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Dark Side of the Neutron?*

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Abstract. We discuss our recently proposed interpretation of the discrepancy between the bottle and beam neutron lifetime experiments as a sign of a dark sector. The difference between the outcomes of the two types of measurements is explained by the existence of a neutron dark decay channel with a branching fraction 1%. Phenomenologically consistent particle physics models for the neutron dark decay can be constructed and they involve a strongly self-interacting dark sector. We elaborate on the theoretical developments around this idea and describe the efforts undertaken to verify it experimentally.

1 Neutron Lifetime Discrepancy

Although the neutron has been known for almost a century, the latest experimental results suggest that it may still be hiding a deep secret. In the currently established framework of particle physics, the Standard Model, the neutron decays almost exclusively through beta decays, involving

$$n \rightarrow p + e^- + \bar{\nu}_e \quad (1)$$

and radiative corrections to this process. A calculation of the neutron lifetime in the Standard Model yields [2]

$$\tau_n^{\text{SM}} = \frac{4908.7(1.9) \text{ s}}{|V_{ud}|^2(1 + 3g_A^2)}, \quad (2)$$

where g_A is the axial-vector coefficient in beta decay, i.e., $\mathcal{M} = \frac{1}{\sqrt{2}} G_F V_{ud} g_V [\bar{p} \gamma_\mu n - g_A \bar{p} \gamma_5 \gamma_\mu n] [\bar{e} \gamma^\mu (1 - \gamma_5) \nu]$. By using the average values of V_{ud} and g_A extracted from experiments and adopted by the Particle Data Group (PDG) [3], one arrives at the neutron lifetime in the range $875.3 \text{ s} < \tau_n < 891.2 \text{ s}$ within 3σ . In turn, a recent lattice QCD calculation of g_A [4, 5] gave $\tau_n = 885 \pm 15 \text{ s}$.

There are two qualitatively different approaches to measuring the neutron lifetime: the bottle experiments and the beam experiments.

The bottle method relies on trapping neutrons in a container and counting them at several points in time. The decaying exponential

$$N_n(t) = N_n(0) \exp(-t/\tau_n) \quad (3)$$

is then fit to the data points $N_n(t)$, and τ_n^{bottle} is read off. Such measurement yields the total neutron lifetime and is independent of the actual decay channels. The average

bottle result quoted by the PDG and based on five experiments [6–10] is

$$\tau_n^{\text{bottle}} = 879.6 \pm 0.6 \text{ s}. \quad (4)$$

The two most recent bottle experiments [11, 12] provided values for τ_n within 2σ of this average.

A different approach has been implemented in beam experiments, where the neutron lifetime is determined by counting the protons (N_p) resulting from neutron decays. Estimating also the number of neutrons in the beam (N_n) that those protons originate from, τ_n^{beam} is given by

$$\tau_n^{\text{beam}} = -\frac{N_n}{dN_p/dt} = \frac{\tau_n}{\text{Br}(n \rightarrow p + \text{anything})}. \quad (5)$$

In the Standard Model $\text{Br}(n \rightarrow p + \text{anything}) = 100\%$, implying the two lifetimes are the same, $\tau_n^{\text{beam}} = \tau_n^{\text{bottle}}$. This equality no longer hold if other, beyond Standard Model neutron decay channels not involving a proton in the final state are allowed. In such a case the branching fraction $\text{Br}(n \rightarrow p + \text{anything}) < 100\%$ and, given Eq. (5),

$$\tau_n^{\text{beam}} > \tau_n^{\text{bottle}}. \quad (6)$$

The average based on two beam experiments [13, 14] (see also Ref. [15] for the original data used in Ref. [14]) and adopted by the PDG is

$$\tau_n^{\text{beam}} = 888.0 \pm 2.0 \text{ s}. \quad (7)$$

This represents a 4.0σ discrepancy with τ_n^{bottle} and hints that the inequality in Eq. (6) might actually hold [16].

