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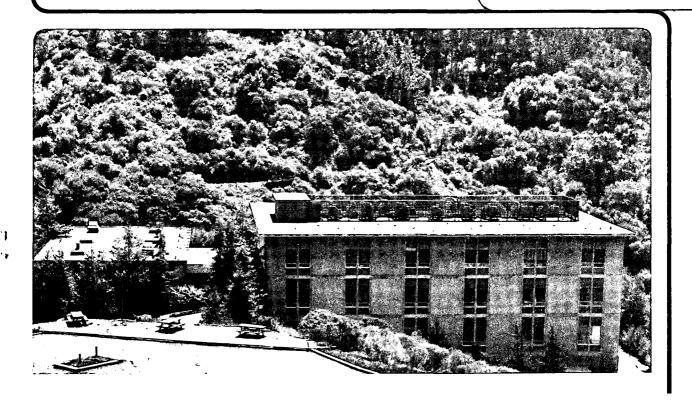
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# THRESHOLD PHOTOELECTRON SPECTRUM OF THE ARGON 3s SATELLITES

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### **ABSTRACT**

Lately a variety of techniques <sup>1,2</sup> have studied the electron correlation satellites with binding energies between the Argon 3s binding energy (29.24 eV) and the 2p<sup>-2</sup> ionization potential (43.38 eV). One of these techniques, Threshold Photoelectron Spectroscopy<sup>3</sup>, with ~90 meV FWHM resolution, revealed at least 25 individual electronic states. All of these could contribute to any other satellite spectrum, and this helped explain some discrepancies between previous measurements. This technique has been applied to the same region with higher resolution (<60 meV at the Ar 3s<sup>-1</sup>peak). In this higher resolution spectrum at least 29 individual electronic states are present. In some cases the multiplet splitting is observed.

### INTRODUCTION

The Argon inner valence satellites have been studied, recently, by a variety of techniques with a variety of incident photon energies. 1-6 By measuring the asymmetry parameter Adam et. al. 1 showed that the lower binding energy satellites have symmetries other than 2S, and at ~40 eV photon energy they also saw evidence for satellites with lower binding energies than had been seen in the preceding X-ray work. 7 Other subsequent investigations 4,5 concentrated upon improving the resolution at higher incident photon energies. Each time the resolution improved more final states emerged. Two other 3,6 studies of these satellites used intrinsically high resolution techniques. One 6 used broad band synchrotron radiation and detected the dispersed fluorescent radiation from the excited states to the ground state of the ion; the other 3 varied the incident photon energy of

synchrotron radiation and detected 0eV electrons. Both of these techniques revealed many new 3s satellites.

### **EXPERIMENT**

In this paper we report another measurement of the argon  $3s^{-1}$  satellites region using the same technique as reference 3. In this case the FWHM of the argon  $3s^{-1}$  peak was less than 60meV, which is an improvement of more than 30% over the previous measurement. The experiment was performed on the new grasshopper beamline at SSRL, using the 600 l/mm grating. The experimental apparatus has been described before.<sup>8,9</sup> In this instance a 1000Å silicon window separated the ultra high vacuum ( $10^{-10}$  torr) of the monochromator from the experimental chamber ( $10^{-5}$  torr). The amount of second order light was determined using the known cross section of neon  $2p^{-1}$ , and the cross sections of the satellites were corrected with this. The reported errors are statistical and do not reflect any error in this calibration procedure. The energy scale was calibrated using the  $3s^{-1}$  argon peak and setting it equal to  $29.24\text{eV}.^{10}$  The uncertainty with the energy scale is  $\pm$  0.02 eV at the lower energies, but probably as high as  $\pm$  0.05 eV close to the double ionization threshold.

### **RESULTS**

Figure 1 contains the Zero Kinetic Energy spectrum of argon from 28.5 to 44.5 eV. It is apparent from the spectrum that the satellites with lower binding energies are more prominent in the zero kinetic energy spectrum, in agreement with the results of Adam et. al. <sup>1</sup> The more congested region from 33.75 to 40.0 eV is presented in Figure 2 with a finer point density than in Figure 1, and the results from these spectra are summarized in Table 1.

A comparison with the higher energy, high resolution work of Brion et. al.<sup>4</sup> indicates that most of the satellites are present at all incident photon energies, with the exception being some of the lower binding energy quartet states, which are visibly present only at threshold. Some of these states are almost degenerate in energy with doubly excited states of the type 3s<sup>2</sup>3p<sup>4</sup>nl n'l',<sup>11</sup> and it has been shown previously that they are enhanced on resonance.<sup>9</sup> Recent work by Schartner et. al.<sup>12</sup> showed that some of the doubly excited states decay intermediately to the 3s<sup>-1</sup> state. It then seems quite possible that some also, or alternatively, decay through 3s<sup>2</sup>3p<sup>4</sup>nl states.

