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Search for Eccentric Binary Black Hole Mergers with Advanced LIGO and Advanced Virgo during Their First and Second Observing Runs

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Search for Eccentric Binary Black Hole Mergers with Advanced LIGO and Advanced Virgo  
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 M. KASPRZACK,<sup>1</sup> S. KATSANEVAS,<sup>29</sup> E. KATSAVOUNIDIS,<sup>14</sup> W. KATZMAN,<sup>7</sup> S. KAUFER,<sup>10</sup> K. KAWABE,<sup>47</sup> N. V. KEERTHANA,<sup>3</sup>  
 F. KÉFÉLIAN,<sup>140</sup> D. KEITEL,<sup>140</sup> R. KENNEDY,<sup>112</sup> J. S. KEY,<sup>148</sup> F. Y. KHALILI,<sup>61</sup> I. KHAN,<sup>16,32</sup> S. KHAN,<sup>9,10</sup>  
 E. A. KHAZANOV,<sup>149</sup> N. KHETAN,<sup>16,17</sup> M. KHURSHEED,<sup>60</sup> N. KIJBUNCHOO,<sup>8</sup> CHUNGLEE KIM,<sup>150</sup> J. C. KIM,<sup>151</sup> K. KIM,<sup>94</sup>  
 W. KIM,<sup>57</sup> W. S. KIM,<sup>152</sup> Y.-M. KIM,<sup>153</sup> C. KIMBALL,<sup>58</sup> P. J. KING,<sup>47</sup> M. KINLEY-HANLON,<sup>46</sup> R. KIRCHHOFF,<sup>9,10</sup>  
 J. S. KISSEL,<sup>47</sup> L. KLEYBOLTE,<sup>134</sup> J. H. KLIKA,<sup>24</sup> S. KLIMENKO,<sup>30</sup> T. D. KNOWLES,<sup>39</sup> P. KOCH,<sup>9,10</sup>  
 S. M. KOEHLLENBECK,<sup>9,10</sup> G. KOEKOEK,<sup>37,154</sup> S. KOLEY,<sup>37</sup> V. KONDRASHOV,<sup>1</sup> A. KONTOS,<sup>155</sup> N. KOPER,<sup>9,10</sup>  
 M. KOROBKO,<sup>134</sup> W. Z. KORTH,<sup>1</sup> M. KOVALAM,<sup>66</sup> D. B. KOZAK,<sup>1</sup> C. KRÄMER,<sup>9,10</sup> V. KRINGEL,<sup>9,10</sup> N. KRISHNENDU,<sup>31</sup>  
 A. KRÓLAK,<sup>156,157</sup> N. KRUPINSKI,<sup>24</sup> G. KUEHN,<sup>9,10</sup> A. KUMAR,<sup>144</sup> P. KUMAR,<sup>158</sup> RAHUL KUMAR,<sup>47</sup> RAKESH KUMAR,<sup>111</sup>  
 L. KUO,<sup>91</sup> A. KUTYNIA,<sup>156</sup> S. KWANG,<sup>24</sup> B. D. LACKEY,<sup>76</sup> D. LAGHI,<sup>20,21</sup> K. H. LAI,<sup>94</sup> T. L. LAM,<sup>94</sup> M. LANDRY,<sup>47</sup>  
 B. B. LANE,<sup>14</sup> R. N. LANG,<sup>159</sup> J. LANGE,<sup>62</sup> B. LANTZ,<sup>51</sup> R. K. LANZA,<sup>14</sup> A. LARTEUX-VOLLARD,<sup>28</sup> P. D. LASKY,<sup>6</sup>  
 M. LAXEN,<sup>7</sup> A. LAZZARINI,<sup>1</sup> C. LAZZARO,<sup>54</sup> P. LEACI,<sup>119,33</sup> S. LEAVEY,<sup>9,10</sup> Y. K. LECOEUICHE,<sup>47</sup> C. H. LEE,<sup>97</sup> H. K. LEE,<sup>160</sup>  
 H. M. LEE,<sup>161</sup> H. W. LEE,<sup>151</sup> J. LEE,<sup>96</sup> K. LEE,<sup>46</sup> J. LEHMANN,<sup>9,10</sup> A. K. LENON,<sup>39</sup> N. LEROY,<sup>28</sup> N. LETENDRE,<sup>34</sup>  
 Y. LEVIN,<sup>6</sup> A. LI,<sup>94</sup> J. LI,<sup>84</sup> K. J. L. LI,<sup>94</sup> T. G. F. LI,<sup>94</sup> X. LI,<sup>48</sup> F. LIN,<sup>6</sup> F. LINDE,<sup>162,37</sup> S. D. LINKER,<sup>133</sup>  
 T. B. LITTENBERG,<sup>163</sup> J. LIU,<sup>66</sup> X. LIU,<sup>24</sup> M. LLORENS-MONTEAGUDO,<sup>22</sup> R. K. L. LO,<sup>94,1</sup> L. T. LONDON,<sup>14</sup>  
 A. LONGO,<sup>164,165</sup> M. LORENZINI,<sup>16,17</sup> V. LORIETTE,<sup>166</sup> M. LORMAND,<sup>7</sup> G. LOSURDO,<sup>21</sup> J. D. LOUGH,<sup>9,10</sup> C. O. LOUSTO,<sup>62</sup>  
 G. LOVELACE,<sup>27</sup> M. E. LOWER,<sup>167</sup> H. LÜCK,<sup>10,9</sup> D. LUMACA,<sup>86,32</sup> A. P. LUNDGREN,<sup>140</sup> R. LYNCH,<sup>14</sup> Y. MA,<sup>48</sup> R. MACAS,<sup>106</sup>  
 S. MACFOY,<sup>25</sup> M. MACINNIS,<sup>14</sup> D. M. MACLEOD,<sup>106</sup> A. MACQUET,<sup>65</sup> I. MAGAÑA HERNANDEZ,<sup>24</sup> F. MAGAÑA-SANDOVAL,<sup>30</sup>  
 R. M. MAGEE,<sup>90</sup> E. MAJORANA,<sup>33</sup> I. MAKSIMOVIC,<sup>166</sup> A. MALIK,<sup>60</sup> N. MAN,<sup>65</sup> V. MANDIC,<sup>43</sup> V. MANGANO,<sup>46,119,33</sup>  
 G. L. MANSELL,<sup>47,14</sup> M. MANSKE,<sup>24</sup> M. MANTOVANI,<sup>29</sup> M. MAPELLI,<sup>53,54</sup> F. MARCHESONI,<sup>52,41</sup> F. MARION,<sup>34</sup> S. MÁRKA,<sup>105</sup>  