The tension between the two types of experiments might arise from underestimated systematic errors, but it may also be an actual sign of new physics. We focus on the latter case. Assuming that the discrepancy between the experimental results originates from an incomplete understanding of the physics behind neutron decay, the results of the two types of experiments can be reconciled if

$$\text{Br}(n \rightarrow p + \text{anything}) \approx 99\%, \quad (8)$$

while the remaining 1% arises from *neutron dark decays*, involving at least one dark sector particle in the final state.

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2 Neutron Dark Decay

To investigate how such decays could have gone unnoticed in other experiments, let us consider a general scenario of a neutron decaying to a final state f with the sum of final state particle masses equal to M_f .

Of course, for the neutron to undergo a dark decay, M_f has to be smaller than the neutron mass, i.e., $M_f < m_n$. The lower bound on M_f is provided by experiments looking for neutron disappearance inside a nucleus. A neutron dark decay inside a nucleus (Z, A) could produce a daughter nucleus in an excited state $(Z, A - 1)^*$, leading to its subsequent de-excitation with the emission of secondary particles, e.g. gamma rays. A search for such signatures has been conducted by the SNO experiment [17] and the KamLAND experiment [18], placing a constraint of $\tau_{n \rightarrow \text{invisible}} > 5.8 \times 10^{29}$ years, adopted by the PDG as the bound on the neutron invisible decay channel.

However, if the condition $M_f > m_n - S_n$ is fulfilled, with S_n being the neutron separation energy in a given nucleus, then the decay $(Z, A) \rightarrow (Z, A - 1) + f$ is kinematically forbidden, while the neutron dark decay $n \rightarrow f$ is still allowed. Among all stable nuclei, the nucleus with the smallest neutron separation energy is ${}^9\text{Be}$, with $S_n({}^9\text{Be}) = 1.664$ MeV. Thus, the requirement of ${}^9\text{Be}$ stability enforces $M_f > m_n - 1.664$ MeV, which leads to the condition

$$937.900 \text{ MeV} < M_f < 939.565 \text{ MeV} . \quad (9)$$

Since $937.9 \text{ MeV} > m_p - m_e$, the requirement in Eq. (9) also assures that proton would not undergo a dark decay.

This opens the way to a whole new class of possible neutron decay channels:

$$n \rightarrow \chi \gamma, \quad n \rightarrow \chi \phi, \quad n \rightarrow \chi e^+ e^-, \quad \dots,$$

where χ is a dark fermion, ϕ is a dark scalar or a dark vector, and the ellipsis denotes other final states involving additional dark particles, photons and neutrinos. We now analyze the first two cases in more detail.

2.1 Neutron \rightarrow dark particle + photon

This simplest case involves only one dark fermion χ and a monochromatic photon in the final state. The allowed range of masses for χ , governed by Eq. (9), is

$$937.900 \text{ MeV} < m_\chi < 939.565 \text{ MeV} . \quad (10)$$

The energy of the corresponding monochromatic photon falls therefore within the range

$$0 < E_\gamma < 1.664 \text{ MeV} . \quad (11)$$

In the limit $m_\chi \rightarrow m_n$, the photon energy $E_\gamma \rightarrow 0$.

The dark fermion χ could be a dark matter particle, in which case its stability would require $m_\chi < m_p + m_e$, so that χ does not undergo beta decay through an off-shell neutron. In this dark matter case the allowed energy range for the photon reduces to $0.782 \text{ MeV} < E_\gamma < 1.664 \text{ MeV}$.

An effective Lagrangian for the decay $n \rightarrow \chi \gamma$ is

$$\begin{aligned} \mathcal{L}_1^{\text{eff}} = & \bar{n} (i\partial - m_n + \frac{g_n e}{2m_n} \sigma^{\mu\nu} F_{\mu\nu}) n \\ & + \bar{\chi} (i\partial - m_\chi) \chi + \varepsilon (\bar{n} \chi + \bar{\chi} n) , \end{aligned} \quad (12)$$

where g_n is the g -factor of the neutron and ε is a model-dependent parameter with mass dimension one that governs the mixing between χ and n . The Lagrangian in Eq. (12) gives a neutron dark decay rate of

$$\Delta\Gamma_{n \rightarrow \chi\gamma} = \frac{g_n^2 e^2}{8\pi} \left(1 - \frac{m_\chi^2}{m_n^2}\right)^3 \frac{m_n \varepsilon^2}{(m_n - m_\chi)^2} . \quad (13)$$

To explain the discrepancy between bottle and beam neutron lifetime experiments, $\Delta\Gamma_{n \rightarrow \chi\gamma} \approx \Gamma_n/100$, where Γ_n is the total neutron decay rate in the Standard Model. A phenomenologically viable particle physics model for the case $n \rightarrow \chi \gamma$ is discussed in Sec. 3.1 (Model 1).