One of the most notable features of this spectrum is that two  $^2P$  states are well resolved into their individual J=1/2 and J=3/2 components. The intensity ratio  $I_3/2/I_{1/2}$  is  $4.5 \pm 1.0$  for the  $^2P$  state with binding energies 32.90 and 33.02 eV and 2.88  $\pm$  0.24 for

the 33.70 and 33.82 eV peaks. This indicates that the ratio is greater than expected statistically in both cases. The error in the intensity, however, does not include possible systematic errors due to the correction for second order light. Therefore these relative intensities must be regarded as tentative.

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## Figure Captions:

- Fig 1. Zero kinetic energy spectrum of argon from the 3s threshold to the double ionization limit. Data points are uncorrected for second order light. The smooth line is a least squares fit to the data.
- Fig. 2. Zero kinetic energy spectrum of argon over the most congested region of fig. 1 (33.2 to 40.0 eV) with a finer point density. Data points are uncorrected for second order light. The smooth line is a least squares fit to the data.

Table 1 Argon 3s Satellites

Argon 3s Satelli	ites					
Assignment	Optical Energy <sup>a</sup>	Energyb	σ(Mb) <sup>b</sup>	Energy <sup>C</sup>	σ(Mb) <sup>C</sup>	Energy <sup>d</sup>
3s <sup>-1</sup>	29.24		0.8		,	
3p <sup>-2</sup> 3d <sup>4</sup> D	32.18	32.17	0.033(4)	32.21	0.045(4)	32.21
$3p^{-2}4s^{4}P$	32.46	32.47	0.027(4)	32.52	0.030(3)	
$3p^{-2}4s^{2}P_{3/2}$	32.90	32.88	0.076(7)	32.93	0.063(5)	
$3p^{-2}4s^{2}P_{1/2}$	33.02			33.04	0.014(3)	33.05
3p <sup>2</sup> 3d <sup>4</sup> F	33.45	33.54	0.071(6)	33.52	0.024(2)	33.56
				33.60	0.016(2)	
$3p^{-2}3d^{2}P_{1/2}$	33.70			33.72	0.069(5)	
$3p^{-2}3d^{2}P_{3/2}$	33.82	33.78	0.300(12)	33.84	0.248(13)	
$3p^{-2}3d^{4}P$	34.06	34.04	0.044(7)	34.05	0.018(2)	34.00
$3p^{-2}4s  ^2D_{3/2}$	34.18	34.20	0.196(11)	34.11	0.020(2)	34.16
3p <sup>-2</sup> 4p <sup>2</sup> D5/2	34.21			34.24	0.177(10)	
$3p^{-2}3d^{2}F$	34.30	34.39	0.088(5)			
$3p^{-2}3d^{2}D$	34.49			34.43	0.060(4)	34.44
3p <sup>-2</sup> 3d <sup>2</sup> G	34.88	34.91	0.039(4)	34.92	0.017(2)	34.83
3p <sup>-2</sup> 4p <sup>4</sup> P	35.01			35.02	0.012(2)	
3p <sup>-2</sup> 4p <sup>4</sup> D	35.30	35.28		35.36	0.023(2)	35.27
3p <sup>-2</sup> 4p <sup>2</sup> D	35.44	35.51	0.160(13)	35.50	0.051(4)	
$3p^{-2}4p^{2}P_{1/2}$	35.56			35.57	0.034(3)	35.59
$3p^{-2}4p^{2}P_{3/2}$	35.62			35.66	0.033(3)	
$3p^{-2}3d^{2}F$	36.03	36.00	0.094(7)	36.06	0.115(7)	36.02
3p-24s 2S	36.50	36.51	0.040(6)	36.55	0.016(2)	36.50
3p <sup>-2</sup> 4p <sup>2</sup> F	36.89	36.90	0.051(7)	36.96	0.038(3)	36.86
3p <sup>-2</sup> 4p <sup>2</sup> P	37.11	37.11	0.061(4)			37.14
$3p^{-2}3d^{2}D$	37.18	37.23	0.066(7)	37.18	0.054(4)	
3p <sup>-2</sup> 4p <sup>2</sup> D	37.25			37.29	0.070(5)	
$3p^{-2}3d^{2}P$	37.40	37.42	0.045(2)	37.48	0.024(3)	37.33
$3p^{-2}3d^{2}P$	38.04	38.03	0.044(2)	38.10	0.043(4)	38.01
3p <sup>-2</sup> 5s <sup>4</sup> P	38.32	38.33	0.007(1)			
3p <sup>-2</sup> 3d <sup>2</sup> S	38.58	38.59	0.038(2)	38.66	0.019(2)	38.56
$3p^{-2}3d^{2}D$	38.91	38.90	0.011(1)	38.98	0.014(20)	38.80
3p <sup>-2</sup> 5p <sup>2</sup> P	39.33	39.35	0.038(2)			39.32
3p <sup>-2</sup> 5p <sup>2</sup> D	39.38					
3p <sup>-2</sup> 4d <sup>2</sup> P	39.39			39.45	0.028(2)	
<sub>3p</sub> -2 <sub>4p</sub> 2 <sub>P</sub>	39.57	39.61	0.028(2)			39.60
$3p^{-2}4d^{2}D$	39.64			39.70	0.022(2)	
$_{3p}^{-2}_{5s}^{2}_{D}$	40.04	40.03	0.025(3)			39.91
3p <sup>-2</sup> 4f	40.07					
3p-24d 2S	41.20	41.16	0.023(4)	41.24	0.027(3)	41.15
sat. total			1.68		1.49	

a. Reference 10.

b. Reference 9.

c. This work.

d. Reference 4.

