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#### Injection, Flow, and Mixing of CO2 in Porous Media with Residual Gas

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#### Abstract

Geologic structures associated with depleted natural gas reservoirs are desirable targets for geologic carbon sequestration (GCS) as evidenced by numerous pilot and industrial-scale GCS projects in these environments world-wide. One feature of these GCS targets that may affect injection is the presence of residual CH<sub>4</sub>. It is well known that CH<sub>4</sub> drastically alters supercritical CO<sub>2</sub> density and viscosity. Furthermore, residual gas of any kind affects the relative permeability of the liquid and gas phases, with relative permeability of the gas phase strongly dependent on the time-history of imbibition or drainage, i.e., dependent on hysteretic relative permeability. In this study, the effects of residual CH<sub>4</sub> on supercritical CO<sub>2</sub> injection were investigated by numerical simulation in an idealized one-dimensional system under three scenarios: (1) with no residual gas; (2) with residual supercritical CO<sub>2</sub>; and (3) with residual CH<sub>4</sub>. We further compare results of simulations that use non-hysteretic and hysteretic relative permeability functions. The primary effect of residual gas is to decrease injectivity by decreasing liquid-phase relative permeability. Secondary effects arise from injected gas effectively incorporating residual gas and thereby extending the mobile-gas plume relative to cases with no residual gas. Third-order effects arise from gas mixing and associated compositional effects on density that effectively create a larger plume per unit mass. Non-hysteretic models of relative permeability can be used to approximate some parts of the behavior of the system, but fully hysteretic formulations are needed to accurately model the entire system.

#### Introduction

Several geologic carbon sequestration (GCS) projects around the world involve  $CO_2$  injection into depleted gas reservoirs. In some cases, injection is directly into the depleted reservoir (e.g., K12-B, van der Meer et al., 2007), while in others it is into the water leg of the reservoir (e.g., Otway Basin Pilot Project (Sharma et al., 2007); In Salah (Ringrose et al., 2009)). Regardless of the details of individual projects, injection into depleted reservoirs may involve injection into porous media that are filled predominantly with brine and residual methane (CH<sub>4</sub>) gas.

There are numerous potential effects of residual  $CH_4$  gas on supercritical  $CO_2$  injection. First, residual gas of any kind will decrease the mobility of the brine, all other things being equal. Second, residual  $CH_4$  provides an initial gas-phase saturation that can be mobilized and incorporated into the

injected plume, thereby increasing its size. Third, the mixing of  $CH_4$  into supercritical  $CO_2$  causes a large decrease in gas mixture density and viscosity that can affect the injectivity and mobility of the gas. In this study, the effects of residual gas on injection of supercritical  $CO_2$  are investigated through modeling of an idealized one-dimensional radial system using TOUGH2/EOS7C (Oldenburg et al., 2004a) enhanced with hysteretic capillary pressure and relative permeability curves (Doughty, 2007). To evaluate and compare results, we refer to the concept of injectivity, defined loosely here as the pressure rise at the well for a given mass injection rate.

# Motivation

# **Prior Work**

There are numerous papers on the concept and evaluation of general feasibility of  $CO_2$  injection into natural gas reservoirs for GCS and for enhanced gas recovery (Blok et al., 1992; Koide et al., 1992; van der Burght, 1992; Oldenburg et al., 2001; Oldenburg et al., 2004). In addition, studies of thermal effects of injecting high-pressure  $CO_2$  into low-pressure depleted reservoirs have been made (Oldenburg, 2007; Maloney and Briceno, 2009; Mathias et al., 2010). The drastic changes in  $CO_2$  properties arising from mixing with  $CH_4$  are well known (e.g., Oldenburg et al., 2004b) and the associated potential beneficial uses have been noted (Oldenburg, 2007). None of these prior studies has investigated the effects of residual gas saturation in general, or the compositional effects associated with mixing supercritical  $CO_2$  and residual  $CH_4$  gas in particular. Furthermore, studies that compare hysteretic and non-hysteretic capillary pressure and relative permeability curves have focused on saline formations and  $CO_2$  trapping (Doughty, 2007) rather than on depleted gas-reservoirs and  $CO_2$  injectivity.

# Effect of Residual Gas on Injection

Residual gas may affect an injection process in many different ways, e.g., inhibiting the displacement of brine through reduced relative permeability of brine, enhancing injection of gas as residual gas becomes mobile at the leading edge of the injected gas plume, or changing the composition of the gas and therefore gas-phase properties as injected gas mixes with residual gas of different composition. The purpose of this study is to investigate the processes of interaction between injected supercritical  $CO_2$  and residual gas to understand potential implications for gas-reservoir injection projects in general. A simplified one-dimensional radial geometry is used to focus on mobility and gas composition effects. The idealized system was chosen to have a constant-pressure boundary condition at a radius of approximately 1 km to model injection into the water leg of a depleted gas-reservoir system with a nearby gas cap that would tend to moderate pressure rise due to injection.