2.2 Neutron \rightarrow two dark particles

A neutron dark decay with the final state consisting of only dark particles is realized by $n \rightarrow \tilde{\chi}^* \rightarrow \chi \phi$, where χ and $\tilde{\chi}$ are dark fermions and ϕ is a dark scalar (ϕ could also be a dark vector). In this case the requirement in Eq. (9) takes the form

$$937.900 \text{ MeV} < m_\chi + m_\phi < 939.565 \text{ MeV} . \quad (14)$$

Since this condition involves only the sum of the χ and ϕ masses, m_χ does not need to be close to m_n , e.g. a scenario where $m_\chi \approx m_\phi \approx m_n/2$ is allowed. However, nuclear stability requires that the mass of the intermediate $\tilde{\chi}$ satisfy

$$m_{\tilde{\chi}} > 937.9 \text{ MeV} \quad (15)$$

to prevent ${}^9\text{Be} \rightarrow {}^8\text{Be} + \tilde{\chi}$. If, in addition, $|m_\chi - m_\phi| < m_p + m_e$, then both χ and ϕ cannot undergo beta decays.

An effective Lagrangian describing $n \rightarrow \chi \phi$ is

$$\begin{aligned} \mathcal{L}_2^{\text{eff}} = & \mathcal{L}_1^{\text{eff}}(\chi \rightarrow \tilde{\chi}) + (\lambda_\phi \bar{\chi} \tilde{\chi} \phi + \text{h.c.}) \\ & + \bar{\chi} (i\partial - m_\chi) \chi + \partial_\mu \phi^* \partial^\mu \phi - m_\phi^2 |\phi|^2 , \end{aligned} \quad (16)$$

resulting in the neutron dark decay rate

$$\Delta\Gamma_{n \rightarrow \chi\phi} = \frac{|\lambda_\phi|^2}{16\pi} \sqrt{f(x, y)} \frac{m_n \varepsilon^2}{(m_n - m_\chi)^2} , \quad (17)$$

where $f(x, y) = [(1-x)^2 - y^2][(1+x)^2 - y^2]^3$, $x = m_\chi/m_n$ and $y = m_\phi/m_n$.

For $m_{\tilde{\chi}} > m_n$ the only available neutron dark decay channel is $n \rightarrow \chi \phi$ and $\Delta\Gamma_{n \rightarrow \chi\phi} \approx \Gamma_n/100$ is needed to explain the neutron lifetime discrepancy. In the case $m_{\tilde{\chi}} < m_n$, the decay channel $n \rightarrow \tilde{\chi} \gamma$ is also allowed. The ratio of the corresponding dark decay rates is

$$\frac{\Delta\Gamma_{n \rightarrow \tilde{\chi}\gamma}}{\Delta\Gamma_{n \rightarrow \chi\phi}} = \frac{2g_n^2 e^2 (1 - \tilde{x}^2)^3}{|\lambda_\phi|^2 \sqrt{f(x, y)}} , \quad (18)$$

where $\tilde{x} = m_{\tilde{\chi}}/m_n$. To account for the experimental discrepancy, $\Delta\Gamma_{n \rightarrow \chi\phi} + \Delta\Gamma_{n \rightarrow \tilde{\chi}\gamma} \approx \Gamma_n/100$. A viable model for the decay $n \rightarrow \chi \phi$ is provided in Sec. 3.2 (Model 2).

3 Particle Physics Models

We emphasize that our neutron dark decay proposal is very general and the models presented below serve only as an illustration of the simplest scenarios. Theories with a more complex dark sector remain to be explored and, as discussed in Sec. 4, already their minimal realizations can solve several outstanding problems in astrophysics.

