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# Topological surface states of semimetal TaSb<sub>2</sub>

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# Abstract

Topological surface states, protected by the global symmetry of the materials, are the keys to understanding various novel electrical, magnetic, and optical properties. TaSb<sub>2</sub> is a newly discovered topological material with unique transport phenomena, including negative magnetoresistance and resistivity plateau, whose microscopic understanding is yet to be reached. In this study, we investigate the electronic band structure of TaSb<sub>2</sub> using angle-resolved photoemission spectroscopy and density functional theory. Our analyses reveal distinct bulk and surface states in TaSb<sub>2</sub>, providing direct evidence of its topological nature. Notably, surface states predominate the electronic contribution near the Fermi level, while bulk bands are mostly located at higher binding energies. Our study underlines the importance of systematic investigations into the electronic structures of topological materials, offering insights into their fundamental properties and potential applications in future technologies.

Keywords Topological materials, Electronic structures, TaSb<sub>2</sub>, Resistivity plateau, Extremely large magnetoresistance

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# **1** Introduction

The topological classifications of electronic structure have led to a significant discovery of topology-driven phenomena in condensed matter physics. In many of these topological materials, the non-trivial topology is manifested with the surface states, such as linear band crossings protected by time-reversal symmetry in topological insulators and Fermi arcs in Weyl semimetals, giving rise to the gapless metallic states. These properties are determined by bulk wavefunctions over the whole Brillouin zone with discrete topological indices.

Various experimental probes have been used to detect nontrivial band topologies. Angle-resolved photoemission spectroscopy (ARPES) offers a direct method for detecting surface states and their dependence on surface termination [1-6]. Furthermore, transport measurements can be employed to identify topological surface states, with characteristic features such as low-temperature resistivity plateaus and negative magnetoresistance [7-15]. The angular dependence of the quantum oscillation measurements is also used to detect surface states, providing insights into the dimensionality of the states around the Fermi energy [9, 11, 15, 16].

In the case that both bulk and topological surface states exist around the Fermi energy, a multifaceted analysis is required to unambiguously determine the topology of the band structures since both states contribute to the response from the external perturbation. In this case, the ARPES measurements combined with first-principles calculations could be an effective strategy in which the theoretical band structures can be used as a guide to distinguish the surface and bulk electronic structures.

Recently discovered TaSb<sub>2</sub> exhibits intriguing transport properties, including resistivity plateaus, negative magnetoresistance (MR), extremely large MR (XMR), and non-trivial Berry phase in Shubnikov-de-Hass (SdH) oscillations, indicating the bulk states non-zero Berry curvatures. It is argued that the unusual magnetoresistance is due to the small electron and hole pockets or from magnetic field-induced Weyl points [9, 16]. The firstprinciples density functional theory (DFT) calculations show the nodal lines, which are gapped in the presence of the spin-orbit coupling (SOC). The calculated topological indices predict the weak topological insulating phase where there are bulk electron and hole pockets crossing the Fermi energy [17], resulting in the topological semimetallic phase. However, there has been no direct observation of the surface states that confirms the proposed topological insulating phase where the topological surface states give rise to the low-temperature resistivity plateaus. This calls for a systematic study of the electronic band structures of TaSb<sub>2</sub>.

In this paper, we investigate the electronic band structure of TaSb<sub>2</sub> using ARPES and DFT calculations. Our analyses reveal bulk and surface bands in TaSb<sub>2</sub>, providing direct evidence for the existence of the topological surface states. Particularly, most bands near the Fermi level are identified as surface states, while bulk bands are located at relatively higher binding energies (E- $E_F < -0.5$  eV, where  $E_F$  is Fermi energy), indicating clear topological properties of TaSb<sub>2</sub> that govern the unique transport phenomena such as the resistivity plateau and XMR. Our study delivers comprehensive investigations into the electronic structures of a complex topological quantum material, providing essential insights into the close correlation between the electronic band structure and transport properties for future technological applications [3, 18-23].

## 2 Results and discussions

## 2.1 Electron band structures of TaSb<sub>2</sub>

The atomic structure of bulk TaSb<sub>2</sub> is presented in Fig. 1a. It has a monoclinic unit cell with space group C2/m (No. 12), in which each Ta site is surrounded by eight Sb atoms. The calculated lattice constants are  $a_c = 10.354$  Å,  $b_c = 3.700$  Å, and  $c_c = 8.384$  Å in good agreement with experimental data (see SM for detail) [9]. The frontier orbitals around the Fermi energy  $(E_F)$  are Ta-d and Sb-p with substantial hybridization between them, as shown in the partial density of states (PDOS) in Fig. 1b. We find a V-shaped density of states around the Fermi energy with a small but finite value at the Fermi energy, consistent with the previous report and metallic transport behavior [9, 17]. The finite density of states at the Fermi energy consists of small electron and hole pockets with band crossings observed along high symmetry lines, as in Fig. 1d, in accordance with the previous DFT calculations [9, 17]. Upon the inclusion of the SOC, the band crossings are fully gapped, but there are electron and hole pockets crossing the Fermi energy, showing a semimetallic ground state.