Three different situations are shown in Figure 1 to illustrate the question being addressed. The base case is the injection of  $CO_2$  into a system with no residual gas (Figure 1a). Figure 1b shows the case of injection of supercritical  $CO_2$  into an aquifer with supercritical  $CO_2$  at residual saturation. Lastly, Figure 1c shows the case of injection of supercritical  $CO_2$  into a system with  $CH_4$  at residual saturation. The specific question being addressed in this study is: what is the effect of residual gas and its composition on the injection of supercritical  $CO_2$ ?



Figure 1. Three alternative cases for injection of  $CO_2$ . (a) zero residual gas. (b) 20% residual gas consisting of  $CO_2$  (white stipple). (c) 20% residual gas consisting of  $CH_4$  (dark stipple).

The factors that make this problem potentially interesting are illustrated in Figures 2 and 3. Figure 2 shows a plot of density and viscosity for gas-mixture compositions between pure  $CO_2$  and pure  $CH_4$  as a function of pressure at reservoir temperatures of 40 °C (left-hand side) and 90 °C (right-hand side). Density in the figures is calculated using WebGasEOS (Reagan and Oldenburg, 2006), a publicly available web-based tool for calculating gas-mixture properties, which implements the Peng-Robinson equation of state for density, and the method of Chung et al. (1988) for viscosity. As shown, the properties of the mixture vary strongly with composition at all pressures and particularly at pressures above the critical pressure of  $CO_2$  (7.4 MPa). The reason for this is that pure  $CO_2$  compresses readily into its supercritical form, but this compressibility is significantly diminished by admixed components such as  $CH_4$ , creating the strong compositional dependence on density. One can imagine a scenario in which injected supercritical  $CO_2$  encounters residual  $CH_4$  resulting in a drastic expansion of the gas mixture.

Figures 3a and 3b show typical non-hysteretic curves (van Genuchten (1980) for capillary pressure and liquid relative permeability, and Corey (1954) for gas relative permeability (see Table 1)). As shown, gas relative permeability increases with gas saturation (liquid saturation decreasing), while water relative permeability decreases with increasing gas saturation. The presence of residual gas thus causes the liquid relative permeability to be lower, all other things being equal. Because the injection of large quantities of  $CO_2$  for GCS requires both the injected  $CO_2$  and the native aqueous fluids (brine) to move, the permeability of the formation to both the injected gas and the brine influences injection pressure, along with the size and character of the injected  $CO_2$  plume. Figures 3c and 3d show hysteretic curves using parameters in Table 1, for several possible turning points (i.e., the saturation at which flow process changes from drainage to imbibition or vice versa). As shown, drainage and imbibition branches are distinct from one another, a feature that models the dependence of capillary pressure and relative permeability on the flow process (drainage or imbibition) that is occurring as well as on the saturation history (Doughty, 2007). Moreover, imbibition branches with different turning points have different values of residual gas saturation  $S_{gr}$ . In Figures 3c and 3d, the drainage branches (red) and the imbibition branches with turning points at  $S_{lr}$  (blue) form an envelope in which other imbibition branches lie. Values of  $S_{gr}$  range from 0 for the drainage branch to  $S_{grmax} = 0.2$  for the imbibition branch with turning point at  $S_{lr}$ .

Prior to injection of  $CO_2$  into depleted gas reservoirs, we envision a history of imbibition during the gas-production phase, i.e., as gas is produced, water flows into the formerly gas-filled regions of the reservoir, entrapping residual gas. By this argument, our simulations of  $CO_2$  injection into the water leg of a depleted reservoir should begin on the imbibition branch of the hysteretic curves with turning point  $S_{lr}$ . However, the injection of the gas (CO<sub>2</sub>) itself is a drainage process. In short, the part of the domain near the well experiencing  $CO_2$  injection (gas saturation increasing) has capillary pressure defined by the second-order drainage branch of the capillary pressure curve (Figure 3c), while the region where brine is being displaced with unchanged gas saturation remains on the imbibition branch.

