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A Perspective on Interfacial Engineering of Lithium Metal Anodes and Beyond

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Abstract

This perspective reviews interfacial engineering of lithium metal anodes. Critical issues and open scientific questions related to coatings on the lithium metal anode are discussed. Essential features for ideal coatings, especially those that can potentially enable lithium plating underneath the coating, are highlighted. While most existing approaches use kinetic control to regulate coating thickness, here we offer a new perspective on thermodynamically-controlled interfacial engineering, focusing on spontaneously-formed 2D interfacial phases (also known as “complexions”). This approach has been applied to other battery systems but has yet to be realized for the lithium metal anodes.

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Introduction

The development of the lithium-ion battery (LIB) has played an indispensable role in the advancement of modern-day electronics. From the very beginning, lithium (Li) has been viewed as the ideal choice for the anode, as demonstrated with the Li-TiS₂ battery developed at Exxon in the 1970s¹. However, the formation of Li dendrites after repeated plating and stripping produced safety concerns that directed research toward alternative anode materials. Today, commercially available Li-ion batteries depend on intercalation chemistries where lithium ions are shuttled back and forth between a graphite anode and a layered transition metal oxide cathode such as LiCoO₂. While traditional “rocking chair” batteries demonstrate excellent cycle life and stability, next-generation Li-ion batteries must push beyond intercalation chemistries in order to attain higher energy densities. To do this, two factors are critical: 1) specific capacity, which refers to how much charge can be stored per mass, and 2) the working voltage window, which is defined by the potential difference between the cathode and anode redox pair. With these two primary considerations, Li metal has re-emerged as the anode of choice.

In order to address the issues surrounding the use of Li, it is important to revisit its past performance. Fig. 1(a) shows the main mechanisms for degradation and failure in Li metal anodes. Li deposition typically exhibits a dendritic morphology that has been shown to penetrate across the cell and reach the cathode, resulting in an internal short and subsequently thermal runaway of the battery²⁻⁶. In addition, the low reduction potential of lithium makes it inherently unstable in conventional electrolytes, resulting in the reduction of the electrolyte to form a passivation layer on the Li surface, commonly known as the “solid-electrolyte interphase (SEI)”⁷. Li₂O, Li₂CO₃, LiF, and Li alkylcarbonate are some of the important components that build up the SEI layer depending on the electrolyte systems^{8,9}. Defects and cracks in the SEI can result in localized Li-ion flux and trigger dendrite growth¹⁰. The formation of Li dendrites increases the surface area of the Li anode resulting in an irreversible consumption of active Li to form new SEI. In addition, the dendrites can also be pinched off from the bulk Li, which is another avenue in irreversible loss of active lithium. The consequence of these two coupled behaviors leads to poor Coulombic efficiency. Numerous strategies have been reported to mitigate these issues, as shown in Fig. 1(b), including electrolyte additives for create stable SEI^{11,12}, 3D lithiophilic host structures for guided Li plating^{13,14}, separator engineering for dendrite detection^{15,16} or anion regulation^{17,18}, and Li-electrolyte interfacial engineering^{19,20}.

This perspective details the use of alloy type surface modifications and coatings to 1) analyze the dendrite suppression mechanisms and 2) discuss insight and strategies for the interfacial stability of the Li metal anode. Fig. 2 shows various thermodynamically and kinetically controlled approaches for both surface modifications and coatings on Li metal anodes with increasing effective interfacial thickness. Hypotheses including mechanical suppression of Li dendrites, selective Li ion diffusion through the coating, and thermodynamic aspects of lithiophilic substrates are discussed below, along with the role of ionic conductivity and Li nucleation with interfacial coatings. Finally, a future outlook on thermodynamically-controlled interfacial engineering is discussed.

Li Plating on Top of the Coating

The idea of using Li-rich alloys for Li dendrite suppression was suggested by Huggins in the 1980s.^{21,22} Improved Li deposition morphology on Li-alloys was hypothesized to be the result of fast Li surface diffusion. Note that the term “Li-alloys” in the battery literature includes intermetallics or composites, while the term “alloys” usually refers only to solid solutions in physical metallurgy. A similar argument based on surface diffusion was proposed by Shi et al.^{23,24} from studies of Li-Mg and Li-Zn alloys fabricated through physical vapor deposition (thermal evaporation; Fig. 2(d)). More recently, studies from Archer’s group showed improved Li plating morphology and Coulombic efficiency using Li-In²⁵ and Li-Sn²⁶ coatings formed on the lithium surface through the spontaneous reduction of In(TFSI)₃ and SnTFSI (TFSI = bis(trifluoromethylsulfonyl)imide) salt solutions, respectively (Fig. 2(e)). Similar beneficial effects were also observed for Li-Si coating fabricated by sputtering Si on Li (Fig. 2(c))²⁷, and Li-Al coating made by laminating an Al foil onto Li (Fig. 2(g))²⁸.

Improvements in plated Li morphology are generally attributed to faster Li surface diffusion in the “alloyed” layer. Chemical diffusion coefficients of Li in the surface layers are obtained through galvanostatic intermittent titration technique²⁹, and the diffusion coefficient of Li in Li-rich alloy phases^{21,22,24} is generally more than two orders of magnitude higher than Li self-diffusion in bulk Li³⁰. Based on density functional theory (DFT) calculations, Archer et al. showed that Li ions exhibit fast surface transport on indium, which helps to explain the observed uniform Li deposition on the indium coating layer²⁵.

Moreover, Li nucleation shows zero overpotential on metal substrates where Li has at least

some solubility and is able to form an alloy, such as Au, Ag, and Mg³¹. This phenomenon can be explained from a nucleation point of view. Plating Li on top of a metal with negligible solid solubility would require heterogeneous nucleation, and that means the system has to overcome a positive nucleation barrier that results in an overpotential³². On the other hand, the room temperature solubility of Li in Au, Ag and Mg enables the formation of a surface solid solution layer upon lithiation at the interface without the need for heterogeneous nucleation, thereby eliminating the overpotential associated with the nucleation process. This surface layer also alters the morphology of the deposited Li³³⁻³⁵. Furthermore, studies have shown ‘dead Li’ formation depends more on the nucleation process, rather than the subsequent growth³⁶⁻³⁸. Therefore, it can be inferred that a substrate with preferred nucleation can decrease the dead Li formation and improve the cycling Coulombic efficiency.

Li Plating Underneath the Coating

Another, and perhaps better, approach is to enable the nucleation of Li underneath the coating. Nazar et al. reported uniform Li plating below a ~10 μm composite layer consisting of a lithium-rich intermetallic (Li-In, Li-As, Li-Bi, and Li-Zn) and metal chloride species³⁹. The proposed mechanism in this particular study focused on the role of the high Li-diffusivity in the intermetallic component of the coating, which is hypothesized to transport Li solely. It was further postulated that the presence of lithium chloride, a by-product of the reaction used to form the layer, serves to impart electronic resistivity to the composite layer to inhibit the reduction of lithium on the outer surface of the coating. Subsequent papers on Li-Ge coating from GeCl_4 -THF steam treatment⁴⁰ and Li-Sb coating from SbI_3 solution⁴¹ reported similar phenomena. For these systems, it is argued that fast Li diffusion in Li-rich compounds plays a critical role in enabling lithium transport through the layer instead of on the surface.