 Z. MÁRKA,<sup>105</sup> C. MARKAKIS,<sup>19</sup> A. S. MARKOSYAN,<sup>51</sup> A. MARKOWITZ,<sup>1</sup> E. MAROS,<sup>1</sup> A. MARQUINA,<sup>104</sup> S. MARSAT,<sup>26</sup>  
 F. MARTELLI,<sup>63,64</sup> I. W. MARTIN,<sup>46</sup> R. M. MARTIN,<sup>36</sup> V. MARTINEZ,<sup>79</sup> D. V. MARTYNOV,<sup>13</sup> H. MASALEHDAN,<sup>134</sup>  
 K. MASON,<sup>14</sup> E. MASSERA,<sup>112</sup> A. MASSEROT,<sup>34</sup> T. J. MASSINGER,<sup>1</sup> M. MASSO-REID,<sup>46</sup> S. MASTROGIOVANNI,<sup>26</sup> A. MATAS,<sup>76</sup>  
 F. MATICHARD,<sup>1,14</sup> L. MATONE,<sup>105</sup> N. MAVALVALA,<sup>14</sup> J. J. MCCANN,<sup>66</sup> R. MCCARTHY,<sup>47</sup> D. E. MCCLELLAND,<sup>8</sup>  
 S. MCCORMICK,<sup>7</sup> L. MCCULLER,<sup>14</sup> S. C. MCGUIRE,<sup>168</sup> C. MCISAAC,<sup>140</sup> J. MCIVER,<sup>1</sup> D. J. MCMANUS,<sup>8</sup> T. MCRRAE,<sup>8</sup>  
 S. T. MCWILLIAMS,<sup>39</sup> D. MEACHER,<sup>24</sup> G. D. MEADORS,<sup>6</sup> M. MEHMET,<sup>9,10</sup> A. K. MEHTA,<sup>18</sup> J. MEIDAM,<sup>37</sup>  
 E. MEJUTO VILLA,<sup>116,69</sup> A. MELATOS,<sup>100</sup> G. MENDELL,<sup>47</sup> R. A. MERCER,<sup>24</sup> L. MERENI,<sup>23</sup> K. MERFELD,<sup>72</sup> E. L. MERILH,<sup>47</sup>  
 M. MERZOUGUI,<sup>65</sup> S. MESHKOV,<sup>1</sup> C. MESSENGER,<sup>46</sup> C. MESSICK,<sup>90</sup> F. MESSINA,<sup>44,45</sup> R. METZDORFF,<sup>73</sup> P. M. MEYERS,<sup>100</sup>  
 F. MEYLAHN,<sup>9,10</sup> A. MIANI,<sup>117,118</sup> H. MIAO,<sup>13</sup> C. MICHEL,<sup>23</sup> H. MIDDLETON,<sup>100</sup> L. MILANO,<sup>80,5</sup> A. L. MILLER,<sup>30,119,33</sup>  
 M. MILLHOUSE,<sup>100</sup> J. C. MILLS,<sup>106</sup> M. C. MILOVICH-GOFF,<sup>133</sup> O. MINAZZOLI,<sup>65,169</sup> Y. MINENKOV,<sup>32</sup> A. MISHKIN,<sup>30</sup>  
 C. MISHRA,<sup>170</sup> T. MISTRY,<sup>112</sup> S. MITRA,<sup>3</sup> V. P. MITROFANOV,<sup>61</sup> G. MITSELMAKHER,<sup>30</sup> R. MITTLEMAN,<sup>14</sup> G. MO,<sup>98</sup>  
 D. MOFFA,<sup>121</sup> K. MOGUSHI,<sup>87</sup> S. R. P. MOHAPATRA,<sup>14</sup> M. MOLINA-RUIZ,<sup>141</sup> M. MONDIN,<sup>133</sup> M. MONTANI,<sup>63,64</sup>  
 C. J. MOORE,<sup>13</sup> D. MORARU,<sup>47</sup> F. MORAWSKI,<sup>56</sup> G. MORENO,<sup>47</sup> S. MORISAKI,<sup>83</sup> B. MOURS,<sup>34</sup> C. M. MOW-LOWRY,<sup>13</sup>  
 F. MUCIACCIA,<sup>119,33</sup> ARUNAVA MUKHERJEE,<sup>9,10</sup> D. MUKHERJEE,<sup>24</sup> S. MUKHERJEE,<sup>108</sup> SUBROTO MUKHERJEE,<sup>111</sup>  
 N. MUKUND,<sup>9,10,3</sup> A. MULLAVEY,<sup>7</sup> J. MUNCH,<sup>57</sup> E. A. MUÑOZ,<sup>42</sup> M. MURATORE,<sup>35</sup> P. G. MURRAY,<sup>46</sup> I. NARDECCHIA,<sup>86,32</sup>  
 L. NATICCHIONI,<sup>119,33</sup> R. K. NAYAK,<sup>171</sup> B. F. NEIL,<sup>66</sup> J. NELSON,<sup>116,69</sup> G. NELEMANS,<sup>67,37</sup> T. J. N. NELSON,<sup>7</sup> M. NERY,<sup>9,10</sup>  
 A. NEUNZERT,<sup>137</sup> L. NEVIN,<sup>1</sup> K. Y. NG,<sup>14</sup> S. NG,<sup>57</sup> C. NGUYEN,<sup>26</sup> P. NGUYEN,<sup>72</sup> D. NICHOLS,<sup>142,37</sup> S. A. NICHOLS,<sup>2</sup>  
 S. NISSANKE,<sup>142,37</sup> F. NOCERA,<sup>29</sup> C. NORTH,<sup>106</sup> L. K. NUTTALL,<sup>140</sup> M. OBERGAULINGER,<sup>22,172</sup> J. OBERLING,<sup>47</sup>  
 B. D. O'BRIEN,<sup>30</sup> G. OGANESYAN,<sup>16,17</sup> G. H. OGIN,<sup>173</sup> J. J. OH,<sup>152</sup> S. H. OH,<sup>152</sup> F. OHME,<sup>9,10</sup> H. OHTA,<sup>83</sup> M. A. OKADA,<sup>15</sup>  
 M. OLIVER,<sup>101</sup> P. OPPERMAN,<sup>9,10</sup> RICHARD J. ORAM,<sup>7</sup> B. O'REILLY,<sup>7</sup> R. G. ORMISTON,<sup>43</sup> L. F. ORTEGA,<sup>30</sup>  
 R. O'SHAUGHNESSY,<sup>62</sup> S. OSSOKINE,<sup>76</sup> D. J. OTTAWAY,<sup>57</sup> H. OVERMIER,<sup>7</sup> B. J. OWEN,<sup>85</sup> A. E. PACE,<sup>90</sup> G. PAGANO,<sup>20,21</sup>  
 M. A. PAGE,<sup>66</sup> G. PAGLIAROLI,<sup>16,17</sup> A. PAI,<sup>130</sup> S. A. PAI,<sup>60</sup> J. R. PALAMOS,<sup>72</sup> O. PALASHOV,<sup>149</sup> C. PALOMBA,<sup>33</sup> H. PAN,<sup>91</sup>