Figure 2 shows the natural cleavage plane measured by ARPES. We find that the  $(1\overline{11})$  plane with respect to the primitive lattice vectors are parallel to the layered atomic planes obtained by disconnecting one of the eight Ta-Sb bonds that has the largest distance (3.01 Å). The corresponding surface Brillouin zone (BZ) is in an elongated hexagonal shape. Moreover, the surface BZ has matching periodicity compared with the surface bands measured by ARPES. Therefore, the  $(1\overline{11})$  plane can be considered as the cleavage plane, consistent with the distance between the Ta and Sb ions. We note that the identification of this unusual cleavage plane is challenging, as it requires a meticulous comparison between the periodicity of the surface band structures and those obtained by



**Fig. 1** Atomic configuration and electronic structures of TaSb<sub>2</sub>. **a** Perspective and top view of the atomic configurations. The lattice vectors with subscript *c* correspond to the conventional unit cell. **b** The Ta-*d* and Sb-*p* orbital-projected DOS, along with the total DOS. **c** High symmetry points in the Brillouin zone where **b**<sub>1</sub>, **b**<sub>2</sub>, and **b**<sub>3</sub> are the reciprocal lattice vectors of the primitive unit cell. **d** Band structures without spin–orbit coupling (SOC) along the high-symmetry lines shown with the red solid lines in (**c**). **e** Band structures with SOC

considering all possible cleavage planes. Our successful identification of the cleavage plane enables further analysis of the electronic band structure related to the topological nature.

## 2.2 Topological surface states of TaSb<sub>2</sub>

To investigate the electronic band structure of TaSb<sub>2</sub>, we performed ARPES measurements and DFT calculations on TaSb<sub>2</sub> single crystals. ARPES intensity plots of the constant energy contours were stacked with energy values ranging from 0, -0.2, -0.4, -0.6, and -0.8 eV, respectively (Fig. 3a). The bands at  $E-E_F=0$  eV were observed, confirming the metallic property of TaSb<sub>2</sub>, consistent with other transport results (Fig. 3b) [9, 16]. The electron-like pocket ( $\alpha$ ) at the X point disappears with emergence of a wave-shaped bands as binding energy increases (Fig. 3a). The ripple-shaped band ( $\zeta$ ) elongates along the  $\overline{\Gamma}$ - $\overline{Y}$  direction, exhibiting 1D-like chain structures across all stacked binding energies. This feature implies the presence of an open Fermi surface feature, suggesting that the open-orbit fermiology may be a contributing factor to the XMR [24, 25]. The experimental Fermi surface (Fig. 3b), exhibits electron ( $\alpha$ ) and hole ( $\beta$ )-like pockets, ripple-shaped features ( $\zeta$ ), and single dot points  $(\gamma)$ .

From the DFT calculations showing both the surface and bulk band structures, we obtained similar Fermi surface features (Fig. 3c) to the ARPES result (Fig. 3b): open

pockets at the  $\overline{X}$  and  $\overline{M}$  points, four closed pockets near the  $\overline{\Gamma}$  point, and a long ripple-shaped feature along the  $\overline{\Gamma}$ - $\overline{Y}$  high symmetry direction, which shows a guasi-1D like shape. The open electron pockets ( $\alpha$ ) at the X points from DFT results are centered at the  $\overline{X}$  points, whereas they are slightly closer to the  $\overline{\Gamma}$  points in ARPES results. With this exception noted, the overall band structures are in agreement with both DFT and ARPES results. We note that most of the bands near the Fermi energy in the SBZ are derived from the surface states, as the bulk states are mostly pushed into the higher binding energy about 20 meV below the  $E_F$  for this particular natural cleavage plane. These surface states are from the weak topological insulating phase of TaSb<sub>2</sub> consistent with the calculated  $Z_2$  classification of (0; 111), where the topological surface states emerge with the band inversion induced by gapped band crossing by SOC, consistent with the previous reports [17].

