. Noguchi⁸⁶, I. Nomidis¹³⁵, M.A. Nomura²⁹, M. Nordberg³⁶,
 J. Novak⁹², T. Novak⁹², O. Novgorodova⁴⁸, R. Novotny¹⁴¹, L. Nozka¹³⁰, K. Ntekas¹⁷¹, E. Nurse⁹⁵,
 F.G. Oakham^{34,am}, H. Oberlack¹¹⁵, J. Ocariz¹³⁵, A. Ochi⁸³, I. Ochoa³⁹, J.P. Ochoa-Ricoux^{146a},
 K. O'Connor²⁶, S. Oda⁸⁸, S. Odaka⁸², S. Oerdek⁵³, A. Ogrodnik^{84a}, A. Oh¹⁰¹, C.C. Ohm¹⁵⁴, H. Oide¹⁶⁵,
 M.L. Ojeda¹⁶⁷, H. Okawa¹⁶⁹, Y. Okazaki⁸⁶, M.W. O'Keefe⁹¹, Y. Okumura¹⁶³, A. Olariu^{27b},
 L.F. Oleiro Seabra^{139a}, S.A. Olivares Pino^{146a}, D. Oliveira Damazio²⁹, J.L. Oliver¹, M.J.R. Olsson¹⁷¹,
 A. Olszewski⁸⁵, J. Olszowska⁸⁵, Ö.O. Öncel²⁴, D.C. O'Neil¹⁵², A.P. O'Neill¹³⁴, A. Onofre^{139a,139e},
 P.U.E. Onyisi¹¹, H. Oppen¹³³, R.G. Oreamuno Madriz¹²¹, M.J. Oreglia³⁷, G.E. Orellana⁸⁹,
 D. Orestano^{75a,75b}, N. Orlando¹⁴, R.S. Orr¹⁶⁷, V. O'Shea⁵⁷, R. Ospanov^{60a}, G. Otero y Garzon³⁰,
 H. Otono⁸⁸, P.S. Ott^{61a}, G.J. Ottino¹⁸, M. Ouchrif^{35c}, J. Ouellette²⁹, F. Ould-Saada¹³³, A. Ouraou^{144,*},
 Q. Ouyang^{15a}, M. Owen⁵⁷, R.E. Owen¹⁴³, V.E. Ozcan^{12c}, N. Ozturk⁸, J. Pacalt¹³⁰, H.A. Pacey³²,
 K. Pachal⁴⁹, A. Pacheco Pages¹⁴, C. Padilla Aranda¹⁴, S. Pagan Griso¹⁸, G. Palacino⁶⁶, S. Palazzo⁵⁰,
 S. Palestini³⁶, M. Palka^{84b}, P. Palni^{84a}, C.E. Pandini⁵⁴, J.G. Panduro Vazquez⁹⁴, P. Pani⁴⁶, G. Panizzo^{67a,67c},
 L. Paolozzi⁵⁴, C. Papadatos¹¹⁰, K. Papageorgiou^{9,g}, S. Parajuli⁴², A. Paramonov⁶, C. Paraskevopoulos¹⁰,
 D. Paredes Hernandez^{63b}, S.R. Paredes Saenz¹³⁴, B. Parida¹⁸⁰, T.H. Park¹⁶⁷, A.J. Parker³¹, M.A. Parker³²,
 F. Parodi^{55b,55a}, E.W. Parrish¹²¹, J.A. Parsons³⁹, U. Parzefall⁵², L. Pascual Dominguez¹³⁵, V.R. Pascuzzi¹⁸,
 J.M.P. Pasner¹⁴⁵, F. Pasquali¹²⁰, E. Pasqualucci^{73a}, S. Passaggio^{55b}, F. Pastore⁹⁴, P. Pasuwan^{45a,45b},
 S. Pataria¹⁰⁰, J.R. Pater¹⁰¹, A. Pathak^{181,i}, J. Patton⁹¹, T. Pauly³⁶, J. Parkes¹⁵³, B. Pearson¹¹⁵,
 M. Pedersen¹³³, L. Pedraza Diaz¹¹⁹, R. Pedro^{139a}, T. Peiffer⁵³, S.V. Peleganchuk^{122b,122a}, O. Penc¹⁴⁰,
 H. Peng^{60a}, B.S. Peralva^{81a}, M.M. Perego⁶⁵, A.P. Pereira Peixoto^{139a}, L. Pereira Sanchez^{45a,45b},
 D.V. Perepelitsa²⁹, E. Perez Codina^{168a}, F. Peri¹⁹, L. Perini^{69a,69b}, H. Pernegger³⁶, S. Perrella³⁶,
 A. Perrevoort¹²⁰, K. Peters⁴⁶, R.F.Y. Peters¹⁰¹, B.A. Petersen³⁶, T.C. Petersen⁴⁰, E. Petit¹⁰², V. Petousis¹⁴¹,
 A. Petridis¹, C. Petridou¹⁶², F. Petrucci^{75a,75b}, M. Pettee¹⁸³, N.E. Pettersson¹⁰³, K. Petukhova¹⁴²,
 A. Peyaud¹⁴⁴, R. Pezoa^{146d}, L. Pezzotti^{71a,71b}, T. Pham¹⁰⁵, P.W. Phillips¹⁴³, M.W. Phipps¹⁷³,
 G. Piacquadio¹⁵⁵, E. Pianori¹⁸, A. Picazio¹⁰³, R.H. Pickles¹⁰¹, R. Piegaia³⁰, D. Pietreanu^{27b}, J.E. Pilcher³⁷,
 A.D. Pilkington¹⁰¹, M. Pinamonti^{67a,67c}, J.L. Pinfeld³, C. Pitman Donaldson⁹⁵, M. Pitt¹⁶¹,
 L. Pizzimento^{74a,74b}, A. Pizzini¹²⁰, M.-A. Pleier²⁹, V. Plesanovs⁵², V. Pleskot¹⁴², E. Plotnikova⁸⁰,