3.1 Model 1 ($n \rightarrow \chi \gamma$)

The minimal model for the neutron dark decay requires only two new particles: a Standard Model singlet Dirac fermion χ and a scalar Φ , chosen to be an $SU(3)_c$ triplet, $SU(2)_L$ doublet and carrying hypercharge $Y = -1/3$. The Lagrangian of such a model is

$$\mathcal{L}_1 = \left[\lambda_q \epsilon^{ijk} \bar{u}_{Li}^c d_{Rj} \Phi_k + \lambda_\chi \Phi^{*i} \bar{\chi} d_{Ri} + \text{h.c.} \right] - M_\Phi^2 |\Phi|^2 - m_\chi \bar{\chi} \chi, \quad (19)$$

where u_L^c is the charge conjugate of u_R . Assigning $B_\chi = 1$ and $B_\Phi = -2/3$, the theory conserves baryon number. A diagram for $n \rightarrow \chi \gamma$ in this model is presented in Fig. 1.

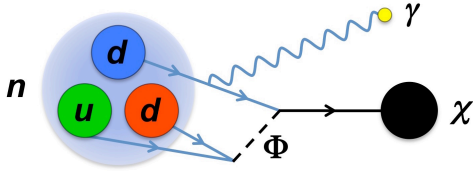


Figure 1. Neutron dark decay $n \rightarrow \chi \gamma$ in Model 1.

The neutron dark decay rate is obtained by matching the Lagrangian in Eq. (12) with that in Eq. (19). The result is given by Eq. (13) with $\varepsilon = \beta \lambda_q \lambda_\chi / M_\Phi^2$, where β is defined through $\langle 0 | \epsilon^{ijk} (\bar{u}_{Li}^c d_{Rj}) d_{Rk}^p | n \rangle = \beta (1 + \gamma_5)_\sigma^\rho u^\sigma$, with u being the neutron spinor. Lattice calculations give $\beta \approx 0.014 \text{ GeV}^3$ [19].

There is a large parameter space available for which $\Delta\Gamma_{n \rightarrow \chi \gamma} \approx \Gamma_n/100$. For example, if one takes the mass of χ to be at the lower end of the allowed range specified in Eq. (10), i.e., $m_\chi = 937.9 \text{ MeV}$, then the mass of Φ and the couplings in the model need to satisfy the relation

$$\frac{M_\Phi}{\sqrt{|\lambda_q \lambda_\chi|}} \approx 400 \text{ TeV}. \quad (20)$$

Therefore, Φ easily avoids all collider bounds provided that $M_\Phi \gtrsim 1 \text{ TeV}$. In addition, since χ is a Dirac fermion, it escapes the stringent constraints arising from neutron-antineutron oscillation [20] and dinucleon decay [21] searches.

3.2 Model 2 ($n \rightarrow \chi \phi$)

The entirely dark decay of the neutron, involving two dark particles in the final state, requires adding four fields to the Standard Model: the Dirac fermions χ and $\tilde{\chi}$, a scalar ϕ and the colored heavy scalar Φ introduced in the previous case. The Lagrangian of the model resembles the one for Model 1 with χ substituted by $\tilde{\chi}$ and an additional interaction term between $\tilde{\chi}$, χ and ϕ , i.e.,

$$\mathcal{L}_2 = \mathcal{L}_1(\chi \rightarrow \tilde{\chi}) + (\lambda_\phi \tilde{\chi} \chi \phi + \text{h.c.}) - m_\phi^2 |\phi|^2 - m_\chi \bar{\chi} \chi. \quad (21)$$

Baryon number is conserved upon assigning $B_{\tilde{\chi}} = B_\phi = 1$ and $B_\chi = 0$. The diagram for the neutron dark decay $n \rightarrow \chi \phi$ in this model is shown schematically in Fig. 2.

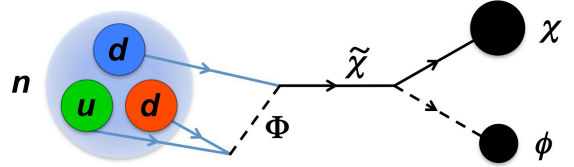


Figure 2. Neutron dark decay $n \rightarrow \chi \phi$ in Model 2.