P. K. PANDA,<sup>144</sup> P. T. H. PANG,<sup>94,37</sup> C. PANKOW,<sup>58</sup> F. PANNARALE,<sup>119,33</sup> B. C. PANT,<sup>60</sup> F. PAOLETTI,<sup>21</sup> A. PAOLI,<sup>29</sup>  
A. PARIDA,<sup>3</sup> W. PARKER,<sup>7,168</sup> D. PASCUCCI,<sup>46,37</sup> A. PASQUALETTI,<sup>29</sup> R. PASSAQUIETI,<sup>20,21</sup> D. PASSUELLO,<sup>21</sup> M. PATIL,<sup>157</sup>  
B. PATRICELLI,<sup>20,21</sup> E. PAYNE,<sup>6</sup> B. L. PEARLSTONE,<sup>46</sup> T. C. PECHSIRI,<sup>30</sup> A. J. PEDERSEN,<sup>42</sup> M. PEDRAZA,<sup>1</sup>  
R. PEDURAND,<sup>23,174</sup> A. PELE,<sup>7</sup> S. PENN,<sup>175</sup> A. PEREGO,<sup>117,118</sup> C. J. PEREZ,<sup>47</sup> C. PÉRIGOIS,<sup>34</sup> A. PERRECA,<sup>117,118</sup>  
J. PETERMANN,<sup>134</sup> H. P. PFEIFFER,<sup>76</sup> M. PHELPS,<sup>9,10</sup> K. S. PHUKON,<sup>3</sup> O. J. PICCINNI,<sup>119,33</sup> M. PICHOT,<sup>65</sup>  
F. PIERGIOVANNI,<sup>63,64</sup> V. PIERRO,<sup>116,69</sup> G. PILLANT,<sup>29</sup> L. PINARD,<sup>23</sup> I. M. PINTO,<sup>116,69,89</sup> M. PIRELLO,<sup>47</sup> M. PITKIN,<sup>46</sup>  
W. PLASTINO,<sup>164,165</sup> R. POGGIANI,<sup>20,21</sup> D. Y. T. PONG,<sup>94</sup> S. PONRATHNAM,<sup>3</sup> P. POPOLIZIO,<sup>29</sup> E. K. PORTER,<sup>26</sup>  
J. POWELL,<sup>167</sup> A. K. PRAJAPATI,<sup>111</sup> J. PRASAD,<sup>3</sup> K. PRASAI,<sup>51</sup> R. PRASANNA,<sup>144</sup> G. PRATTEN,<sup>101</sup> T. PRESTEGARD,<sup>24</sup>  
M. PRINCIPE,<sup>116,89,69</sup> G. A. PRODI,<sup>117,118</sup> L. PROKHOROV,<sup>13</sup> M. PUNTURO,<sup>41</sup> P. PUPPO,<sup>33</sup> M. PÜRRE, <sup>76</sup> H. QI,<sup>106</sup>  
V. QUETSCHKE,<sup>108</sup> P. J. QUINONEZ,<sup>35</sup> F. J. RAAB,<sup>47</sup> G. RAAJMAKERS,<sup>142,37</sup> H. RADKINS,<sup>47</sup> N. RADULESCO,<sup>65</sup> P. RAFFAI,<sup>110</sup>  
S. RAJA,<sup>60</sup> C. RAJAN,<sup>60</sup> B. RAJBHANDARI,<sup>85</sup> M. RAKHMANOV,<sup>108</sup> K. E. RAMIREZ,<sup>108</sup> A. RAMOS-BUADES,<sup>101</sup> JAVED RANA,<sup>3</sup>  
K. RAO,<sup>58</sup> P. RAPAGNANI,<sup>119,33</sup> V. RAYMOND,<sup>106</sup> M. RAZZANO,<sup>20,21</sup> J. READ,<sup>27</sup> T. REGIMBAU,<sup>34</sup> L. REI,<sup>59</sup> S. REID,<sup>25</sup>  
D. H. REITZE,<sup>1,30</sup> P. RETTEGNO,<sup>127,176</sup> F. RICCI,<sup>119,33</sup> C. J. RICHARDSON,<sup>35</sup> J. W. RICHARDSON,<sup>1</sup> P. M. RICKER,<sup>19</sup>  
G. RIEMENSCHNEIDER,<sup>176,127</sup> K. RILES,<sup>137</sup> M. RIZZO,<sup>58</sup> N. A. ROBERTSON,<sup>1,46</sup> F. ROBINET,<sup>28</sup> A. ROCCHI,<sup>32</sup> L. ROLLAND,<sup>34</sup>  
J. G. ROLLINS,<sup>1</sup> V. J. ROMA,<sup>72</sup> M. ROMANELLI,<sup>71</sup> R. ROMANO,<sup>4,5</sup> C. L. ROMEL,<sup>47</sup> J. H. ROMIE,<sup>7</sup> C. A. ROSE,<sup>24</sup> D. ROSE,<sup>27</sup>  
K. ROSE,<sup>121</sup> D. ROSIŃSKA,<sup>74</sup> S. G. ROSOFSKY,<sup>19</sup> M. P. ROSS,<sup>177</sup> S. ROWAN,<sup>46</sup> A. RÜDIGER,<sup>9,10,\*</sup> P. RUGGI,<sup>29</sup> G. RUTINS,<sup>132</sup>  
K. RYAN,<sup>47</sup> S. SACHDEV,<sup>90</sup> T. SADECKI,<sup>47</sup> M. SAKELLARIADOU,<sup>146</sup> O. S. SALAFIA,<sup>178,44,45</sup> L. SALCONI,<sup>29</sup> M. SALEEM,<sup>31</sup>  