Our hysteretic formulation (Doughty, 2009) for relative permeability does not distinguish between first-order imbibition and second-order drainage, as indicated by the two-way arrows on Figure 3d. This approximation has worked well for previous problems involving CO<sub>2</sub> injection into saline formations. However, for the present problem, this approximation will not adequately represent the second-order drainage accompanying injection of CO<sub>2</sub> into residual gas. Therefore, being guided by the form of the capillary pressure curve, in which the second-order drainage branch rapidly approaches the primary drainage branch as  $S_g$  increases (green dashed line in Figure 3c), for relative permeability we approximate the second-order drainage branch by the first-order imbibition branch with turning point at  $S_l = 0.6$  (green dashed line in Figure 3d), but use a slightly increased value of  $S_{grmax} = 0.265$ , which makes  $S_{gr} = 0.2$  and assures that the initial conditions represent immobile gas. Although admittedly ad hoc, this approach captures the essential features of the problem: it preserves the zero gas relative permeability far from the injection well while enabling a rapidly-increasing gas relative permeability near the injection well. Of course, it would be preferable to incorporate the second-order drainage branches directly in the relative permeability functions, and such an effort is planned as part of continuing code development.

This study is aimed at understanding the role of residual gas, from both the phase-interference and compositional perspectives, in controlling  $CO_2$  injection processes. Fluid flow problems in which drainage and imbibition occur in different regions at different times cannot be fully captured using non-hysteretic curves, but it is of interest to see if any relevant features of the problem can be represented using the simpler non-hysteretic forms.



Figure 2. Variation of gas density and viscosity as a function of  $CO_2$ -CH<sub>4</sub> mole fraction and pressure at (a) 40 °C and (b) 90 °C.



Figure 3. Non-hysteretic variation of (a) capillary pressure and (b) gas and liquid relative permeability as a function of liquid saturation using functions given in Table 1 and two different  $S_{gr}$  values. Hysteretic variation of (c) capillary pressure and (d) gas and liquid relative permeability using functions given in Table 1, for several possible turning points (dots). The drainage branches (red) and the imbibition branches with turning points at  $S_{lr}$  (blue) form an envelope in which other imbibition branches lie. Arrows identify whether drainage ( $S_g$  increasing) or imbibition ( $S_l$  increasing) occurs on a given branch.

#### Methods

Injection simulations were carried out using TOUGH2/EOS7C (Pruess et al., 1999; Oldenburg et al., 2004) to accurately model multiphase flow and multicomponent  $CO_2$  and  $CH_4$  gas-mixture properties. TOUGH2/EOS7C models five components (water, brine,  $CO_2$ , a tracer, and  $CH_4$ ). TOUGH2/EOS7C uses the Peng-Robinson equation of state to calculate gas-mixture density, and the method of Chung et al. (1988) for calculating mixture viscosity. EOS7C uses a fugacity equilibrium approach for gas-mixture solubility that is well-suited for deep reservoir environments (Oldenburg et al., 2004). Salinity effects on gas solubility in the aqueous phase are not modeled, and we have used pure water (zero salinity) as the aqueous phase in the results shown below. To investigate the phase

interference effects of the presence of residual gas on injection in a simplified geometry, we have used both hysteretic and non-hysteretic models as implemented in iTOUGH2/EOS7C (Finsterle, 2004, 2008; Doughty, 2007, 2009).

For this problem, a one-dimensional radial grid was used to emphasize the residual-gas effects being investigated without the complications added by a more complicated geometry. The discretization of the one-dimensional radial problem is shown in Figure 4, emphasizing the details of the fine discretization around the well. Beyond the near-well region, grid spacing is uniform (2 m) out to a radial distance of 300 m (the maximum extent of CO<sub>2</sub> during the 2-year simulation period), beyond which the grid gradually coarsens. Preliminary studies using a grid that gradually coarsened over all r showed that fronts can be significantly smeared out by numerical dispersion. The constantpressure boundary condition at r = 1074 m was chosen to model the presence of a gas cap somewhere in the system that would tend to moderate pressure changes during water-leg  $CO_2$ injection. The injection rate is 100 t  $CO_2/day$  (1.16 kg/s) into a 30-m thick layer with permeability of 500 mD ( $5 \times 10^{-13}$  m<sup>2</sup>) and porosity equal to 0.20. Two different reservoir scenarios are considered: (1) a 1 km-deep reservoir with initial pressure  $P = 1 \times 10^7$  Pa, T = 40 °C; and (2) a 2 kmdeep system with initial pressure  $P = 2 \times 10^7$  Pa, T = 90 °C. All simulations are isothermal. The residual gas saturation  $S_{gr}$  is set to 0.20 or 0.01 for the non-hysteretic case, to mimic various aspects of the hysteretic case, in which  $S_{gr} = 0$  during drainage and  $S_{gr}$  varies from 0 to  $S_{grmax} = 0.2$  during imbibition. See Table 1 for specific values of parameters, and Figure 3 for the capillary pressure and relative permeability curves used in the simulations. Note that a small amount of CH<sub>4</sub> is co-injected in all the simulations for an efficient numerical solution, but it does not materially affect the results.