After matching the Lagrangians in Eqs. (16) and (21), the rate for the neutron dark decay $n \rightarrow \chi \phi$ is given by Eq. (17) with $\varepsilon = \beta \lambda_q \lambda_{\tilde{\chi}} / M_\Phi^2$. The condition $\Delta\Gamma_{n \rightarrow \chi \phi} \approx \Gamma_n/100$, required to explain the neutron lifetime discrepancy when $m_{\tilde{\chi}} > m_n$, is satisfied for a wide range of parameters. In particular, adopting $m_\chi = 937.9 \text{ MeV}$, $m_\phi \approx 0$ and $m_{\tilde{\chi}} = 2m_n$, the mass of Φ and the couplings of the model have to satisfy

$$\frac{M_\Phi}{\sqrt{|\lambda_q \lambda_{\tilde{\chi}} \lambda_\phi|}} \approx 300 \text{ TeV}, \quad (22)$$

again consistent with collider, neutron-antineutron oscillation and dinucleon decay constraints.

In the case $m_{\tilde{\chi}} < m_n$, the additional neutron decay channel $n \rightarrow \tilde{\chi} \gamma$ is also available and the sum of the two rates should add up to $\approx \Gamma_n/100$, with their ratio governed by Eq. (18).

4 Theoretical Developments

Our work inspired several theoretical efforts to explore further implications of neutron dark decays. This involved studying the physics of neutron stars in the presence of the new neutron decay channels, constructing neutron dark decay models with a more complex dark sector including self-interactions, building models with dark decays of mesons, as well as inventing alternative, although related, ways to explain the neutron lifetime discrepancy. We discuss those theoretical ideas below.

4.1 Neutron star constraints

The impact of neutron dark decays on neutron stars was considered in Refs. [22–24]. The resulting production of dark particles changes the energy density and pressure inside a neutron star, modifying its equation of state. This in turn changes the predictions for the maximum allowed neutron star masses, since they are derived from integrating the Tolman-Oppenheimer-Volkoff equation that explicitly depends on the equation of state.

It was shown that the observed neutron star masses ($2M_\odot$ for the heaviest neutron stars discovered) are allowed if strong repulsive self-interactions are present in the dark sector of our models. Such interactions are easily introduced in the representative Models 1 and 2 discussed in Sec. 3 by simply adding a dark vector boson coupled strongly to the dark particle χ .

Interestingly, a strongly self-interacting dark sector lies along the lines of the self-interacting dark matter paradigm, which was introduced two decades ago [25] to solve the core-cusp and missing satellite problem of the Λ CDM model.

4.2 Models with a self-interacting dark sector

A model of this type was constructed in Ref. [26], where a neutron dark decay involving a dark fermion and a dark photon in the final state was considered, i.e., $n \rightarrow \chi A'$. The effective Lagrangian is

$$\begin{aligned} \mathcal{L}^{\text{eff}} = & \bar{n} (i\not{D} - m_n + \frac{g_n e}{2m_n} \sigma^{\mu\nu} F_{\mu\nu}) n \\ & + \bar{\chi} (i\not{D} - m_\chi) \chi + \varepsilon (\bar{n} \chi + \bar{\chi} n) \\ & - \frac{1}{4} F'_{\mu\nu} F'^{\mu\nu} - \frac{\delta}{2} F_{\mu\nu} F'^{\mu\nu} - \frac{1}{2} m_{A'}^2 A'_\mu A'^\mu, \end{aligned} \quad (23)$$

where the covariant derivative $D_\mu = \partial_\mu - i g' A'_\mu$. It was shown that the strength of the dark photon coupling to the dark particle χ , governed by the parameter g' and resulting in repulsive interactions between the χ particles, can be chosen such that the neutron lifetime discrepancy is explained and, at the same time, all astrophysical bounds are satisfied, including constraints from neutron stars, galaxy clusters, cosmic microwave background, Big Bang nucleosynthesis and supernovae. If the dark particle χ in this model is stable, it can contribute to the dark matter in the universe, but cannot account for all of the dark matter.