When lithium plating occurs underneath a coating layer, the performance improvements are hypothesized to be associated with the mechanical suppression of dendrite growth and selective Li-ion transport³⁹. Theoretical models have shown that a coating material under conformal compressive stress with a shear modulus that is higher than Li, can cause preferential deposition of Li at concave regions (rather than at peaks), thereby reducing the surface roughness of the deposited Li^{42,43}. Interestingly, these composite coatings are proposed to exhibit selective Li-ion transport across the coating while blocking the anions, thus preventing the occurrence of side

reactions between Li and the electrolyte³⁹. The fast Li transport properties in Li-rich alloys were also argued to provide a uniform Li flux across the layer, resulting in deposition underneath the coating³⁹. Table 1 summarizes the Li-ion diffusion coefficients for various Li-rich intermetallic compounds. Using Li₂₂Sn₅ as an example, room temperature Li diffusion coefficients were found to be in the range between 1.9×10^{-7} and 5.9×10^{-7} cm² s⁻¹ based on electrochemical measurements⁴⁴. This is significantly higher than the diffusion coefficient of Li in bulk Li metal (7.65×10^{-11} cm² s⁻¹)³⁰.

Critical Issues and Open Questions

Although many studies have reported improvements on plating morphology and cycling stability, Li-metal anodes are still far from commercialization. More studies are needed to understand cell failure mechanisms with these applied coatings. In the case of Li plating underneath the coating layer, the proposed working mechanisms (mechanical suppression and selective ion diffusion) assume that the coating can remain undamaged upon cycling. However, this assumption requires further experimental validation. Cell stacking pressure means that deformation can occur along with Li plating. Moreover, stripping can easily cause cracks to form on the coating. As shown in Fig. 3, dense, conformal, and micrometer-thick coatings consisting of Li salts or intermetallic compounds are brittle⁴⁵ and do not allow for the plastic deformation that results from the Li volume change on cycling. A similar argument has also been proposed for LiPON protection layer on Li metal anodes. LiPON, which is commonly used as a solid-state electrolyte, produces an ion-conducting and relatively brittle coating. The cycling stability of LiPON protected Li was enhanced at low plating/stripping capacity. Unfortunately, for high energy density cells, cracks were observed due to the substantial volume changes in the Li metal, and dendrites eventually formed^{46,47}. For interfacial coatings that favor Li plating underneath, there likely exists a current density threshold for effective dendrite suppression. Providing a unified testing protocol or computation model to identify and predict areal capacity/current density limitations exceeds the scope of this perspective. However, it is worth noting that studies are needed to identify this threshold and factors that determine it, with a goal of eventually increasing this threshold value.

Although the use of a lithiophilic substrate, such as Li-In²⁵ and Li-Sn³¹ interfacial coatings, has been shown to alter the initial nucleation process, their beneficial effects on the subsequent

growth and cycling are questionable. As shown in Fig. 3, after the nuclei cover a significant area on the substrate, the plated Li becomes the substrate for the incoming Li-ion, and the deposition morphology mainly depends on Li-ion mass transfer in the liquid electrolyte⁴⁸. Under high Li plating current densities, Li-ion diffusion in the electrolyte becomes the limiting process, and a concentration gradient builds up near the deposition surface. The resulting electric field causes fractal dendrites to grow³⁸. Moreover, since the plated Li is in direct contact with the liquid electrolyte, side reactions between Li and the electrolyte continue to occur and consume the electrolyte. Therefore, while lithiophilic substrates demonstrate improved morphology and growth of Li deposits initially, there is a limitation on the effectiveness of this approach once the surface of the coating is fully (or partially) covered by lithium.

To fully comprehend the mechanisms of coating and further improve it, the microstructure, ion conduction mechanisms, and Li nucleation behaviors all require further studies. Specifically, Li conduction through a micrometer-thick composite coating on Li metal consisting of a Li-rich intermetallic compound and Li salts, as shown in Fig. 2(f), is complex. In the case of a fully dense coating, the only ion pathway for the observed deposition of Li underneath the coating is through the intermetallic compound network. Since the Li-ion diffusion coefficient in the organic electrolyte is usually higher than the typical values in intermetallic compounds^{22,49}, the rate-limiting step is governed by the coating layer as Li-ions diffuse from the liquid electrolyte, across the coating layer, and nucleate underneath. There will be a kinetic limitation for the maximum ion flux across the coating, and therefore a current threshold above which it becomes favorable for Li to nucleate on the top of the coating. The inherent ionic diffusivity of the coating layer and the coating thickness are important parameters to tune this current threshold. Previous studies have also pointed out the desolvation of Li ions from the solvent molecule determines the activation energies of Li-ion transport across the liquid/solid interface^{50,51}. The solvation ability of the electrolyte solvent can be tuned with different Li salts. A recent study of Li-ion transport at the $\text{LiNi}_{0.33}\text{Mn}_{0.33}\text{Co}_{0.33}\text{O}_2$ /electrolyte interface reported that Lithium bis(trifluoromethane)sulfonimide (LiTFSI) based electrolyte can have ~100-fold increase in exchange current density at the solid-liquid electrolyte interface in comparison to the LiPF_6 based electrolyte⁵². This behavior was attributed to the lower desolvation energy of Li^+ with TFSI.

Moreover, composite films are almost never fully dense, e.g., the microstructures of the

composite coatings like those shown in Fig. 2(f) that enable underneath Li plating. Thus, more careful characterization is required. Porosity may exist in such a composite coating layer due to salt dissolution and inherent defects; if the pores are connected above a percolation threshold, Li conduction can occur through the percolated electrolyte within the coating. Here, characterization via cryogenic electron microscopy (cryoEM) can potentially provide more detailed information on the coating microstructure without severe beam damage⁵³. Evaluation of changes in the coating structure and in the location of Li nucleation with different current densities can be done using a variety of methods, including in-situ X-ray tomography⁵⁴ and ex-situ cryo-focused ion beam-scanning electron microscope (Cryo-FIB-SEM).⁵⁵ Direct imaging methods such as these can be helpful in understanding the coating degradation mechanism, and to provide insights for the future design of protection coatings.

Desirable Features of Coatings

Fig. 4 depicts several essential attributes of one ideal interfacial protection layer for the Li metal anode^{3,6}. A coating that allows for lithium deposition to take place underneath, as is likely achieved in the case shown in Fig. 2(f), is preferred for a dendrite-free anode. Unlike Li plating on top, the effectiveness of suppressing Li dendrites focuses less on the nucleation process and more on the effective transport and lithium flux across the coating layer. Having fast, selective lithium transport through the layer is a critical requirement. For example, prior studies already suggested that solid electrolyte type coatings can allow selective Li-ion diffusion while no anion species can diffuse across the layer^{50,51}. Li transport across the layer must be kinetically faster than the process of plating on top of the layer. Sluggish transport kinetics could result in the layer transport process becoming a rate-limiting step, leading to concentration gradients within the layer and added polarization. Additionally, having high electronic resistance in the layer is beneficial, as it provides a less favorable environment for lithium reduction on the surface of the coating layer, effectively reducing the kinetics for lithium plating on top of the layer and thus allowing Li transport through the films to better compete kinetically. In the case of Nazar et al.'s work³⁹, the composite nature of the layer with a chlorine species fulfilled this need, but potentially reduced the Li transport rate due to the large “inactive” component to the film. In addition to selective and fast Li-ion diffusion, the coating layer should be thermodynamically stable with Li metal. Specifically, solid electrolyte coatings have good Li-ion conductivity, but many of them are not stable in contact with Li metal. One such example is $\text{Li}_{1+x}\text{Al}_x\text{Ge}_{2-x}(\text{PO}_3)_3$

(LAGP), where Ge^{4+} can be reduced by Li metal⁵⁶. More discussion about solid electrolyte coatings and their stability can be found elsewhere^{57,58}.