P. Podberezko^{122b,122a}, R. Poettgen⁹⁷, R. Poggi⁵⁴, L. Poggioli¹³⁵, I. Pogrebnyak¹⁰⁷, D. Pohl²⁴, I. Pokharel⁵³,
 G. Polesello^{71a}, A. Poley^{152,168a}, A. Policicchio^{73a,73b}, R. Polifka¹⁴², A. Polini^{23b}, C.S. Pollard⁴⁶,
 V. Polychronakos²⁹, D. Ponomarenko¹¹², L. Pontecorvo³⁶, S. Popa^{27a}, G.A. Popeneciu^{27d}, L. Portales⁵,
 D.M. Portillo Quintero⁵⁸, S. Pospisil¹⁴¹, K. Potamianos⁴⁶, I.N. Potrap⁸⁰, C.J. Potter³², H. Potti¹¹,
 T. Poulsen⁹⁷, J. Poveda¹⁷⁴, T.D. Powell¹⁴⁹, G. Pownall⁴⁶, M.E. Pozo Astigarraga³⁶, A. Prades Ibanez¹⁷⁴,
 P. Pralavorio¹⁰², M.M. Prapa⁴⁴, S. Prell⁷⁹, D. Price¹⁰¹, M. Primavera^{68a}, M.L. Proffitt¹⁴⁸, N. Proklova¹¹²,
 K. Prokofiev^{63c}, F. Prokoshin⁸⁰, S. Protopopescu²⁹, J. Proudfoot⁶, M. Przybycien^{84a}, D. Pudzha¹³⁷,
 A. Puri¹⁷³, P. Puzo⁶⁵, D. Pyatiizbyantseva¹¹², J. Qian¹⁰⁶, Y. Qin¹⁰¹, A. Quadt⁵³, M. Queitsch-Maitland³⁶,
 M. Racko^{28a}, F. Ragusa^{69a,69b}, G. Rahal⁹⁸, J.A. Raine⁵⁴, S. Rajagopalan²⁹, A. Ramirez Morales⁹³,
 K. Ran^{15a,15d}, D.M. Rauch⁴⁶, F. Rauscher¹¹⁴, S. Rave¹⁰⁰, B. Ravina⁵⁷, I. Ravinovich¹⁸⁰, J.H. Rawling¹⁰¹,
 M. Raymond³⁶, A.L. Read¹³³, N.P. Readioff¹⁴⁹, M. Reale^{68a,68b}, D.M. Rebutzi^{71a,71b}, G. Redlinger²⁹,
 K. Reeves⁴³, D. Reikher¹⁶¹, A. Reiss¹⁰⁰, A. Rej¹⁵¹, C. Rembser³⁶, A. Renardi⁴⁶, M. Renda^{27b},
 M.B. Rendel¹¹⁵, A.G. Rennie⁵⁷, S. Resconi^{69a}, E.D. Resseguie¹⁸, S. Rettie⁹⁵, B. Reynolds¹²⁷,
 E. Reynolds²¹, O.L. Rezanova^{122b,122a}, P. Reznicek¹⁴², E. Ricci^{76a,76b}, R. Richter¹¹⁵, S. Richter⁴⁶,
 E. Richter-Was^{84b}, M. Ridel¹³⁵, P. Rieck¹¹⁵, O. Rifki⁴⁶, M. Rijssenbeek¹⁵⁵, A. Rimoldi^{71a,71b},
 M. Rimoldi⁴⁶, L. Rinaldi^{23b}, T.T. Rinn¹⁷³, G. Ripellino¹⁵⁴, I. Riu¹⁴, P. Rivadeneira⁴⁶,
 J.C. Rivera Vergara¹⁷⁶, F. Rizatdinova¹²⁹, E. Rizvi⁹³, C. Rizzi³⁶, S.H. Robertson^{104,ab}, M. Robin⁴⁶,
 D. Robinson³², C.M. Robles Gajardo^{146d}, M. Robles Manzano¹⁰⁰, A. Robson⁵⁷, A. Rocchi^{74a,74b},
 E. Rocco¹⁰⁰, C. Roda^{72a,72b}, S. Rodriguez Bosca¹⁷⁴, A. Rodriguez Rodriguez⁵², A.M. Rodríguez Vera^{168b},
 S. Roe³⁶, J. Roggel¹⁸², O. Røhne¹³³, R. Röhrig¹¹⁵, R.A. Rojas^{146d}, B. Roland⁵², C.P.A. Roland⁶⁶,
 J. Roloff²⁹, A. Romaniouk¹¹², M. Romano^{23b,23a}, N. Rompotis⁹¹, M. Ronzani¹²⁵, L. Roos¹³⁵, S. Rosati^{73a},
 G. Rosin¹⁰³, B.J. Rosser¹³⁶, E. Rossi⁴⁶, E. Rossi^{75a,75b}, E. Rossi^{70a,70b}, L.P. Rossi^{55b}, L. Rossini⁴⁶,
 R. Rosten¹⁴, M. Rotaru^{27b}, B. Rottler⁵², D. Rousseau⁶⁵, G. Rovelli^{71a,71b}, A. Roy¹¹, D. Roy^{33e},
 A. Rozanov¹⁰², Y. Rozen¹⁶⁰, X. Ruan^{33e}, T.A. Ruggeri¹, F. Rühr⁵², A. Ruiz-Martinez¹⁷⁴, A. Rummler³⁶,
 Z. Rurikova⁵², N.A. Rusakovich⁸⁰, H.L. Russell¹⁰⁴, L. Rustige^{38,47}, J.P. Rutherford⁷, E.M. Rüttinger¹⁴⁹,
 M. Rybar¹⁴², G. Rybkin⁶⁵, E.B. Rye¹³³, A. Ryzhov¹²³, J.A. Sabater Iglesias⁴⁶, P. Sabatini⁵³,
 L. Sabetta^{73a,73b}, S. Sacerdoti⁶⁵, H.F.W. Sadrozinski¹⁴⁵, R. Sadykov⁸⁰, F. Safai Tehrani^{73a},
 B. Safarzadeh Samani¹⁵⁶, M. Safdari¹⁵³, P. Saha¹²¹, S. Saha¹⁰⁴, M. Sahinsoy¹¹⁵, A. Sahu¹⁸²,
 M. Saimpert³⁶, M. Saito¹⁶³, T. Saito¹⁶³, H. Sakamoto¹⁶³, D. Salamani⁵⁴, G. Salamanna^{75a,75b},
 A. Salnikov¹⁵³, J. Salt¹⁷⁴, A. Salvador Salas¹⁴, D. Salvatore^{41b,41a}, F. Salvatore¹⁵⁶, A. Salvucci^{63a,63b,63c},
 A. Salzburger³⁶, J. Samarati³⁶, D. Sammel⁵², D. Sampsonidis¹⁶², D. Sampsonidou¹⁶², J. Sánchez¹⁷⁴,
 A. Sanchez Pineda^{67a,36,67c}, H. Sandaker¹³³, C.O. Sander⁴⁶, I.G. Sanderswood⁹⁰, M. Sandhoff¹⁸²,
 C. Sandoval^{22b}, D.P.C. Sankey¹⁴³, M. Sannino^{55b,55a}, Y. Sano¹¹⁷, A. Sansoni⁵¹, C. Santoni³⁸,
 H. Santos^{139a,139b}, S.N. Santpur¹⁸, A. Santra¹⁷⁴, K.A. Saoucha¹⁴⁹, A. Sapronov⁸⁰, J.G. Saraiva^{139a,139d},
 O. Sasaki⁸², K. Sato¹⁶⁹, F. Sauerburger⁵², E. Sauvan⁵, P. Savard^{167,am}, R. Sawada¹⁶³, C. Sawyer¹⁴³,
 L. Sawyer^{96,ag}, I. Sayago Galvan¹⁷⁴, C. Sbarra^{23b}, A. Sbrizzi^{67a,67c}, T. Scanlon⁹⁵, J. Schaarschmidt¹⁴⁸,
 P. Schacht¹¹⁵, D. Schaefer³⁷, L. Schaefer¹³⁶, S. Schaepe³⁶, U. Schäfer¹⁰⁰, A.C. Schaffer⁶⁵, D. Schaile¹¹⁴,
 R.D. Schamberger¹⁵⁵, E. Schanet¹¹⁴, C. Scharf¹⁹, N. Scharmberg¹⁰¹, V.A. Schegelsky¹³⁷, D. Scheirich¹⁴²,
 F. Schenck¹⁹, M. Schernau¹⁷¹, C. Schiavi^{55b,55a}, L.K. Schildgen²⁴, Z.M. Schillaci²⁶, E.J. Schioppa^{68a,68b},
 M. Schioppa^{41b,41a}, K.E. Schleicher⁵², S. Schlenker³⁶, K.R. Schmidt-Sommerfeld¹¹⁵, K. Schmieden³⁶,
 C. Schmitt¹⁰⁰, S. Schmitt⁴⁶, L. Schoeffel¹⁴⁴, A. Schoening^{61b}, P.G. Scholer⁵², E. Schopf¹³⁴, M. Schott¹⁰⁰,
 J.F.P. Schouwenberg¹¹⁹, J. Schovancova³⁶, S. Schramm⁵⁴, F. Schroeder¹⁸², A. Schulte¹⁰⁰,
 H-C. Schultz-Coulon^{61a}, M. Schumacher⁵², B.A. Schumm¹⁴⁵, Ph. Schune¹⁴⁴, A. Schwartzman¹⁵³,
 T.A. Schwarz¹⁰⁶, Ph. Schwemling¹⁴⁴, R. Schwienhorst¹⁰⁷, A. Sciandra¹⁴⁵, G. Sciolla²⁶,
 M. Scornajenghi^{41b,41a}, F. Scuri^{72a}, F. Scutti¹⁰⁵, L.M. Scyboz¹¹⁵, C.D. Sebastiani⁹¹, P. Seema¹⁹,
 S.C. Seidel¹¹⁸, A. Seiden¹⁴⁵, B.D. Seidlitz²⁹, T. Seiss³⁷, C. Seitz⁴⁶, J.M. Seixas^{81b}, G. Sekhniadze^{70a},
 S.J. Sekula⁴², N. Semprini-Cesari^{23b,23a}, S. Sen⁴⁹, C. Serfon²⁹, L. Serin⁶⁵, L. Serkin^{67a,67b}, M. Sessa^{60a},