A. SAMAJDAR,<sup>37</sup> L. SAMMUT,<sup>6</sup> E. J. SANCHEZ,<sup>1</sup> L. E. SANCHEZ,<sup>1</sup> N. SANCHIS-GUAL,<sup>179</sup> J. R. SANDERS,<sup>180</sup>  
K. A. SANTIAGO,<sup>36</sup> E. SANTOS,<sup>65</sup> N. SARIN,<sup>6</sup> B. SASSOLAS,<sup>23</sup> O. SAUTER,<sup>137,34</sup> R. L. SAVAGE,<sup>47</sup> P. SCHALE,<sup>72</sup> M. SCHEEL,<sup>48</sup>  
J. SCHEUER,<sup>58</sup> P. SCHMIDT,<sup>13,67</sup> R. SCHNABEL,<sup>134</sup> R. M. S. SCHOFIELD,<sup>72</sup> A. SCHÖNBECK,<sup>134</sup> E. SCHREIBER,<sup>9,10</sup>  
B. W. SCHULTE,<sup>9,10</sup> B. F. SCHUTZ,<sup>106</sup> J. SCOTT,<sup>46</sup> S. M. SCOTT,<sup>8</sup> E. SEIDEL,<sup>19</sup> D. SELLERS,<sup>7</sup> A. S. SENGUPTA,<sup>181</sup>  
N. SENNETT,<sup>76</sup> D. SENTENAC,<sup>29</sup> V. SEQUINO,<sup>59</sup> A. SERGEEV,<sup>149</sup> Y. SETYAWATI,<sup>9,10</sup> D. A. SHADDOCK,<sup>8</sup> T. SHAFFER,<sup>47</sup>  
M. S. SHAHRIAR,<sup>58</sup> M. B. SHANER,<sup>133</sup> A. SHARMA,<sup>16,17</sup> P. SHARMA,<sup>60</sup> P. SHAWHAN,<sup>77</sup> H. SHEN,<sup>19</sup> R. SHINK,<sup>182</sup>  
D. H. SHOEMAKER,<sup>14</sup> D. M. SHOEMAKER,<sup>78</sup> K. SHUKLA,<sup>141</sup> S. SHYAMSUNAR,<sup>60</sup> K. SIELLEZ,<sup>78</sup> M. SIENIAWSKA,<sup>56</sup> D. SIGG,<sup>47</sup>  
L. P. SINGER,<sup>81</sup> D. SINGH,<sup>90</sup> N. SINGH,<sup>74</sup> A. SINGHAL,<sup>16,33</sup> A. M. SINTES,<sup>101</sup> S. SITMUKHAMBETOV,<sup>108</sup> V. SKLIRIS,<sup>106</sup>  
B. J. J. SLAGMOLEN,<sup>8</sup> T. J. SLAVEN-BLAIR,<sup>66</sup> J. R. SMITH,<sup>27</sup> R. J. E. SMITH,<sup>6</sup> S. SOMALA,<sup>183</sup> E. J. SON,<sup>152</sup> S. SONI,<sup>2</sup>  
B. SORAZU,<sup>46</sup> F. SORRENTINO,<sup>59</sup> T. SOURADEEP,<sup>3</sup> E. SOWELL,<sup>85</sup> A. P. SPENCER,<sup>46</sup> M. SPERA,<sup>53,54</sup> A. K. SRIVASTAVA,<sup>111</sup>  
V. SRIVASTAVA,<sup>42</sup> K. STAATS,<sup>58</sup> C. STACHIE,<sup>65</sup> M. STANDKE,<sup>9,10</sup> D. A. STEER,<sup>26</sup> M. STEINKE,<sup>9,10</sup> J. STEINLECHNER,<sup>134,46</sup>  
S. STEINLECHNER,<sup>134</sup> D. STEINMEYER,<sup>9,10</sup> S. P. STEVENSON,<sup>167</sup> D. STOCKS,<sup>51</sup> R. STONE,<sup>108</sup> D. J. STOPS,<sup>13</sup> K. A. STRAIN,<sup>46</sup>  
G. STRATTA,<sup>184,64</sup> S. E. STRIGIN,<sup>61</sup> A. STRUNK,<sup>47</sup> R. STURANI,<sup>185</sup> A. L. STUVER,<sup>186</sup> V. SUDHIR,<sup>14</sup> T. Z. SUMMERSCALES,<sup>187</sup>  
L. SUN,<sup>1</sup> S. SUNIL,<sup>111</sup> A. SUR,<sup>56</sup> J. SURESH,<sup>83</sup> P. J. SUTTON,<sup>106</sup> B. L. SWINKELS,<sup>37</sup> M. J. SZCZEPAŃCZYK,<sup>35</sup> M. TACCA,<sup>37</sup>  
S. C. TAIT,<sup>46</sup> C. TALBOT,<sup>6</sup> D. B. TANNER,<sup>30</sup> D. TAO,<sup>1</sup> M. TÁPAI,<sup>131</sup> A. TAPIA,<sup>27</sup> J. D. TASSON,<sup>98</sup> R. TAYLOR,<sup>1</sup>  
R. TENORIO,<sup>101</sup> L. TERKOWSKI,<sup>134</sup> M. THOMAS,<sup>7</sup> P. THOMAS,<sup>47</sup> S. R. THONDAPU,<sup>60</sup> K. A. THORNE,<sup>7</sup> E. THRANE,<sup>6</sup>  
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## ABSTRACT