Figure 4. Discretization of the one-dimensional radial problem showing detail around the well at r = 0 m. All boundaries are closed except the right-hand side which is held at constant pressure.

Thickness	30 m
Radius	1 km (open right-hand side boundary, see Fig. 4)
Porosity $(\phi)$	0.2
Pressure (Pa)	Shallow reservoir case: $1 \times 10^7$ Pa
	Deep reservoir case: $2 \times 10^7$ Pa
Temperature $(T)$	Shallow reservoir case: 40 °C
	Deep reservoir case: 90 °C
Permeability ( <i>k</i> )	$5 \times 10^{-13} \text{ m}^2$
Capillary Pressure $(P_c)$	van Genuchten <sup>1,2</sup>
	$\lambda = 0.63, S_{lr} = 0.19, \alpha = 1.5 \text{ x } 10^{-3} \text{ Pa}^{-1}, P_{max} = 1 \text{ x } 10^{5} \text{ Pa},$
$\mathbf{P}_{1}$	$S_{ls} = 1.$
Relative permeability $(k_r)$	Non-nysteretic liquid: van Genuchten '; Gas:
	Corey
	$\lambda = 0.63, S_{lr} = 0.21, S_{gr} = 0.20 \text{ or } 0.01$
	Hysteretic liquid and gas relative permeability <sup>4</sup>
	$\lambda = 0.63, S_{lr} = 0.21, S_{grmax} = 0.20, \gamma = 2, \lambda_{gas} =$
	0.5
Molec. diffusivity	Liquid: $10^{-10} \text{ m}^2 \text{ s}^{-1}$
coefficients $(d_{\beta}^{\kappa})^{1}$	Gas: $10^{-5} \text{ m}^2 \text{ s}^{-1}$
	$\theta = 0.0, P_0 = 10^5 \text{ Pa}$
Tortuosity ( $\tau_0$ )	0.25
Saturation-dependent	Equal to relative permeability
tortuosity ( $\tau_{\beta}$ )	

Table 1. Properties of the idealized one-dimensional radial model reservoir.

<sup>1</sup>Pruess et al. (1999)

<sup>2</sup> $\lambda$  is *m* in van Genuchten (1980) <sup>3</sup>Corey (1954)

<sup>4</sup>Doughty (2007, 2009)

### Results

Pressure (*P*), liquid saturation (*S*<sub>*l*</sub>), mass fraction of CO<sub>2</sub> in the gas ( $X_g^{CO2}$ ), gas density ( $\rho_g$ ), and mass fraction of CH<sub>4</sub> in the gas ( $X_g^{CH4}$ ) for the three cases considered are shown in Figures 5–7 for a shallow reservoir ( $P = 1 \ge 10^7$  Pa, T = 40 °C) and for a deep reservoir ( $P = 2 \ge 10^7$  Pa, T = 90 °C). Results from both hysteretic and non-hysteretic relative permeability curves (with two different values of residual gas saturation) are presented for *P* and *S*<sub>*l*</sub>, the differences between which will be discussed first for each of the three cases.

In particular, for the case of no residual gas (Figure 5) the pressure curves reveal differences for hysteretic and non-hysteretic results. Pressure rise for the case of non-hysteretic  $S_{gr} = 0.20$  is larger than for the cases of non-hysteretic  $S_{gr} = 0.01$  and hysteretic relative permeability. The reason for the difference is that non-hysteretic  $S_{gr} = 0.20$  requires the gas saturation in the injection zone to exceed 0.20 before CO<sub>2</sub> becomes mobile, hence the pressure builds up more before flow occurs. For non-hysteretic  $S_{gr} = 0.01$ , CO<sub>2</sub> with gas saturation exceeding 0.01 is mobile, thus creating less

pressure build up. The hysteretic case involves gas relative permeability following the drainage branch as  $CO_2$  is injected, which matches closely the non-hysteretic  $S_{gr} = 0.01$  case. For locations beyond the gas front, pressures are nearly equal for non-hysteretic and hysteretic cases because there is no gas saturation to cause phase interference. Note in this first scenario the  $CO_2$  plume extends to r = 140 m after two years.