H. Severini¹²⁸, S. Sevova¹⁵³, F. Sforza^{55b,55a}, A. Sfyrta⁵⁴, E. Shabalina⁵³, J.D. Shahinian¹⁴⁵, N.W. Shaikh^{45a,45b}, D. Shaked Renous¹⁸⁰, L.Y. Shan^{15a}, M. Shapiro¹⁸, A. Sharma¹³⁴, A.S. Sharma¹, P.B. Shatalov¹²⁴, K. Shaw¹⁵⁶, S.M. Shaw¹⁰¹, M. Shehade¹⁸⁰, Y. Shen¹²⁸, A.D. Sherman²⁵, P. Sherwood⁹⁵, L. Shi⁹⁵, C.O. Shimmin¹⁸³, Y. Shimogama¹⁷⁹, M. Shimojima¹¹⁶, J.D. Shinner⁹⁴, I.P.J. Shipsey¹³⁴, S. Shirabe¹⁶⁵, M. Shiyakova^{80,z}, J. Shlomi¹⁸⁰, A. Shmeleva¹¹¹, M.J. Shochet³⁷, J. Shojaii¹⁰⁵, D.R. Shope¹⁵⁴, S. Shrestha¹²⁷, E.M. Shrif^{33e}, M.J. Shroff¹⁷⁶, E. Shulga¹⁸⁰, P. Sicho¹⁴⁰, A.M. Sickles¹⁷³, E. Sideras Haddad^{33e}, O. Sidiropoulou³⁶, A. Sidoti^{23b,23a}, F. Siegert⁴⁸, Dj. Sijacki¹⁶, M.Jr. Silva¹⁸¹, M.V. Silva Oliveira³⁶, S.B. Silverstein^{45a}, S. Simion⁶⁵, R. Simoniello¹⁰⁰, C.J. Simpson-allsoy²¹, S. Simsek^{12b}, P. Sinervo¹⁶⁷, V. Sinetckii¹¹³, S. Singh¹⁵², M. Sioli^{23b,23a}, I. Siral¹³¹, S.Yu. Sivoklov¹¹³, J. Sjölin^{45a,45b}, A. Skaf⁵³, E. Skorda⁹⁷, P. Skubic¹²⁸, M. Slawinska⁸⁵, K. Sliwa¹⁷⁰, R. Slovak¹⁴², V. Smakhtin¹⁸⁰, B.H. Smart¹⁴³, J. Smiesko^{28b}, N. Smirnov¹¹², S.Yu. Smirnov¹¹², Y. Smirnov¹¹², L.N. Smirnova^{113,r}, O. Smirnova⁹⁷, E.A. Smith³⁷, H.A. Smith¹³⁴, M. Smizanska⁹⁰, K. Smolek¹⁴¹, A. Smykiewicz⁸⁵, A.A. Snesarev¹¹¹, H.L. Snoek¹²⁰, I.M. Snyder¹³¹, S. Snyder²⁹, R. Sobie^{176,ab}, A. Soffer¹⁶¹, A. Søggaard⁵⁰, F. Sohns⁵³, C.A. Solans Sanchez³⁶, E.Yu. Soldatov¹¹², U. Soldevila¹⁷⁴, A.A. Solodkov¹²³, A. Soloshenko⁸⁰, O.V. Solovyanov¹²³, V. Solovyev¹³⁷, P. Sommer¹⁴⁹, H. Son¹⁷⁰, A. Sonay¹⁴, W. Song¹⁴³, W.Y. Song^{168b}, A. Sopczak¹⁴¹, A.L. Soppio⁹⁵, F. Sopkova^{28b}, S. Sottocornola^{71a,71b}, R. Soualah^{67a,67c}, A.M. Soukharev^{122b,122a}, D. South⁴⁶, S. Spagnolo^{68a,68b}, M. Spalla¹¹⁵, M. Spangenberg¹⁷⁸, F. Spanò⁹⁴, D. Sperlich⁵², T.M. Spieker^{61a}, G. Spigo³⁶, M. Spina¹⁵⁶, D.P. Spiteri⁵⁷, M. Spousta¹⁴², A. Stabile^{69a,69b}, B.L. Stamas¹²¹, R. Stamen^{61a}, M. Stamenkovic¹²⁰, A. Stampekis²¹, E. Stanecka⁸⁵, B. Stanislaus¹³⁴, M.M. Stanitzki⁴⁶, M. Stankaityte¹³⁴, B. Stapf¹²⁰, E.A. Starchenko¹²³, G.H. Stark¹⁴⁵, J. Stark⁵⁸, P. Staroba¹⁴⁰, P. Starovoitov^{61a}, S. Stärz¹⁰⁴, R. Staszewski⁸⁵, G. Stavropoulos⁴⁴, M. Stegler⁴⁶, P. Steinberg²⁹, A.L. Steinhebel¹³¹, B. Stelzer^{152,168a}, H.J. Stelzer¹³⁸, O. Stelzer-Chilton^{168a}, H. Stenzel⁵⁶, T.J. Stevenson¹⁵⁶, G.A. Stewart³⁶, M.C. Stockton³⁶, G. Stoicea^{27b}, M. Stolarski^{139a}, S. Stonjek¹¹⁵, A. Straessner⁴⁸, J. Strandberg¹⁵⁴, S. Strandberg^{45a,45b}, M. Strauss¹²⁸, T. Strebler¹⁰², P. Strizenec^{28b}, R. Ströhmer¹⁷⁷, D.M. Strom¹³¹, R. Stroynowski⁴², A. Strubig⁵⁰, S.A. Stucci²⁹, B. Stugu¹⁷, J. Stupak¹²⁸, N.A. Styles⁴⁶, D. Su¹⁵³, W. Su^{60c,148}, X. Su^{60a}, V.V. Sulin¹¹¹, M.J. Sullivan⁹¹, D.M.S. Sultan⁵⁴, S. Sultansoy^{4c}, T. Sumida⁸⁶, P. Sun⁵⁷, S. Sun¹⁰⁶, X. Sun¹⁰¹, C.J.E. Suster¹⁵⁷, M.R. Sutton¹⁵⁶, S. Suzuki⁸², M. Svatos¹⁴⁰, M. Swiatlowski^{168a}, S.P. Swift², T. Swirski¹⁷⁷, A. Sydorenko¹⁰⁰, I. Sykora^{28a}, M. Sykora¹⁴², T. Sykora¹⁴², D. Ta¹⁰⁰, K. Tackmann^{46,x}, J. Taenzer¹⁶¹, A. Taffard¹⁷¹, R. Tafirout^{168a}, E. Tagiev¹²³, R. Takashima⁸⁷, K. Takeda⁸³, T. Takeshita¹⁵⁰, E.P. Takeva⁵⁰, Y. Takubo⁸², M. Talby¹⁰², A.A. Talyshev^{122b,122a}, K.C. Tam^{63b}, N.M. Tamir¹⁶¹, J. Tanaka¹⁶³, R. Tanaka⁶⁵, S. Tapia Araya¹⁷³, S. Tapprogge¹⁰⁰, A. Tarek Abouelfadl Mohamed¹⁰⁷, S. Tarem¹⁶⁰, K. Tariq^{60b}, G. Tarna^{27b,d}, G.F. Tartarelli^{69a}, P. Tas¹⁴², M. Tasevsky¹⁴⁰, E. Tassi^{41b,41a}, A. Tavares Delgado^{139a}, Y. Tayalati^{35f}, A.J. Taylor⁵⁰, G.N. Taylor¹⁰⁵, W. Taylor^{168b}, H. Teagle⁹¹, A.S. Tee⁹⁰, R. Teixeira De Lima¹⁵³, P. Teixeira-Dias⁹⁴, H. Ten Kate³⁶, J.J. Teoh¹²⁰, K. Terashi¹⁶³, J. Terron⁹⁹, S. Terzo¹⁴, M. Testa⁵¹, R.J. Teuscher^{167,ab}, S.J. Thais¹⁸³, N. Themistokleous⁵⁰, T. Thevenaux-Pelzer⁴⁶, F. Thiele⁴⁰, D.W. Thomas⁹⁴, J.O. Thomas⁴², J.P. Thomas²¹, E.A. Thompson⁴⁶, P.D. Thompson²¹, E. Thomson¹³⁶, E.J. Thorpe⁹³, R.E. Ticse Torres⁵³, V.O. Tikhomirov^{111,ai}, Yu.A. Tikhonov^{122b,122a}, S. Timoshenko¹¹², P. Tipton¹⁸³, S. Tisserant¹⁰², K. Todome^{23b,23a}, S. Todorova-Nova¹⁴², S. Todt⁴⁸, J. Tojo⁸⁸, S. Tokár^{28a}, K. Tokushuku⁸², E. Tolley¹²⁷, R. Tombs³², K.G. Tomiwa^{33e}, M. Tomoto^{82,117}, L. Tompkins¹⁵³, P. Tornambe¹⁰³, E. Torrence¹³¹, H. Torres⁴⁸, E. Torró Pastor¹⁷⁴, M. Toscani³⁰, C. Tosciri¹³⁴, J. Toth^{102,aa}, D.R. Tovey¹⁴⁹, A. Traeet¹⁷, C.J. Treado¹²⁵, T. Trefzger¹⁷⁷, F. Tresoldi¹⁵⁶, A. Tricoli²⁹, I.M. Trigger^{168a}, S. Trincaz-Duvoid¹³⁵, D.A. Trischuk¹⁷⁵, W. Trischuk¹⁶⁷, B. Trocme⁵⁸, A. Trofymov⁶⁵, C. Troncon^{69a}, F. Trovato¹⁵⁶, L. Truong^{33c}, M. Trzebinski⁸⁵, A. Trzupek⁸⁵, F. Tsai⁴⁶, J.C-L. Tseng¹³⁴, P.V. Tsiarshka^{108,af}, A. Tsirigotis^{162,u}, V. Tsiskaridze¹⁵⁵, E.G. Tskhadadze^{159a}, M. Tsopoulou¹⁶², I.I. Tsukerman¹²⁴, V. Tsulaia¹⁸, S. Tsuno⁸², D. Tsybychev¹⁵⁵, Y. Tu^{63b}, A. Tudorache^{27b}, V. Tudorache^{27b}, T.T. Tulbure^{27a}, A.N. Tuna⁵⁹, S. Turchikhin⁸⁰, D. Turgeman¹⁸⁰, I. Turk Cakir^{4b,s},