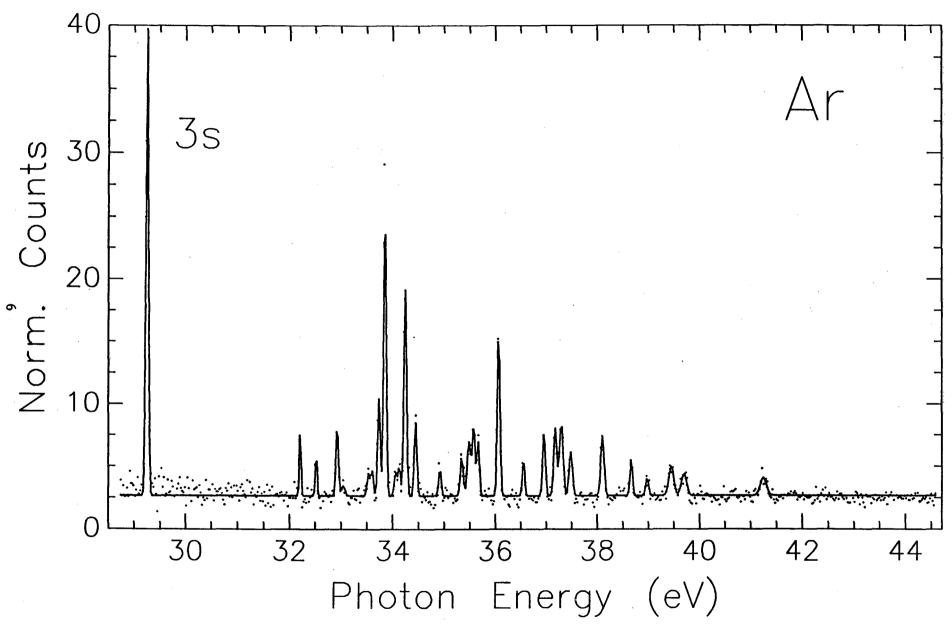
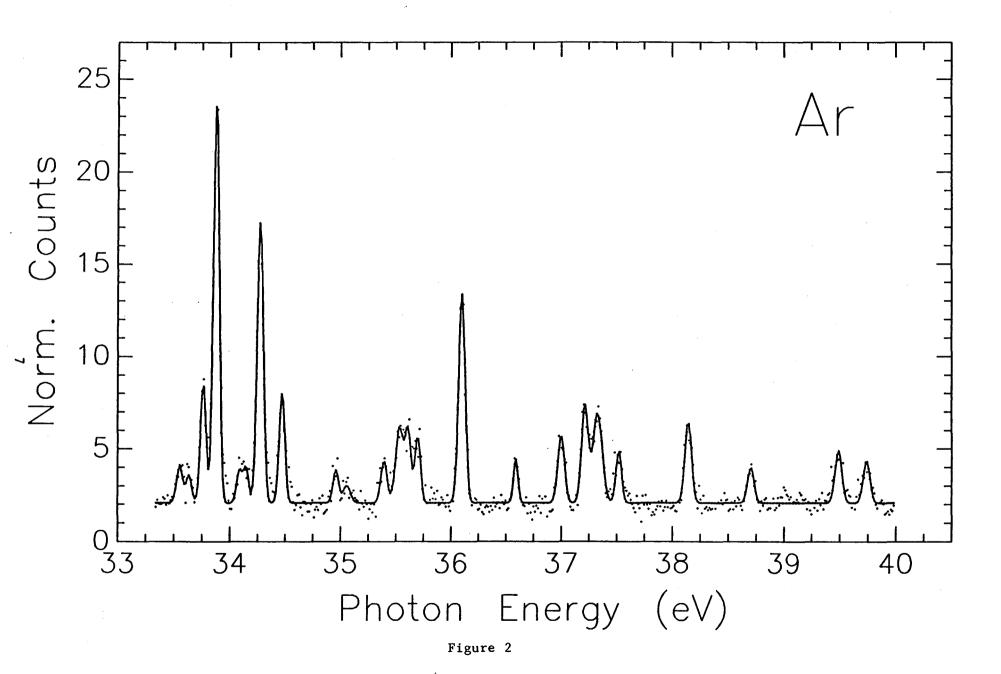


Figure 1



**3** 

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