Figures 6 and 7 show results for the cases of  $CO_2$  injection into the system with gas at residual saturation. The results for pressure show large differences between hysteretic and non-hysteretic curves and for the two different values of  $S_{gr}$ . The small pressure increase obtained for nonhysteretic curves with  $S_{gr} = 0.01$  arises because gas is mobile at nearly all saturations. In contrast, for  $S_{gr} = 0.20$ , the gas relative permeability near the well and the liquid relative permeability beyond the injected gas plume are lower than in the  $S_{gr} = 0.01$  case causing greater pressure increase. The hysteretic pressure profile is intermediate between these extremes as it models drainage conditions near the well and imbibition conditions beyond the injected plume front. The hysteretic liquid saturation profile is also intermediate between the non-hysteretic profiles over most of the plume. At low liquid saturations (small r), the hysteretic profile follows the non-hysteretic profile with  $S_{gr}$  = 0.01. Starting at  $S_l = 0.6$  it then parallels the non-hysteretic profile with  $S_{gr} = 0.20$ . This behavior reflects the gas-phase relative permeability curve shown in Figure 3d, in which the dashed line used to represent the second-order drainage branch coincides with the primary drainage branch for  $S_l \leq$ 0.6, but more closely resembles the first-order imbibition branch for  $S_l > 0.6$ . Comparing Figure 3d to Figure 3b indicates that the primary drainage and first-order imbibition branches are identical to the non-hysteretic gas relative permeability curves for  $S_{gr} = 0.01$  and  $S_{gr} = 0.20$ , respectively.

Considering now just the hysteretic results of Figures 5 and 6, we see the first-order observation that the pressure increase at the well is larger for the case with residual  $CO_2$  gas (Figure 6) than for the case of fully saturated conditions with zero residual gas (Figure 5). This occurs because of the decreased mobility of brine in the cases where residual gas is present. By this mechanism, rather than enhancing injectivity as might be expected by the presence of gas at the outset, residual gas inhibits injectivity by limiting the mobility of the brine that must be displaced in order for injection to occur. The second observation from comparison of results of Figures 5 and 6 is that the plume radius as defined by the region in which gas is mobile is larger in Figure 6 than in Figure 5. Specifically, in the case of zero residual gas (Figure 5), the injected  $CO_2$  plume extends approximately 140 m as defined by the gas saturation. For the case with residual  $CO_2$  gas (Figures 6), the gas saturation front extends about 230 m. Apparently, the injected gas is augmented by the initial residual gas which is incorporated into the mobile gas plume. Some of this increase in plume extent occurs because a small fraction of the injected  $CO_2$  dissolves when no  $CO_2$  is initially present. Without this dissolution, simulation results (not shown) indicate that the plume in the no-residual-gas case would extend to about 160 m.

We now consider the most interesting case, that of supercritical  $CO_2$  injection into a system with residual  $CH_4$  gas. Comparing only the hysteretic results, we observe the pressure rise is slightly lower for the case of supercritical  $CO_2$  injection into the system with residual  $CH_4$  (Figure 7) than into the system with residual  $CO_2$  (Figure 6). The reason for this appears to be the lower viscosity of  $CH_4$  gas relative to supercritical  $CO_2$ . Apparently the displacement of  $CH_4$  ahead of the  $CO_2$ injection front requires less pressure than does the displacement of  $CO_2$ . The drastic reduction in gas-mixture density as supercritical  $CO_2$  mixes with  $CH_4$  reveals itself as a third-order effect in the liquid saturation curve which shows the larger gas plume size (r = 260 m for the shallow reservoir and r = 240 m for the deep reservoir) (Figure 7) relative to the gas plume size (r = 230 m) in Figure 6. The effect is stronger for the lower pressure and temperature case than for the higher pressure and temperature case. This large decrease in density causes a volume expansion that is also reflected in the shape of the liquid saturation curve. In contrast to the smooth increase in liquid saturation with radial distance shown in Figure 6, the liquid saturation curve in Figure 7 shows two distinct regions: one from r = 0 to r = 140 m (where gas is predominantly CO<sub>2</sub>) and the other from r = 140 m to r =260 m (where gas is predominantly CH<sub>4</sub>) with a notable break in slope at the transition from CO<sub>2</sub> to CH<sub>4</sub>. Note as an aside that the local maximum in liquid saturation (minimum in gas saturation) near the plume front for the non-hysteretic case ( $S_{gr} = 0.01$ ) is an artifact of the non-hysteretic relative permeability model and does not occur for the hysteretic case. In short, the simulations show that supercritical CO<sub>2</sub> mixes with residual CH<sub>4</sub> causing the total volume of gas to increase, which creates a larger gas plume relative to the residual-CO<sub>2</sub> gas case.