It is also worth acknowledging that at higher currents, there exists a threshold where plating on top of the layer will be more kinetically favorable. Further studies are needed to design a layer where the threshold can be pushed further towards high power usage.

Despite its potential advantages, plating underneath a coating layer results in volume expansion and stress build-up from the plated lithium that must be accommodated. Therefore, the mechanical properties of the layer should also be considered as the layer material must be able to withstand the stress fluctuations generated by the plating and stripping process. To better accommodate the volume expansion, one strategy is to combine inorganic coating materials with various polymers to form a composite coating with increased elasticity^{59,60}. Alternatively, the interfacial coating can be combined with a 3D porous Li-host structure to alleviate the volume changes during Li plating and stripping²⁰.

Lastly, self-healing, where a coating can repair itself or rebuild during cycling, would be a highly desirable feature for any protective coating. Some of the self-healing functions may be achieved in thermodynamically-controlled coatings that form spontaneously, as discussed in the next section.

Thermodynamically-Controlled Interfacial Engineering: A New Perspective

As described in the preceding sections, current research on interfacial engineering of the Li metal anode is generally based on kinetically-controlled methods. Here, the thickness of the coating is controlled by the supply of the coating-forming material, e.g., via sputtering (Fig. 2(c)), vapor deposition (Fig. 2(d)), and reactions with vapor (Fig. 2(e)) or solution (Fig. 2(f)), or the thickness of the initial metal foil (Fig. 2(g)). In other words, the thickness of the coating is not self-limiting. Another issue with any kinetically-controlled coating method is the relatively high possibility of forming defects in the coating.

As the field moves forward, we may further exploit interfacial engineering based on thermodynamic approaches⁶¹. One example, albeit not the focus of this perspective, is represented by the adsorption of cations such as Cs^+ and Rb^{2+} from additives in the liquid electrolytes on the Li metal surface (Fig. 2(a))⁶², which will form adsorbates that serve as a

“protective coating” on the lithium anode. Such adsorption can be viewed to form spontaneously at thermodynamic equilibrium. Therefore, it may rebuild (though it can be inhibited by SEI formation) upon Li plating and stripping to self-heal and provide uniform Li deposition. An in-depth discussion of this approach is beyond the scope of this perspective, but can be found in other reviews of electrolyte additives^{2,9}.

A good example of thermodynamically-controlled interfacial engineering is to alloy (dope) Li with a surface-active or segregating metal element M to promote the surface segregation of M in a Li- M system as a spontaneously formed protective “coating”. Equivalently in thermodynamics, segregation is the same as adsorption at a thermodynamic equilibrium; a surface phase can form via segregation/adsorption of a dopant (alloying element) or spreading of a prewetting precursor film⁶³. In a broader context, Fig. 5 shows that various 2D interfacial phases can form at thermodynamic equilibria.^{64–66} These interfacial phases form spontaneously (which may lead to less defects in comparison with kinetically-controlled coating processing) when chemical segregation couples with interfacial structural transitions (e.g., reconstruction) or disordering (e.g., formation of premelting like interfacial layers).⁶⁴ Such interfacial phases, which are 2D thermodynamically (that is, no degree of freedom in the third dimension perpendicular to the interface), have been named as “complexions”⁶⁴ to differentiate them from thin layers of bulk (3D) phases present at interfaces (with arbitrary or “kinetically-controlled” thicknesses). In other words, a 2D interfacial phase (or complexion) should exhibit a thermodynamically-determined or self-selecting thickness (independent of kinetic factors, once a thermodynamic equilibrium is achieved)⁶⁴. While it is successfully applied to improve the performance of various cathodes (Figs. 2(b) and 5(d)) and other anodes such as TiO₂ (Fig. 5(c)), this new perspective for interfacial engineering has yet to be explored for Li metal anodes.

Let us briefly explain the underlying physics of the formation of 2D interfacial phases in several steps, as follows. First, we consider that the β phase can perfectly wet an α - γ interface if:

$$\gamma_{\alpha\beta} + \gamma_{\beta\gamma} < \gamma_{\alpha\gamma}^{(0)}, \quad (1)$$

where $\gamma_{\alpha\beta}$, $\gamma_{\beta\gamma}$, and $\gamma_{\alpha\gamma}^{(0)}$ are the interfacial energies of α - β , β - γ , and α - γ interfaces, respectively, and the superscript “(0)” in $\gamma_{\alpha\gamma}^{(0)}$ denotes it is a hypothetical “clean” interface

without the wetting or adsorption (while the equilibrium $\gamma_{\alpha\gamma} \equiv \gamma_{\alpha\beta} + \gamma_{\beta\gamma}$ for the case of perfect wetting)⁶³. Second, if the β phase is not stable as a (3D) bulk phase, a β -like (2D) interfacial phase (of microscopic thickness h) can still be stabilized at the α - γ interface thermodynamically by the reduction of interfacial energies ($-\Delta\gamma$ defined below), if:

$$-\Delta\gamma \equiv \gamma_{\alpha\gamma}^{(0)} - (\gamma_{\alpha\beta} + \gamma_{\beta\gamma}) < \Delta G_{\beta}^{(vol)} \cdot h, \quad (2)$$

where $\Delta G_{\beta}^{(vol)}$ is the volumetric free energy penalty for forming the metastable β phase⁶⁷⁻⁶⁹. This phenomenon is called “prewetting” in physics, which refers to “wetting” occurring when the phase does the wetting is not yet a stable bulk phase^{63,70,71}. Third, in a phenomenological thermodynamic theory, the interfacial excess grand potential of this β -like interfacial phase can be written as^{67,68}:

$$\Phi^x = \gamma_{\alpha\gamma}^{(0)} + \Delta\gamma \cdot f(h) + \Delta G_{\beta}^{(vol)} \cdot h. \quad (3)$$

Here, $f(h)$ is a dimensionless interfacial coefficient that describes the details of the thickness-dependent interfacial interactions, which should satisfy the boundary conditions: $f(0) = 0$ and $f(+\infty) = 1$. Minimization of the interfacial excess grand potential in Eq. (3) defines the equilibrium interfacial configuration (with an “equilibrium” or thermodynamically-determined thickness h_{EQ}). In a continuum approximation, the interfacial coefficient depends on thickness exponentially if a short-range interaction dominates or quadratically if an unretarded London dispersion force dominates⁶⁷⁻⁶⁹. For an idealized “hard-sphere” liquid, an additional oscillatory structural (solvation) interaction⁷² arises, producing an idealized series of discrete complexions⁶⁴: i.e., clean, monolayer, bilayer, trilayer, nanolayer, and wetting layer complexions, as shown in Fig. 5(b). In other words, these six complexions correspond to: $h_{EQ} = 0, 1\sigma, 2\sigma, 3\sigma, x$ (~ 1 nm), and $+\infty$, respectively, where σ is the thickness of a monolayer of adsorbates⁷³. This series of complexions have been discovered by Dillon and Harmer⁷⁴ at Al_2O_3 grain boundaries with different dopants, but they can also exist in other types of interfaces. A more rigorous thermodynamic model can be found at Ref. 75, with further discussion in an overview article⁶⁴. For Li-M systems, however, the interfacial coefficient $f(h)$ should be more complex than the overly simplified “hard spheres” model due to strong Li-M bonding, so we do not expect a simple series of Dillon-Harmer complexions to occur in a sequence. Here, reconstruction like