Many of the astrophysical constraints are alleviated if one assumes non-thermal dark matter production. This was shown in Ref. [27], where a model for the neutron dark decay $n \rightarrow \chi \phi$ was constructed, based on our Model 2, but with a dark boson introduced to mediate large self-interactions of χ . The Lagrangian for the dark sector is

$$\begin{aligned} \mathcal{L}_D = & g \bar{\chi} \not{Z}_D \chi + (\lambda_\phi \bar{\chi} \chi \phi + \text{h.c.}) \\ & - i g Z_D^\mu (\phi^* \partial_\mu \phi - \phi \partial_\mu \phi^*). \end{aligned} \quad (24)$$

There exists a choice of parameters for which this model satisfies neutron star constraints, remains consistent with all other astrophysical bounds and χ makes up all of the dark matter in the universe. In addition, due to the self-interactions of χ , the model is shown to solve the small-scale structure problems of the Λ CDM model.

4.3 Hadron dark decays

The idea of dark decays can be applied also to other neutral hadrons. In Ref. [28] it was argued that the mesons K_L^0 and B^0 can decay to dark sector particles at measurable rates. An explicit model was constructed with a dark sector consisting of several families of dark fermions. An analogous mechanism that prevents neutron beta decays in neutron stars, i.e., Pauli blocking, also forbids neutron dark decays inside a neutron star in this model.

4.4 Baryogenesis

It has recently been shown that the model addressing the neutron lifetime puzzle based on the Lagrangian in Eq. (23) provides a successful framework for low-scale baryogenesis [29]. In addition, a model very similar to our Model 2, with couplings of $\bar{\chi}$ to other quark flavors and a Majorana (instead of Dirac) fermion χ , has been proposed in the context of low-scale baryogenesis as well [30].

4.5 Related solutions

Taking into consideration only the experimental data for g_A from experiments performed after the year 2002, the bottle neutron lifetime is favored [31]. Based on this observation, explanations of the neutron lifetime discrepancy have been put forward in which it is the bottle lifetime that is equivalent to the Standard Model prediction for τ_n . The difference in outcomes of the bottle and beam measurements is explained via neutron-mirror neutron oscillations resonantly enhanced in large magnetic fields thus affecting only beam measurements [32], or by invoking a sizable Fierz interference term canceling the dark decay contribution to the neutron decay rate [33].

5 Experimental Searches

Several experimental efforts have been undertaken directly after our results were announced, searching specifically for the signatures we proposed.

5.1 Neutron \rightarrow dark matter + photon

Within the first few weeks after our results became public, a dedicated experiment was performed at the Los Alamos UCN facility looking for the monochromatic photon in the neutron dark decay $n \rightarrow \chi \gamma$ [34]. The search was sensitive to final state photons with energies $0.782 \text{ MeV} < E_\gamma < 1.664 \text{ MeV}$ and challenged the case $\text{Br}(n \rightarrow \chi \gamma) \approx 1\%$ at a significance level of 2.2σ . The remaining photon energy range, i.e., $E_\gamma < 0.782 \text{ MeV}$, is left to be explored.

5.2 Neutron \rightarrow dark particle + e^+e^-

Another dedicated experiment, also performed at the Los Alamos UCN facility, looked for e^+e^- pairs from the neutron dark decay $n \rightarrow \chi e^+e^-$ [35]. This search excluded the case with $\text{Br}(n \rightarrow \chi e^+e^-) \approx 1\%$ for the electron-positron energy range $E_{e^+e^-} \gtrsim 2m_e + 100 \text{ keV}$ with a confidence of nearly 100%. The remaining 100 keV energy window was beyond experimental sensitivity.