R.J. Turner²¹, R. Turra^{69a}, P.M. Tuts³⁹, S. Tzamarias¹⁶², E. Tzovara¹⁰⁰, K. Uchida¹⁶³, F. Ukegawa¹⁶⁹, G. Unal³⁶, M. Unal¹¹, A. Undrus²⁹, G. Unel¹⁷¹, F.C. Ungaro¹⁰⁵, Y. Unno⁸², K. Uno¹⁶³, J. Urban^{28b}, P. Urquijo¹⁰⁵, G. Usai⁸, Z. Uysal^{12d}, V. Vacek¹⁴¹, B. Vachon¹⁰⁴, K.O.H. Vadla¹³³, T. Vafeiadis³⁶, A. Vaidya⁹⁵, C. Valderanis¹¹⁴, E. Valdes Santurio^{45a,45b}, M. Valente⁵⁴, S. Valentinetti^{23b,23a}, A. Valero¹⁷⁴, L. Valéry⁴⁶, R.A. Vallance²¹, A. Vallier³⁶, J.A. Valls Ferrer¹⁷⁴, T.R. Van Daalen¹⁴, P. Van Gemmeren⁶, S. Van Stroud⁹⁵, I. Van Vulpen¹²⁰, M. Vanadia^{74a,74b}, W. Vandelli³⁶, M. Vandembroucke¹⁴⁴, E.R. Vandewall¹²⁹, A. Vaniachine¹⁶⁶, D. Vannicola^{73a,73b}, R. Vari^{73a}, E.W. Varnes⁷, C. Varni^{55b,55a}, T. Varol¹⁵⁸, D. Varouchas⁶⁵, K.E. Varvell¹⁵⁷, M.E. Vasile^{27b}, G.A. Vasquez¹⁷⁶, F. Vazeille³⁸, D. Vazquez Furelos¹⁴, T. Vazquez Schroeder³⁶, J. Veatch⁵³, V. Vecchio¹⁰¹, M.J. Veen¹²⁰, L.M. Veloce¹⁶⁷, F. Veloso^{139a,139c}, S. Veneziano^{73a}, A. Ventura^{68a,68b}, A. Verbytskyi¹¹⁵, V. Vercesi^{71a}, M. Verducci^{72a,72b}, C.M. Vergel Infante⁷⁹, C. Vergis²⁴, W. Verkerke¹²⁰, A.T. Vermeulen¹²⁰, J.C. Vermeulen¹²⁰, C. Vernieri¹⁵³, P.J. Verschuuren⁹⁴, M.C. Vetterli^{152,am}, N. Viaux Maira^{146d}, T. Vickey¹⁴⁹, O.E. Vickey Boeriu¹⁴⁹, G.H.A. Viehhauser¹³⁴, L. Vigani^{61b}, M. Villa^{23b,23a}, M. Villaplana Perez³, E.M. Villhauer⁵⁰, E. Vilucchi⁵¹, M.G. Vinciter³⁴, G.S. Virdee²¹, A. Vishwakarma⁵⁰, C. Vittori^{23b,23a}, I. Vivarelli¹⁵⁶, M. Vogel¹⁸², P. Vokac¹⁴¹, S.E. von Buddenbrock^{33e}, E. Von Toerne²⁴, V. Vorobel¹⁴², K. Vorobev¹¹², M. Vos¹⁷⁴, J.H. Vossebeld⁹¹, M. Vozak¹⁰¹, N. Vranjes¹⁶, M. Vranjes Milosavljevic¹⁶, V. Vrba¹⁴¹, M. Vreeswijk¹²⁰, N.K. Vu¹⁰², R. Vuillermet³⁶, I. Vukotic³⁷, S. Wada¹⁶⁹, P. Wagner²⁴, W. Wagner¹⁸², J. Wagner-Kuhr¹¹⁴, S. Wahdan¹⁸², H. Wahlberg⁸⁹, R. Wakasa¹⁶⁹, V.M. Walbrecht¹¹⁵, J. Walder¹⁴³, R. Walker¹¹⁴, S.D. Walker⁹⁴, W. Walkowiak¹⁵¹, V. Wallangen^{45a,45b}, A.M. Wang⁵⁹, A.Z. Wang¹⁸¹, C. Wang^{60a}, C. Wang^{60c}, F. Wang¹⁸¹, H. Wang¹⁸, H. Wang³, J. Wang^{63a}, P. Wang⁴², Q. Wang¹²⁸, R.-J. Wang¹⁰⁰, R. Wang^{60a}, R. Wang⁶, S.M. Wang¹⁵⁸, W.T. Wang^{60a}, W. Wang^{15c}, W.X. Wang^{60a}, Y. Wang^{60a}, Z. Wang¹⁰⁶, C. Wanotayaroj⁴⁶, A. Warburton¹⁰⁴, C.P. Ward³², R.J. Ward²¹, N. Warrack⁵⁷, A.T. Watson²¹, M.F. Watson²¹, G. Watts¹⁴⁸, B.M. Waugh⁹⁵, A.F. Webb¹¹, C. Weber²⁹, M.S. Weber²⁰, S.A. Weber³⁴, S.M. Weber^{61a}, A.R. Weidberg¹³⁴, J. Weingarten⁴⁷, M. Weirich¹⁰⁰, C. Weiser⁵², P.S. Wells³⁶, T. Wenaus²⁹, B. Wendland⁴⁷, T. Wengler³⁶, S. Wenig³⁶, N. Wermes²⁴, M. Wessels^{61a}, T.D. Weston²⁰, K. Whalen¹³¹, A.M. Wharton⁹⁰, A.S. White¹⁰⁶, A. White⁸, M.J. White¹, D. Whiteson¹⁷¹, B.W. Whitmore⁹⁰, W. Wiedenmann¹⁸¹, C. Wiel⁴⁸, M. Wielers¹⁴³, N. Wieseotte¹⁰⁰, C. Wiglesworth⁴⁰, L.A.M. Wiik-Fuchs⁵², H.G. Wilkens³⁶, L.J. Wilkins⁹⁴, H.H. Williams¹³⁶, S. Williams³², S. Willocq¹⁰³, P.J. Windischhofer¹³⁴, I. Wingerter-Seez⁵, E. Winkels¹⁵⁶, F. Winklmeier¹³¹, B.T. Winter⁵², M. Wittgen¹⁵³, M. Wobisch⁹⁶, A. Wolf¹⁰⁰, R. Wölker¹³⁴, J. Wollrath⁵², M.W. Wolter⁸⁵, H. Wolters^{139a,139c}, V.W.S. Wong¹⁷⁵, N.L. Woods¹⁴⁵, S.D. Worm⁴⁶, B.K. Wosiek⁸⁵, K.W. Woźniak⁸⁵, K. Wraight⁵⁷, S.L. Wu¹⁸¹, X. Wu⁵⁴, Y. Wu^{60a}, J. Wuerzinger¹³⁴, T.R. Wyatt¹⁰¹, B.M. Wynne⁵⁰, S. Xella⁴⁰, L. Xia¹⁷⁸, J. Xiang^{63c}, X. Xiao¹⁰⁶, X. Xie^{60a}, I. Xioidis¹⁵⁶, D. Xu^{15a}, H. Xu^{60a}, H. Xu^{60a}, L. Xu²⁹, T. Xu¹⁴⁴, W. Xu¹⁰⁶, Y. Xu^{15b}, Z. Xu^{60b}, Z. Xu¹⁵³, B. Yabsley¹⁵⁷, S. Yacoob^{33a}, D.P. Yallup⁹⁵, N. Yamaguchi⁸⁸, Y. Yamaguchi¹⁶⁵, A. Yamamoto⁸², M. Yamatani¹⁶³, T. Yamazaki¹⁶³, Y. Yamazaki⁸³, J. Yan^{60c}, Z. Yan²⁵, H.J. Yang^{60c,60d}, H.T. Yang¹⁸, S. Yang^{60a}, T. Yang^{63c}, X. Yang^{60b,58}, Y. Yang¹⁶³, Z. Yang^{60a}, W.-M. Yao¹⁸, Y.C. Yap⁴⁶, E. Yatsenko^{60c}, H. Ye^{15c}, J. Ye⁴², S. Ye²⁹, I. Yeletsikh⁸⁰, M.R. Yexley⁹⁰, E. Yigitbasi²⁵, P. Yin³⁹, K. Yorita¹⁷⁹, K. Yoshihara⁷⁹, C.J.S. Young³⁶, C. Young¹⁵³, J. Yu⁷⁹, R. Yuan^{60b,h}, X. Yue^{61a}, M. Zaazoua^{35f}, B. Zabinski⁸⁵, G. Zacharis¹⁰, E. Zaffaroni⁵⁴, J. Zahreddine¹³⁵, A.M. Zaitsev^{123,ah}, T. Zakareishvili^{159b}, N. Zakharchuk³⁴, S. Zambito³⁶, D. Zanzi³⁶, S.V. Zeiβner⁴⁷, C. Zeitnitz¹⁸², G. Zemaityte¹³⁴, J.C. Zeng¹⁷³, O. Zenin¹²³, T. Ženiš^{28a}, D. Zerwas⁶⁵, M. Zgubič¹³⁴, B. Zhang^{15c}, D.F. Zhang^{15b}, G. Zhang^{15b}, J. Zhang⁶, K. Zhang^{15a}, L. Zhang^{15c}, L. Zhang^{60a}, M. Zhang¹⁷³, R. Zhang¹⁸¹, S. Zhang¹⁰⁶, X. Zhang^{60c}, X. Zhang^{60b}, Y. Zhang^{15a,15d}, Z. Zhang^{63a}, Z. Zhang⁶⁵, P. Zhao⁴⁹, Z. Zhao^{60a}, A. Zhemchugov⁸⁰, Z. Zheng¹⁰⁶, D. Zhong¹⁷³, B. Zhou¹⁰⁶, C. Zhou¹⁸¹, H. Zhou⁷, M.S. Zhou^{15a,15d}, M. Zhou¹⁵⁵, N. Zhou^{60c}, Y. Zhou⁷, C.G. Zhu^{60b}, C. Zhu^{15a,15d}, H.L. Zhu^{60a}, H. Zhu^{15a}, J. Zhu¹⁰⁶, Y. Zhu^{60a}, X. Zhuang^{15a}, K. Zhukov¹¹¹, V. Zhulanov^{122b,122a}, D. Zieminska⁶⁶, N.I. Zimine⁸⁰, S. Zimmermann^{52,*}, Z. Zinonos¹¹⁵, M. Ziolkowski¹⁵¹, L. Živković¹⁶, G. Zobernig¹⁸¹, A. Zoccoli^{23b,23a}, K. Zoch⁵³, T.G. Zorbass¹⁴⁹, R. Zou³⁷, L. Zwalinski³⁶.