those observed in the Ni-Bi^{65,66}, Cu-Bi⁷⁶, Al-Cu⁷⁷, Si-Au^{73,78}, or transition metal doped WC⁷⁹ are expected (with possibly even more complexity and diversity). Disordered interfacial phases (nanoscale amorphous-like intergranular films) similar to those observed in W-Ni^{80,81}, Mo-Ni⁸², Ni-W⁸³, Ni-S⁸⁴, Cu-Zr,⁸⁵ and Cu-Zr-Hf⁸⁶ may also exist in certain Li-M systems. Noting that the above thermodynamic framework is applicable surfaces, grain boundaries, and other types of interfaces. The specific atomic configurations for 2D interfacial phases in Li-M systems, which should also depend on the specific metal M, have yet to be characterized experimentally.

Different from previously discussed kinetically-controlled coatings where the thickness is determined by the amount of the coating material supplied or the processing time (Fig. 2 (c-f)), a thermodynamically-controlled 2D surface phase can form spontaneously with a self-limiting thickness at a thermodynamic equilibrium⁶¹. Prior studies have demonstrated that Li₃PO₄ based surface amorphous films (SAFs), which are a type of 2D surface phases with self-selecting or “equilibrium” thickness on the order of one nanometer, can form and improve cycling stability and rate capability of various cathode materials, including LiNi_{0.5}Mn_{1.5}O₄, LiCoO₂, and Li_{1.33}Ni_{0.3}Mn_{0.57}O₄ (see one example in Fig. 2(b))⁸⁷⁻⁸⁹. WO₃-based surface phases have also been used to improve the discharge capacities of Li_{1.33}Ni_{0.3}Mn_{0.57}O₄ by decreasing the surface Ni/Mn ratio and changing the surface valence state⁸⁸. Surface nitridation has been used to form a TiO_xN_y-based (or TiN-like) surface phase (Fig. 5(c), i.e. a “surface precursor” to the bulk TiN phase) to increase the rate capability of TiO₂ anode by increasing the surface electronic conductivity (resembling the properties of the conductive TiN)⁹⁰. It should be noted that most of these 2D surface phases formed in facile mixing and annealing at elevated temperatures. For low-melting lithium, if an equilibrium surface phase can form near room temperature, it may also be able to rebuild during cycling and serve as a self-healing surface coating – an intriguing direction that needs to be explored in future studies.

In addition to interfacial engineering of the Li metal anode, this thermodynamic approach also leads to other potential opportunities. Such interfacial complexions can also exist at grain boundaries to provide potential benefits^{91,92}. For example, analogous phosphate-based “amorphous” films of equilibrium thicknesses on the order of one nanometer can form at both surfaces and grain boundaries of sintered LiFePO₄ cathode particles to provide fast Li⁺ conduction pathways (Fig. 5(d))^{93,94}. Specifically, for the Li metal anode, premelting-like grain boundaries, which are the interfacial precursor to the liquid phase formed below the bulk melting

temperature,⁷⁰ may form near room temperature. If so, this could provide super-plasticity (via promoting grain boundary sliding and Coble creep at room temperature), which would be a very desirable property. This effect represents yet another new perspective that needs to be verified and realized in a future study. Yet another possibility is to use complexions to tailor the interfaces in (kinetically-controlled) thick protective coatings, e.g., to improve the interfacial ionic conductivity of the composite coatings similar to that shown in Fig. 2(f).

Advanced characterization is needed to elucidate the mechanism of interfacial phases. Various transmission electron microscopy (TEM) techniques including aberration-corrected scanning transmission electron microscopy (AC-STEM), energy dispersive X-ray spectroscopy (EDS), and electron energy loss spectroscopy (EELS) can be used to study the atomic level structure^{65,95}, elemental composition⁶⁶, and oxidation state of the interfacial phases⁷³, respectively. In-situ TEM setups^{96,97} have been recently developed, which enables the characterization of dynamic changes of the interfacial phases under a real working environment. Due to the high reactivity of Li metal, TEM characterization needs to be conducted under cryogenic temperature to prevent electron beam damage on Li metal^{53,98,99}. Care must be taken to prevent air-exposure during the sample transfer as Li can easily react with oxygen, nitrogen, and water vapor. Cryo-FIB-SEM can also be used to study the cross-section microstructure and Li deposition morphology⁵⁵. Electrochemical impedance spectroscopy (EIS) combined with equivalent circuit model fitting can also be used to quantify bulk and interfacial charge transport properties⁵⁰.

In general, 2D interfacial phases or complexions, which have structures that are neither observed in, nor necessarily stable as, conventional 3D bulk phases, can potentially help to obtain exceptional properties that many bulk phases cannot achieve alone, including ionic conductivity, interfacial stability, and cyclability. Some of these ideas have already been demonstrated for various battery systems, as discussed in a recent review⁶¹. Specifically, the feasibility of using 2D interface phases spontaneously formed at thermodynamic equilibria to tailor the lithium metal anode offers potential new opportunities, that must be further explored.

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Data Availability Statement

Data sharing is not applicable to this perspective article as no new data were created or analyzed in this study.

TABLE 1. A summary of Li diffusivity in various Li-rich alloy type coatings. In the method column, the following abbreviations are used: nuclear magnetic relaxation (NMR) and galvanostatic intermittent titration technique (GITT).

Material	Li Diffusivity ($\text{cm}^2 \text{s}^{-1}$)	Temperature	Method	Ref
Li (self-diffusion)	7.65×10^{-11}	25 °C	NMR	100
	6.12×10^{-11}	25 °C	Thermodynamic Model	30
Li₁₃Sn₅	5.01×10^{-5} to 7.59×10^{-4}	415 °C	GITT	101
Li₂₂Sn₅	1.9×10^{-7} to 5.9×10^{-7}	25 °C	GITT	44
	6.58×10^{-5} to 1.91×10^{-4}	415 °C	GITT	101
Li₂₂Si₅	5.13×10^{-5} to 7.24×10^{-5}	415 °C	GITT	102
Li₃Sb	2.0×10^{-4} to 4.0×10^{-4}	360 °C	GITT	103
Li₃Bi	1.0×10^{-6} to 3.0×10^{-6}	25 °C	GITT	104
LiZn	4.0×10^{-7} to 4.0×10^{-8}	25 °C	GITT	24,44
Li-In (47-62 at % Li)	4.73×10^{-7} to 3.98×10^{-5}	415 °C	GITT	105
Li_xAg (x=4.7-5.0)	0.12×10^{-8} to 4.0×10^{-8}	25 °C	GITT	106
Li-Mg (β-phase)	$\approx 10^{-8}$	25 °C	GITT	23
	$\approx 10^{-11}$	25 °C	Neutron Tomography	107
	2.3×10^{-11}	25 °C	Sand's Equation	108