When formed through dynamical interactions, stellar-mass binary black holes may retain eccentric orbits ( $e > 0.1$  at 10 Hz) detectable by ground-based gravitational-wave detectors. Eccentricity can therefore be used to differentiate dynamically-formed binaries from isolated binary black hole mergers. Current template-based gravitational-wave searches do not use waveform models associated to eccentric orbits, rendering the search less efficient to eccentric binary systems. Here we present results of a search for binary black hole mergers that inspiral in eccentric orbits using data from the first and second observing runs (O1 and O2) of Advanced LIGO and Advanced Virgo. We carried out the search with the coherent WaveBurst algorithm, which uses minimal assumptions on the signal morphology and does not rely on binary waveform templates. We show that it is sensitive to binary mergers with a detection range that is weakly dependent on eccentricity for all bound systems. Our search did not identify any new binary merger candidates. We interpret these results in light of eccentric binary formation models. We rule out formation channels with rates  $\gtrsim 100 \text{ Gpc}^{-3} \text{ yr}^{-1}$  for  $e > 0.1$  assuming a black hole mass spectrum with power-law index  $\lesssim 2$ .

*Keywords:* eccentric BBH, LIGO-Virgo

## 1. INTRODUCTION

In their first two observing runs, the Advanced LIGO and Advanced Virgo detectors discovered 10 binary black hole (BBH) mergers and a binary neutron star merger (Abbott et al. 2018b). These detections have already provided a wealth of information on cosmic processes, including the rate, mass, spin and redshift distribution of BBH mergers (Abbott et al. 2016f, 2018a), constraints on general relativity (Abbott et al. 2016e, 2019c), estimates of the Hubble constant (Abbott et al. 2017a; Soares-Santos et al. 2019; Abbott et al. 2019b), and constraints on multi-messenger emission from the mergers (Adrián-Martínez et al. 2016; Abbott et al.

2016d; Albert et al. 2017; Burns et al. 2018; Abbott et al. 2019a, 2008).

A key question that remains unanswered is how BBHs are formed. Viable formation channels include isolated binary evolution (e.g. Bethe & Brown 1998; Belczynski et al. 2002, 2014, 2016; Dominik et al. 2013; Menekens & Vanbeveren 2014; Spera et al. 2015; Eldridge & Stanway 2016; Marchant et al. 2016; Mandel & de Mink 2016; Mapelli et al. 2017; Mapelli & Giacobbo 2018; Stevenson et al. 2017; Giacobbo & Mapelli 2018; Kruckow et al. 2018; Barrett et al. 2018) and dynamical encounters in dense stellar environments, such as globular clusters (e.g. Portegies Zwart & McMillan 2000; O’Leary et al. 2006; Sadowski et al. 2008; Downing et al. 2010, 2011; Rodriguez et al. 2015, 2016a,b; Rodriguez & Loeb 2018; Askar et al. 2017; Samsing 2018; Samsing

\* Deceased, July 2018.



et al. 2018; Fragione & Kocsis 2018; Zevin et al. 2019), young star clusters (e.g. Banerjee et al. 2010; Ziosi et al. 2014; Mapelli 2016; Banerjee 2017, 2018; Di Carlo et al. 2019; Kumamoto et al. 2018) and galactic nuclei (e.g. O’Leary et al. 2009; Antonini & Perets 2012; Antonini & Rasio 2016; Petrovich & Antonini 2017; Stone et al. 2017b,a; Rasskazov & Kocsis 2019). Moreover, the dynamical process known as Kozai–Lidov (KL) resonance (Kozai 1962; Lidov 1962) can trigger the merger of a BBH, even if the BBH has not been formed in a dense star cluster. In fact, if the BBH is orbited by a tertiary body (i.e. the BBH is the inner binary of a stable hierarchical triple system), the KL mechanism triggers oscillations of the BBH’s eccentricity, which might speed up the merger by gravitational-wave emission. Each channel is expected to produce black hole mergers with different mass and spin distributions (Abbott et al. 2018a; Mandel & O’Shaughnessy 2010; Rodriguez et al. 2016c; Abbott et al. 2016a; Farr et al. 2017). The limited statistics from the low number of systems detected through gravitational waves and model uncertainties so far do not allow strong constraints on the formation channels.

Orbital eccentricity is a distinguishing feature of dynamical formation channels. Gravitational-wave emission acts to circularize binary orbits by the time they reach orbital frequencies to which Advanced LIGO and Advanced Virgo are sensitive ( $\gtrsim 10$  Hz). Eccentric orbits in the Advanced LIGO and Advanced Virgo band indicate either that the binary was formed with small orbital separation and therefore did not have time to circularize, or that some dynamical process increased the eccentricity. For example, KL-induced mergers are expected to be associated with high eccentricities (Antonini et al. 2017; Fragione & Bromberg 2019; Fragione et al. 2019a; Fragione & Kocsis 2019). The detection of gravitational waves from an eccentric binary would suggest that binary systems can form dynamically, and could help distinguish between different dynamical formation scenarios (KL oscillations in triple systems or dynamical encounters in dense stellar clusters) (Lower et al. 2018).