For the conditions considered here, the radial extent of the plume increases by almost a factor of two when residual CH<sub>4</sub> is present (260 m, Figure 7) compared to injection into a formation with no residual gas (140 m, Figure 5). This finding raises the question of how much, if at all, plume extent would increase for CO<sub>2</sub> injection into a liquid-saturated formation containing dissolved CH<sub>4</sub>, which is considered a likely scenario for CO<sub>2</sub> injection into the water leg of a depleted gas reservoir or a saline formation located in the vicinity of a petroleum resource. Figure 8 shows simulation results for a case in which the formation initially contains dissolved CH<sub>4</sub> just below the solubility limit. Comparison of the saturation profiles for Figures 5 and 8 shows that they are identical right up to the leading edge, where the plume for the dissolved CH<sub>4</sub> case extends about 10 m farther, to r = 150 m. The mass fraction profile indicates that the plume is nearly pure CO<sub>2</sub> out to r = 140, then nearly pure CH<sub>4</sub> at the leading edge, indicating that as CO<sub>2</sub> moves out into the formation, creating a gas phase and partially dissolving into the liquid phase, the initially-dissolved CH<sub>4</sub> exsolves immediately, and is then pushed ahead of the growing CO<sub>2</sub> plume. Because the mass of CH<sub>4</sub> initially present for the dissolved case (Figure 8) is much smaller than for the residual gas case (Figure 7), the increase in radial extent of the plume is correspondingly much smaller.

In summary, the three main effects observed in the preceding simulations are (1) the reduction in injectivity caused by decreased brine mobility due to the presence of residual gas regardless of composition, (2) the larger extent of the gas plume caused by incorporation of dissolved or residual gas into the injected gas plume, and (3) the expansion of the gas plume caused by mixing of injected  $CO_2$  with residual CH<sub>4</sub>. Hysteretic relative permeability is needed to model properly all of the variables over the whole domain, although non-hysteretic models can be used for modeling parts of the system.



Figure 5. Results after two years of injection of  $CO_2$  into a formation saturated with brine, showing pressure, saturation, gas density (kg m<sup>-3</sup>), and mass fractions for the shallow (a) and deep (b) reservoirs. In the upper frames, where the dash-dot line is not visible, it coincides with the solid line.



Figure 6. Results after two years of injection of  $CO_2$  into a formation containing 20% residual gas consisting mainly of  $CO_2$ , showing pressure, saturation, gas density (kg m<sup>-3</sup>), and mass fractions for the shallow (a) and deep (b) reservoirs.



Figure 7. Results after two years of injection of  $CO_2$  into a formation containing 20% residual gas consisting mainly of  $CH_4$ , showing pressure, saturation, gas density (kg m<sup>-3</sup>), and mass fractions for the shallow (a) and deep (b) reservoirs.



Figure 8. Results after two years of injection of  $CO_2$  into a formation containing dissolved  $CH_4$  just below the solubility limit, showing pressure, saturation, gas density (kg m<sup>-3</sup>), and mass fractions for the shallow (a) and deep (b) reservoirs. In the upper frames, where the dash-dot line is not visible, it coincides with the solid line.

### Effects of Alternative Outer Radial Boundary Conditions

To examine the impact of the constant-pressure boundary approximately 1 km from the injection well, several cases with an infinite-acting reservoir were also modeled. For these cases, the grid extends 100 km, farther than any of the pressure responses propagate during the two-year injection period. Figure 9a shows results for a shallow formation containing dissolved  $CH_4$  (compare to Figure 8a) and Figure 9b shows results for a shallow formation containing 20% residual  $CH_4$  gas (compare to Figure 7a). The only noticeable differences are the pressure profiles, which are about 0.05 MPa higher for the infinite-acting reservoir. The resulting density increase is barely noticeable (as expected based on Figure 2a), and the saturation and mass fraction profiles are essentially unchanged. Simulation results for other initial conditions are comparable. Thus we can conclude that as far as plume development goes, the open reservoir model is equivalent to an infinite reservoir model.