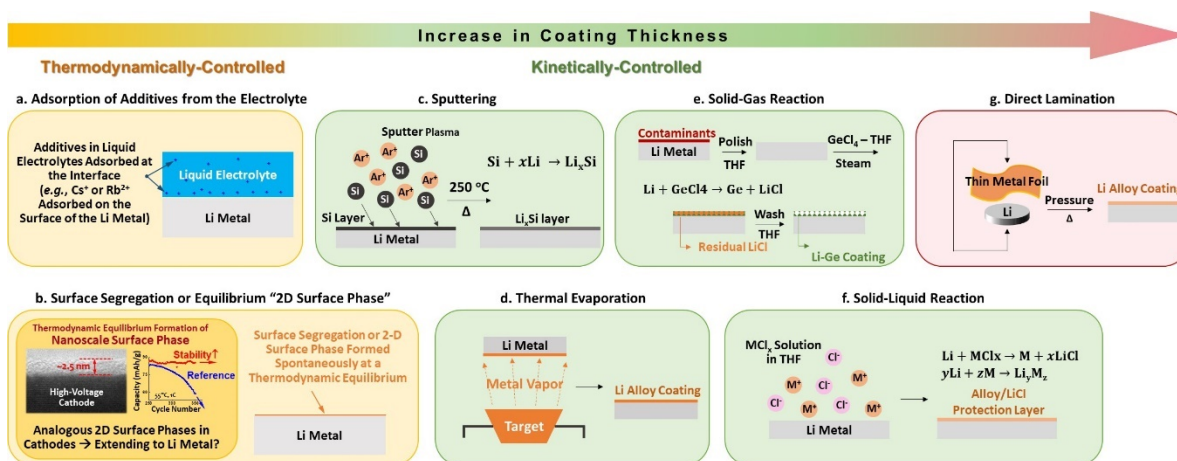


FIG. 2. Schematic illustration of various methods for surface modifications and coatings with increasing effective interfacial thickness. (a) It should be noted that the adsorption of cations⁶² (e.g., Cs⁺ and Rb²⁺) from the additives in liquid electrolytes on the surface of liquid metals represents a case of thermodynamically-controlled surface modification (albeit it is not the focus of this perspective). (b) Likewise, we propose to utilize surface segregation (or spreading of a wetting or prewetting precursor film⁶³) or equilibrium formation of 2-D surface phases, which have been proven feasible and useful to improve the stability of cathode surfaces⁸⁹, as a potentially new surface engineering method for lithium metal anodes (yet to be explored). More commonly adopted, various kinetically-controlled surface coatings to improve lithium method anodes include: (c) Li_xSi layer fabricated from sputtering and annealing²⁷, (d) alloy coatings formed by metal evaporation onto Li metal surface^{23,24}, (e) Li-Ge based surface modification synthesized from spontaneous reaction between Li and GeCl₄-THF vapor⁴⁰, and (f) alloy-LiCl composite protection layer synthesized from reaction between Li and metal chloride solution³⁹. (g) Finally, direct lamination of metal foils on lithium with controlled pressure and temperature represents another method to make even thick alloy type interfacial layers²⁸.

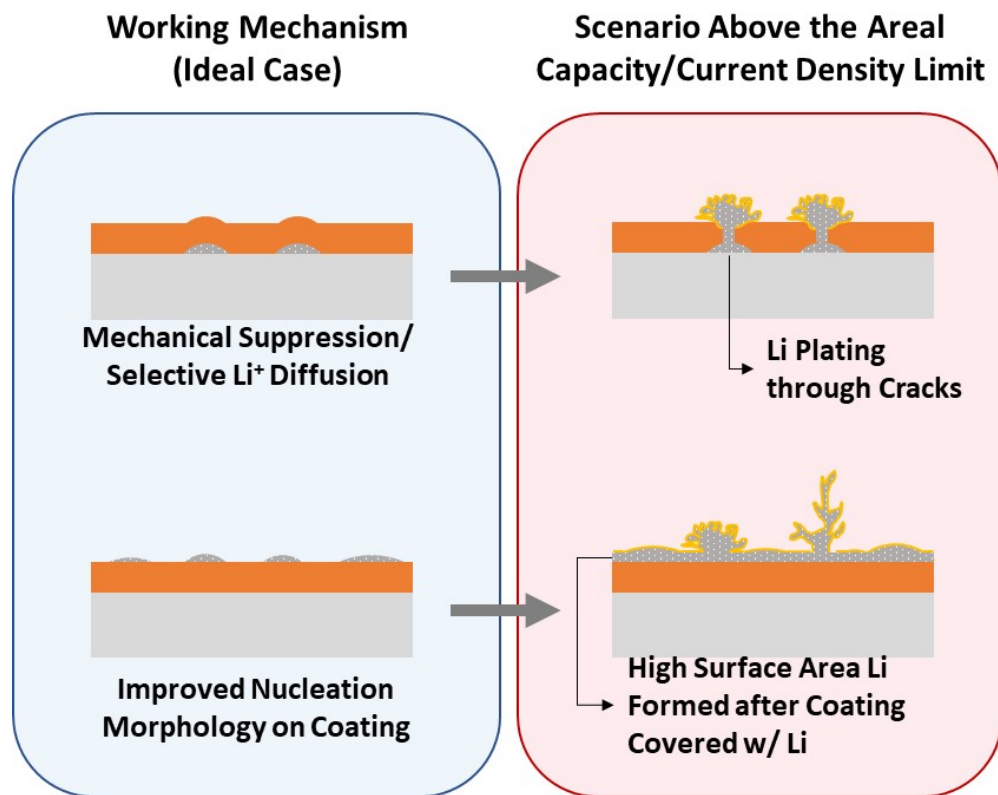


FIG. 3. Schematics of possible failure mechanisms of alloy type coatings after the areal capacity or current density limit is exceeded.