In the following we define eccentricity at the time when the gravitational-wave frequency of the binary is at 10 Hz (Peters & Mathews 1963). Eccentricity constantly evolves during the inspiral.

Template-based gravitational-wave searches used by Advanced LIGO and Advanced Virgo currently do not include eccentric orbital templates (Abbott et al. 2018b). Quasi-circular waveform templates are able to detect binaries with small eccentricities ( $e \lesssim 0.1$ ), but are inefficient at extracting moderately to highly eccentric binaries (Brown & Zimmerman 2010). Multiple

efforts for generating the full inspiral merger ringdown (IMR) waveforms for the binaries with eccentric orbits are underway (Cao & Han 2017; Hinder et al. 2018; Ireland et al. 2019; Huerta et al. 2018; Hinderer & Babak 2017). However, the lack of a reliable and complete waveform model prevents the implementation of a matched-filtering search at this time, and led to the development of alternative search methods (Tiwari et al. 2016; Coughlin et al. 2015; Lower et al. 2018).

Here we report the results of a search for eccentric binary black hole mergers with the coherent WaveBurst (cWB) algorithm that does not rely on binary system waveforms. cWB is sensitive to binaries of any eccentricity, and in particular to high-mass black holes. The search has been carried out over data from Advanced LIGO and Advanced Virgo’s O1 and O2 observing runs, and found no evidence of eccentric binary signals. This paper evaluates the sensitivity of cWB to eccentric binary mergers, and infers constraints from non-detection on the rate of eccentric mergers.

## 2. DETECTORS AND ANALYSIS METHOD

### 2.1. *Advanced LIGO and Advanced Virgo*

The Advanced LIGO detectors began their first observing run O1 on September 12, 2015, which lasted until January 19, 2016 (Abbott et al. 2016b). During this time they accumulated  $T_{\text{obs},1} = 48$  days of coincident data during which both LIGO Hanford and LIGO Livingston detectors were operating. The second observing run O2 started on November 30, 2016 and lasted until August 25, 2017, resulting in  $T_{\text{obs},2} = 118$  days of coincident data (Abbott et al. 2018b). Advanced Virgo joined the Advanced LIGO detectors on August 1, 2017. The detectors’ sensitivity was not uniform during these runs, and there was a marked sensitivity increase from O1 to O2 (Abbott et al. 2018c). As adding Advanced Virgo data was not improving the sensitivity of the search, this analysis only uses data from the Advanced LIGO detectors.

### 2.2. *Search description*

The search for eccentric binary black hole mergers uses the same configuration of the cWB pipeline (Klimenko et al. 2008, 2016) as the binary black hole merger search reported in Abbott et al. (2018b). An early version of the search is described in (Tiwari et al. 2016). cWB is designed to search for transient signals, without specifying a waveform model. It identifies coherent excess power in multi-resolution time-frequency representations of the detectors’ strain data, for signal frequencies up to 1 kHz and duration up to a few seconds. The excess power is collected in the time-frequency plane

assuming monotonically increasing frequency for better collection of the signal energy from binary black holes. The search identifies events that are coherent in multiple detectors and reconstructs the source sky location and signal waveforms by using the constrained maximum likelihood method.

The cWB detection statistic  $\rho$  is based on the coherent energy  $E_c$  obtained by cross-correlating the signal waveforms reconstructed in the network of detectors. It is proportional to the coherent network signal-to-noise ratio. The estimation of statistical significance of an event is done by ranking the  $\rho$  of the event against the  $\rho$  distribution for background events obtained by repeating the analysis on time-shifted data. To exclude astrophysical events from the background sample, the time shifts are much larger than the expected signal delay between the detectors. Each cWB event was assigned a False Alarm Rate (FAR) based on the rate of background triggers with  $\rho$  higher than that of the event.

To increase the robustness against non stationary detector noise generating glitches, cWB uses signal-independent vetoes: the network correlation  $c_c = E_c/(E_c + E_n)$ , where  $E_n$  is the residual noise energy estimated after the reconstructed signal pixels are subtracted from the data. For a gravitational-wave signal we expected  $c_c \approx 1$  while for glitches  $c_c \ll 1$ . Events with  $c_c < 0.7$  are rejected.

Detector characterization studies are also carried out to ensure that candidate events are not due to instrumental or environmental artifacts. We have rejected the times where significant instrumental artifacts make the data unusable (Abbott et al. 2016c).

### 2.3. Simulated astrophysical signals

In order to estimate the sensitivity of our search, we simulated eccentric BBH signals, injected them into detector data and searched for them using cWB. We used a BH mass range of  $5 M_\odot - 50 M_\odot$  (Abbott et al. 2018a), and eccentricities in the  $e \in [0, 0.99]$  range. We assumed that BHs have zero spin. These simulations were carried out to quantify the search sensitivity for individual binaries. Below we considered different mass distributions to characterize our sensitivity.

At the time of the analysis only one set of templates was available for the generation of full inspiral-merger-ringdown eccentric binary waveforms including generic spin configurations by East et al. (2013). It uses a prescription based on the equations of motion of a geodesic in a Kerr spacetime, coupled with the quadrupole formula for the gravitational radiation. The model defines an effective Kerr spacetime whose mass and spin parameters are set equal to the total mass and orbital angular

momentum of the binary. The binary is evolved based on the behavior of a timelike geodesic in the effective Kerr spacetime, but the mass and angular momentum of this spacetime are changed at each time step based on the emitted energy and angular momentum calculated in the quadrupole approximation. This approach reproduces the correct orbital dynamics in the Newtonian limit and general-relativistic test particle limit. This model also incorporates strong-field features such as pericenter precession, frame dragging, and the existence of unstable orbits and related zoom-whirl dynamics (East et al. 2013). The inspiral waveforms obtained using the above treatment are stitched to a merger model that was developed for quasicircular mergers but also performs well for eccentric mergers with little modification (Baker et al. 2008; Kelly et al. 2011). In Fig. 1 we show this waveform for the case of circular and eccentric ( $e = 0.5$ ) orbits.