- ¹Department of Physics, University of Adelaide, Adelaide; Australia.
- ²Physics Department, SUNY Albany, Albany NY; United States of America.
- ³Department of Physics, University of Alberta, Edmonton AB; Canada.
- ⁴(^a)Department of Physics, Ankara University, Ankara; (^b)Istanbul Aydin University, Application and Research Center for Advanced Studies, Istanbul; (^c)Division of Physics, TOBB University of Economics and Technology, Ankara; Turkey.
- ⁵LAPP, Université Grenoble Alpes, Université Savoie Mont Blanc, CNRS/IN2P3, Annecy; France.
- ⁶High Energy Physics Division, Argonne National Laboratory, Argonne IL; United States of America.
- ⁷Department of Physics, University of Arizona, Tucson AZ; United States of America.
- ⁸Department of Physics, University of Texas at Arlington, Arlington TX; United States of America.
- ⁹Physics Department, National and Kapodistrian University of Athens, Athens; Greece.
- ¹⁰Physics Department, National Technical University of Athens, Zografou; Greece.
- ¹¹Department of Physics, University of Texas at Austin, Austin TX; United States of America.
- ¹²(^a)Bahcesehir University, Faculty of Engineering and Natural Sciences, Istanbul; (^b)Istanbul Bilgi University, Faculty of Engineering and Natural Sciences, Istanbul; (^c)Department of Physics, Bogazici University, Istanbul; (^d)Department of Physics Engineering, Gaziantep University, Gaziantep; Turkey.
- ¹³Institute of Physics, Azerbaijan Academy of Sciences, Baku; Azerbaijan.
- ¹⁴Institut de Física d'Altes Energies (IFAE), Barcelona Institute of Science and Technology, Barcelona; Spain.
- ¹⁵(^a)Institute of High Energy Physics, Chinese Academy of Sciences, Beijing; (^b)Physics Department, Tsinghua University, Beijing; (^c)Department of Physics, Nanjing University, Nanjing; (^d)University of Chinese Academy of Science (UCAS), Beijing; China.
- ¹⁶Institute of Physics, University of Belgrade, Belgrade; Serbia.
- ¹⁷Department for Physics and Technology, University of Bergen, Bergen; Norway.
- ¹⁸Physics Division, Lawrence Berkeley National Laboratory and University of California, Berkeley CA; United States of America.
- ¹⁹Institut für Physik, Humboldt Universität zu Berlin, Berlin; Germany.
- ²⁰Albert Einstein Center for Fundamental Physics and Laboratory for High Energy Physics, University of Bern, Bern; Switzerland.
- ²¹School of Physics and Astronomy, University of Birmingham, Birmingham; United Kingdom.
- ²²(^a)Facultad de Ciencias y Centro de Investigaciones, Universidad Antonio Nariño, Bogotá; (^b)Departamento de Física, Universidad Nacional de Colombia, Bogotá, Colombia; Colombia.
- ²³(^a)INFN Bologna and Università di Bologna, Dipartimento di Fisica; (^b)INFN Sezione di Bologna; Italy.
- ²⁴Physikalisches Institut, Universität Bonn, Bonn; Germany.
- ²⁵Department of Physics, Boston University, Boston MA; United States of America.
- ²⁶Department of Physics, Brandeis University, Waltham MA; United States of America.
- ²⁷(^a)Transilvania University of Brasov, Brasov; (^b)Horia Hulubei National Institute of Physics and Nuclear Engineering, Bucharest; (^c)Department of Physics, Alexandru Ioan Cuza University of Iasi, Iasi; (^d)National Institute for Research and Development of Isotopic and Molecular Technologies, Physics Department, Cluj-Napoca; (^e)University Politehnica Bucharest, Bucharest; (^f)West University in Timisoara, Timisoara; Romania.
- ²⁸(^a)Faculty of Mathematics, Physics and Informatics, Comenius University, Bratislava; (^b)Department of Subnuclear Physics, Institute of Experimental Physics of the Slovak Academy of Sciences, Kosice; Slovak Republic.
- ²⁹Physics Department, Brookhaven National Laboratory, Upton NY; United States of America.
- ³⁰Departamento de Física, Universidad de Buenos Aires, Buenos Aires; Argentina.
- ³¹California State University, CA; United States of America.