A more extreme change in the outer radial boundary condition is to consider a closed reservoir. This is easily modeled by replacing the constant-pressure grid block at r = 1074 m in the original grid (Figure 4) with a normal grid block. Figure 10 shows results for a closed shallow reservoir after two years for the cases with dissolved CH<sub>4</sub> initially present (Figure 10a; compare to Figures 8a and 9a) and 20% residual CH<sub>4</sub> initially present (Figure 10b; compare to Figures 7a and 9b). For the closed right-hand side boundary condition, there is a large pressure increase accompanying CO<sub>2</sub> injection when no gas is initially present (Figure 10a) because liquid compressibility is small, enabling the pressure response to propagate rapidly to the closed boundary. The large pressure increase (about 7.6 MPa, compared to less than 0.1 MPa for the open and infinite-acting cases) causes CO<sub>2</sub> density to increase dramatically (from 570 kg/m<sup>3</sup> to 795 kg/m<sup>3</sup>), resulting in a more compact CO<sub>2</sub> plume (a plume extent of r = 120 m at two years rather than the r = 150 m previously obtained). In contrast, when residual gas is initially present, fluid compressibility is significantly larger, and the closed boundary has only a small effect (Figure 10b), with pressure increasing about 0.15 MPa more than for the open case and 0.10 MPa more than for the infinite-acting case. This completely closed reservoir is probably not a realistic choice for actual CO<sub>2</sub> storage, but it serves as a limiting case for the more promising "semi-closed" reservoirs, which are bounded radially, but allow significant pressure release through under- and overlying shale layers (Zhou et al., 2009).



Figure 9. Results after two years of injection of  $CO_2$  into an infinite-acting shallow reservoir containing (a) dissolved  $CH_4$  and (b) 20% residual gas consisting mainly of  $CH_4$ , showing pressure, saturation, gas density (kg m<sup>-3</sup>), and mass fractions. In the upper frames, where the dash-dot line is not visible, it coincides with the solid line.



Figure 10. Results after two years of injection of  $CO_2$  into a closed shallow reservoir containing (a) dissolved  $CH_4$  and (b) 20% residual gas consisting mainly of  $CH_4$ , showing pressure, saturation, gas density (kg m<sup>-3</sup>), and mass fractions. In the upper frames, where the dash-dot line is not visible, it coincides with the solid line. Note the two different pressure scales for (a) and (b); initial reservoir pressure is 10 MPa for both.

### Displaying Results as Transients for Observation Well Locations

Displaying simulation results as a function of radial distance for a given time provides an excellent way to visualize the multi-phase flow and transport processes accompanying  $CO_2$  injection, but it is not necessarily the optimal approach for designing field experiments or interpreting field data. Rather than having a complete picture of pressure, saturation, and gas content throughout the subsurface at a given time, we typically have a transient record from a limited number of observation wells. Figure 11 shows the simulation results for four years of  $CO_2$  injection into open and closed shallow reservoirs for initial conditions of brine nearly saturated with dissolved  $CH_4$ , as the transient response at three observation wells located 50, 100, and 150 m away from the injection well. Figure 12 shows the same results for initial conditions of brine containing 20% residual  $CH_4$ .

The open-reservoir case satisfies the conditions required for a similarity-variable  $r^2/t$  to exist: onedimensional radial geometry, essentially infinite radial extent, spatially uniform material properties and initial conditions, and steady flow rate at the injection well. Thus, the transient responses at each observation well can be collapsed to a single curve by dividing each time by the  $r^2$  value of that well. Such an exercise can provide useful insight into grid discretization errors, and in fact was used for that purpose here, to confirm that the 2-m grid spacing within the domain of the CO<sub>2</sub> plume is adequate. However, plotting versus  $t/r^2$  does not illustrate the actual timing of CH<sub>4</sub> and CO<sub>2</sub> arrival as well as plotting versus t does. For example, Figure 11a illustrates clearly that as observation well distance increases, the duration of the CH<sub>4</sub> pulse at the leading edge of the CO<sub>2</sub> plume increases.

The difference in pressure response for the open and closed reservoirs is strikingly displayed in the transient plots. When plotted on a scale big enough to show the entire range of pressures reached for the closed reservoir, as is done in Figure 11, the responses for the different observation wells collapse to a single curve. That is, the pressure over the entire model is roughly uniform. For the open reservoir (Figure 11a), this uniform pressure is essentially constant, so the gas density is purely a function of gas-phase composition. In contrast, for the closed reservoir (Figure 11b), as pressure increases,  $CO_2$  density increases sharply and  $CH_4$  density increases modestly (consistent with Figure 2a), and this variation is superposed on the gas-phase composition dependence as  $CH_4$  and  $CO_2$  arrive at each observation well. The overall density increase makes the plume much more compact, greatly lengthening the arrival time at more distant observation wells.

When residual gas is present (Figure 12), the difference between open and closed reservoir cases is much smaller, as the overall pressure increase for the closed reservoir case is small. Here the transient response highlights how long a time delay there is between the arrival of the plume (as manifested by a liquid saturation decrease) and the arrival of  $CO_2$  at an observation well, and how this delay grows significantly as observation well distance from the injection well increases. In contrast to the dissolved  $CH_4$  initial condition, where the  $CH_4$  pulse is rather narrow (Figure 11), when residual  $CH_4$  is initially present,  $CO_2$  arrival may not occur for a considerable time after the initial decrease in  $S_l$  (Figure 12). For example, for the 100-m observation well, with dissolved  $CH_4$  initially present,  $CO_2$  arrival more than 200 days between initial  $S_l$  decline and  $CO_2$  arrival.