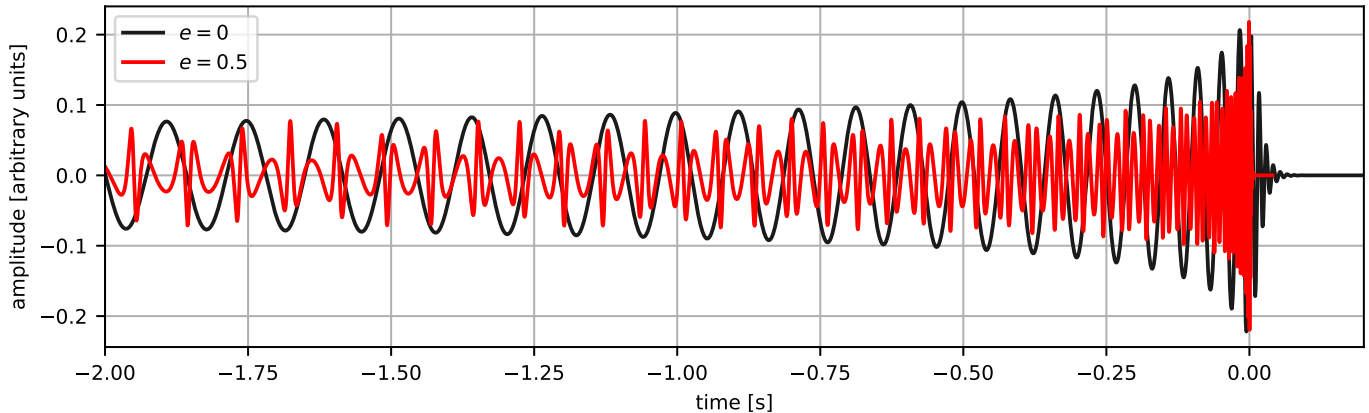
The waveforms we used here to simulate gravitational waves from eccentric binaries are approximate. Compared to gravitational waveforms obtained using general-relativistic numerical simulations, the waveforms differ in overlap by up to a few percent (East et al. 2013). However, we found that cWB can detect the waveforms used here with equal sensitivity to precise waveforms obtained using numerical simulations, making their use appropriate to characterize search sensitivity. This is due to the fact that cWB does not rely on the precise gravitational waveform, making a few percent difference negligible.

## 3. RESULTS

This search has detected 7 of the 10 BBH events that were identified by template based searches (its sensitivity compared to template based searches is higher for higher mass binaries; see Table 1 in Abbott et al. 2018b). We considered these events to have no eccentricity. Our search did not detect any gravitational-wave event beyond these. Therefore, we concluded that no eccentric BBH merger has been detected. Below we present our search sensitivity to interpret this non-detection.

We note that the detection of only 7 out of 10 BBH events by cWB is consistent with its relative sensitivity compared to template-based searches, which are more sensitive in particular at low black hole masses (see Fig. 2). Our interpretation below does not depend on, or make use, of this fractional overlap between the two search types.

We further note that the detection by cWB and the less-confident template-based detection of an event would not necessarily mean that the event was an eccentric binary (see, e.g., Abbott et al. 2018b). The ec-



**Figure 1.** Examples of gravitational waveforms for a  $10 M_{\odot} - 10 M_{\odot}$  BBH system with eccentricities 0 (black) and 0.5 (red).

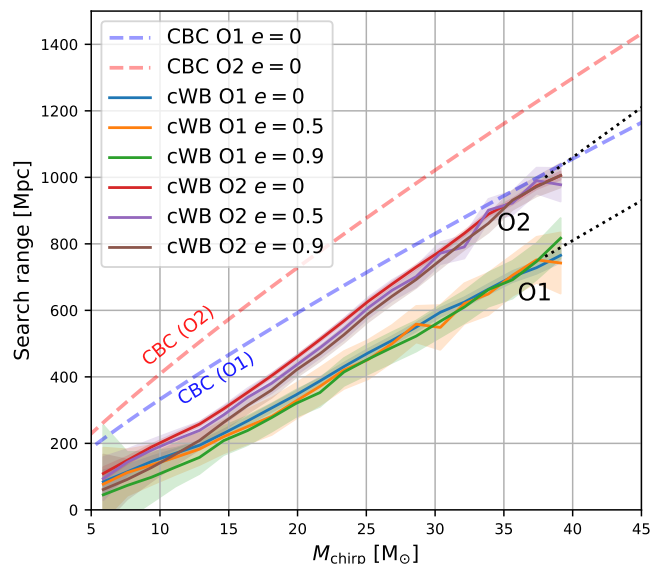
centricity of a detected event would need to be independently measured (e.g., Lower et al. 2018).

### 3.1. Sensitivity to eccentric mergers

We characterized the sensitivity of our search by calculating its *range*—the distance, averaged over observation time, sky location and orientation, within which a BBH can be detected with false alarm rate  $\leq 10^{-2} \text{ yr}^{-1}$ . For this calculation we adopted a standard cosmological model with Hubble parameter  $H_0 = 67.9 \text{ km s}^{-1} \text{ Mpc}^{-1}$  and  $\Omega_M = 0.3065$  (Ade et al. 2016). The range depends on the black hole masses, and is different for the O1 and O2 observing runs. In particular it depends on the chirp mass  $\mathcal{M}$  of the binary, where  $\mathcal{M} \equiv (m_1 m_2)^{3/5} (m_1 + m_2)^{-1/5}$  for black hole masses  $m_1$  and  $m_2$ . We find that cWB range is independent of the eccentricity for the whole mass range considered (see also Tiwari et al. 2016). Our ranges, using the eccentric waveforms described in Section 2.3, are shown in Fig. 2. We additionally see that the sensitive range of cWB grows faster with chirp mass than the range of template-based searches, making cWB additionally useful for circular binaries at higher masses (see also Abbott et al. 2017b).