- ³²Cavendish Laboratory, University of Cambridge, Cambridge; United Kingdom.
- ³³(^a)Department of Physics, University of Cape Town, Cape Town; (^b)iThemba Labs, Western Cape; (^c)Department of Mechanical Engineering Science, University of Johannesburg, Johannesburg; (^d)University of South Africa, Department of Physics, Pretoria; (^e)School of Physics, University of the Witwatersrand, Johannesburg; South Africa.
- ³⁴Department of Physics, Carleton University, Ottawa ON; Canada.
- ³⁵(^a)Faculté des Sciences Ain Chock, Réseau Universitaire de Physique des Hautes Energies - Université Hassan II, Casablanca; (^b)Faculté des Sciences, Université Ibn-Tofail, Kénitra; (^c)Faculté des Sciences Semlalia, Université Cadi Ayyad, LPHEA-Marrakech; (^d)Moroccan Foundation for Advanced Science Innovation and Research (MAScIR), Rabat; (^e)Faculté des Sciences, Université Mohamed Premier and LPTPM, Oujda; (^f)Faculté des sciences, Université Mohammed V, Rabat; Morocco.
- ³⁶CERN, Geneva; Switzerland.
- ³⁷Enrico Fermi Institute, University of Chicago, Chicago IL; United States of America.
- ³⁸LPC, Université Clermont Auvergne, CNRS/IN2P3, Clermont-Ferrand; France.
- ³⁹Nevis Laboratory, Columbia University, Irvington NY; United States of America.
- ⁴⁰Niels Bohr Institute, University of Copenhagen, Copenhagen; Denmark.
- ⁴¹(^a)Dipartimento di Fisica, Università della Calabria, Rende; (^b)INFN Gruppo Collegato di Cosenza, Laboratori Nazionali di Frascati; Italy.
- ⁴²Physics Department, Southern Methodist University, Dallas TX; United States of America.
- ⁴³Physics Department, University of Texas at Dallas, Richardson TX; United States of America.
- ⁴⁴National Centre for Scientific Research "Demokritos", Agia Paraskevi; Greece.
- ⁴⁵(^a)Department of Physics, Stockholm University; (^b)Oskar Klein Centre, Stockholm; Sweden.
- ⁴⁶Deutsches Elektronen-Synchrotron DESY, Hamburg and Zeuthen; Germany.
- ⁴⁷Lehrstuhl für Experimentelle Physik IV, Technische Universität Dortmund, Dortmund; Germany.
- ⁴⁸Institut für Kern- und Teilchenphysik, Technische Universität Dresden, Dresden; Germany.
- ⁴⁹Department of Physics, Duke University, Durham NC; United States of America.
- ⁵⁰SUPA - School of Physics and Astronomy, University of Edinburgh, Edinburgh; United Kingdom.
- ⁵¹INFN e Laboratori Nazionali di Frascati, Frascati; Italy.
- ⁵²Physikalisches Institut, Albert-Ludwigs-Universität Freiburg, Freiburg; Germany.
- ⁵³II. Physikalisches Institut, Georg-August-Universität Göttingen, Göttingen; Germany.
- ⁵⁴Département de Physique Nucléaire et Corpusculaire, Université de Genève, Genève; Switzerland.
- ⁵⁵(^a)Dipartimento di Fisica, Università di Genova, Genova; (^b)INFN Sezione di Genova; Italy.
- ⁵⁶II. Physikalisches Institut, Justus-Liebig-Universität Giessen, Giessen; Germany.
- ⁵⁷SUPA - School of Physics and Astronomy, University of Glasgow, Glasgow; United Kingdom.
- ⁵⁸LPSC, Université Grenoble Alpes, CNRS/IN2P3, Grenoble INP, Grenoble; France.
- ⁵⁹Laboratory for Particle Physics and Cosmology, Harvard University, Cambridge MA; United States of America.
- ⁶⁰(^a)Department of Modern Physics and State Key Laboratory of Particle Detection and Electronics, University of Science and Technology of China, Hefei; (^b)Institute of Frontier and Interdisciplinary Science and Key Laboratory of Particle Physics and Particle Irradiation (MOE), Shandong University, Qingdao; (^c)School of Physics and Astronomy, Shanghai Jiao Tong University, KLPPAC-MoE, SKLPPC, Shanghai; (^d)Tsung-Dao Lee Institute, Shanghai; China.
- ⁶¹(^a)Kirchhoff-Institut für Physik, Ruprecht-Karls-Universität Heidelberg, Heidelberg; (^b)Physikalisches Institut, Ruprecht-Karls-Universität Heidelberg, Heidelberg; Germany.
- ⁶²Faculty of Applied Information Science, Hiroshima Institute of Technology, Hiroshima; Japan.
- ⁶³(^a)Department of Physics, Chinese University of Hong Kong, Shatin, N.T., Hong Kong; (^b)Department of Physics, University of Hong Kong, Hong Kong; (^c)Department of Physics and Institute for Advanced