We further investigate the experimental electronic band structures of  $TaSb_2$  along various high symmetry directions, comparing them to the DFT calculations. In Fig. 4, calculated bulk band structures with SOC are presented in black solid lines. They exist mostly away from the Fermi level in higher binding energy in this particular cleavage plane (111). SOC separates the valence and conduction bands, leading to no band crossing points for the bulk bands, thereby classifying  $TaSb_2$  as a topological material with weak topological invariants, which



**Fig. 2** The cleaved atomic plane and surface Brillouin zone. **a** Primitive unit cell with  $(1\overline{11})$  plane shown as a red-colored area. **b** The surface Brillouin zone (BZ) corresponding to the  $(1\overline{11})$  plane, denoted as a red-colored area with newly defined high symmetry points,  $\overline{\Gamma}$ ,  $\overline{X}$ ,  $\overline{Y}$ , and  $\overline{M}$ . The red dotted arrows (**b**<sub>1s</sub> and **b**<sub>2s</sub>) are the reciprocal lattice vectors of the surface BZ

is consistent with the previous report [26]. All surface bands, except for the  $\overline{\Gamma}$ - $\overline{Y}$  direction, exhibit band crossings that create nodal lines in line with the previous report [17]. However, bulk bands with SOC gap out the nodal lines and separate the valence and conduction bands. Consequently, TaSb<sub>2</sub> exhibits nearly compensated semimetal behavior and universally possesses surface states near Fermi level, which leads to TaSb<sub>2</sub> as a topological insulator [17, 26].

# 2.3 Topological nature of TaSb<sub>2</sub>

Transition-metal dipnictides,  $TPn_2$  (T=Nb, Ta and Pn=Sb, As), exhibit unique topological characters. They are weak topological insulators in zero magnetic fields, but under external magnetic fields, can be categorized as Type-II Weyl materials [27]. Typically, it is suggested that the resistivity plateau observed in these materials

arises from the competition between the insulating bulk state and metallic surface states. The initial increase in conductivity with decreasing temperature is followed by a resistivity plateau where the conductance of the metallic surface state saturates the resistivity of the insulating bulk state. On the other hand, recent research suggests that classical magneto-resistance theories offer an alternative explanation for these resistivity plateaus observed in materials, without necessarily relying on topological surface states. This alternative explanation considers factors such as impurity scattering, field induced metalinsulator transition, electron-phonon interactions, and electron-electron interactions [11–16].

 $TaSb_2$  exhibits various interesting transport properties, including positive extreme magnetoresistance (XMR) and high mobility. In the low-temperature regime, it demonstrates both the negative MR and resistivity plateau when the applied field is parallel and perpendicular to the current, respectively [9].

There has been considerable interest in understanding the microscopic mechanisms of XMR and identifying novel XMR materials, such as transition-metal dipnictides (TmPn<sub>2</sub>) MoAs<sub>2</sub>, and W<sub>2</sub>As<sub>3</sub> [12, 13, 28–30]. Proposed mechanisms to explain XMR include nontrivial band topology, electron-hole compensation, open-orbit Fermi surface (FS) topology, and forbidden backscattering at zero field. In our analysis of the electronic band structure, we observed a distinct non-closing band feature in the FS (Fig. 3a, b), a characteristic consistently present in XMR semimetals of the TmPn<sub>2</sub> family with the C12/m1 space group. This finding suggests that openorbit fermiology, together with electron-hole compensation, may play a key role in the XMR behavior of TmPn<sub>2</sub> materials, including TaSb<sub>2</sub> [24, 25, 31].

The resistivity plateau observed in TaSb<sub>2</sub> at low temperatures is a result of a magnetic field-induced resistivity plateau with the broken time-reversal symmetry. This differs from the typical resistivity plateau in topological insulators with time-reversal symmetry, such as NbSb<sub>2</sub>, NbAs<sub>2</sub>, TaAs<sub>2</sub>, and WTe<sub>2</sub> [15, 28, 31–34]. These properties are attributed to field-induced metal-insulator transition or Kohler's rule. On the other hand, calculations of the bulk electronic band structure suggest that TaSb<sub>2</sub> has weak topological properties, potentially leading to the presence of surface states contributing to the observed resistivity plateau [17]. In addition, due to the presence of electron and hole pockets at the Fermi level, this resistivity plateau cannot be elucidated by Kohler's rule based on single scattering process. Given that most of the transport measurements are done along the [110] direction that is included in the  $(1\overline{11})$  cleavage plane, the observed numerous surface states near the Fermi surface need also be considered as a major source of the resistivity plateau



**Fig. 3** Electronic band structure of TaSb<sub>2</sub>. **a** Constant energy contours in the energy range from  $E - E_F = 0$  eV to -0.8 eV. Experimental (**b**) and calculated (**c**) Fermi surface (FS) with hv = 55 eV. Red lines denote the surface Brillouin zone (SBZ) with high-symmetry direction labeled. The band energies are shifted to the higher binding energy of about 20 meV for better comparison. **d** Calculated band structure with bulk (black solid line) and surface bands (blue to yellow color scale)

in addition to the contribution from the compensating bulk electron and hole pockets. We note that our theoretical calculation of another (001) cleavage plane containing the [110] direction show substantial surface states (See SM for details), supporting the robust surface state contributions in the transport along the [110] direction.