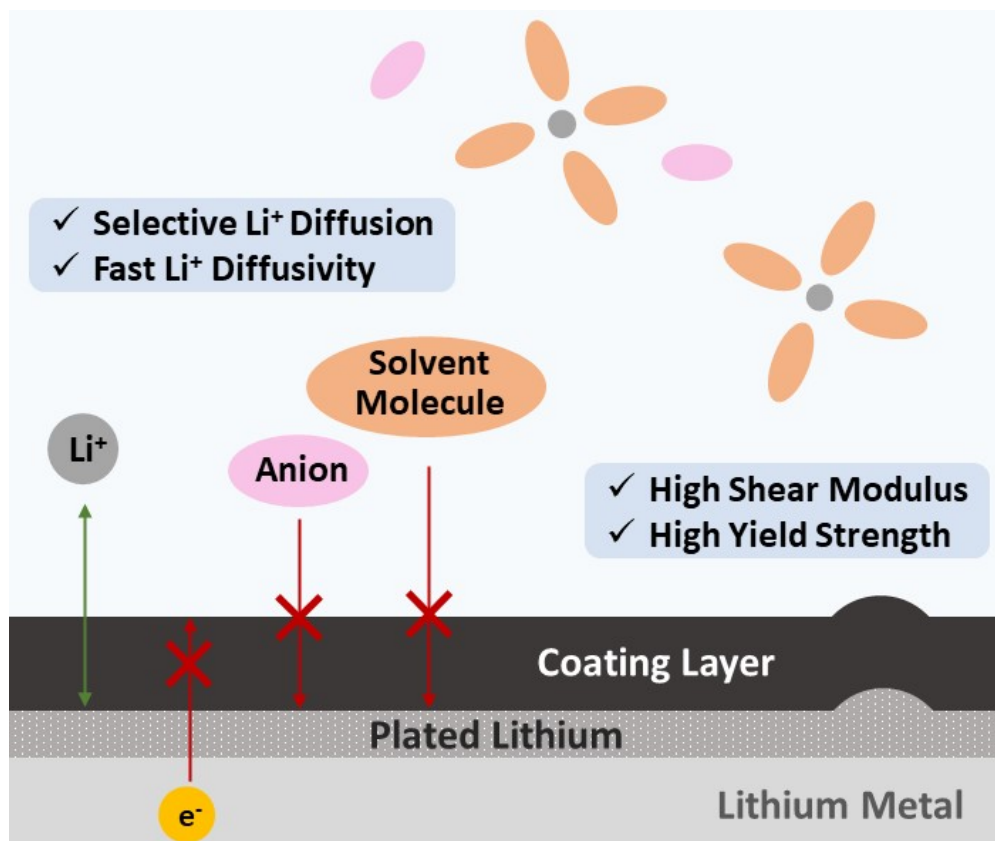


FIG. 4. Schematic illustration of ideal features of the interfacial protection layer on lithium metal^{3,6}.

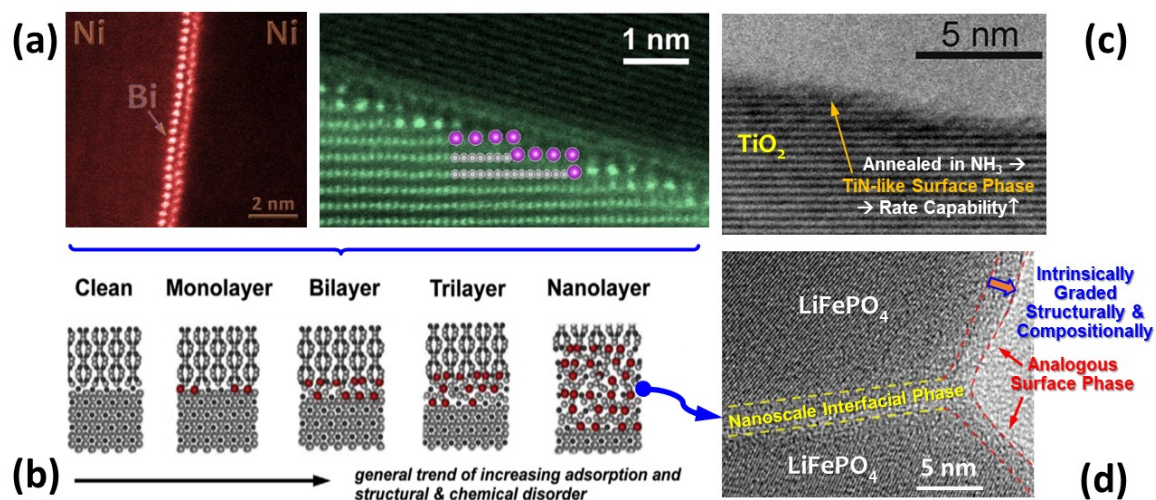


FIG. 5. Perspectives on thermodynamically-controlled interfacial engineering. (a) Example of ordered bismuth adsorption at nickel general grain boundaries^{65,66}. (b) Schematic illustration of a series of 2D interfacial phases at grain boundaries⁶⁴ and presumably also hetero-phase interfaces. (c) Surface nitridation of TiO₂ Li-ion battery anode material to form TiO_xN_y-based (TiN-like) surface phase to increase the rate capability and surface electronic conductivity⁹⁰. (d) Enhanced Li⁺ conductivity from utilizing analogous 2D interfacial phases at surfaces and grain boundaries in LiFePO₄ cathode material⁹³.

References

- ¹ M.S. WHITTINGHAM, *Science* (80-.). **192**, 1126 LP (1976).
- ² X.B. Cheng, R. Zhang, C.Z. Zhao, and Q. Zhang, *Chem. Rev.* **117**, 10403 (2017).
- ³ H. Zhou, S. Yu, H. Liu, and P. Liu, *J. Power Sources* **450**, 227632 (2020).
- ⁴ D. Lin, Y. Liu, and Y. Cui, *Nat. Nanotechnol.* **12**, 194 (2017).
- ⁵ L. Ma, M.S. Kim, and L.A. Archer, *Chem. Mater.* **29**, 4181 (2017).
- ⁶ R. Xu, X.B. Cheng, C. Yan, X.Q. Zhang, Y. Xiao, C.Z. Zhao, J.Q. Huang, and Q. Zhang, *Matter* **1**, 317 (2019).
- ⁷ J.B. Goodenough and Y. Kim, *Chem. Mater.* **22**, 587 (2010).
- ⁸ K. Xu, (2004).
- ⁹ K. Xu, *Chem. Rev.* **114**, 11503 (2014).
- ¹⁰ R.L. Sacci, N.J. Dudney, K.L. More, L.R. Parent, I. Arslan, N.D. Browning, and R.R. Unocic, *Chem. Commun.* **50**, 2104 (2014).
- ¹¹ X.Q. Zhang, X.B. Cheng, X. Chen, C. Yan, and Q. Zhang, *Adv. Funct. Mater.* **27**, 1 (2017).
- ¹² D. Aurbach, K. Gamolsky, B. Markovsky, Y. Gofer, M. Schmidt, and U. Heider, *Electrochim. Acta* **47**, 1423 (2002).
- ¹³ Y. Liu, D. Lin, Z. Liang, J. Zhao, K. Yan, and Y. Cui, *Nat. Commun.* **7**, 1 (2016).
- ¹⁴ D. Lin, Y. Liu, Z. Liang, H.W. Lee, J. Sun, H. Wang, K. Yan, J. Xie, and Y. Cui, *Nat. Nanotechnol.* **11**, 626 (2016).
- ¹⁵ H. Wu, D. Zhuo, D. Kong, and Y. Cui, *Nat. Commun.* **5**, 1 (2014).
- ¹⁶ M.S. Gonzalez, Q. Yan, J. Holoubek, Z. Wu, H. Zhou, N. Patterson, V. Petrova, H. Liu, and P. Liu, *Adv. Mater.* **1906836**, 1 (2020).
- ¹⁷ S. Bai, X. Liu, K. Zhu, S. Wu, and H. Zhou, *Nat. Energy* **1**, (2016).
- ¹⁸ S. Bai, Y. Sun, J. Yi, Y. He, Y. Qiao, and H. Zhou, *Joule* **2**, 2117 (2018).
- ¹⁹ E. Cha, M.D. Patel, J. Park, J. Hwang, V. Prasad, K. Cho, and W. Choi, *Nat. Nanotechnol.* **13**, 337 (2018).
- ²⁰ Y. Gao, Z. Yan, J.L. Gray, X. He, D. Wang, T. Chen, Q. Huang, Y.C. Li, H. Wang, S.H. Kim, T.E. Mallouk, and D. Wang, *Nat. Mater.* **18**, 384 (2019).
- ²¹ R.A. Huggins, *J. Power Sources* **22**, 341 (1988).
- ²² R.A. Huggins, *J. Power Sources* **26**, 109 (1989).
- ²³ Z. Shi, M. Liu, D. Naik, and J.L. Gole, *J. Power Sources* **92**, 70 (2001).
- ²⁴ Z. Shi, M. Liu, and J.L. Gole, *Electrochem. Solid-State Lett.* **3**, 312 (2000).