5.3 Nuclear dark decays

There exists a number of unstable nuclei for which the neutron separation energy is smaller than for ${}^9\text{Be}$, i.e., $S_n < 1.664$ MeV. Those include ${}^7\text{H}$, ${}^{11}\text{Li}$, ${}^{11}\text{Be}$, ${}^{13}\text{Li}$, ${}^{14}\text{B}$, ${}^{15}\text{C}$, ${}^{16}\text{Be}$, ${}^{17}\text{B}$, ${}^{17}\text{C}$, ${}^{19}\text{B}$, ${}^{19}\text{C}$, ${}^{22}\text{C}$, ${}^{22}\text{N}$, as well as heavier ones. For these particular nuclei a neutron dark decay can lead to nuclear dark decays if the final state dark particle mass m_χ falls within the range

$$937.9 \text{ MeV} < m_\chi < m_n - S_n. \quad (25)$$

We proposed to search for such nuclear dark decays in our original paper [1], focusing on the corresponding signatures for ${}^{11}\text{Li}$, for which $S_n({}^{11}\text{Li}) = 0.396$ MeV. In that case the decay chain ${}^{11}\text{Li} \rightarrow {}^{10}\text{Li} + \chi \rightarrow {}^9\text{Li} + n + \chi$ is allowed and the ${}^9\text{Li}$ long lifetime could be used to discriminate against the background from ${}^{11}\text{Li}$ beta decays. However, ${}^9\text{Li}$ can be produced also in beta-delayed deuteron emission [36, 37] and the distinction between this and the dark channel would be extremely difficult.

It was argued in Ref. [38] that, from an experimental point of view, there is a much better candidate: ${}^{11}\text{Be}$, for which $S_n({}^{11}\text{Be}) = 0.502$ MeV. It was also suggested that the presence of an unexpectedly high number of ${}^{10}\text{Be}$ in ${}^{11}\text{Be}$ decays described in Ref. [39] might in fact be a sign of the neutron dark decay $n \rightarrow \chi \phi$ like in our Model 2, leading to the nuclear dark decay

$${}^{11}\text{Be} \rightarrow {}^{10}\text{Be} + \tilde{\chi}^* \rightarrow {}^{10}\text{Be} + \chi + \phi, \quad (26)$$

and not necessarily, as initially conjectured, due to an enhanced βp channel resulting from an unknown resonance.

In addition, it was shown in Ref. [40] that the nuclear dark decay in Eq. (26) is consistent with the observed Standard Model decay rates of ${}^{11}\text{Be}$ as long as $m_{\tilde{\chi}} > m_n - S_n({}^{11}\text{Be})$, i.e.,

$$m_{\tilde{\chi}} > 939.064 \text{ MeV}. \quad (27)$$

This condition is obviously satisfied in the model with a self-interacting dark sector of Ref. [27], where the $\tilde{\chi}$ mass was chosen to be $m_{\tilde{\chi}} = 800$ GeV.

Very recently, an experiment at the CERN-ISOLDE laboratory was performed [38, 41] with the goal of determining whether the final state of ${}^{11}\text{Be}$ decays contains protons in the final state or not. The results have not yet been published.

5.4 Ongoing beam measurements

There are currently two operating beam experiments measuring the neutron lifetime, the first one at the National Institute of Standards and Technology (NIST) [42, 43] and the second one at the Japan Proton Accelerator Research Complex (J-PARC) [44, 45]. If those experiments provide results consistent with the current beam average, the tension between bottle and beam measurements will increase, supporting the viability of models presented here.

5.5 Expanding the scope of bottle experiments

Perhaps the most straightforward, although technically challenging way to tackle the neutron lifetime puzzle would be to modify the existing bottle experimental setup. Including a proton detection system in bottle experiments would enable measuring the branching fraction $\text{Br}(n \rightarrow p + \text{anything})$ independently of the beam experiment. Such modification would enable a direct test of the premise that the difference of outcomes between the bottle and beam measurements is due to neutron decays that do not produce a proton, without any dependence on the specific model realization of the non-proton final state.

6 Final Remarks

Given the theoretical and experimental developments related to our proposal, Model 2 with a self-interacting dark sector seems like a very promising candidate theory for explaining the neutron lifetime discrepancy. This model is not only consistent with all current experimental constraints, but it is also interesting from a theoretical perspective, with its solution to the small-scale structure problem and perhaps a novel mechanism for baryogenesis.

Even if the neutron lifetime puzzle gets resolved by future higher precision bottle and beam measurements, dark decays of the neutron at a smaller rate will still be allowed and certainly interesting to consider. It would be incredible if the good old neutron became the key to unraveling the mystery of the dark side of our universe.

Acknowledgements

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