# **3** Conclusions

In conclusion, our investigation into the electronic band structure of TaSb<sub>2</sub> has provided valuable insights into its topological nature and unique transport phenomena. Through a combination of ARPES and DFT calculations, we have identified both bulk and surface bands in TaSb<sub>2</sub>, offering direct evidence of its topological properties. Our results imply that the presence of an open Fermi surface may be a shared characteristic in XMR materials with the C12/m1 space group, potentially working in synergy with electron-hole compensation to elucidate the origin of the XMR effect. In addition, a significant proportion of the bands near the Fermi level are identified as surface states, while bulk bands are situated at relatively higher binding energies. This observation underscores the clear topological properties of TaSb<sub>2</sub>, which can be a key factor in understanding unique transport phenomena such as resistivity plateau. While classical magneto-resistance theories offer alternative explanations for resistivity plateaus, our investigation stresses the critical role of topological surface states in shaping the transport properties of TaSb<sub>2</sub>, suggesting avenues for further exploration of its topological properties.

# 4 Methods

## 4.1 Single crystal growth

Single crystals of  $TaSb_2$  were synthesized using chemical vapor transport methods as described previously [9, 16, 35, 36].

# 4.2 ARPES measurement

ARPES measurements were performed at the HERS endstation of the Beamline 10.0.1, Advanced Light Source, Lawrence Berkeley National Laboratory. The ARPES system is equipped with a Scienta R4000 electron analyzer and has base pressure  $3 \times 10^{-11}$  Torr. The photon energy was set at 55 eV with energy and angular resolution of 25 meV and 0.1 degree. Measurements were made at 15 K.

## 4.3 First-principles calculations

First-principles DFT calculations were performed using the Vienna ab initio simulation package (VASP) [37, 38]. The generalized gradient approximation with Perdew–Burke–Ernzerhof parameterizations [39] was used for the exchange–correlation functional. The projector augmented wave method [40] was used with an energy cut-off of 500 eV. The  $\Gamma$ -point centered  $8 \times 8 \times 5$  k-point grid was used. Convergence was reached if the consecutive energy difference was less than  $10^{-6}$  eV. The



**Fig. 4** Experimental and theoretical electronic band structures of TaSb<sub>2</sub>. **a–f** ARPES intensity plots taken at hv = 55 eV (left), corresponding second-derivatives ARPES spectra for enhanced visibility (middle), and calculated band structure (right) along high symmetry directions,  $\overline{\Gamma} \cdot \overline{X}$  (**a**, **c**),  $\overline{Y} \cdot \overline{M}$  (**b**, **d**),  $\overline{\Gamma} \cdot \overline{Y}$  (**e**), and  $\overline{\Gamma} \cdot \overline{M}$  (**f**). In the calculated band structure (right panels), black solid lines represent bulk states, while yellow dispersions (blue to yellow color scale) denote surface bands

atomic structures were relaxed with a force threshold of 0.001 eV Å<sup>-1</sup>. For the calculation of the surface band structures and the  $Z_2$  index set, Wannier90 code [41] and WannierTools were used [42].

#### Abbreviations

/ issic futions	
ARPES	Angle-resolved photoemission spectroscopy
MR	Magnetoresistance
XMR	Extremely large magnetoresistance
SdH	Shubnikov-de-Hass
DFT	Density functional theory
SOC	Spin–orbit coupling
E <sub>F</sub>	Fermi energy
PDOS	Partial density of states
ΒZ	Brillouin zone
SBZ	Surface Brillouin zone
XMR	Extreme magnetoresistance
VSAP	Vienna ab initio simulation package

# **Supplementary Information**

The online version contains supplementary material available at https://doi.org/10.1186/s40580-024-00457-y.

Supplementary Material 1: Table S1. Comparison between the calculated and experimental lattice parameters. Figure S1.  $k_z$  dispersion of TaSb<sub>2</sub>. Figure S2. The cleaved planes parallel to the b<sub>c</sub> direction and corresponding surface BZs.

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## Author contributions

J.-E. L., S. Y. P., H. R., and S. -K. M. proposed and designed the research. Y. L. and C. P. performed single-crystal growth. J.-E. L., J. H., and H. R. carried out the ARPES measurements and analyzed the ARPES data with assistance from C. H. and S. -K. M.; S. Y. P. carried out the density functional calculations and provided theoretical support. J. -E. L., S. Y. P., H. R., and S. -K. M. wrote the manuscript and revised it with assistance from all other authors. All authors contributed to the scientific planning and discussions.

#### Data availability

The data generated during the current study are available from the corresponding authors upon reasonable request.

## Declarations

## **Competing interests**

The authors declare that they have no competing interests.

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