- ²⁵ S. Choudhury, Z. Tu, S. Stalin, D. Vu, K. Fawole, D. Gunceler, R. Sundararaman, and L.A. Archer, *Angew. Chemie* **129**, 13250 (2017).
- ²⁶ Z. Tu, S. Choudhury, M.J. Zachman, S. Wei, K. Zhang, L.F. Kourkoutis, and L.A. Archer, *Nat. Energy* **3**, 310 (2018).
- ²⁷ W. Tang, X. Yin, S. Kang, Z. Chen, B. Tian, S.L. Teo, X. Wang, X. Chi, K.P. Loh, H.W. Lee, and G.W. Zheng, *Adv. Mater.* **30**, 1 (2018).
- ²⁸ H. Kim, J.T. Lee, D.C. Lee, M. Oschatz, W. Il Cho, S. Kaskel, and G. Yushin, *Electrochem. Commun.* **36**, 38 (2013).
- ²⁹ W. Weppner, *J. Electrochem. Soc.* **124**, 1569 (1977).
- ³⁰ E. Dologlou, *Glas. Phys. Chem.* **36**, 570 (2010).
- ³¹ K. Yan, Z. Lu, H.W. Lee, F. Xiong, P.C. Hsu, Y. Li, J. Zhao, S. Chu, and Y. Cui, *Nat. Energy* **1**, (2016).
- ³² R.W. Balluffi, S.M. Allen, and W.C. Carter, *Kinetics of Materials* (John Wiley & Sons, 2005).
- ³³ D.R. Ely and R.E. García, *J. Electrochem. Soc.* **160**, A662 (2013).
- ³⁴ A. Pei, G. Zheng, F. Shi, Y. Li, and Y. Cui, *Nano Lett.* **17**, 1132 (2017).
- ³⁵ K.S. Nagy, S. Kazemiabnavi, K. Thornton, and D.J. Siegel, *ACS Appl. Mater. Interfaces* **11**, 7954 (2019).
- ³⁶ A.J. Sanchez, E. Kazyak, Y. Chen, K.H. Chen, E.R. Pattison, and N.P. Dasgupta, *ACS Energy Lett.* 994 (2020).
- ³⁷ K.H. Chen, K.N. Wood, E. Kazyak, W.S. Lepage, A.L. Davis, A.J. Sanchez, and N.P. Dasgupta, *J. Mater. Chem. A* **5**, 11671 (2017).
- ³⁸ K.N. Wood, M. Noked, and N.P. Dasgupta, *ACS Energy Lett.* **2**, 664 (2017).
- ³⁹ X. Liang, Q. Pang, I.R. Kochetkov, M.S. Sempere, H. Huang, X. Sun, and L.F. Nazar, *Nat. Energy* **2**, (2017).
- ⁴⁰ K. Liao, S. Wu, X. Mu, Q. Lu, M. Han, P. He, Z. Shao, and H. Zhou, *Adv. Mater.* **30**, 1 (2018).
- ⁴¹ T. Chen, W. Kong, P. Zhao, H. Lin, Y. Hu, R. Chen, W. Yan, and Z. Jin, *Chem. Mater.* **31**, 7565 (2019).
- ⁴² C. Monroe and J. Newman, *J. Electrochem. Soc.* **151**, A880 (2004).
- ⁴³ C. Monroe and J. Newman, *J. Electrochem. Soc.* **152**, A396 (2005).
- ⁴⁴ A. Anani and R.A. Huggins, **134**, 3098 (1987).
- ⁴⁵ A.M. Russell, *Adv. Eng. Mater.* **5**, 629 (2003).
- ⁴⁶ W. Liu, R. Guo, B. Zhan, B. Shi, Y. Li, H. Pei, Y. Wang, W. Shi, Z. Fu, and J. Xie, *ACS Appl. Energy Mater.* **1**, 1674 (2018).
- ⁴⁷ A.C. Kozen, C.F. Lin, O. Zhao, S.B. Lee, G.W. Rubloff, and M. Noked, *Chem. Mater.* **29**,

6298 (2017).

⁴⁸ J. Xiao, *Science* (80-.). **366**, 426 (2019).

⁴⁹ M.T. Ong, O. Vernali, E.W. Draeger, A.C.T. Van Duin, V. Lordi, and J.E. Pask, *J. Phys. Chem. B* **119**, 1535 (2015).

⁵⁰ M.R. Busche, T. Drossel, T. Leichtweiss, D.A. Weber, M. Falk, M. Schneider, M.L. Reich, H. Sommer, P. Adelhelm, and J. Janek, *Nat. Chem.* **8**, 426 (2016).

⁵¹ Y. Yamada, F. Sagane, Y. Iriyama, T. Abe, and Z. Ogumi, *J. Phys. Chem. C* **113**, 14528 (2009).

⁵² B. Wen, Z. Deng, P.C. Tsai, Z.W. Lebens-Higgins, L.F.J. Piper, S.P. Ong, and Y.M. Chiang, *Nat. Energy* (2020).

⁵³ Y. Li, Y. Li, A. Pei, K. Yan, Y. Sun, C.L. Wu, L.M. Joubert, R. Chin, A.L. Koh, Y. Yu, J. Perrino, B. Butz, S. Chu, and Y. Cui, *Science* (80-.). **358**, 506 (2017).

⁵⁴ K.J. Harry, D.T. Hallinan, D.Y. Parkinson, A.A. MacDowell, and N.P. Balsara, *Nat. Mater.* **13**, 69 (2014).

⁵⁵ J.Z. Lee, T.A. Wynn, M.A. Schroeder, J. Alvarado, X. Wang, K. Xu, and Y.S. Meng, *ACS Energy Lett.* **4**, 489 (2019).

⁵⁶ Y. Zhu, X. He, and Y. Mo, *ACS Appl. Mater. Interfaces* **7**, 23685 (2015).

⁵⁷ A. Banerjee, X. Wang, C. Fang, E.A. Wu, and Y.S. Meng, (2020).

⁵⁸ Y. Xiao, Y. Wang, S.H. Bo, J.C. Kim, L.J. Miara, and G. Ceder, *Nat. Rev. Mater.* **5**, 105 (2020).

⁵⁹ H. Zhang, C. Li, M. Piszcz, E. Coya, T. Rojo, L.M. Rodriguez-Martinez, M. Armand, and Z. Zhou, *Chem. Soc. Rev.* **46**, 797 (2017).

⁶⁰ Z. Jiang, L. Jin, Z. Han, W. Hu, Z. Zeng, Y. Sun, and J. Xie, *Angew. Chemie - Int. Ed.* **58**, 11374 (2019).

⁶¹ J. Luo, *Energy Storage Mater.* **21**, 50 (2019).

⁶² F. Ding, W. Xu, G.L. Graff, J. Zhang, M.L. Sushko, X. Chen, Y. Shao, M.H. Engelhard, Z. Nie, J. Xiao, X. Liu, P. V. Sushko, J. Liu, and J.G. Zhang, *J. Am. Chem. Soc.* **135**, 4450 (2013).

⁶³ P.G. De Gennes, *Rev. Mod. Phys.* **57**, 827 (1985).