- Study, Hong Kong University of Science and Technology, Clear Water Bay, Kowloon, Hong Kong; China.
- ⁶⁴Department of Physics, National Tsing Hua University, Hsinchu; Taiwan.
- ⁶⁵IJCLab, Université Paris-Saclay, CNRS/IN2P3, 91405, Orsay; France.
- ⁶⁶Department of Physics, Indiana University, Bloomington IN; United States of America.
- ⁶⁷(^a)INFN Gruppo Collegato di Udine, Sezione di Trieste, Udine; (^b)ICTP, Trieste; (^c)Dipartimento Politecnico di Ingegneria e Architettura, Università di Udine, Udine; Italy.
- ⁶⁸(^a)INFN Sezione di Lecce; (^b)Dipartimento di Matematica e Fisica, Università del Salento, Lecce; Italy.
- ⁶⁹(^a)INFN Sezione di Milano; (^b)Dipartimento di Fisica, Università di Milano, Milano; Italy.
- ⁷⁰(^a)INFN Sezione di Napoli; (^b)Dipartimento di Fisica, Università di Napoli, Napoli; Italy.
- ⁷¹(^a)INFN Sezione di Pavia; (^b)Dipartimento di Fisica, Università di Pavia, Pavia; Italy.
- ⁷²(^a)INFN Sezione di Pisa; (^b)Dipartimento di Fisica E. Fermi, Università di Pisa, Pisa; Italy.
- ⁷³(^a)INFN Sezione di Roma; (^b)Dipartimento di Fisica, Sapienza Università di Roma, Roma; Italy.
- ⁷⁴(^a)INFN Sezione di Roma Tor Vergata; (^b)Dipartimento di Fisica, Università di Roma Tor Vergata, Roma; Italy.
- ⁷⁵(^a)INFN Sezione di Roma Tre; (^b)Dipartimento di Matematica e Fisica, Università Roma Tre, Roma; Italy.
- ⁷⁶(^a)INFN-TIFPA; (^b)Università degli Studi di Trento, Trento; Italy.
- ⁷⁷Institut für Astro- und Teilchenphysik, Leopold-Franzens-Universität, Innsbruck; Austria.
- ⁷⁸University of Iowa, Iowa City IA; United States of America.
- ⁷⁹Department of Physics and Astronomy, Iowa State University, Ames IA; United States of America.
- ⁸⁰Joint Institute for Nuclear Research, Dubna; Russia.
- ⁸¹(^a)Departamento de Engenharia Elétrica, Universidade Federal de Juiz de Fora (UFJF), Juiz de Fora; (^b)Universidade Federal do Rio De Janeiro COPPE/EE/IF, Rio de Janeiro; (^c)Instituto de Física, Universidade de São Paulo, São Paulo; Brazil.
- ⁸²KEK, High Energy Accelerator Research Organization, Tsukuba; Japan.
- ⁸³Graduate School of Science, Kobe University, Kobe; Japan.
- ⁸⁴(^a)AGH University of Science and Technology, Faculty of Physics and Applied Computer Science, Krakow; (^b)Marian Smoluchowski Institute of Physics, Jagiellonian University, Krakow; Poland.
- ⁸⁵Institute of Nuclear Physics Polish Academy of Sciences, Krakow; Poland.
- ⁸⁶Faculty of Science, Kyoto University, Kyoto; Japan.
- ⁸⁷Kyoto University of Education, Kyoto; Japan.
- ⁸⁸Research Center for Advanced Particle Physics and Department of Physics, Kyushu University, Fukuoka ; Japan.
- ⁸⁹Instituto de Física La Plata, Universidad Nacional de La Plata and CONICET, La Plata; Argentina.
- ⁹⁰Physics Department, Lancaster University, Lancaster; United Kingdom.
- ⁹¹Oliver Lodge Laboratory, University of Liverpool, Liverpool; United Kingdom.
- ⁹²Department of Experimental Particle Physics, Jožef Stefan Institute and Department of Physics, University of Ljubljana, Ljubljana; Slovenia.
- ⁹³School of Physics and Astronomy, Queen Mary University of London, London; United Kingdom.
- ⁹⁴Department of Physics, Royal Holloway University of London, Egham; United Kingdom.
- ⁹⁵Department of Physics and Astronomy, University College London, London; United Kingdom.
- ⁹⁶Louisiana Tech University, Ruston LA; United States of America.
- ⁹⁷Fysiska institutionen, Lunds universitet, Lund; Sweden.
- ⁹⁸Centre de Calcul de l'Institut National de Physique Nucléaire et de Physique des Particules (IN2P3), Villeurbanne; France.
- ⁹⁹Departamento de Física Teórica C-15 and CIAFF, Universidad Autónoma de Madrid, Madrid; Spain.
- ¹⁰⁰Institut für Physik, Universität Mainz, Mainz; Germany.

- ¹⁰¹School of Physics and Astronomy, University of Manchester, Manchester; United Kingdom.
- ¹⁰²CPPM, Aix-Marseille Université, CNRS/IN2P3, Marseille; France.
- ¹⁰³Department of Physics, University of Massachusetts, Amherst MA; United States of America.
- ¹⁰⁴Department of Physics, McGill University, Montreal QC; Canada.
- ¹⁰⁵School of Physics, University of Melbourne, Victoria; Australia.
- ¹⁰⁶Department of Physics, University of Michigan, Ann Arbor MI; United States of America.
- ¹⁰⁷Department of Physics and Astronomy, Michigan State University, East Lansing MI; United States of America.
- ¹⁰⁸B.I. Stepanov Institute of Physics, National Academy of Sciences of Belarus, Minsk; Belarus.
- ¹⁰⁹Research Institute for Nuclear Problems of Byelorussian State University, Minsk; Belarus.
- ¹¹⁰Group of Particle Physics, University of Montreal, Montreal QC; Canada.
- ¹¹¹P.N. Lebedev Physical Institute of the Russian Academy of Sciences, Moscow; Russia.
- ¹¹²National Research Nuclear University MEPhI, Moscow; Russia.
- ¹¹³D.V. Skobeltsyn Institute of Nuclear Physics, M.V. Lomonosov Moscow State University, Moscow; Russia.
- ¹¹⁴Fakultät für Physik, Ludwig-Maximilians-Universität München, München; Germany.
- ¹¹⁵Max-Planck-Institut für Physik (Werner-Heisenberg-Institut), München; Germany.
- ¹¹⁶Nagasaki Institute of Applied Science, Nagasaki; Japan.
- ¹¹⁷Graduate School of Science and Kobayashi-Maskawa Institute, Nagoya University, Nagoya; Japan.
- ¹¹⁸Department of Physics and Astronomy, University of New Mexico, Albuquerque NM; United States of America.
- ¹¹⁹Institute for Mathematics, Astrophysics and Particle Physics, Radboud University/Nikhef, Nijmegen; Netherlands.
- ¹²⁰Nikhef National Institute for Subatomic Physics and University of Amsterdam, Amsterdam; Netherlands.
- ¹²¹Department of Physics, Northern Illinois University, DeKalb IL; United States of America.
- ¹²²^(a) Budker Institute of Nuclear Physics and NSU, SB RAS, Novosibirsk; ^(b) Novosibirsk State University Novosibirsk; Russia.
- ¹²³Institute for High Energy Physics of the National Research Centre Kurchatov Institute, Protvino; Russia.
- ¹²⁴Institute for Theoretical and Experimental Physics named by A.I. Alikhanov of National Research Centre "Kurchatov Institute", Moscow; Russia.
- ¹²⁵Department of Physics, New York University, New York NY; United States of America.
- ¹²⁶Ochanomizu University, Otsuka, Bunkyo-ku, Tokyo; Japan.
- ¹²⁷Ohio State University, Columbus OH; United States of America.
- ¹²⁸Homer L. Dodge Department of Physics and Astronomy, University of Oklahoma, Norman OK; United States of America.
- ¹²⁹Department of Physics, Oklahoma State University, Stillwater OK; United States of America.
- ¹³⁰Palacký University, RCPTM, Joint Laboratory of Optics, Olomouc; Czech Republic.
- ¹³¹Institute for Fundamental Science, University of Oregon, Eugene, OR; United States of America.
- ¹³²Graduate School of Science, Osaka University, Osaka; Japan.
- ¹³³Department of Physics, University of Oslo, Oslo; Norway.
- ¹³⁴Department of Physics, Oxford University, Oxford; United Kingdom.
- ¹³⁵LPNHE, Sorbonne Université, Université de Paris, CNRS/IN2P3, Paris; France.
- ¹³⁶Department of Physics, University of Pennsylvania, Philadelphia PA; United States of America.
- ¹³⁷Konstantinov Nuclear Physics Institute of National Research Centre "Kurchatov Institute", PNPI, St. Petersburg; Russia.
- ¹³⁸Department of Physics and Astronomy, University of Pittsburgh, Pittsburgh PA; United States of

America.

¹³⁹(*a*) Laboratório de Instrumentação e Física Experimental de Partículas - LIP, Lisboa; (*b*) Departamento de Física, Faculdade de Ciências, Universidade de Lisboa, Lisboa; (*c*) Departamento de Física, Universidade de Coimbra, Coimbra; (*d*) Centro de Física Nuclear da Universidade de Lisboa, Lisboa; (*e*) Departamento de Física, Universidade do Minho, Braga; (*f*) Departamento de Física Teórica y del Cosmos, Universidad de Granada, Granada (Spain); (*g*) Dep Física and CEFITEC of Faculdade de Ciências e Tecnologia, Universidade Nova de Lisboa, Caparica; (*h*) Instituto Superior Técnico, Universidade de Lisboa, Lisboa; Portugal.

¹⁴⁰Institute of Physics of the Czech Academy of Sciences, Prague; Czech Republic.

¹⁴¹Czech Technical University in Prague, Prague; Czech Republic.

¹⁴²Charles University, Faculty of Mathematics and Physics, Prague; Czech Republic.

¹⁴³Particle Physics Department, Rutherford Appleton Laboratory, Didcot; United Kingdom.

¹⁴⁴IRFU, CEA, Université Paris-Saclay, Gif-sur-Yvette; France.

¹⁴⁵Santa Cruz Institute for Particle Physics, University of California Santa Cruz, Santa Cruz CA; United States of America.