Figure 11. Transient response at 50 m (solid), 100 m (long dash), and 150 m (short dash) away from the injection well for CO<sub>2</sub> injection into (a) open and (b) closed shallow reservoirs containing brine with dissolved CH<sub>4</sub>. When no gas phase exists ( $S_l = 1$ ), gas mass fractions are shown as zero and  $\rho_g$  shows the density a gas phase would have if it were present.



Figure 12. Transient response at 50 m (solid), 100 m (long dash), and 150 m (short dash) away from the injection well for  $CO_2$  injection into (a) open and (b) closed shallow reservoirs containing brine with 20% residual CH<sub>4</sub>.

Figure 13 summarizes all the effects of initial and boundary conditions on plume development by plotting plume arrival time at each observation well for all cases modeled. For an open or infiniteacting reservoir, the arrival time of the plume at a given observation well decreases sharply as one goes from a brine-saturated formation, to a formation containing dissolved gas, to a formation containing residual gas. For example, for an observation well located 100 m from the injection well, the plume would be expected 372 days after injection commenced for a brine-saturated formation, at 303 days for a formation with dissolved  $CH_4$ , and at 98 days for a formation with residual  $CH_4$ . For a closed reservoir, arrival times for cases with no residual gas initially present are significantly longer, due to the density increase accompanying pressure increase. The effect of the composition of the gas initially present is relatively small. For dissolved gas, arrival time is slightly earlier when  $CO_2$  is initially present because none of the injected  $CO_2$  can dissolve (when  $CH_4$  is initially present, some CO<sub>2</sub> dissolves into the aqueous phase and some CH<sub>4</sub> exsolves, but the net effect is more dissolution and a later plume arrival). For residual gas, arrival time is slightly later when CO<sub>2</sub> is initially present, because the greater density of CO<sub>2</sub> results in a more compact plume.



Effect of Initial and Boundary Conditions on Plume Arrival

Figure 13. Plume arrival time at observation wells located 50, 100, and 150 m from the injection well, for various initial and boundary conditions. Symbol color identifies the composition of the gas initially present.

### Summary, Conclusions, and Future Work

Applications of TOUGH2/EOS7C with and without hysteretic relative permeability were carried out to investigate the effects of residual gas on  $CO_2$  injection in depleted gas reservoirs. Results suggest that residual gas reduces injectivity by reducing the mobility of the brine that has to be displaced. A secondary effect is the tendency for the gas plume to extend farther from the injection well as residual gas is incorporated into the mobile gas plume. The extent is greater when the residual gas is  $CH_4$ , with its much lower density than supercritical  $CO_2$ , which strongly decreases plume density, resulting in a larger plume volume. When a closed reservoir is considered, pressure can increase significantly during injection, which in turn increases  $CO_2$  density, creating a plume that grows more slowly as time goes on.

The maximum increase in radial extent of the injected gas plume is about a factor of two in the onedimensional system investigated here, and could be somewhat larger in a two-dimensional (r,z)system where buoyancy flow can be included, since the lower gas density caused by the mobilized CH<sub>4</sub> will make the plume more buoyant. While a factor of two increase in plume extent is certainly non-trivial, it is relatively modest compared to the potential effects of other poorly known system parameters, such as vertical permeability, which controls buoyancy flow and gravity override, and horizontal permeability anisotropy (e.g., due to a contribution from fracture permeability), which controls the direction of plume movement, and other permeability heterogeneity that cause flow channelization.

This preliminary study has motivated future work in several areas. First, using a two-dimensional rz model geometry will enable the interplay of buoyancy flow and initial conditions to be examined, and using a three-dimensional model will allow treatment of dipping formations and incorporation of heterogeneous hydrologic properties, which could promote preferential flow. Second, the impact of varying the capillary pressure and relative permeability parameters should be investigated to address different formation types and associated parameter uncertainty. Laboratory data (e.g., Bachu and Bennion, 2008) suggest that parameters such as residual phase saturations and end-point relative permeability vary over a large range among different rock samples, and this variation is expected to strongly impact both the pressure and saturation responses to CO<sub>2</sub> injection. Finally, a more rigorous treatment of hysteresis is necessary, in which the higher-order scanning curves (2<sup>nd</sup>-order drainage and 3<sup>rd</sup>-order imbibition) already incorporated in the capillary pressure function are added to the gas- and liquid-relative permeability formulations.

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