### 3.2. Astrophysical constraints

In order to compare our results to astrophysical source populations, we calculated the volume-time (VT) probed by our search. VT depends on the mass distribution of the BBH population. Dynamical formation channels are expected to result in different BH mass and mass ratio distributions than BBH mergers from field binaries (Rodriguez et al. 2018; Kimpson et al. 2016). We considered a BBH mass distribution such that the mass of the more massive BH,  $m_1$ , follows a power-law distribution  $m^{-\beta}$  within the range  $[5 M_{\odot}, 50 M_{\odot}]$  for different  $\beta$  values (see below), while the second BH’s mass,  $m_2$  is uniformly distributed within the range  $[5 M_{\odot}, m_1]$ .



**Figure 2.** Range of the cWB analysis to BBH mergers as a function of the binary’s chirp mass, separately for the O1 and O2 observing runs, and for different orbital eccentricities (see legend). The shaded regions represent  $1\sigma$  uncertainties. The dotted lines are linear fits on the ranges at chirp masses  $> 30 M_{\odot}$  for  $e = 0$ . For comparison, we show the sensitive ranges for template-based searches for compact binary coalescence (CBC), assuming  $e = 0$ , for O1 and O2 (Abbott et al. 2018c). Masses are given in source-frame.

The mass distribution of BBH mergers detected by Advanced LIGO and Advanced Virgo so far is somewhat different from this assumed distribution (Abbott et al. 2018b,a). However, eccentric BBH merger channels are likely responsible for only a subset of these observations and therefore they do not fully determine the overall spectrum. With this mass distribution model, we find that  $\text{VT}(\beta) \approx \{6.6, 2.4, 0.75\} \times 10^{-2} \text{ Gpc}^3 \text{ yr}$  for  $\beta = \{1, 2, 3\}$ , respectively.

The BBH merger rate density for processes that can lead to eccentric orbits in the Advanced LIGO and Advanced Virgo band is mostly predicted to be up to a few  $\text{Gpc}^{-3}\text{yr}^{-1}$  (Antonini & Rasio 2016; Rodriguez et al. 2016b; Silsbee & Tremaine 2017; Bartos et al. 2017; Petrovich & Antonini 2017; Hoang et al. 2018; Hamers et al. 2018; Yang et al. 2019; Kocsis et al. 2006), while some more extreme models predict merger rate densities up to  $100\text{Gpc}^{-3}\text{yr}^{-1}$  (VanLandingham et al. 2016; Rasskazov & Kocsis 2019). The fraction of mergers from these processes that have high eccentricity ( $e \gtrsim 0.1$ ) ranges from  $\sim 1\%$  (Randall & Xianyu 2018b,a) to close to all mergers (Petrovich & Antonini 2017; Gondán et al. 2018).

In order to understand the astrophysical rate density constraints of our results, we considered a dynamical formation channel that produces BBH mergers at rate density  $R_{\text{dyn}}$ , with a mass power-law index  $\beta$  (see above). We assumed that a fraction  $f_{\text{ecc}}$  of mergers from this channel have eccentricities  $e > 0.1$ , and that this BBH sub-population follows the mass distribution considered here. We further assumed that all BBH mergers detected by Advanced LIGO and Advanced Virgo so far have eccentricities  $e < 0.1$ . The expected number of eccentric mergers ( $e > 0.1$ ) from this model detected by cWB during O1 and O2 is then

$$\langle N_{\text{cWB,ecc}} \rangle = R_{\text{dyn}} f_{\text{ecc}} \text{VT}(\beta). \quad (1)$$

Given that no such eccentric merger was detected, the Neyman 90% confidence-level upper limit is

$$R_{\text{ul,ecc}} = \frac{2.3}{f_{\text{ecc}} \text{VT}(\beta)}. \quad (2)$$

where  $\beta = \{1, 2, 3\}$ . We obtained

$$R_{\text{ul,ecc}} = \{30, 90, 300\} f_{\text{ecc}}^{-1} \text{Gpc}^{-3} \text{yr}^{-1} \quad (3)$$

for  $\beta = \{1, 2, 3\}$ , respectively. The quoted approximate values were rounded to the first significant digit. We found that this result does not depend on the eccentricity distribution of the source population as our search sensitivity only weakly depends on eccentricity.

Our results rule out models predicting  $\gtrsim 100\text{Gpc}^{-3}\text{yr}^{-1}$  merger rate densities (VanLandingham et al. 2016; Rasskazov & Kocsis 2019) for  $\beta \lesssim 2$  if the majority of mergers in the given model have eccentricities  $e > 0.1$ , while the results are consistent with a number of other models (e.g., Antonini et al. 2017; Fragione et al. 2019b; Antonini & Perets 2012).

#### 4. CONCLUSION

We searched for eccentric BBH mergers using the cWB algorithm. We showed that the sensitivity of our method

is independent of the eccentricity at the time the binary enters Advanced LIGO and Advanced Virgo's frequency band at  $\sim 10\text{Hz}$ .

Our search only uncovered binaries that have also been found by template-based searches that do not appear to have eccentric orbits. We interpreted this non-detection in light of the expected merger rate density of BBH formation channels that can produce eccentric orbits, and the fraction of these mergers that have eccentricities  $\gtrsim 0.1$ . Our results rule out the highest end of the rate density predictions ( $\gtrsim 100\text{Gpc}^{-3}\text{yr}^{-1}$ ) assuming that the majority of the binaries from these channels have  $e > 0.1$ , and that the power-law index of the BH mass spectrum is  $\lesssim 2$ .

Future observing runs by Advanced LIGO, Advanced Virgo and KAGRA (Aso et al. 2013) will provide substantially improved sensitivity to probe formation mechanisms resulting in eccentric binaries (Abbott et al. 2018c).

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