¹⁴⁶(*a*) Departamento de Física, Pontificia Universidad Católica de Chile, Santiago; (*b*) Universidad Andres Bello, Department of Physics, Santiago; (*c*) Instituto de Alta Investigación, Universidad de Tarapacá; (*d*) Departamento de Física, Universidad Técnica Federico Santa María, Valparaíso; Chile.

¹⁴⁷Universidade Federal de São João del Rei (UFSJ), São João del Rei; Brazil.

¹⁴⁸Department of Physics, University of Washington, Seattle WA; United States of America.

¹⁴⁹Department of Physics and Astronomy, University of Sheffield, Sheffield; United Kingdom.

¹⁵⁰Department of Physics, Shinshu University, Nagano; Japan.

¹⁵¹Department Physik, Universität Siegen, Siegen; Germany.

¹⁵²Department of Physics, Simon Fraser University, Burnaby BC; Canada.

¹⁵³SLAC National Accelerator Laboratory, Stanford CA; United States of America.

¹⁵⁴Physics Department, Royal Institute of Technology, Stockholm; Sweden.

¹⁵⁵Departments of Physics and Astronomy, Stony Brook University, Stony Brook NY; United States of America.

¹⁵⁶Department of Physics and Astronomy, University of Sussex, Brighton; United Kingdom.

¹⁵⁷School of Physics, University of Sydney, Sydney; Australia.

¹⁵⁸Institute of Physics, Academia Sinica, Taipei; Taiwan.

¹⁵⁹(*a*) E. Andronikashvili Institute of Physics, Iv. Javakhishvili Tbilisi State University, Tbilisi; (*b*) High Energy Physics Institute, Tbilisi State University, Tbilisi; Georgia.

¹⁶⁰Department of Physics, Technion, Israel Institute of Technology, Haifa; Israel.

¹⁶¹Raymond and Beverly Sackler School of Physics and Astronomy, Tel Aviv University, Tel Aviv; Israel.

¹⁶²Department of Physics, Aristotle University of Thessaloniki, Thessaloniki; Greece.

¹⁶³International Center for Elementary Particle Physics and Department of Physics, University of Tokyo, Tokyo; Japan.

¹⁶⁴Graduate School of Science and Technology, Tokyo Metropolitan University, Tokyo; Japan.

¹⁶⁵Department of Physics, Tokyo Institute of Technology, Tokyo; Japan.

¹⁶⁶Tomsk State University, Tomsk; Russia.

¹⁶⁷Department of Physics, University of Toronto, Toronto ON; Canada.

¹⁶⁸(*a*) TRIUMF, Vancouver BC; (*b*) Department of Physics and Astronomy, York University, Toronto ON; Canada.

¹⁶⁹Division of Physics and Tomonaga Center for the History of the Universe, Faculty of Pure and Applied Sciences, University of Tsukuba, Tsukuba; Japan.

¹⁷⁰Department of Physics and Astronomy, Tufts University, Medford MA; United States of America.

- ¹⁷¹Department of Physics and Astronomy, University of California Irvine, Irvine CA; United States of America.
- ¹⁷²Department of Physics and Astronomy, University of Uppsala, Uppsala; Sweden.
- ¹⁷³Department of Physics, University of Illinois, Urbana IL; United States of America.
- ¹⁷⁴Instituto de Física Corpuscular (IFIC), Centro Mixto Universidad de Valencia - CSIC, Valencia; Spain.
- ¹⁷⁵Department of Physics, University of British Columbia, Vancouver BC; Canada.
- ¹⁷⁶Department of Physics and Astronomy, University of Victoria, Victoria BC; Canada.
- ¹⁷⁷Fakultät für Physik und Astronomie, Julius-Maximilians-Universität Würzburg, Würzburg; Germany.
- ¹⁷⁸Department of Physics, University of Warwick, Coventry; United Kingdom.
- ¹⁷⁹Waseda University, Tokyo; Japan.
- ¹⁸⁰Department of Particle Physics and Astrophysics, Weizmann Institute of Science, Rehovot; Israel.
- ¹⁸¹Department of Physics, University of Wisconsin, Madison WI; United States of America.
- ¹⁸²Fakultät für Mathematik und Naturwissenschaften, Fachgruppe Physik, Bergische Universität Wuppertal, Wuppertal; Germany.
- ¹⁸³Department of Physics, Yale University, New Haven CT; United States of America.
- ^a Also at Borough of Manhattan Community College, City University of New York, New York NY; United States of America.
- ^b Also at Centro Studi e Ricerche Enrico Fermi; Italy.
- ^c Also at CERN, Geneva; Switzerland.
- ^d Also at CPPM, Aix-Marseille Université, CNRS/IN2P3, Marseille; France.
- ^e Also at Département de Physique Nucléaire et Corpusculaire, Université de Genève, Genève; Switzerland.
- ^f Also at Departament de Física de la Universitat Autònoma de Barcelona, Barcelona; Spain.
- ^g Also at Department of Financial and Management Engineering, University of the Aegean, Chios; Greece.
- ^h Also at Department of Physics and Astronomy, Michigan State University, East Lansing MI; United States of America.
- ⁱ Also at Department of Physics and Astronomy, University of Louisville, Louisville, KY; United States of America.
- ^j Also at Department of Physics, Ben Gurion University of the Negev, Beer Sheva; Israel.
- ^k Also at Department of Physics, California State University, East Bay; United States of America.
- ^l Also at Department of Physics, California State University, Fresno; United States of America.
- ^m Also at Department of Physics, California State University, Sacramento; United States of America.
- ⁿ Also at Department of Physics, King's College London, London; United Kingdom.
- ^o Also at Department of Physics, St. Petersburg State Polytechnical University, St. Petersburg; Russia.
- ^p Also at Department of Physics, University of Fribourg, Fribourg; Switzerland.
- ^q Also at Dipartimento di Matematica, Informatica e Fisica, Università di Udine, Udine; Italy.
- ^r Also at Faculty of Physics, M.V. Lomonosov Moscow State University, Moscow; Russia.
- ^s Also at Giresun University, Faculty of Engineering, Giresun; Turkey.
- ^t Also at Graduate School of Science, Osaka University, Osaka; Japan.
- ^u Also at Hellenic Open University, Patras; Greece.
- ^v Also at IJCLab, Université Paris-Saclay, CNRS/IN2P3, 91405, Orsay; France.
- ^w Also at Institutio Catalana de Recerca i Estudis Avancats, ICREA, Barcelona; Spain.
- ^x Also at Institut für Experimentalphysik, Universität Hamburg, Hamburg; Germany.
- ^y Also at Institute for Mathematics, Astrophysics and Particle Physics, Radboud University/Nikhef, Nijmegen; Netherlands.
- ^z Also at Institute for Nuclear Research and Nuclear Energy (INRNE) of the Bulgarian Academy of Sciences, Sofia; Bulgaria.

- aa* Also at Institute for Particle and Nuclear Physics, Wigner Research Centre for Physics, Budapest; Hungary.
- ab* Also at Institute of Particle Physics (IPP); Canada.
- ac* Also at Institute of Physics, Azerbaijan Academy of Sciences, Baku; Azerbaijan.
- ad* Also at Instituto de Fisica Teorica, IFT-UAM/CSIC, Madrid; Spain.
- ae* Also at Istanbul University, Dept. of Physics, Istanbul; Turkey.
- af* Also at Joint Institute for Nuclear Research, Dubna; Russia.
- ag* Also at Louisiana Tech University, Ruston LA; United States of America.
- ah* Also at Moscow Institute of Physics and Technology State University, Dolgoprudny; Russia.
- ai* Also at National Research Nuclear University MEPhI, Moscow; Russia.
- aj* Also at Physics Department, An-Najah National University, Nablus; Palestine.
- ak* Also at Physikalisches Institut, Albert-Ludwigs-Universität Freiburg, Freiburg; Germany.
- al* Also at The City College of New York, New York NY; United States of America.
- am* Also at TRIUMF, Vancouver BC; Canada.
- an* Also at Università di Napoli Parthenope, Napoli; Italy.
- ao* Also at University of Chinese Academy of Sciences (UCAS), Beijing; China.
- * Deceased