⁶⁴ P.R. Cantwell, M. Tang, S.J. Dillon, J. Luo, G.S. Rohrer, and M.P. Harmer, *Acta Mater.* **62**, 1 (2014).

⁶⁵ J. Luo, H. Cheng, K.M. Asl, C.J. Kiely, and M.P. Harmer, *Science* (80-.). **333**, 1730 (2011).

⁶⁶ Z. Yu, P.R. Cantwell, Q. Gao, D. Yin, Y. Zhang, N. Zhou, G.S. Rohrer, M. Widom, J. Luo, and M.P. Harmer, *Science* (80-.). **358**, 97 (2017).

⁶⁷ J. Luo, *Crit. Rev. Solid State Mater. Sci.* **32**, 67 (2007).

- ⁶⁸ J. Luo and Y.-M. Chiang, *Annu. Rev. Mater. Res.* **38**, 227 (2008).
- ⁶⁹ J. Luo, *J. Am. Ceram. Soc.* **95**, 2358 (2012).
- ⁷⁰ J.W. Cahn, *J. Chem. Phys.* **66**, 3667 (1977).
- ⁷¹ D. Bonn, J. Eggers, J. Indekeu, and J. Meunier, *Rev. Mod. Phys.* **81**, 739 (2009).
- ⁷² J.N. Israelachvili, *Intermolecular and Surface Forces* (Academic press, 2011).
- ⁷³ S. Ma, P.R. Cantwell, T.J. Pennycook, N. Zhou, M.P. Oxley, D.N. Leonard, S.J. Pennycook, J. Luo, and M.P. Harmer, *Acta Mater.* **61**, 1691 (2013).
- ⁷⁴ S.J. Dillon, M. Tang, W.C. Carter, and M.P. Harmer, *Acta Mater.* **55**, 6208 (2007).
- ⁷⁵ J. Luo, *Appl. Phys. Lett.* **95**, 10 (2009).
- ⁷⁶ A. Kundu, K.M. Asl, J. Luo, and M.P. Harmer, *Scr. Mater.* **68**, 146 (2013).
- ⁷⁷ P. Parajuli, D. Romeu, V. Hounkpati, R. Mendoza-Cruz, J. Chen, M.J. Yacamán, J. Flowers, and A. Ponce, *Acta Mater.* **181**, 216 (2019).
- ⁷⁸ C. Hu and J. Luo, *Scr. Mater.* **158**, 11 (2019).
- ⁷⁹ Z. Luo, C. Hu, L. Xie, H. Nie, C. Xiang, X. Gu, J. He, W. Zhang, Z. Yu, and J. Luo, *Mater. Horizons* **7**, 173 (2020).
- ⁸⁰ V.K. Gupta, D.H. Yoon, H.M. Meyer, and J. Luo, *Acta Mater.* **55**, 3131 (2007).
- ⁸¹ J. Luo, V.K. Gupta, D.H. Yoon, and H.M. Meyer, *Appl. Phys. Lett.* **87**, 1 (2005).
- ⁸² X. Shi and J. Luo, *Phys. Rev. Lett.* **105**, 1 (2010).
- ⁸³ J.D. Schuler, O.K. Donaldson, and T.J. Rupert, *Scr. Mater.* **154**, 49 (2018).
- ⁸⁴ T. Hu, S. Yang, N. Zhou, Y. Zhang, and J. Luo, *Nat. Commun.* **9**, (2018).
- ⁸⁵ A. Khalajhedayati, Z. Pan, and T.J. Rupert, *Nat. Commun.* **7**, (2016).
- ⁸⁶ C.M. Grigorian and T.J. Rupert, *Acta Mater.* **179**, 172 (2019).
- ⁸⁷ H. Liu, J. Huang, D. Qian, S. Hy, C. Fang, J. Luo, and Y.S. Meng, *J. Electrochem. Soc.* **163**, A971 (2016).
- ⁸⁸ J. Huang, H. Liu, T. Hu, Y.S. Meng, and J. Luo, *J. Power Sources* **375**, 21 (2018).
- ⁸⁹ J. Huang and J. Luo, *Phys. Chem. Chem. Phys.* **16**, 7786 (2014).
- ⁹⁰ M. Samiee and J. Luo, *J. Power Sources* **245**, 594 (2014).
- ⁹¹ D.H.S. Tan, A. Banerjee, Z. Chen, and Y.S. Meng, *Nat. Nanotechnol.* **15**, 170 (2020).
- ⁹² J. Luo, *J. Mater.* **1**, 22 (2015).
- ⁹³ A. Kayyar, H. Qian, and J. Luo, *Appl. Phys. Lett.* **95**, 95 (2009).
- ⁹⁴ B. Kang and G. Ceder, *Nature* **458**, 190 (2009).

- ⁹⁵ T. Meiners, T. Frolov, R.E. Rudd, G. Dehm, and C.H. Liebscher, *Nature* **579**, 375 (2020).
- ⁹⁶ X.H. Liu, L. Zhong, L.Q. Zhang, A. Kushima, S.X. Mao, J. Li, Z.Z. Ye, J.P. Sullivan, and J.Y. Huang, *Appl. Phys. Lett.* **98**, 1 (2011).
- ⁹⁷ R.R. Unocic, X.G. Sun, R.L. Sacci, L.A. Adamczyk, D.H. Alsem, S. Dai, N.J. Dudney, and K.L. More, *Microsc. Microanal.* **20**, 1029 (2014).
- ⁹⁸ X. Wang, M. Zhang, J. Alvarado, S. Wang, M. Sina, B. Lu, J. Bouwer, W. Xu, J. Xiao, J.-G. Zhang, J. Liu, and Y.S. Meng, *Nano Lett.* **17**, 7606 (2017).
- ⁹⁹ M.J. Zachman, Z. Tu, S. Choudhury, L.A. Archer, and L.F. Kourkoutis, *Nature* **560**, 345 (2018).
- ¹⁰⁰ R. Messer and F. Noack, *Appl. Phys.* **6**, 79 (1975).
- ¹⁰¹ C.J. Wen and R.A. Huggins, *J. Solid State Chem.* **35**, 376 (1980).
- ¹⁰² C.J. Wen and R.A. Huggins, *J. Solid State Chem.* **37**, 271 (1981).
- ¹⁰³ W. Weppner and R.A. Huggins, *J. Solid State Chem.* **22**, 297 (1977).
- ¹⁰⁴ M. Hiratani, K. Miyauchi, and T. Kudo, *Solid State Ionics* **28–30**, 1406 (1988).
- ¹⁰⁵ C.J. Wen and R.A. Huggins, *Mater. Res. Bull.* **15**, 1225 (1980).
- ¹⁰⁶ S. Jin, Y. Ye, Y. Niu, Y. Xu, H. Jin, J. Wang, Z. Sun, A. Cao, X. Wu, Y. Luo, H. Ji, and L.-J. Wan, *J. Am. Chem. Soc.* **142**, 8818 (2020).
- ¹⁰⁷ Y. Zhang, K.S.R. Chandran, M. Jagannathan, H.Z. Bilheux, and J.C. Bilheux, *J. Electrochem. Soc.* **164**, A28 (2017).
- ¹⁰⁸ T. Krauskopf, B. Mogwitz, C. Rosenbach, W.G. Zeier, and J. Janek, *Adv. Energy Mater.* **9**, (2019).